

$$-\frac{dn}{dt} = \beta(T_0) n_b n_a n_b + 3HT$$

$$n = rs = r \cdot \frac{2\pi^2 g_s}{45} T^3 \quad dn = dr \cdot s + \frac{3rsdT}{T}$$

$$-dt = \frac{dT}{HT} \quad \text{if } g_s = \text{const}$$

$$\frac{sdr + 3rs \frac{dT}{T}}{\frac{1}{HT} dr} = n_a \cdot n_b \cdot n_b \beta(T_0) + 3HT$$

$$sHT \frac{dr}{dT} + 3(rs) = n_a \cdot n_b^2 \beta(T_0) + 3HT$$

$$\frac{dr}{dT} = \frac{n_a n_b^2 s^2}{HT} \beta(T_0) \sim \frac{1}{T^5}$$

$$T_b = T_a \approx \frac{T^2}{T_{\text{eq}}}$$

$$T_{\text{eq}} = \frac{T_{\text{eq}}}{\sqrt{h} 2^{1/3}}$$

$$T_g = 2^{1/3} T$$

$$2l = \frac{g_{\text{rel}}(T)}{g_{\text{rel}}(T_{\text{rec}})}$$

$$n_a = n_b \Rightarrow r_a = r_b$$

$$\frac{dr}{dT} = \frac{r^2 s^2}{HT} \cdot \frac{25m}{M} e^{10} \cdot \frac{1}{T_a^{3/2}} (1 - \cos \theta)$$

$$\frac{dr}{r^3} = \frac{s^2}{H} \cdot e^{10} \frac{25m}{M} (1 - \cos \theta) \cdot \frac{1}{T} \cdot \frac{T_{\text{eq}}^{3/2}}{T^3}$$

$$\text{RD: } \frac{dr}{r^3} = \frac{\left(\frac{2\pi^2 g_s}{45}\right)^2 T^6}{h_{\text{RD}} m_{\text{Pl}}^2} e^{10} \frac{25m}{M} (1 - \cos \theta) \frac{T_{\text{eq}}^{3/2}}{T^{4+6}}$$

$$d\left(\frac{1}{r^2}\right) = \frac{\left(\frac{2\pi^2 g_s}{45}\right)^2 e^{10} \frac{25m}{M} (1 - \cos \theta) T_{\text{eq}}^{3/2}}{h_{\text{RD}} m_{\text{Pl}}^2} \frac{1}{T^5}$$

$$\frac{1}{2r^2} = \frac{D_{\text{RD}}}{5T^3} + C$$

$$C = \frac{1}{2r_{\text{rec}}^2} - \frac{D_{\text{RD}}}{5T_{\text{rec}}^3}$$

$$\frac{1}{r^2(T)} = \frac{1}{r_{\text{rec}}^2} + \frac{2D_{\text{RD}}}{3} \left(\frac{1}{T^3} - \frac{1}{T_{\text{rec}}^3} \right)$$

$$r^2(T) = \frac{1 + \frac{2D_{\text{RD}}}{3} \left(\frac{1}{T^3} - \frac{1}{T_{\text{rec}}^3} \right) r_{\text{rec}}^2}{1}$$

$$r(T)_{\text{RD}} = r_{\text{rec}} \sqrt{\frac{1}{1 + \frac{2D_{\text{RD}}}{3} \left(\frac{1}{T^3} - \frac{1}{T_{\text{rec}}^3} \right)}}$$

$$n = rs = r \frac{2\pi^2 g_s}{45} T^3,$$

$$g_s = \Sigma_{\text{bos}} (T_{\text{bos}}/T)^3 + 7/8 \Sigma_{\text{ferm}} (T_{\text{ferm}}/T)^3, \quad (14)$$

where sums are over sorts of ultrarelativistic bosons and fermions, T is the temperature of photons. Hubble constant on radiation and matter dominated (RD and MD) stages can be represented as following

$$H = \begin{cases} \sqrt{\frac{4\pi^3 g_s}{45}} \frac{T^2}{m_{\text{Pl}}} = h_{\text{RD}} \frac{T^2}{m_{\text{Pl}}} & \text{at RD stage,} \\ \frac{2}{3t_{\text{mod}}} \left(\frac{T}{2.7\text{K}} \right)^{3/2} = h_{\text{MD}} \frac{T^{3/2}}{m_{\text{Pl}}^{1/2}} & \text{at MD stage.} \end{cases} \quad (15)$$

Here m_{Pl} is the Plank mass, $t_{\text{mod}} \approx 4.4 \cdot 10^{17}$ s is the modern age of Universe, $h_{\text{RD}} \approx 5.5 \sqrt{g_s/11}$, $h_{\text{MD}} \approx 3.08 \cdot 10^{-14}$ and analogously to g_s of Eq.(14)

$$g_s = \Sigma_{\text{bos}} (T_{\text{bos}}/T)^4 + 7/8 \Sigma_{\text{ferm}} (T_{\text{ferm}}/T)^4. \quad (16)$$

Value $T = T_{\text{RM}} \approx 1.2$ eV provides good transition between given dependences of H on RD- and MD-stages, connected them with factor $\sqrt{\frac{T}{T_{\text{RM}}}}$. For time intervals when $g_{s,s} \approx \text{const}$ it is just the following transition between differentials

$$\frac{1}{T^3} = \frac{1}{T_{\text{rec}}^3} + \dots$$

$$r(T_{\text{rec}}) = r_{\text{rec}}$$

$$\frac{1}{2} \left(\frac{1}{r^2} - \frac{1}{r_{\text{rec}}^2} \right) = \frac{D_{\text{RD}}}{5} \left(\frac{1}{T^3} - \frac{1}{T_{\text{rec}}^3} \right)$$