

Drift current: An electric field has the net movement of charge due to an electric field is called drift. The net drift of charge gives rise to a drift current.

$$\textcircled{1} - \frac{q}{m} = qE = qb^2(92) = pb^2$$

Drift current density:

If we have a positive volume density of moving at an average drift velocity per unit area

$$J_{\text{drift}} = (p^2 d) \text{ (area)} = ab^2 q = pb^2$$

If the volume charge density is due to positively charged holes, then

$$J_{p \text{ drift}} = (e P) v_{dp}$$

v_{dp} = Average

resultant drift velocity of holes

The equation of motion of a positively charged hole into the presence of an electric field is nonzero

m_p^* = effective mass

$F = m_p^* a$ and $E = \frac{F}{q}$

$a = \text{magnitude of acceleration}$

$t = \text{time taken to move from left to right}$

$d = \text{distance moved by hole}$



Average drift velocity is \propto to electric field

$$v_{dp} = \mu_p E \quad \text{where mobility}$$

is proportional to $\frac{1}{m_p}$ factor

$$J_{pdif} = (\mu_p) v_{dp} = e \mu_p p E \quad \text{---(1)}$$

The same discussion of drift applies to electrons

$$J_{ndif} = p v_{dn} = (-e n) v_{dn} \\ = (-e n) (-\mu_n E) = p n \mu_n E$$

$$\text{at slab is } J_{ndif} = e \mu_n n E \quad | \quad v_{dn} = -\mu_n E$$

$$J_{drift} = (e) \text{ current up } E$$

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$$q = q_0 A$$

Corriess diffusion

Diffusion is the process whereby particles move from a region of high concentration toward a region of low concentration

$$J_{diffusion} = \frac{q}{t} \frac{dN}{dx}$$

If the gas molecules were electrically charged, the net flow of charge would result in a diffusion current.

Diffusion current Density :

The net rate of electron from F_n .

$$F_n = \frac{1}{2} n(-e) v_{th} - \frac{1}{2} n(+e) v_{th} = \frac{1}{2} v_{th} [n(-e) - n(+e)]$$

$$F_n = \frac{1}{2} n(-e) v_{th} - \frac{1}{2} n(+e) v_{th} \quad \text{--- (1)}$$

If we expand equation (1) into Taylor series.

$$F_n = \frac{1}{2} v_{th} \left\{ \left[n(0) - e \frac{dn}{dx} \right] - \left[n(0) + e \frac{dn}{dx} \right] \right\}$$

$$= -v_{th} e \frac{dn}{dx}$$

(each electron has charge $-e$) So the current is

$$J = -e F_n = +e v_{th} e \frac{dn}{dx}$$

$$J_{ndrift} = e D_n \frac{dn}{dx}, D_n \text{ is called electron diffusion coefficient}$$

$$J_{pdif} = -e D_p \frac{dp}{dx}, D_p \text{ hole diffusion coefficient}$$

[5.4]

$T = 300 \text{ K}$, electron concentration varies

linearly from $1 \times 10^{18} \text{ cm}^{-3}$ over

distance of 0.10 cm to $7 \times 10^{17} \text{ cm}^{-3}$ over

$$D_n = 225 \text{ cm}^2/\text{s}$$

Set:

+ (a) The diffusion current density is given by

$$J_{\text{diff}} = e D_n \frac{dn}{dx_b}$$

$$\approx e D_n \frac{\Delta n}{l}$$

$$= 1.6 \times 10^{-19} \times 225 \times \frac{(6 \times 10^{18} - 7 \times 10^{17})}{0.10}$$

$$= 108 \text{ A/cm}^2 = E$$

better if $\frac{nb}{xb} \propto l$ = first

$\frac{nb}{xb} q b l = \text{first } E$

$nb \propto D_n, q b$

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→ Total current density

→ Induced electric field

→ example 5.5

Total current density

We now have four possible independent current mechanisms in a semiconductor. These components are electrons or electron drift and diffusion currents and holes drift and diffusion currents.

The total current density is the sum of these four components, or for the one-dimensional case.

$$J_{cb} = e \mu n E_x + e \mu p E_x + e D_n \frac{dn}{dx} - e D_p \frac{dp}{dx}$$

This equation may be generalized to three dimensions.

$$J = e \mu n E + e \mu p E + e D_n \nabla n - e D_p \nabla p$$

The electron mobility gives an indication of how well an electron moves in a semiconductor as a result of the force of an electric field.

Induced electric field

The separation of positive and negative charges induces an electric field that is in a direction opposite to the diffusion process.

The electric potential ϕ is related to potential energy by the charge (-e). So

$$\phi = +\frac{1}{e} (E_F - E_{F_i})$$

The electric field for the one-dimensional situation is defined as

$$E_x = - \frac{d\phi}{dx} = - \frac{1}{e} \frac{dE_{F_i}}{dx}$$

$$n_i = n_i \exp \left[\frac{E_F - E_{F_i}}{kT} \right] \approx N_d(x)$$

$$E_F - E_{F_i} = kT \ln \frac{N_d(x)}{n_i}$$

$$\Rightarrow - \frac{dE_{F_i}}{dx} = \frac{kT}{N_d(x)} \frac{dN_d(x)}{dx}$$

$$E_x = - \left(\frac{kT}{e} \right) \frac{1}{N_d(x)} \frac{dN_d(x)}{dx}$$

Since we have an electric field, there will be a potential difference through the semiconductor due to the nonuniform doping.

$$V_{DD} = T$$

$$x \rho_{01} - \delta_{01} = (x) b N$$

$$m_0 V_{DD} x \rho_{01} = T$$

so we can write $\frac{V_{DD}}{x \rho_{01}} = \frac{T}{m_0}$

$$\text{we can also write } \frac{\delta_{01}}{x \rho_{01}} = \frac{(x) b N}{x b} = N$$

the carrier density is

$$\text{then we can write } \frac{\delta_{01}}{x \rho_{01}} = \frac{(x) b N}{x b} = \left(\frac{T}{m_0} \right) = N$$

$$\text{So we get } \frac{\delta_{01}}{x \rho_{01}} = \frac{N}{(x) b N} = \frac{1}{b}$$

If we assume

$$\frac{\delta_{01}}{x \rho_{01}} = \frac{1}{b}$$

$$T = \frac{1}{b}$$

$$b = 2 \times 10^3$$

$$T = 2 \times 10^3 \text{ K}$$

$$m_0 V_{DD} = x T$$

5-6

Electric field $E_x = ?$

$$T = 300 \text{ K}$$

$$N_d(x) = 10^{16} - 10^9 x$$

$$x = 0 \leq x \leq 1 \text{ cm}$$

Solⁿ: Taking the derivative of the donor concentration, we have

$$\frac{dN_d(x)}{dx} = -10^{19}$$

$$E_x = -\left(\frac{kT}{e}\right) \frac{1}{N_d(x)} \frac{dN_d(x)}{dx}$$

$$= -\frac{(0.0259)(-10^{19})}{10^{16} - 10^9 x}$$

at $x = 0$

$$E_x = 25.9 \text{ V/cm}$$

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Notes

- * The Einstein Relation
- * Math
- * Carrier generation and Recombination
- * The semiconductor in equilibrium
- * Excess carrier generation and Recombination
- * Math

The Einstein Relation

We assume there are no electrical connections that the semiconductor is in thermal equilibrium, so the individual electron and hole currents must be zero.

$$J_n = e n \frac{dn}{dx} = 0$$

$$J_n = e n \frac{dn}{dx} = e D_n \frac{dN_d(x)}{dx}$$

If we assume quasi-neutrality so that $n \approx N_d(x)$

$$J_n = 0 = e v_n N_d(x) E_x + e D_n \frac{dN_d(x)}{dx}$$

Putting the value of E_x =

$$0 = -e v_n N_d(x) \frac{kT}{e} \cdot \frac{1}{N_d(x)} \cdot \frac{dN_d(x)}{dx} + e D_n \frac{dN_d(x)}{dx}$$

-①

$$\text{Equation ① valid for } \frac{D_n}{v_n} = \frac{kT}{e} - ①$$

The hole current must also be zero in semiconductor

$$\frac{D_p}{v_p} = \frac{kT}{e} - ②$$

combining ① and ②

$$\frac{D_n}{v_n} = \frac{D_p}{v_p} = \frac{kT}{e}$$

this is the Einstein relation

5.6

mobility of carrier is $1600 \text{ cm}^2/\text{V.s}$

at $T = 300 \text{ K}$

Determine diffusion coefficient

$\zeta \pi$:

$$\frac{(x)bHb}{x^b} \propto \frac{kT}{e} \Rightarrow D = \frac{kT}{e} \cdot \frac{(x)bHb}{x^b} e^{kT} = 0 \text{ at } T = 0$$

$$= 0.0259 \times 1600 \text{ cm}^2/\text{s}$$

$$\frac{(x)bHb}{x^b} = 25.9 \text{ cm}^2/\text{s}$$

Chapter six (6)

* Carrier generation and Recombination

Formation of the electron-hole pairs

These pairs can be generated by
To move into the conduction band or valence
band. This is called carrier generation.
Or the electron hole pair can be formed
destroyed by recombination.

Recombination of the electron-hole pairs

When the electron-hole pair recombines,
the energy released is called recombination
energy. This energy is released in the form
of heat. This heat is called recombination
heat. The recombination heat is
also called Auger heat.

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Chapter - 7

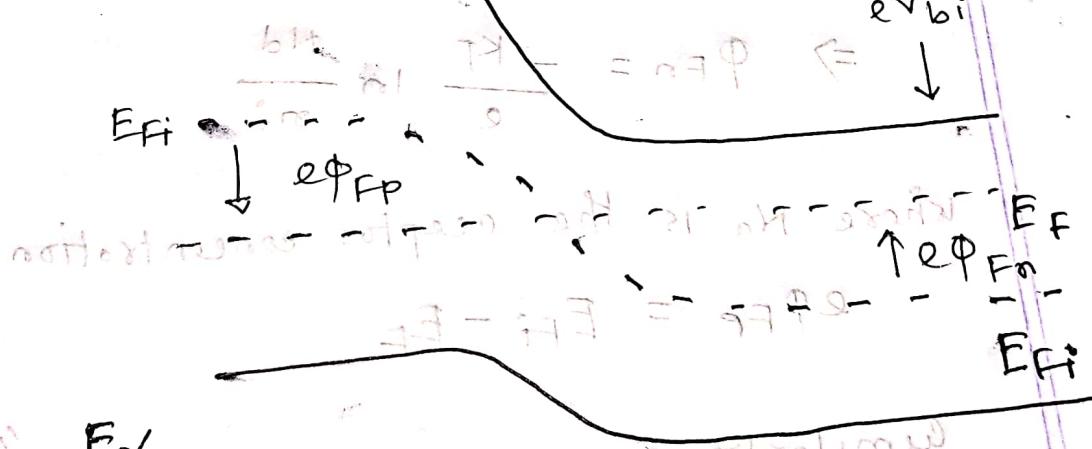
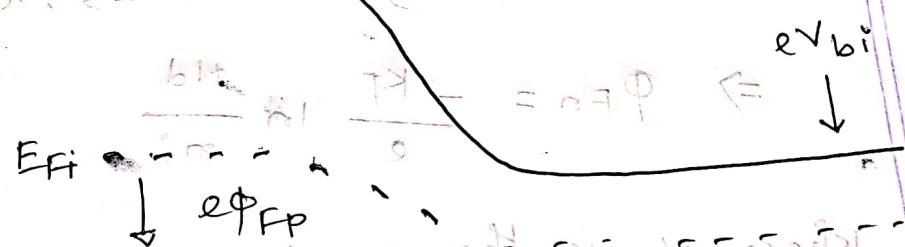
- Basic structure of the P-n Junction
- Built-in potential barrier
- Math 7.1

$$\left[\frac{E_F - E_V}{kT} \right]_{P} = 10^9 \cdot 10^{-10} = 10^8$$

Electrons in the conduction band of the n region see a potential barrier in trying to move into the conduction band of the p region. This potential barrier is referred to as the built-in potential barrier and is denoted by V_{bi} -

$$\left(\frac{E_F - E_V}{kT} \right)_{P} = 10^9 \cdot 10^{-10} = 10^8$$

$bH = 0.2263$ Mod. of formation point



Now we can write $\frac{eH}{kT} V_{bi} = |P_{Fn}| f |Q_{Fn}| - ①$

in then region, the electron concentration in the conduction band is given by

$$n_0 = N_i \exp \left[-\frac{(E_e - E_F)}{KT} \right]$$

which can also be written in the form

$$n_0 = n_i \exp \left[\frac{E_F - E_{F_i}}{KT} \right]$$

at the potential Φ_{Fn} in then region as

$$\text{at } \Phi_{Fn} = E_{F_i} - E_F \text{ and so}$$

$$-\frac{e\Phi_{Fn}}{KT} \text{ before}$$

$$n_0 = n_i \exp \left(\frac{-e\Phi_{Fn}}{KT} \right)$$

Taking natural log both sides, $n_0 = N_d$

$$\Rightarrow \Phi_{Fn} = -\frac{KT}{e} \ln \frac{N_d}{n_i}$$

where N_d is the acceptor concentration

$$-e\Phi_{Fp} = E_{F_i} - E_F$$

Similarly

$$\Phi_{Fp} = +\frac{KT}{e} \ln \frac{N_a}{n_i}$$

$V_{6.1''}$

$\frac{KT}{e} \ln \frac{N_a}{n_i}$

$\frac{KT}{e} \ln \frac{N_a}{n_i}$

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- * Reverse bias
- * Space charge width
- * Math 7.1

[7.1] To calculate the built-in potential in a
pn Junction

$$T = 300\text{K}, N_a = 1 \times 10^{38} \text{ cm}^{-3}, N_d = 1 \times 10^{15} \text{ cm}^{-3}$$

$$n_i = 1.5 \times 10^{10} \text{ cm}^{-3}$$

$$\begin{aligned} \text{Sol: } V_{bi} &= \sqrt{T} \ln \frac{N_a N_d}{n_i^2} \\ &= 0.0259 \ln \frac{10^{18} \times 10^{15}}{(1.5 \times 10^{10})^2} \\ &= .754 \text{ V} \end{aligned}$$

* Space charge width: