MAT257 Notes

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This document is a collection of notes for the course MAT257: Analysis II, as taught by Professor Edward Bierstone in 2018 at the University of Toronto. The notes are a combination of notes I made in class (which can be found in their original form in the notes folder in this repository) and scans of handwritten notes which Professor Bierstone has generously given me the permission to use.

1 Introduction

TODO: this

2 Differentiation

TODO: this

3 Integration

3.1 The (Riemann) Integral Over a Rectangle

TODO: this

3.2 Integrals Over More General Bounded Sets

TODO: this

3.3 Fubini's Theorem

TODO: rewrite

We are now going to talk about Fubini's theorem, which is about how to integrate over a high-dimensional rectangle by repeatedly performing individual integrals. Let's start with an example. Suppose we want to integrate over a rectangle

$$A = [a, b] \times [c, d] \subseteq \mathbb{R}^2 \tag{1}$$

and let's suppose, to make it easy, we have a continuous, non-negative function f, defined on this rectangle, say

$$z = f(x, y) > 0 \tag{2}$$

The idea is that if we fix a point on the x-axis, say x, we can consider the "slice" in the y-direction determined by this point. We could find, for example, the area of that slice, and it's reasonable to expect that the integral of f on the rectangle, we could obtain by integrating the area of that slice along the length of the rectangle.

The idea then is to try to find the area of such a slice, which would be in this case, for some x,

$$h(x) = \int_{c}^{d} g_x(y)dy = \int_{c}^{d} f(x,y)dy$$
(3)

where $g_x(y) = f(x, y)$ fixing x. It's reasonable to expect that

$$\int_{A} f = \int_{a}^{b} h = \int_{a}^{b} \left(\int_{c}^{d} f(x, y) dy \right) dx \tag{4}$$

So that's the idea of Fubini's theorem: that we should be able to integrate a function over a rectangle by repeated one-dimensional integrals. Of course, we're not interested in integrating only continuous functions, we want to look at more general, integrable functions, but if you think about that, supposing f is integrable, this could run into a problem: one of the functions g_x might not be integrable on [c, d]. After all, it's set of discontinuities could be $x_0 \times [c, d]$ for some x_0 , so $\int_c^d g_{x_0}$ makes no sense!

So we'll have to formulate something maybe a little bit more technical, but it's to capture that problem. Suppose we just have $f:A\to\mathbb{R}$ bounded. The function may or may not be integrable, meaning the supremum of lower sums is equal to the infimum of lower sums. Whether the function is integrable or not, we can still look at the supermum of lower sums and the infimum of lower sums, which is what we'll do.

We'll define the <u>lower</u> and <u>upper integrals</u> of f on A, $L \int_A f$ and $U \int_A f$ respectively, to be the supremum of all the lower sums $L(f, \mathcal{P})$ and infimum of all the upper sums $U(f, \mathcal{P})$ respectively. We can now write down our theorem as the previous formula, but taking into account that we don't know that the function g_x mentioned before, we just replace it by the lower or upper integral:

Theorem 1. Suppose $A \subset \mathbb{R}^m$ and $B \subset \mathbb{R}^n$ are closed rectangles and $f: A \times B \to \mathbb{R}$ is integrable. For all $x \in A$, define

$$g_x: B \to \mathbb{R}, g_x(y) = f(x, y)$$
 (5)

Set

$$\mathcal{L}(x) = \mathcal{L} \int_{B} g_{x} = \mathcal{L} \int_{B} f(x, y) dy, \mathcal{U}(x) = \mathcal{U} \int_{B} g_{x} = \mathcal{U} \int_{B} f(x, y) dy$$
 (6)

Then $\mathcal{L}(x)$, $\mathcal{U}(x)$ are integrable on A and

$$\int_{A \times B} f = \int_{A} \mathcal{L} = \int_{A} \mathcal{U} = \int_{A} \left(L \int_{B} f(x, y) dy \right) dx = \int_{A} \left(U \int_{B} f(x, y) dy \right) dx \tag{7}$$

Before we prove this, some remarks:

1. We also have

$$\int_{A\times B} f = \int_{B} \left(\mathcal{L} \int_{A} (f(x,y) dx) \right) dy = \int_{B} \left(\mathcal{U} \int_{A} (f(x,y) dx) \right) dy \tag{8}$$

2. If g_x is integrable on B for every $x \in A$, then

$$\int_{A \times B} f = \int_{A} \left(\int_{B} f(x, y) dy \right) dx \tag{9}$$

3. If

$$A = [a_1, b_1] \times [a_2, b_2] \times \dots \times [a_n, b_n]$$
(10)

we can apply Fubini's theorem repeatedly to obtain

$$\int_{A} f = \int_{a_{n}}^{b_{n}} \left(\dots \int_{a_{2}}^{b_{2}} \left(\int_{a_{1}}^{b_{1}} f(x_{1}, \dots, x_{n}) dx_{1} \right) dx_{2} \dots \right) dx_{n}$$
(11)

if f is "sufficiently nice" (otherwise we'll have to sprinkle in some L's or U's).

Now here's an application to $\int_C f$ when f is integrable on a Jordan-measurable set $C \subset \mathbb{R}^n$. What we're going to use is Fubini's theorem in a rectangle with the formula

$$\int_{C} f = \int_{A} f \chi_{C} \tag{12}$$

where $A \supseteq C$ is a closed rectangle. Let's do some quick examples:

1. Let's say we were integrating on the region

$$C = [-1, 1] \times [-1, 1] \setminus \{x \in \mathbb{R}^2, x_1^2 + x_2^2 < 1\}$$
(13)

i.e. the unit square with the unit circle removed from it. We have

$$\int_C f = \int_{[-1,1]^2} f \chi_C \tag{14}$$

We can write

$$\chi_C = \begin{cases} 1 & \text{if } -1 \le y \le -\sqrt{1-x^2} \text{ or } \sqrt{1-x^2} \le y \le 1\\ 0 & \text{otherwise} \end{cases}$$
 (15)

We can write

$$\int_{-1}^{1} f(x,y)\chi_C(x,y)dy = \int_{-1}^{-\sqrt{1-x^2}} f(x,y)dy + \int_{\sqrt{1-x^2}}^{1} f(x,y)dy$$
 (16)

i.e. the limits of integration correspond to the boundary of C. So we get

$$\int_{C} f = \int_{-1}^{1} \left(\int_{-1}^{-\sqrt{1-x^{2}}} f(x,y) dy + \int_{\sqrt{1-x^{2}}}^{1} f(x,y) dy \right) dx \tag{17}$$

2. Say we want to integrate over a triangle C, under the line from (0,0) to (a,a). We have

$$\int_{C} f = \int_{0}^{a} \left(\int_{y}^{a} f(x, y) dx \right) dy \tag{18}$$

We could have done it the other way, integrating first with respect to y and then with respect to x, and get

$$\int_C f = \int_0^a \left(\int_0^x f(x, y) dy \right) dx \tag{19}$$

This shows the form might not be the same in different directions. And it could be important to exploit the difference. For example, what if our function f(x, y) only depends on one of the variables, say y (i.e. f is independent of x). Then

$$\int_{0}^{a} \left(\int_{0}^{x} f(y) dy \right) dx \tag{20}$$

doesn't simplify very much, but

$$\int_0^a \left(\int_y^a f(y) dx \right) dy = \int_0^a (a - y) f(y) dy \tag{21}$$

So this double integral reduces to a single integral with respect to y.

TODO: proof

TODO: polish Before we move on, we will give some more examples of computing integrals, using Fubini's theorem and the change of variables theorems as tools.

Let us begin by recalling some facts about the regions on whuch we can integrate. What are the kinds of region on which we integrate? One of the kinds of regions would be the region between the graphs of two functions that are defined on a Jordan-measurable set. Specifically, if we have a subset $C \subset \mathbb{R}^n$ that is Jordan-measurable, and two integrable functions $\varphi(x) \leq \psi(x)$ defined on C, we're interested in integrating over the region bounded by the graphs of the two functions Let's call this region S, that is, define

$$S = \{(x, y) \in \mathbb{R}^n \times \mathbb{R} : x \in C, y \in [\varphi(x), \psi(x)]\}$$
(22)

We have that S is a Jordan-measurable set. We want to be able to integrate a continuous function on S. If f(x,y) is a bounded continuous on S, then we can compute using Fubini's theorem that

$$\int_{S} f = \int_{C} \left(\int_{\varphi(x)}^{\psi(x)} f(x, y) dy \right) dx \tag{23}$$

We went through the proof of this before, but I want to use it to compute some examples.

(a) Consider

$$\int_{0}^{2} \int_{y/2}^{1} y e^{-x^{3}} dx dy \tag{24}$$

There's no hope in proceeding naively, since e^{-x^3} doesn't even have an elementary primitive, i.e. you cannot write down it's integral with only elementary functions. But let's consider what region we're integrating along: the area under the line between (0,0) and (2,1). So using the above observation, we can rewrite the integral to be

$$\int_0^1 \left(\int_0^{2x} y e^{-x^3} dy \right) dx = \int_0^1 \frac{y^2}{2} e^{-x^3} \Big|_0^{2x} dx = 2 \int_0^1 x^2 e^{-x^3} dx = -\frac{2}{3} e^{-x^3} \Big|_0^1 = \frac{2}{3} (1 - e^{-1}) \quad (25)$$

Here, the change of variable was a useful thing to do, because it enabled us to write down an original integral, which we couldn't do in the original form.

(b) Consider

$$\int_{2}^{4} \int_{4/x}^{\frac{20-4x}{8-x}} (y-4)dydx \tag{26}$$

So here, again, should we just go ahead and do it as it's written? Well then we'll get $\frac{y^2}{2} - 4y$ and we'll substitute those things in and we'll just get a terrible mess. But what's the region we're integrating over? $\frac{4}{x}$ is like a hyperbola, which is 2 when x is 2 and 1 when x is 4. The other function on top, what does it look like? It's also a hyperbola: it's a constant plus something over 8 - x, so it'll open down. And this is the region we're integrating on: the area between a hyperbola open up and another opening down.

If we change the order of integration, we'll be integrating on the outside with respect to y, which goes from 1 to 2, and we'll be integrating on the inside with respect to x with x going from 4/y to... let's solve:

$$y = \frac{20 - 4x}{8 - x} = 4 + \frac{12}{x - 8} \implies x - 8 = \frac{12}{y - 4} \iff x = 8 + \frac{12}{y - 4}$$
 (27)

Hence we can write the above as

$$\int_{1}^{2} \left(\int_{4/y}^{8+12/(y-4)} (y-4) dx \right) dy = \int_{1}^{2} (y-4)x \Big|_{4/y}^{8+12/(y-4)} dy = \int_{1}^{2} (y-4) \left(8 + \frac{12}{y-4} - \frac{4}{y} \right) dy$$
(28)

which is all stuff that's very easy to integrate

(c) Let's see how to integrate a simple function z over a region of 3-space

$$S = \{(x, y, z) \in \mathbb{R}^3 : x^2 + y^2 \le z^2 \land x^2 + y^2 + z^2 \le 1\}$$
 (29)

So, what is this region? The second equation says we're inside the closed ball of radius 1 centered at the origin. The first equation is the region inside two cones, opening up and down from the origin (kind of like an hour glass). So we want the intersection of these two regions. In terms of the theorem that we wrote down at the beginning, the set C is like the discs inside the circles formed by the intersections of the top and bottom of the cones and circles, and we're integrating on the region between the semicircles above and below the plane, and the cones above and below the plane. So what's this going to be? Well, by symmetry, it's going to be zero.

Let's make it a little harder. Let's consider only the top, i.e.

$$S^{+} = \{(x, y, z) \in S : z \ge 0\}$$

So, rewriting our integral to be over the disc C, we obtain

$$\int \int_{x^2+y^2 \le \frac{1}{2}} \left(\int_{\sqrt{x^2+y^2}}^{\sqrt{1-(x^2+y^2)}} z dz \right) dx dy$$

$$= \int \int_{x^2+y^2 \le \frac{1}{2}} \frac{1}{2} (1 - (x^2 + y^2) - (x^2 + y^2)) dx dy$$

$$= \frac{1}{2} \int \int_{x^2+y^2 \le \frac{1}{2}} dx dy - \int \int_{x^2+y^2-\frac{1}{2}} (x^2 + y^2) dx dy = \frac{1}{2} \pi \frac{1}{2} = \frac{\pi}{4}$$
(30)

How do we justify that last, "magical" step? Well, the prettiest way to do so is to change to polar coordinates, but we haven't justified this quite yet, which is the point of this example. You'll remember from first year caclulus that we can write

$$\int_{a}^{b} f(g(x))g'(x)dx = \int_{g(a)}^{g(b)} f(u)du$$
 (31)

This is what we're going to be doing in several variables. But first: polar coordinates? That means writing

$$x = r\cos\theta, y = r\cos\theta, \theta \in [0, 2\pi], r \in [0, \infty)$$
(32)

This gives that

$$\frac{\partial(x,y)}{\partial(r,\theta)} = \begin{pmatrix} \cos\theta & -r\sin\theta\\ \sin\theta & r\cos\theta \end{pmatrix} \implies \det\frac{\partial(x,y)}{\partial(r,\theta)} = r \implies dxdy = drd\theta \tag{33}$$

"generalizing" equation 31. On the other hand, we have that

$$x^2 + y^2 = r^2 (34)$$

Since we have that, for this disc, θ ranges over the whole interval $[0, 2\pi]$ whereas r ranges over $[0, 1/\sqrt{2}]$, we can hence rewrite the above integral as

$$\int_0^{2\pi} \int_0^{\frac{1}{\sqrt{2}}} r^2 \cdot r dr d\theta = \int_0^{2\pi} d\theta \int_0^{\frac{1}{\sqrt{2}}} r^3 dr = 2\pi \frac{r^4}{4} \Big|_0^{\frac{1}{\sqrt{2}}} = \frac{\pi}{8}$$
 (35)

3.4 Partitions of Unity

3.4.1 Bump Functions

A bump function, in the one dimensional case, is a function which goes up to 1 from 0, stays there for a while, and then returns to zero. But we're interested in doing this in n-dimensions. How? Given an open $U \subset \mathbb{R}^n$, and compact $C \subset U$, we want to construct a \mathcal{C}^{∞} function $f: U \to \mathbb{R}$ with $0 \le f \le 1$ such that $\forall x \in C, f(x) = 1$ and f(x) = 0 outside some compact subset of U. So it looks like the 1-dimensional bump function, except over an n-dimensional area C.

How can we construct such a function? We'll do it as a kind of series of exercises, based on one particular function, which you can imagine ought to be the kind of thing that you have to build this function, that is,

$$f(x) = \begin{cases} e^{1/x^2} & \text{if } x \neq 0 \\ 0 & \text{if } x = 0 \end{cases}$$
 (36)

We proceed as follows, building up by steps:

(a) There is a \mathcal{C}^{∞} function $g: \mathbb{R} \to \mathbb{R}$ such that g > 0 on (-1,1), 0 elsewhere. It's easy to write down a formula for such a function using the function given in 36, as follows:

$$g(x) = \begin{cases} e^{-\frac{1}{(x-1)^2}} e^{-\frac{1}{(x+1)^2}} & \text{if } x \in (-1,1) \\ 0 & \text{if } x \notin (-1,1) \end{cases}$$
(37)

This works because $e^{-\frac{1}{(x-1)^2}}$ becomes completely flat at x=1, whereas $e^{-\frac{1}{(x+1)^2}}$ becomes completely flat at x=-1.

(b) Given $a \in \mathbb{R}^n, \delta > 0$, there is a C^{∞} function $g_{a,\delta} : \mathbb{R}^n \to \mathbb{R}$ such that

$$g_{a,\delta} > 0 \text{ on } |x - a| < \delta, g_{a,\delta} = 0 \text{ outside}$$
 (38)

So how could achieve this? We could use the function defined in 37 to write

$$g_{a,\delta} = g\left(\frac{|x-a|^2}{\delta^2}\right) \tag{39}$$

(c) There is a \mathcal{C}^{∞} function $h: U \to \mathbb{R}$ such that h is greater than 0 on C and h = 0 outside a compact subset of U. So how do we do this?

We could cover C with a finite set of open balls $U_1, ..., U_k$ with centers $a_1, ..., a_k$ and radii $\delta_1, ..., \delta_k$ whose closures lie in U. We can do this since C is compact. Once we do this, we can define quite simply

$$h(x) = \sum_{i=1}^{k} g_{a_k, \delta_k}(x) \tag{40}$$

(d) As a fourth thing, going back to the real line, let's show that if we're given any $\epsilon \in \mathbb{R}^+$, we can find a \mathcal{C}^{∞} function ψ_{ϵ} which is zero up to 0, becomes 1 at ϵ , and then remains 1 forever after. How can we do this from what we've already done? We're going to do it starting with a \mathcal{C}^{∞} function like that in equation 37, but instead of on (-1,1), on $(0,\epsilon)$. More generally, actually, we can use any function g such that g > 0 on $(0,\epsilon)$ and is zero outside. We can then define ψ_{ϵ} by

$$\psi_{\epsilon}(x) = \frac{\int_0^x g}{\int_0^{\epsilon} g} \tag{41}$$

(e) Now, putting everything together, how can we define the bump function that we wanted to begin with? We begin with the function h as defined in equation 40, and we're going to modify it using equation 41 as follows: letting $\epsilon = \min_C h > 0$, which exists since C is compact, let

$$f(x) = \psi_{\epsilon} \circ h \tag{42}$$

Now there are other ways of doing it. One of them is to write down one integral formula from the beginning, but it would be harder to see just where this formula comes from in the first place.

3.4.2 Definition

We can now get to defining partitions of unity:

Theorem 2. Given $A \subset \mathbb{R}^n$ and an open cover \mathcal{O} of A, there is a countable collection Φ of \mathcal{C}^{∞} functions φ on \mathbb{R}^n with the following properties:

- (a) $0 \le \varphi \le 1$
- (b) For all $x \in A$, there is an open neighborhood V_x of x in \mathbb{R}^n such that all but finitely many φ vanish on V_x
- (c) For all $x \in A$, $\sum_{\varphi \in \Phi} \varphi(x) = 1$ (sum is finite in a neighborhood of x)
- (d) For all $\varphi \in \Phi$, there is some open $U \in \mathcal{O}$ such that $\varphi = 0$ outside a compact subset of U.

Definition 1. A collection φ satisfying conditions 2a, 2b, 2c above is called a \underline{C}^{∞} partition for A. If condition 2d holds, this partition is called <u>subordinate to \mathcal{O} </u>

This is all \mathcal{C}^{∞} stuff, but we could have made something weaker. We could have asked for a \mathcal{C}^r partition, or even a partition that was merely continous. So suppose we did that, and only cared about a \mathcal{C}^0 partition. Then, we could also, though the proof will have to come next time, we could use functions that look like trapezoids, using |x| instead of $e^{-\frac{1}{x^2}}$. If we wanted a \mathcal{C}^r partition of unity, then of course, we could do the same thing, but only caring about \mathcal{C}^r instead of \mathcal{C}^{∞} , and hence using the function given by

$$f(x) = \begin{cases} |x^{r+1}| & \text{on} \quad [0, \infty) \\ -|x^{r+1}| & \text{on} \quad (-\infty, 0] \end{cases}$$
 (43)

Before we try to construct such a (\mathcal{C}^{∞}) partition, let's discuss why this is an interesting idea. We're going to use this for several things in this course:

- To extend the definition of the integral to more general regions
- To prove the change of variables theorem
- To extend the definition of the integral to manifolds
- To prove Stoke's theorem

So why? Suppose we wanted to extend the definition of the integral to manifolds. Take some shape in 3-space, with a boundary, which is a manifold. Right now, we can integrate over a rectangle. Once we prove the change of variables theorem, we'll be able to integrate over a space diffeomorphic to a rectangle, a sort of deformed rectangle. Now, if we have a coordinate chart, a diffeomorphism from \mathbb{R}^n to some part of the manifold, we could use this to extend the definition of the integral from the rectangle to at least some kind of "rectangular" piece of the manifold. The question remains: how can we extend the definition to the *entire* manifold? A reasonable way to try to do this is to try to cover the entire manifold by a grid of rectangles. If we add up the integrals on each of those pieces...

That was the way this was doen classically: you'd take a manifold and divide it up by some kind of rectangular grid, and just add up the contributions of every piece. This classical approach, however, ran into a big obstacle, which was: how do you know that you can really cover every manfield with such a grid? How can you really divide up a really complex surface into rectangles? That is an extremely hard problem: it's true you can do it, but it's an extremely hard problem, and was solved way after the extension of integrals to manifolds.

So how can we avoid this issue? The way this was solved was, you do something a million times simpler. When you cover the manifold with rectangles, you don't care that they all fit together in a nice grid. Just, for every point, you take a rectangle containing that point, and we don't care whether they overlap or not. How do we exploit that? We use a partition of unity! Because what we'll do is, we'll pick a partition of unity such that every one of these bump functions sits in teh interior of one of these rectangles. Now, if we want to integrate our function f over the manifold, and we have our partition of unity Φ , and we look at φf , it vanishes outside one of these rectangles, so we can just integrate over the rectangle and get the answer for the manifold. We can then write

$$\int f = \sum \int \varphi f \tag{44}$$

So the key is to take a function, and make it's support very small, and then work with those small pieces. And we're going to use this in the same way to prove Stoke's theorem. Stoke's theorem is like the FTC. What does the FTC say? It says that

$$\int_{a}^{b} f' = f(a) - f(b) \tag{45}$$

So the integral of f over the <u>boundary</u> of [a,b] is the same as the integral of the derivative of f over the whole interval. So Stoke's theorem is that the integral of a function over the boundary of a manifold is the same as the integral of the derivative over the whole manifold.

We can prove this for rectangles. Then, if we want to prove it for manifolds, again we need to try to cover the manfolds with a grid like this. If we can do that, we can prove Stoke's theorem for surfaces

that we know can be covered by a grid. But we can do it for general manifolds in a really cheap way, again by multiplying with partitions of unity. So that's whay we're going to do in the rest of the course. But first, we have to prove Theorem 2, because that's what we're going to use to glue together little pieces in a very nice but cheap way.

Proof. The way we're going to do this is by considering a number of cases building up to the general case.

(a) Assume A is compact. That means finitely many $U \in \mathcal{O}$ cover A. Call them $U_1, ..., U_q$. In this case we're going to construct a *finite* (not countable) partition of unity subordinate to $\{U_1, ..., U_q\}$. For each $i \in \{1, ..., q\}$, the first thing that we're going to do is show that we can find a compact set $D_i \subset U_i$ such that the interiors of the D_i cover A. How?

For all $x \in A$, let D_x be a closed ball with centre x lying inside some U_i . Since A is compact, finitely many of the interiors of the D_x cover A. Number them $D_{x_1}, ..., D_{x_p}$. For each $i \in \{1, ..., q\}$, let

$$D_i = \bigcup_{D_{x_j} \subset U_i} D_{x_j} \tag{46}$$

Let ψ_i be a \mathcal{C}^{∞} function $0 \leq \psi_i \leq 1$ such that $\psi_i > 0$ on D_i , and is 0 outside a compact subset of U_i . Notice if we do this, then, since the interior of the D_i 's cover's A,

$$\forall x \in U \supseteq A, \sum_{i=1}^{q} \psi_i(x) > 0 \tag{47}$$

where U is an open neighborhood of A. SO we define our partition of unity as

$$\Phi = \{\varphi_1, ..., \varphi_q\}, \varphi_i = \frac{\psi_i}{\sum_{i=1}^q \psi_i}$$

$$\tag{48}$$

(b) Now assume $A = \bigcup_{i \in \mathbb{N}} A_i$ where each A_i is compact and a subset of the interior of A_{i+1} . Let $B_i = A_i \setminus \text{Int } A_{i-1}$. Let \mathcal{O}_i be covers of B_i by the open sets

$$U \cap (\operatorname{Int} A_{i+1} \setminus A_{i-2}) \tag{49}$$

for each $U \in \mathcal{O}$. By Case 2a there is a finite partition of unity Φ_i for B_i subordinate to \mathcal{O}_i . Consider

$$\sigma = \sum_{i} \sum_{\psi_i \in \Phi_i} \psi_i \tag{50}$$

This is a finite sum in some neighborhood of every $x \in A$. Take for a partition

$$\varphi = \frac{\psi}{\sigma} \tag{51}$$

for each $\psi \in \Phi_i$, for every i.

(c) Assume A is open. Then $A = \bigcup_{i \in \mathbb{N}} A_i$ as in Case 2b. How can we get such a representation? Try

$$A_i = \{x : |x| \le i, d(x, \mathbb{R}^n \setminus A) > i^{-1}\}$$
 (52)

(d) Let A be arbitrary. Apply Case 2c to $\bigcup_{U \in \mathcal{O}} U$. A partition of unity subordinate to \mathcal{O} is also a partition of unity for A subordinate to \mathcal{O}

3.5 Extended Definition of the Integral ("Improper Integral")

Suppose that A is open and f is bounded and continuous. Then under the current definition, the expression

$$\int_{A} f \tag{53}$$

may not exist, since the boundary may not be Jordan-measurable. So we now want to try and rectify that, by attempting to extend the definition of the integral to open subsets $A \subset \mathbb{R}^n$ and *locally* bounded functions $f: A \to \mathbb{R}$, i.e. for every point in A, there is an open neighborhood around A such that f is bounded around A. Furthermore, we assume the set of discontinuities is of measure zero.

Just like the improper integral from first year calculus, it won't always exist. But we'll want to know the value when it does. We'll begin with a useful note: if f vanishes outside of a compact subset C of A, then in fact f is integrable on C (since we can put that compact subset in a big rectangle and multiply by the characteristic function). In fact, we can say that f is integrable on any bounded open subset U of A containing C, for the same reason, and

$$\int_{U} f = \int_{C} f \tag{54}$$

So it makes sense to say

$$\int_{A} f = \int_{C} f \tag{55}$$

This is going to be our starting point. Now to go further, we're going to talk about a partition of unity. Let \mathcal{O} be an open cover of A such that $U \subset A$ for all $U \in \mathcal{O}$ and let Φ be a partition of unity for A subordinate to \mathcal{O} . Note that this is the same as saying let Φ be a partition of unity for A subordinate to $\{A\}$, since that counts as an open covering of A. As in equation 55, we can define, $\forall \varphi \in \Phi$,

$$\int_{A} \varphi |f| \tag{56}$$

We can now extend the definition of the integral as follows:

Definition 2. f is integrable on A if

$$\sum_{\varphi \in \Phi} \int_{A} \varphi |f| \tag{57}$$

converges. Then we define

$$\int_{A} f = \sum_{\varphi \in \Phi} \varphi f \tag{58}$$

Note that equation 58 must converge, and in fact converges absolutely, if equation 57 converges, because

$$\sum_{\varphi \in \Phi} \left| \int_{A} \varphi f \right| \le \sum_{\varphi \in \Phi} \int_{A} |\varphi| |f| = \sum_{\varphi \in \Phi} \int_{A} \varphi f \tag{59}$$

since φ is nonnegative.

Theorem 3. Suppose $A \subset \mathbb{R}$ is open, $f : A \to \mathbb{R}$ is locally bounded and the set of discontinuities of f has measure 0.

(a) Given partitions Φ, Ψ as above, not necessarily subordinate to the same open covering, though this doesn't matter since both will be subordinate to $\{A\}$, then if equation 57 converges for Φ , then it converges for Ψ and the integral defined using Φ and the integral defined using Ψ are equal, i.e.

$$\sum_{\psi \in \Psi} \int_{A} \psi f = \sum_{\varphi \in \Phi} \int_{A} \varphi f \tag{60}$$

Note that equation 60 would not necessarily be true if equation 57 does not converge.

(b) The new definition of the integral always exists if both A and f are bounded, showing it truly generalizes the old definition. Specificially, if both A and f are bounded, whether or not the boundary of A has measure zero,

$$\sum_{\varphi \in \Phi} \int_{A} \varphi |f| \tag{61}$$

(equation 57) converges

(c) If A is Jordan-measurable and f is bounded, then

$$\int_{A} f = \sum_{\varphi \in \Phi} \int_{A} \varphi f \tag{62}$$

where the left is as defined before.

There are other ways of defining the integral, some of which look more like the improper integral from first year calculus. Another thing that you could do is the following: since you can always write an open set as a kind of "expanding union" of compact sets $C_i \subseteq C_{i+1}$, you can actually do that in such a way such that each C_i is Jordan-measurable, meaning the integral on each of these C_i 's exist. We could then define

$$\int_{A} f = \lim_{i \to \infty} \int_{C_{i}} f \tag{63}$$

After we finish proving this theorem, either I'll prove this or stick it on the next problem set.

Proof. (a) $\varphi \cdot f$ vanishes outside a compact set, and only finitely many ψ are nonzero on this set. So

$$\sum_{\psi \in \Psi} \psi = 1 \implies \sum_{\varphi \in \Phi} \int_{A} \varphi f = \sum_{\varphi \in \Phi} \int_{A} \left(\sum_{\psi \in \Psi} \psi \right) \varphi f = \sum_{\varphi \in \Phi} \sum_{\psi \in \Psi} \int_{A} \varphi \psi f \tag{64}$$

A priori, we don't know that we can interchange the order in the double sum of the last equal expression, but we could write down exactly the same stuff with the absolute value in it, which tells us we can. Interchanging them, we get the desired result.

(b) Let B be a big rectangle containing A, and assume $|f| \leq M$ on A. For any finite subset F of Φ ,

$$\sum_{\varphi \in F} \int_{A} \varphi |f| \le M \int_{A} \sum_{\varphi \in F} \varphi \le MV(B) \tag{65}$$

implying the desired result. We still have to show that in the situation that $\int_A f$ is defined according to our earlier definition, we get the same result. We'll do so next time.

Let's recap. What we've shown so far is that

- (a) The existence of the integral and it's value is independent of the parititon of unity chosen to compute it.
- (b) The integral always exists if A and f are both bounded.
- (c) We want to show that if we're in a situation where the integral exists according to our old definition, so that in particular A and f will be bounded, then our new definition agrees with the old one. That is, if moreover, A is Jordan measurable, then

$$\int_{A} f = \sum_{\varphi \in \Phi} \int_{A} \varphi f \tag{66}$$

where the left hand side is computed using the old definition

Proof. For every $\epsilon > 0$, there is a compact Jordan-measurable subset $C \subset A$ such that

$$V(A \setminus C) = \int_{A \setminus C} 1 = \int_{A} \chi_{A \setminus C} < \epsilon \tag{67}$$

So how do we get this C? We're given that A is Jordan measurable, so as A is bounded we can cover the boundary of A by the interiors of finitely many closed balls \mathcal{B} (or rectangles) of total volume $< \epsilon$. So we can take

$$C = A \setminus \text{Int} \bigcup \mathcal{B} \tag{68}$$

(we could also take the union of interiors). Why is C Jordan measurable? Since the boundary of C is a subset of the boundary of $Int \bigcup \mathcal{B}$, which is has measure zero, it follows that C has boundary of measure zero and is hence Jordan measurable.

Since C is compact, only finitely many of our partition functions $\varphi \in \Phi$ are nonzero on C. So we'll take a finite sum, look at the difference, and show that that difference goes to zero: consider any finite subset F of Φ including the ones which are nonzero on C. Consider

$$\left| \int_{A} f - \sum_{\varphi \in F} \int_{A} \varphi f \right| \tag{69}$$

As the absolute value of the integral is less than or equal to the integral of the absolute value, we have that equation 69 is less than or equal to

$$\int_{A} \left| f - \sum_{\varphi \in F} \varphi f \right| \tag{70}$$

Since f is bounded, we can choose M such that $|f| \leq M$. So equation 70 is less than or equal to

$$M \int_{A} \left(1 - \sum_{\varphi \in F} \varphi \right) = M \int_{A} \sum_{\varphi \in \Phi \setminus F} \varphi \tag{71}$$

Since

$$\forall \varphi \in \Phi \setminus F, \forall c \in C, \varphi(c) = 0 \tag{72}$$

we have that equation 71 is less than or equal to

$$MV(A \setminus C) < M\epsilon$$
 (73)

As I mentioned last time, you can do this in another way which perhaps looks more like the improper integral from first year calculus, and that's integrating with bigger and bigger sets. Let me state that, but I won't prove it, and rather put it on the problem set:

Theorem 4. Let A be open, $f: A \to \mathbb{R}$ be locally bounded and let the set of discontinuties of f have measure zero. Write

$$A = \bigcup_{n \in \mathbb{N}} C_n \tag{74}$$

where each C_n is compact and Jordan Measurble and satisfy $C_n \subset \operatorname{Int} C_{n+1}$. Part of the exercise is to show that we can do this. Then f is integrable on A if and only if

$$\left\{ \int_{C_i} |f| : n \in \mathbb{N} \right\} \tag{75}$$

is bounded. Note that if this is the case, the sequence

$$\left\{ \int_{C_i} f : n \in \mathbb{N} \right\} \tag{76}$$

converges absolutely, and we can write

$$\int_{A} f = \lim_{n \to \infty} \int_{C_n} f \tag{77}$$

Proof. Exercise.
$$\Box$$

It's good to appreciate that the improper integral which we defined using a partition of unity can be expressed in this way. It's also really good to do the above exercise as an exercise in using partitions of unity.

There's a little gap in the treatment of partitions of unity last time which I want to fill here. In the first case, where A was compact, we defined

$$\psi_1, \dots, \psi_p \tag{78}$$

such that $\psi_1 + ... \psi_p > 0$ on an open set $U \supseteq A$. We then let

$$\varphi_i = \frac{\psi_i}{\psi_1 + \dots \psi_p} \tag{79}$$

This is all well and good, but individual φ_i might be defined on sets that go outside U, and we have to worry about dividing by zero to guarantee that each bump function is \mathcal{C}^{∞} everywhere. To deal with this, we can simply multiply everything by another bump function (we can find one which is 1 on A and zero outside a compact set U).

Let f(x) be a \mathcal{C}^{∞} bump function such that f=1 on A, f=0 outside a compact subset of U. Then the set of functions

$$\{f \cdot \varphi_i : i \in 1, ..., p\} \tag{80}$$

is a partition of unity as required.

Before we finish, I want to say something about integration by substitution, or change of variables.

3.6 Change of Variables

Recall the single-variable integration by substitution formula

$$\int_{g(a)}^{g(b)} f = \int_{a}^{b} (f \circ g)g' \tag{81}$$

This holds whenever "both sides make sense", e.g. if f is continuous and g is continuously differentiable. The hypotheses, however, could be weaker. When we talk about integration by substitution, our g is going to be the change of variables. But change of variables means like "invertible". So not just any g, but in particular 1-to-1. I wanted to point out how, in terms of our formula, we should think about this in the case where g is one-to-one. So suppose it is so. Then equation 81 can be rewritten as

$$\int_{g([a,b])} f = \int_{[a,b]} (f \circ g)|g'| \tag{82}$$

since otherwise the equaton would not be correct if g was decreasing. So this is what our change-of-variables formula is going to look like. We'll look at that next time.

TODO: rest

3.7 Parametrically Defined Curves

TODO: this

4 Manifolds

4.1 What is a manifold?

TODO: this

4.2 Functions Between Manifolds

TODO: this

4.3 Manifolds with Boundary

TODO: this

4.4 Multilinear Algebra

TODO: this

4.5 Vector Fields and Differential Forms

TODO: this

4.6 The Differential Operator

TODO: this

5 Integration on Manifolds

TODO: this

5.1 Integration of Parametrized Curves

TODO: this

5.2 Integral of a k-form over a k-cube

TODO: this

5.3 Integration of Differential Forms on Manifolds

TODO: this

5.4 Manifolds with Boundary

TODO: this

5.5 Stoke's Theorem on Manifolds

TODO: this