Quantum Mechanics

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Dirac notation

Inner product of two vectors

$$\langle a|b\rangle = \langle b|a\rangle^*$$
 (1.1)
 $\langle a|a\rangle \ge 0$ (1.2)

$$a|a\rangle \ge 0 \tag{1.2}$$

Postulate X: time evolution 1.1

$$i\hbar\frac{\partial}{\partial t}\left|\Psi(t)\right\rangle = \hat{\mathbf{H}}(t)\left|\Psi(t)\right\rangle \tag{1.3}$$

Postulates

2.1 Postulate 1: state

Postulate 1 The state of a physical system is described by a state vector that belongs to a complex vector space V, called the state space of the system.

State space

$$|\Psi_1\rangle + |\Psi_2\rangle = |\Psi_3\rangle \tag{3.1}$$

$$|\Psi_1\rangle + |\Psi_2\rangle = |\Psi_2\rangle + |\Psi_1\rangle \tag{3.2}$$

Operators

4.1 Basic Properties

An operator acting on a ket creates a new ket:

$$\hat{A} |\Psi\rangle = |\Psi'\rangle \tag{4.1}$$

$$|\Psi\rangle\,, |\Psi'\rangle \in V$$
 (4.2)

Operators are linear

$$\hat{A}(a_1 | \Psi_1 \rangle + a_2 | \Psi_2 \rangle) = (a_1 \,\hat{A} | \Psi_1 \rangle + a_2 \,\hat{A} | \Psi_2 \rangle) \tag{4.3}$$

$$|\Psi\rangle, |\Psi'\rangle \in V; a_1, a_2 \in \mathbb{C}$$
 (4.4)

Operators are associative and commutative under addition

$$\hat{A} + (\hat{B} + \hat{C}) = (\hat{A} + \hat{B}) + \hat{C}$$
 (4.5)

$$\hat{A} + \hat{B} = \hat{B} + \hat{A} \tag{4.6}$$

Multiplying operators is interpreted as applying them to kets. It is associative but NOT (in general) commutative.

$$\hat{A}\,\hat{B}\,|\Psi\rangle = \hat{A}(\hat{B}\,|\Psi\rangle) = \hat{A}\,|\Psi'\rangle \tag{4.7}$$

$$\hat{A}(\hat{B}\,\hat{C}) = (\hat{A}\,\hat{B})\,\hat{C} \tag{4.8}$$

$$\hat{A}\,\hat{B} \neq \hat{B}\,\hat{A} \tag{4.9}$$

The lack of commutativeness makes the "commutator" useful

$$[\hat{A}, \hat{B}] = \hat{A} \hat{B} - \hat{B} \hat{A}$$
 (4.10)

(4.11)

+

4.2 Hermitian Operators

+ An operator is called Hermitian (or self-adjoint) if it is it's own conjugate $\hat{A}=\hat{A}^{\dagger}$ +

$$\hat{A} |A\rangle = |B\rangle \rightarrow \langle A| \hat{A}^{\dagger} = \langle B|$$
 (4.12)

$$\hat{A} |A\rangle = |B\rangle \rightarrow \langle A| \hat{A} = \langle B|$$
 (4.13)

(4.14)

A hermitian operator \hat{A} has the following properties

- 1. $\hat{A} |\lambda\rangle = \lambda |\lambda\rangle \rightarrow \lambda \in \mathbb{R}$
- 2. $\langle \hat{A} \rangle = \langle \Psi | \hat{A} | \Psi \rangle \in \mathbb{R}$
- 3. All eigenvectors with different eigenvalues are orthogonal

4.3 Projection Operators

A projection operator is defined

$$[\hat{A}, \hat{B}] = \hat{A} \,\hat{B} - \hat{B} \,\hat{A}$$
 (4.15)

(4.16)

Eigenvectors and Eigenvalues

5.1 Degenerate Eigenvectors

If two eigenvectors have the same eigenvalue:

$$\hat{H} |\lambda_1\rangle = \lambda |\lambda_1\rangle \tag{5.1}$$

$$\hat{H} |\lambda_2\rangle = \lambda |\lambda_2\rangle \tag{5.2}$$

their linear combination is an eigenvector as well:

$$\alpha \,\hat{\mathbf{H}} \,|\lambda_1\rangle = \lambda \alpha \,|\lambda_1\rangle \tag{5.3}$$

$$\beta \,\hat{\mathbf{H}} \,|\lambda_1\rangle = \lambda\beta \,|\lambda_1\rangle \tag{5.4}$$

$$\hat{\mathbf{H}}[\alpha |\lambda_1\rangle + \beta |\lambda_2\rangle] = \lambda[\alpha |\lambda_1\rangle + \beta |\lambda_2\rangle] \tag{5.5}$$

Therefore it is possible to create two orthogonal eigenvectors for this eigenvalue. The probability in the degenerate case is the sum of the probabilities for each eigenvector $|\langle \hat{\mathbf{H}}||\lambda_1\rangle|^2+|\langle \hat{\mathbf{H}}||\lambda_2\rangle|^2$