

18/03/2021

Inductor

→


(DC)

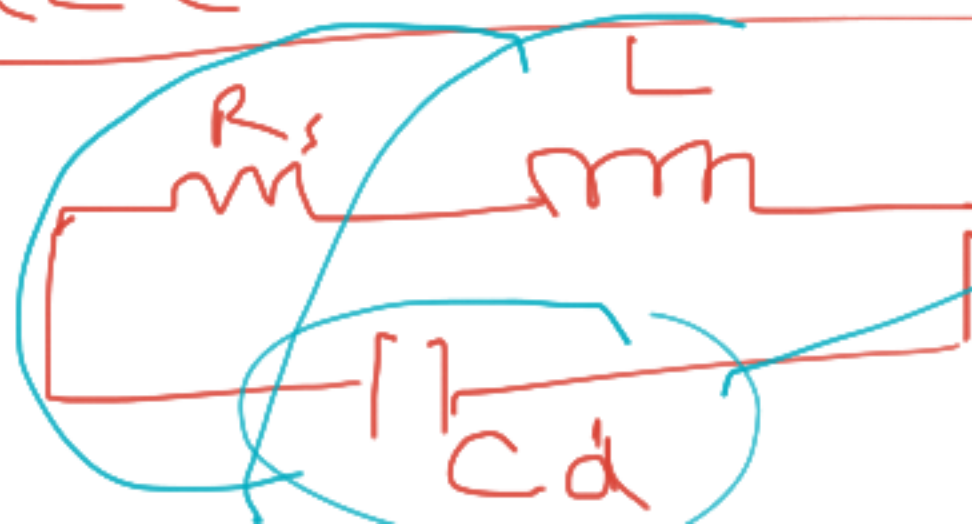


coiled to increase the magnetic flux
linkage



(distributed capacitor)
(C)

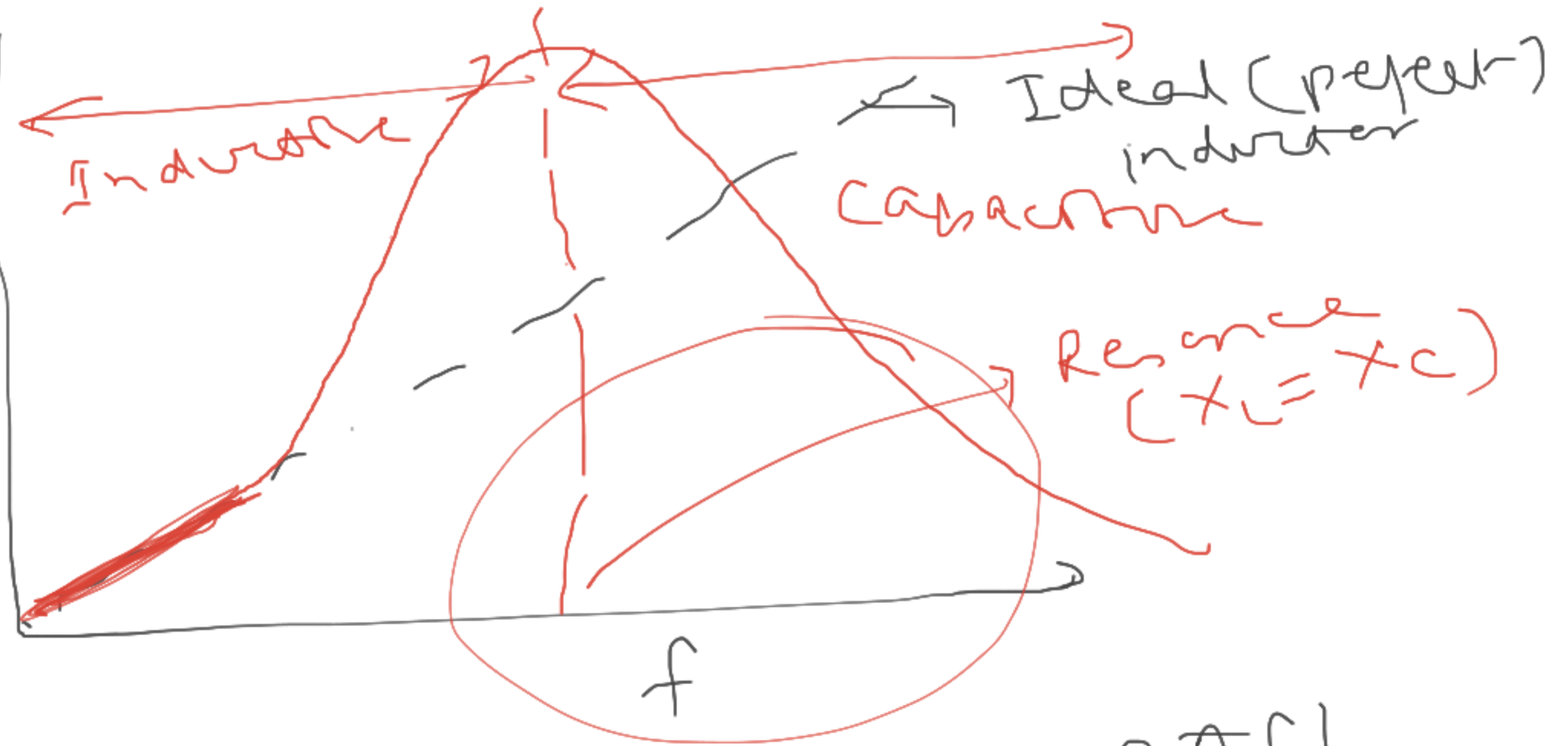
Eqn for AP
inductor at
AC (or RF or high freq)



Stray elements
Parasitic

$L, C \rightarrow$ resonance
 R_s - series resistance

Impedance
of
inductor
 $Z = X_L$
 $= j\omega L$
 $= j2\pi fL$
(Ω)



Quality factor $Q = \frac{X_L}{R_S} = \frac{\omega L}{R_S} = \frac{2\pi fL}{R_S}$

For a perfect inductor, $R_S = 0$
 $\Rightarrow Q = \infty$

To increase θ of L

\rightarrow use a larger diameter
 \Rightarrow resistance decreases $\left[R = \frac{\rho L}{A} \right]$
 \Rightarrow Q increase \rightarrow Area

$\Rightarrow C_d$ increases
 \rightarrow Spread the windings apart (d)
 $\Rightarrow C_d$ decreases $[C = \frac{\epsilon A}{d}]$

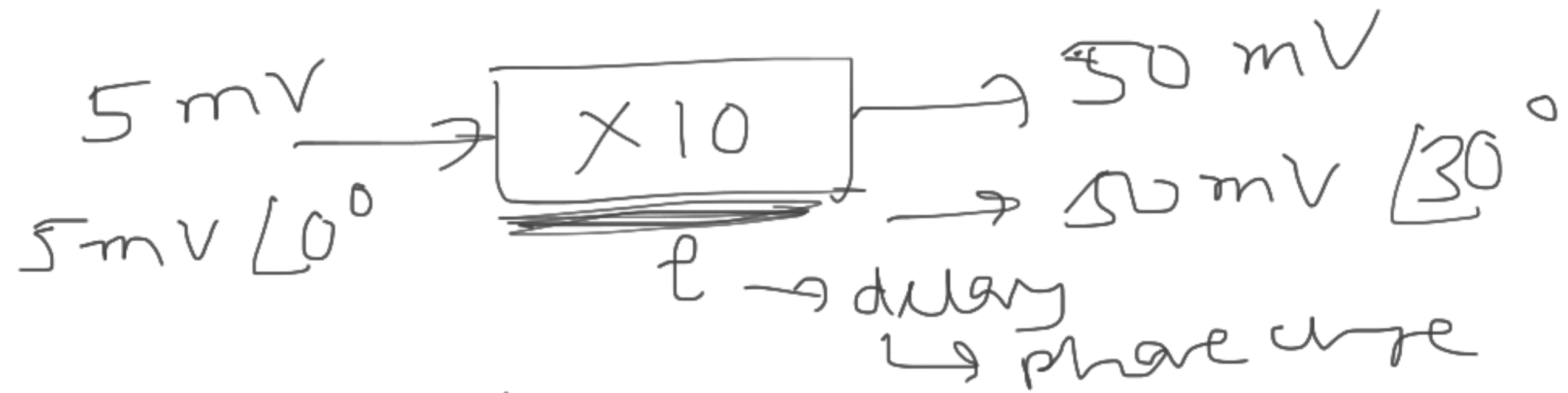
$\Rightarrow \mathcal{O}_S$ increases

$\Rightarrow \Phi$ increases
 \rightarrow increase the permeability of ^{magnetic} flux linkage
 by using a ferrite rod between the
 core $\rightarrow \Phi$ increases



Applications of L

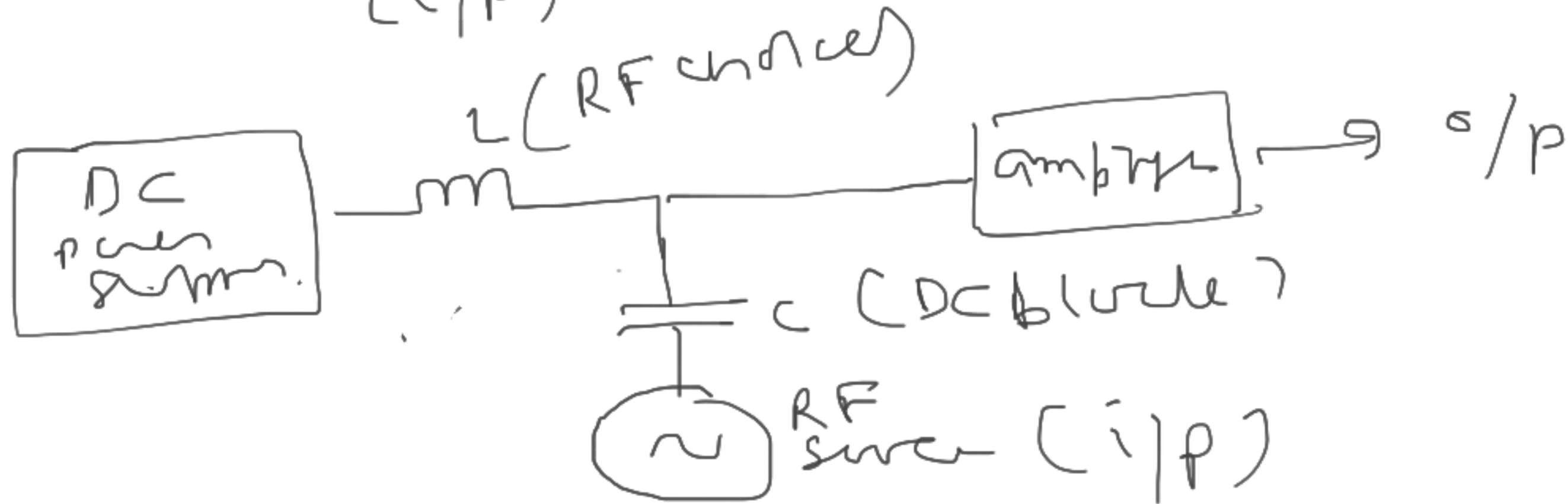
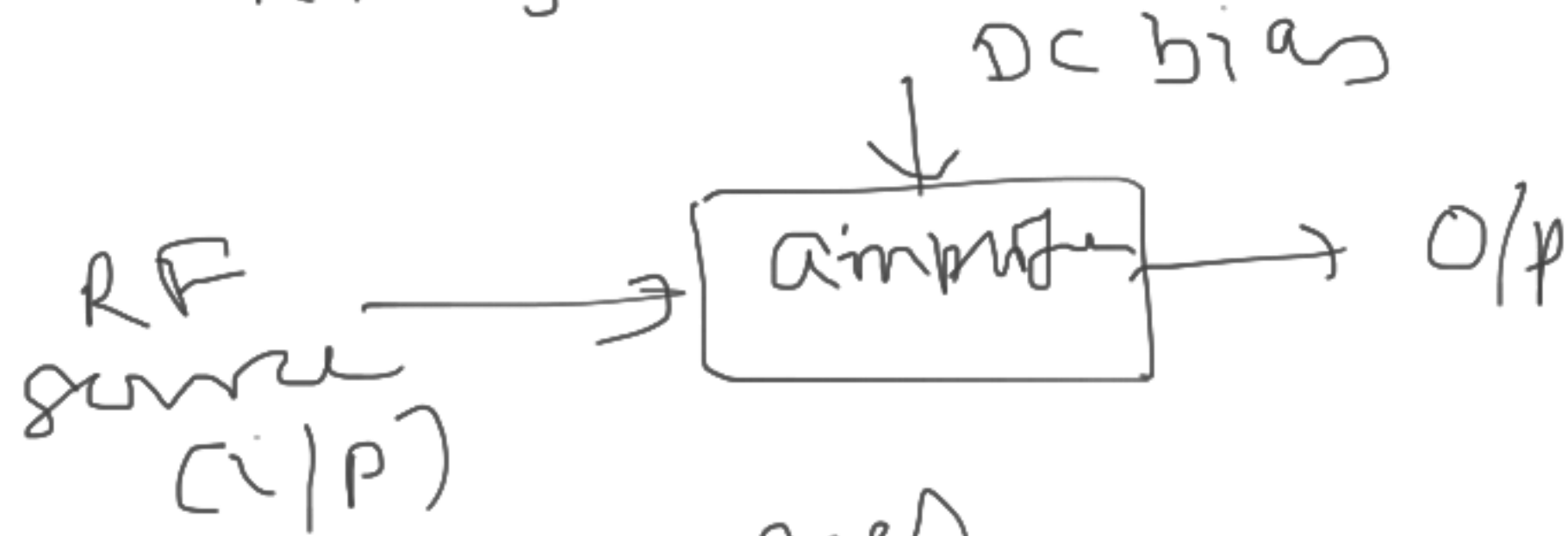
- LC filter
- oscillator $(f = \frac{1}{2\pi\sqrt{LC}})$
- delay (phase change)



- RF choke
(allows dc to pass through)
doesn't allow RF or ac to pass through

Bias Tee (T)

DC power supply (low freq or 2nd freq)
RF generator (high freq) —



Chip resists (no leads)



Size code



0402

0603

0805

1206

1218

L (mils)

40

60

80

120

120

w (mils)

20

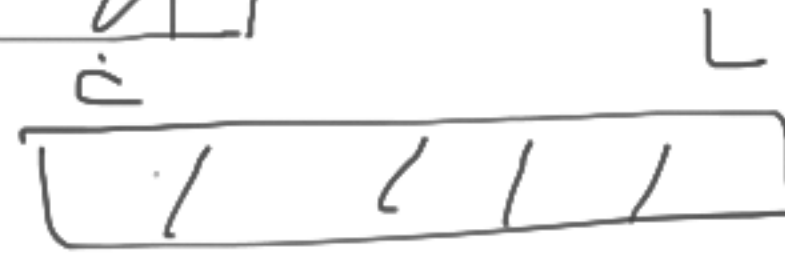
30

50

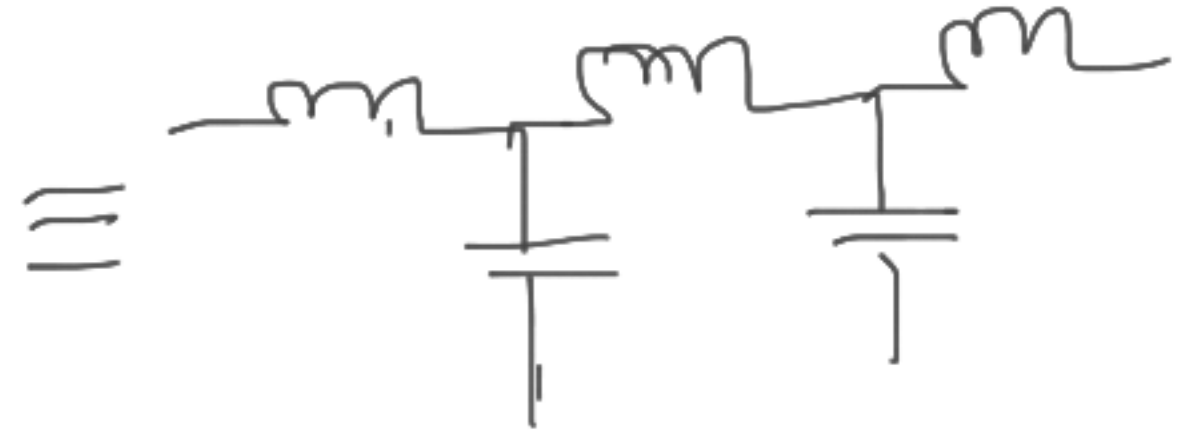
60

180

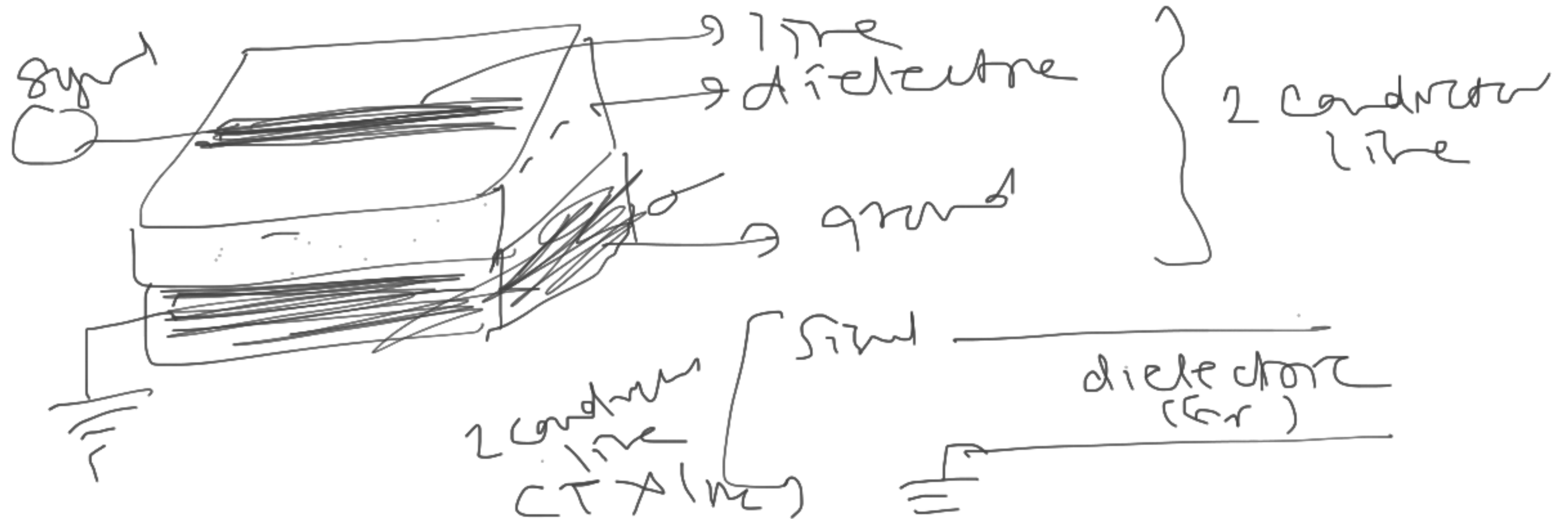
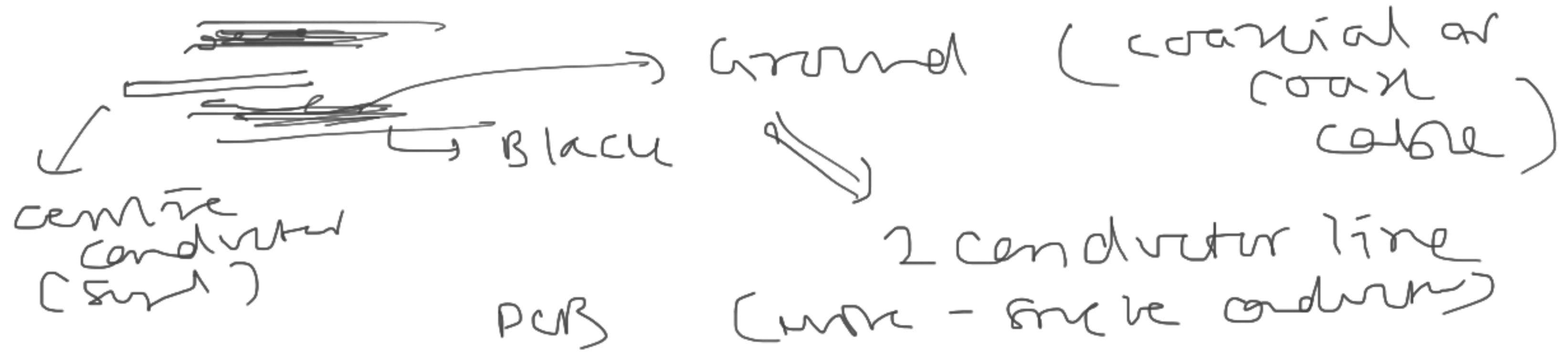
L, C (printed form)



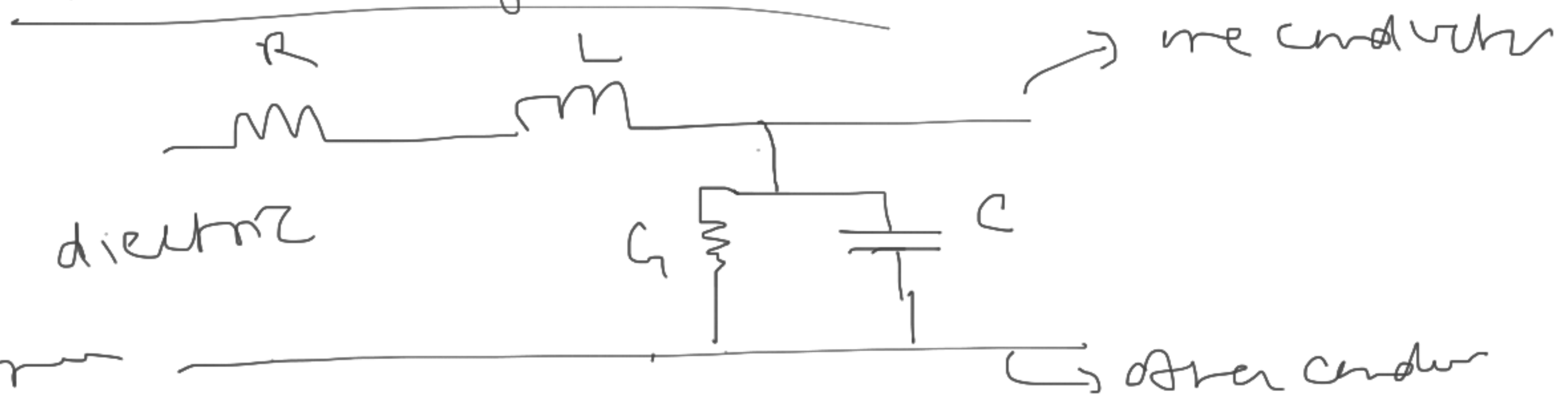
wire
R
□
L



Transmission Line



Eqvt circuit of TX line



- $R \rightarrow$ due to resistance of conductor
- $L \rightarrow$ represents magnetic flux generated by current on TX line
- $C \rightarrow$ capacitance due to dielectric between two conductors
- $G \rightarrow$ represents coupling currents between two conductors because dielectric (conductance) has some conductivity (leakage current)

R, L, G, C (not lumped or fixed values)
 ohm, H \rightarrow \downarrow mhos or ohm^{-1} or Siemens (S)

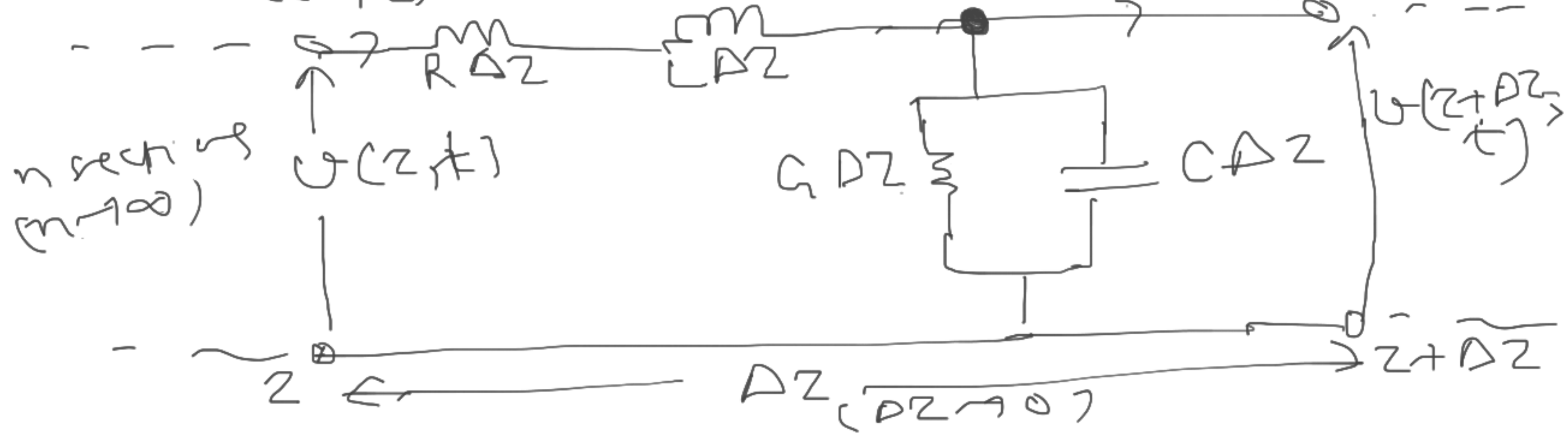
$$G \neq \frac{1}{R}$$

\rightarrow Distributed \rightarrow $R, L, G, C \rightarrow \frac{F}{\text{length}}$
 ohm
 $\frac{\text{length}}{\text{cm}}$

$V \neq IR$
 $\checkmark [KVL \times KCL \times$

Analysis of Transmission (TX) line

We divide TX line of length l into
 no. of small sections ($l = \Delta z$), $\Delta z \rightarrow 0$
 ($n \rightarrow \infty$)



$R \rightarrow \frac{\text{ohm}}{\text{length}}$
 (distributed)

$R \Delta z = \frac{\text{ohm length}}{\text{length}} = \text{fixed}$
 instantaneous (t)

Applying KVL (loop eqn)

→ sum of voltages in a loop = 0

$$\begin{aligned} \varphi(z, t) &= R \Delta z \cdot \bar{I}(z, t) - L \Delta z \cdot \frac{\partial \bar{I}}{\partial z}(z, t) \\ &= \varphi(z + \Delta z, t), \end{aligned}$$

where
 $\varphi(z,t)$ = instantaneous voltage at z
 " " " $(z+\Delta z)$

$$G(Z+DZ, t) =$$

Rearrange & dividing by ΔZ & combining

9 terms, we get

So here, we get

$$- \left[\frac{\psi(z+\Delta z, t) - \psi(z, t)}{\Delta z} \right] = R \hat{\psi}(z, t) + \left[\frac{\partial \hat{\psi}}{\partial t}(z, t) \right]$$

In limit $\Delta Z \rightarrow 0$
(difference goes to differential),
we have

$$\boxed{-\frac{\partial \psi(z,t)}{\partial z} = R \bar{i}(z,t) + L \frac{\partial \bar{i}(z,t)}{\partial t}}$$

(1)


Similarly, applying KCL to node N

→ sum of currents into a node ≥ 0

$$\begin{aligned} i(z,t) - G \Delta Z \cdot \psi(z + \Delta Z, t) - C \Delta Z \cdot \frac{\partial \psi(z + \Delta Z, t)}{\partial t} \\ = i(z + \Delta Z, t) \end{aligned}$$

Rearranging & dividing by ΔZ & combining
i terms, we set a limit $\Delta Z \rightarrow 0$,
we get

$$\frac{-\partial \bar{V}(Z, t)}{\partial Z} = G V(Z, t) + C \frac{\partial V(Z, t)}{\partial t}$$

① & ② general TX line eqns 

(or Telegrapher's eqns)

For harmonic time dependence,
use of phasors simplifies transmission
line eqns to ordinary differential eqns.

For a cosine reference, we write

$$V(z, t) = \operatorname{Re} [V(z) e^{j\omega t}] \quad (3)$$

$$\text{or } i(z, t) = \operatorname{Re} [I(z) e^{j\omega t}] \quad (4)$$

Substituting (3) or (4) in (1) or (2), we get
ordinary differential eqns.

$$\boxed{-\frac{dV(z)}{dz} = (R + j\omega L) I(z)} \quad (5)$$

$$\boxed{-\frac{dI(z)}{dz} = (G + j\omega C) V(z)} \quad (6)$$

Eqn (5) & (6) are time harmonic transmission line eqn.

In these eqn., V & I are complex

Eqn (5) & (6) $\rightarrow V(z), I(z) = ?$

Taking $\frac{d}{dz}$ of one eqn. (say, (6)) & substituting into another eqn. (5)

$$(6) \Rightarrow \frac{d^2 I(z)}{dz^2} = (G + j\omega C) \cdot \frac{dV(z)}{dz}$$

\rightarrow next lecture