

Nuclear matter EoS with minimal length ($\pi(\ell)$ -model)

Minimal length \Rightarrow Komutasi non-local field operator harus bersifat

$$\delta(\vec{x} - \vec{y}) \rightarrow \tilde{\delta}(\vec{x} - \vec{y}) = \frac{1}{\pi(\vec{x} - \vec{y})} \sin \left[\frac{\pi(\vec{x} - \vec{y})}{2\beta} \right]$$

referensi utama:

- 1.) Kh. Nouicer, J. Phys. A: Math. Gen. 38 (2005)
- 2.) Frassino-Panella, PRD 85 (2012), 045030
- 3.) Panella, PRD 76, 095012 (2007)

(1) momentum modified to (to all orders in β):

$$p_i \rightarrow \bar{p}_i(p) = \frac{1}{\sqrt{\beta}} p \left[\tan^{-1}(\sqrt{\beta} p) \right] p_i \quad \dots (1)$$

(2) GUP (+) Dirac equations (in momentum rep):

$$\left(\gamma^\mu \bar{p}_\mu - m^+ \right) \psi(\vec{p}) = 0 \quad \dots (2)$$

(3) State "plane-wave" = maximized localized state:

$$\psi(\vec{x}, t) = A e^{-i(\omega_p t - \vec{p} \cdot \vec{x})} \quad \dots (3)$$

(4) dispersion relation

from (2) one will get:

$$(p_0)^2 - \vec{p}^2 - m^{+2} = 0$$

$$(p_0)^2 = \vec{p}^2 + m^{+2}$$

$$(p_0)^2 = \frac{1}{\beta} \left[\tan^{-1}(\sqrt{\beta} p) \right]^2 + m^{+2} \quad \dots (4)$$

①

(5) Normalisasi Condition $\sum_{\alpha=1}^{\infty} U^+(\vec{p}, \alpha) U(\vec{p}, \alpha)$ disinyal

$$\sum_{\alpha=1}^{\infty} U^+(\vec{p}, \alpha) U(\vec{p}, \alpha) = 1.$$

$$\sum_{\alpha=1}^{\infty} U(\vec{p}, \alpha) U(\vec{p}, \alpha) = \frac{2m}{m_p}$$

(6) Field operator (2nd quantization with GUP)

expand in maximally localized state

$$\psi(\vec{x}, t) = \sum_{\alpha} \int \frac{d^3 p}{(1 + \beta p^2)(2\pi)^3} N_p \left[b(\vec{p}, \alpha) e^{i(-\vec{p} \cdot \vec{x} + \omega_p t)} U(\vec{p}, \alpha) + b^{\dagger}(\vec{p}, \alpha) e^{-i(-\vec{p} \cdot \vec{x} + \omega_p t)} U(\vec{p}, \alpha) \right]$$

demanding the commutation (ref- Greiner & Ryder):

$$\{\psi_a(\vec{x}, t), \psi_b^+(\vec{y}, t)\} = \delta_{ab} \delta(\vec{x} - \vec{y})$$

With Commutation of annihilation & creation operator:

$$\boxed{\{b(\vec{p}, \alpha), b^+(\vec{q}, \beta)\} = (2\pi)^3 \delta_{\alpha\beta} \delta(\vec{p} - \vec{q}) = \{d(\vec{p}, \alpha), d^+(\vec{q}, \beta)\}}$$

left-hand side:

$$\begin{aligned} \{\psi_a(\vec{x}, t), \psi_b^+(\vec{y}, t)\} &= \sum_{\alpha\beta} \int \frac{d^3 p d^3 q N_p N_q}{(2\pi)^6 (1 + \beta p^2)(1 + \beta q^2)} \left[(2\pi)^3 \delta_{\alpha\beta} \delta(\vec{p} - \vec{q}) U_a(\vec{p}, \alpha) U_b^*(\vec{q}, \beta) \right. \\ &\quad \left. e^{i(\vec{q} \cdot \vec{y} - \vec{p} \cdot \vec{x})} e^{i(\omega_p - \omega_q)t} \right. \\ &\quad \left. + (2\pi)^3 \delta_{\alpha\beta} U_a(\vec{p}, \alpha) U_b^*(\vec{q}, \beta) \delta(\vec{p} - \vec{q}) \right. \\ &\quad \left. e^{i(\vec{q} \cdot \vec{x} - \vec{p} \cdot \vec{y})} e^{i(\omega_q - \omega_p)t} \right] \end{aligned}$$

②

$$\{\Psi_a(\vec{x}, t), \Psi_b^\dagger(\vec{y}, t)\} = \sum_{\alpha} \int \frac{d^3 p}{(2\pi)^3} \frac{N_p^2}{(1+\beta p^2)^2} \left[U_a(\vec{p}, \alpha) U_b^\dagger(\vec{p}, \alpha) e^{i \vec{p} \cdot (\vec{x} - \vec{y})} + U_a^\dagger(\vec{p}, \alpha) U_b(\vec{p}, \alpha) e^{i \vec{p} \cdot (\vec{y} - \vec{x})} \right]$$

$$\begin{aligned} &= \int \frac{d^3 p N_p^2 e^{i \vec{p} \cdot \vec{x}}}{(2\pi)^3 (1+\beta p^2)^2} \left[\sum_{\alpha} \left(U_a(\vec{p}, \alpha) U_b^\dagger(\vec{p}, \alpha) + U_a^\dagger(\vec{p}, \alpha) U_b(\vec{p}, \alpha) \right) \right] \\ &= \int \frac{d^3 p N_p^2}{(2\pi)^3 (1+\beta p^2)^2} \left[2 \delta_{ab} \right] e^{i \vec{p} \cdot (\vec{x} - \vec{y})} \\ &= 2 \int \frac{d^3 p N_p^2}{(2\pi)^3 (1+\beta p^2)^2} e^{i \vec{p} \cdot (\vec{x} - \vec{y})} \end{aligned}$$

since r.h.s :

$$\begin{aligned} \{\Psi_a(\vec{x}, t), \Psi_b^\dagger(\vec{y}, t)\} &= \delta_{ab} \int \tilde{S}(\vec{x} - \vec{y}) \\ &= \delta_{ab} \int \frac{d^3 p}{(2\pi)^3 (1+\beta p^2)^2} e^{i \vec{p} \cdot (\vec{x} - \vec{y})} \end{aligned}$$

$$\text{so we have } 2N_p^2 = 1 \quad \text{or} \quad N_p = \frac{1}{\sqrt{2}}$$

accordingly :

$$\begin{aligned} \Psi(\vec{x}, t) &= \sum_{\alpha} \int \frac{d^3 p}{(1+\beta p^2)(2\pi)^3 \sqrt{2}} \left[b(\vec{p}, \alpha) e^{i \vec{p} \cdot \vec{x}_n} U(\vec{p}, \alpha) \right. \\ &\quad \left. + d^\dagger(\vec{p}, \alpha) e^{-i \vec{p} \cdot \vec{x}_n} U^\dagger(\vec{p}, \alpha) \right] \end{aligned}$$

(3)

7.) $\Sigma\omega$ Lagrangian

$$\mathcal{L} = \bar{\Psi} \left[i \gamma^n (\partial_n + ig_\omega \omega_n) - \frac{m}{(m - g_\omega \sigma)} \right] \Psi + \frac{1}{2} \left[(\partial_n \sigma) (\partial^n \sigma) - m_\sigma^2 \sigma^2 \right] - \frac{1}{4} \omega_{n\nu} \omega^{n\nu} + \frac{1}{2} M_\omega^2 \omega_n \omega^n$$

in the mean relativistic scheme:

7.1) Dirac equation for nucleon

$$\left[\gamma_n (i \partial^n - g_\omega \omega_n) - (m - g_\omega \sigma) \right] \psi(x) = 0 \quad \dots (7.1)$$

dirac $\psi(x)$ spinor expand in maximally localized state

7.2) Condition for meson $\Sigma\omega$ (under MRF scheme):

$$M_\sigma^2 \langle \sigma \rangle = g_\sigma \langle \bar{\Psi} \Psi \rangle$$

$$M_\omega^2 \langle \omega_0 \rangle = g_\omega \langle \Psi^\dagger \Psi \rangle \quad \dots (7.2)$$

$$M_\omega^2 \langle \omega_k \rangle = g_\omega \langle \bar{\Psi} \gamma_k \Psi \rangle = 0$$

(g)

8.) Equation of State

8.1) Scalar density $\rho_s = \langle \bar{\psi} \psi \rangle$

$$\begin{aligned}
 \langle \bar{\psi} \psi \rangle &= \sum_{\alpha\beta} \int \frac{d^3 p d^3 q}{(1+\beta p^2)(1+\beta q^2)(2\pi)^3 2} \bar{U}(\vec{p}, \alpha) U(\vec{q}, \beta) e^{i x^\mu (\bar{p}_\mu - \bar{q}_\mu)} \delta(\vec{p} - \vec{q}) \delta_{\alpha\beta} \\
 &= \int \frac{d^3 p}{2(2\pi)^3 (1+\beta p^2)^2} \underbrace{\sum_{\alpha} \bar{U}(\vec{p}, \alpha) U(\vec{p}, \alpha)}_{\frac{2m^*}{\omega_p}} \\
 &= \int \frac{d^3 p}{(2\pi)^3 (1+\beta p^2)^2} \frac{m^*}{\omega_p} \\
 &= \int \frac{d^3 p}{(2\pi)^3 (1+\beta p^2)^2} \frac{m - g_r \Gamma}{\sqrt{\bar{p}^2 + m^*^2}} \\
 &= \int \frac{d^3 p}{(2\pi)^3 (1+\beta p^2)^2} \frac{m - g_r \Gamma}{\left[\frac{1}{\beta} (\tan^{-1} \sqrt{\beta} p)^2 + (m - g_r \Gamma)^2 \right]^{1/2}}
 \end{aligned}$$

$$\rho_s = \langle \bar{\psi} \psi \rangle = \frac{1}{2\pi^2} \int_0^{p_f} \frac{p^2 dp}{(1+\beta p^2)^2} \frac{m - g_r \Gamma}{\left[\frac{1}{\beta} (\tan^{-1} \sqrt{\beta} p)^2 + (m - g_r \Gamma)^2 \right]^{1/2}}$$

- - - (8.1)

Scalar density

(5)

(8.2) Vector density $\rho_0 = \langle \psi^\dagger \psi \rangle$

$$\begin{aligned}
 \langle \psi^\dagger \psi \rangle &= \sum_{\alpha\beta} \int \frac{d^3 p d^3 q}{2(2\pi)^3 (1+\beta p^2)(1+\beta q^2)} U^\dagger(\vec{p}, \alpha) U(\vec{q}, \beta) e^{ix^n(\vec{p}_n - \vec{q}_n)} \delta(\vec{p} \cdot \vec{q}) \delta_{\alpha\beta} \\
 &= \int \frac{d^3 p}{2(2\pi)^3 (1+\beta p^2)^2} \sum_{\alpha} U^\dagger(\vec{p}, \alpha) U(\vec{q}, \alpha) \\
 &= \int \frac{d^3 p}{(2\pi)^3 (1+\beta p^2)^2} \\
 &= \frac{1}{2\pi^2} \int_0^{p_p} \frac{p^2 dp}{(1+\beta p^2)^2} \\
 &= \frac{1}{2\pi^2} \left[\frac{\tan^{-1}[\sqrt{\beta} p_f] - \frac{\sqrt{\beta} p_f}{1+\beta p_f^2}}{2\beta^{3/2}} \right]
 \end{aligned}$$

$$\rho_0 = \langle \psi^\dagger \psi \rangle = \frac{1}{4\pi^2 (1+\beta p_f^2) \beta^{3/2}} \left[(1+\beta p_f^2) \tan^{-1}(\sqrt{\beta} p_f) - \sqrt{\beta} p_f \right]$$

... (8.2)

Vector density

(b)

8.3) Energy density

menggunakan bentuk (7.156) pada Glendenning:

$$\epsilon = -\langle \mathcal{L} \rangle + \langle \bar{\psi} \gamma_0 k_0 \psi \rangle$$

with

$$k_0 = \gamma_\omega \omega_0 + \omega_p = \gamma_\omega \omega_0 + \sqrt{\vec{p}^2 + m^*{}^2}$$

so:

$$(1) \langle \bar{\psi} \gamma_0 k_0 \psi \rangle = \langle \psi^\dagger k_0 \psi \rangle$$

$$\begin{aligned} &= \frac{1}{2\pi^2} \int_0^{p_F} \frac{p^2 dp}{(1 + \beta p^2)^2} \left[\gamma_\omega \omega_0 + \sqrt{\vec{p}^2 + m^*{}^2} \right] \\ &= \gamma_\omega \omega_0 \langle \psi^\dagger \psi \rangle + \frac{1}{2\pi^2} \int_0^{p_F} p^2 \sqrt{\frac{1}{\beta} (\tan^{-1} r_B p)^2 + m^*{}^2} dp \end{aligned}$$

$$(2) \langle \mathcal{L} \rangle = -\frac{1}{2} m_\sigma^2 \vec{r}^2 + \frac{1}{2} M_\omega \omega_0^2$$

accordingly from (7.2):

$$M_\omega^2 \omega_0 = \gamma_\omega \langle \psi^\dagger \psi \rangle$$

$$M_\omega^2 \omega_0^2 = \gamma_\omega \omega_0 \langle \psi^\dagger \psi \rangle$$

so

$$\epsilon = -\langle \mathcal{L} \rangle + \langle \bar{\psi} \gamma_0 k_0 \psi \rangle$$

$$\boxed{\epsilon = \frac{1}{2} m_\sigma^2 \vec{r}^2 + \frac{1}{2} M_\omega^2 \omega_0^2 + \frac{1}{2\pi^2} \int_0^{p_F} p^2 \sqrt{\frac{1}{\beta} (\tan^{-1} r_B p)^2 + (m - \gamma_\omega \vec{r})^2} dp}$$

(7)

energy density

--- (8.3)

8.3) Pressure

$$P = \langle \delta \rangle + \frac{1}{3} \langle \bar{\Psi} \chi; k_i \Psi \rangle$$

$$= -\frac{1}{2} M_\sigma^2 \bar{v}^2 + \frac{1}{2} M_\omega^2 (\bar{\ell})_0^2 + \frac{1}{6\pi^2} \int_0^{P_F} \frac{p^2 dp}{(1+\beta p^2)^2} \cdot \frac{\bar{P}^2}{\sqrt{\bar{E}^2 + M^2}}$$

$$\bar{P} = -\frac{1}{2} M_\sigma^2 \bar{v}^2 + \frac{1}{2} M_\omega^2 (\bar{\ell})_0^2 + \frac{1}{6\pi^2} \int_0^{P_F} \frac{p^2 dp}{(1+\beta p^2)^2} \cdot \frac{\frac{1}{\beta} (\tan^{-1} \sqrt{\beta} p)^2}{\left[\frac{1}{\beta} (\tan^{-1} \sqrt{\beta} p)^2 + M^2 \right]^{\gamma_1}}$$

$$\boxed{\bar{P} = -\frac{1}{2} M_\sigma^2 \bar{v}^2 + \frac{1}{2} M_\omega^2 (\bar{\ell})_0^2 + \frac{1}{6\pi^2} \int_0^{P_F} \frac{p^2 (\tan^{-1}(\sqrt{\beta} p))^2}{\beta (1+\beta p^2)^2 \left[\frac{1}{\beta} (\tan^{-1}(\sqrt{\beta} p))^2 + (M-\frac{1}{2}\bar{v}) \right]^{\gamma_1}}}$$

Pressure

(d)