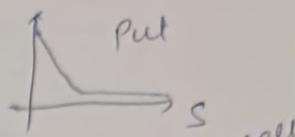
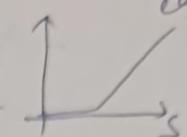


Computational Finance

$$V(\delta S, t) \quad 0 \leq S < \infty$$

$$\frac{\partial V}{\partial t} + \frac{\sigma^2 S^2}{2} \frac{\partial^2 V}{\partial S^2} + (\kappa - \delta) S \frac{\partial V}{\partial S} - rV = 0$$


$$V(S, T) = \begin{cases} (K-S)^+ \\ (S-K)^+ \end{cases}$$


$$V(0, t) = \dots$$

$$V(S_{\text{man}}, t) = \dots$$

1) Finite Domain \rightarrow eqn. more complicated.

2) Infinite domain \rightarrow eqn. simple.

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \quad -\infty < S \leq \infty$$

$$\frac{U_m^{n+1} - U_m^n}{\Delta t} = \frac{U_{m+1}^n - 2U_m^n + U_{m-1}^n}{\Delta x^2} \quad \text{FTCS}$$

$S_{\min} < S \leq S_{\max}$

$$\frac{U_m^n - U_m^{n-1}}{\Delta t} = \frac{U_{m+1}^n - 2U_m^n + U_{m-1}^n}{2\Delta x} \quad \text{BTCS}$$

$$\frac{U_{m+1}^{n+1} - U_m^{n-1}}{\Delta t} = \frac{U_{m+1}^n - 2U_m^n + U_{m-1}^n}{2\Delta x} \quad \text{CTCS}$$

$$S \in [a, b] \quad N = \frac{b-a}{\Delta x} \quad \Delta x = \frac{b-a}{N} = \Delta S$$

$$t \in [0, T] \quad M, \quad \Delta t = \frac{T}{M} = \Delta t$$

$$V(S, t)$$

$$V_m^n \approx V_m(S_m, t_n)$$

$$\text{FTCS} \quad \frac{U_m^{n+1} - U_m^n}{\Delta t} = \frac{U_{m+1}^n - 2U_m^n + U_{m-1}^n}{\Delta x^2}$$

$$\Rightarrow U_m^{n+1} = \alpha U_{m+1}^n + (1-2\alpha) U_m^n + \alpha U_{m-1}^n$$

$$0 < \alpha < 1/2$$

$$BTCS \quad \frac{U_m^{n+1} - U_m^n}{k} = \frac{U_{m+1}^{n+1} - 2U_m^{n+1} + U_{m-1}^{n+1}}{h^2}$$

$$-\lambda U_{m+1}^{n+1} + ((1+2\lambda)U_m^{n+1}) + \lambda U_{m-1}^{n+1} = U_m^n$$

$$AV = F$$

$$U = A^{-1}F$$

$$\text{Crank Nicolson} \quad \frac{1}{2}(PTCS + BTCS)$$

$$\frac{U_m^{n+1} - U_m^n}{k} = \frac{U_{m+1}^{n+1} - 2U_m^{n+1} + U_{m-1}^{n+1} + \lambda U_{m+1}^n - 2U_m^n + U_{m-1}^n}{2h^2}$$

$$\Rightarrow U_{m+1}^{n+1}$$

$$\boxed{\frac{\lambda}{2}(U_{m+1}^n - \lambda U_m^n + \frac{\lambda}{2} U_{m-1}^n) = (1+\lambda)U_m^{n+1} - \lambda U_{m+1}^{n+1} + \frac{\lambda}{2} U_{m-1}^{n+1}}$$

$$\gamma^2 y'' + \gamma y' + y = 0$$

$$\begin{aligned} \frac{dy}{dt} &= t \\ y &= \gamma \log t \end{aligned}$$

$$\frac{dy}{dt} = \log t + \gamma + \frac{1}{t} \frac{dt}{dt}$$

$$\frac{d^2y}{dt^2} = \frac{1}{t} \frac{dt}{dt} + \frac{t}{t} \frac{dt}{dt} + \gamma \left(-\frac{1}{t^2} \right)$$

$\therefore V_t, V_S, V_{SS}$

for BSM. Transformation - ①

$$\begin{cases} Y = \ln S \rightarrow -\infty < Y < \infty \\ T = T-t \\ V(S, T) = e^{-\gamma(T-t)} v(y, z) \end{cases}$$

$$\frac{\partial V}{\partial t} + \frac{\sigma^2}{2} S^2 \frac{\partial^2 V}{\partial S^2} + (\bar{r} - \delta) S \frac{\partial V}{\partial S} - \gamma V = 0$$

$$V_t = \gamma e^{-\gamma(T-t)} v(y, z) - e^{-\gamma(T-t)} v_T(y, z)$$

$$V_S = e^{-\gamma(T-t)} v_y(y, z) \frac{1}{S}$$

$$v_{yy} = e^{-\alpha(T-t)} v_y(y, \tau) \left(-\frac{1}{S^2} \right) + e^{-\alpha(T-t)} v_{yy}(y, \tau) \left(\frac{1}{S^2} \right)$$

$$S^2 v_{yy} = e^{-\alpha(T-t)} (v_{yy}(y, \tau) - v_y(y, \tau))$$

$$S v_y = e^{-\alpha(T-t)} v_y(y, \tau)$$

$$v_t = e^{-\alpha(T-t)} (\alpha v - v_z)$$

$$e^{-\alpha(T-t)} (v_F - v_E) + \frac{\sigma^2}{2} (e^{-\alpha(T-t)} (v_{yy} - v_y))$$

$$+ (\alpha - \delta) (e^{-\alpha(T-t)} v_y) - \alpha e^{-\alpha(T-t)} v = 0$$

$$-v_F + \frac{\sigma^2}{2} (v_{yy} - v_y) + (\alpha - \delta) v_y = 0$$

$$\frac{\sigma^2}{2} v_{yy} + v_y (\alpha - \delta - \frac{\sigma^2}{2}) - v_F = 0 \quad \left| \quad \frac{\partial^2 u}{\partial n^2} = \frac{\partial u}{\partial t} \right.$$

$$v(y, 0) = V_T(e^y)$$

Transformation 2.

$$n = y + \left(\alpha - \delta - \frac{\sigma^2}{2} \right) t$$

$$\tilde{z} = \frac{1}{2} \sigma^2 t$$

$$v(y, t) = u(n, \tilde{z})$$

$$v_y = u_n \psi_n(n, \tilde{z})$$

$$v_{yy} = \psi_{nn}(n, \tilde{z})$$

$$v_z = \frac{\sigma^2}{2} u_{\tilde{z}} + (\alpha - \delta - \frac{\sigma^2}{2}) u_n$$

$$\frac{\sigma^2}{2} u_{nn} + u_n (\alpha - \delta - \frac{\sigma^2}{2}) - \frac{\sigma^2}{2} u_{\tilde{z}} - (\alpha - \delta - \frac{\sigma^2}{2}) u_n = 0$$

$$\boxed{\psi_{nn} = u_{\tilde{z}}}$$

$$-\infty < n < \infty$$

↓ restrict

$$n_{\min} < n < n_{\max}$$

$$\boxed{u(n, 0) = V_T(e^y)}$$

$$S(t) = S_0 e^{\alpha t}$$

$$V(S, t) = e^{-\alpha(T-t)} V(y, t) = e^{-\alpha(T-t)} u(\eta, \tilde{t})$$

Remark: Both the transformation can be combined

$$y = \ln S$$

$$\tilde{t} = (T-t)$$

$$\eta = \ln S + \left(\alpha - \frac{\sigma^2}{2} \right) \tilde{t}$$

$$\tilde{t} = \frac{1}{2} \sigma^2 (T-t)$$

$$V(S, t) = e^{-\alpha(T-t)} u(\eta, \tilde{t})$$

Note: Here we have assumed $k=1$ (strike price)

Remarks

- ① Since the value of the option is given at final time $t=T$ in order to make the initial time $\tilde{t}=T-t$
- ② $y = \ln S$, $S = e^y$ is used to convert the variable coeff. to constant coefficient.
- ③ $V = e^{-\alpha(T-t)} V(y, \tilde{t})$ to eliminate the integrating factor.

$$\frac{dV}{d\tilde{t}} - \alpha V = f.$$

$$\text{I.f. } = e^{-\alpha \tilde{t}}$$

If $\alpha(t)$ is function of t

$$\text{then I.f. } e^{-\int_0^t \alpha(s) ds}.$$

Suppose the $\alpha(t), S(t), \sigma(t)$ are function of t then the transformation will be

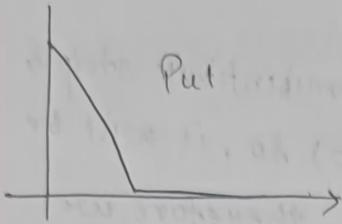
$$\eta = \ln S + \int_t^T \left(\alpha(s) - \delta(s) - \frac{\sigma^2(s)}{2} \right) ds$$

$$\tilde{t} = \frac{1}{2} \int_t^T \sigma^2(s) ds$$

$$V(S, t) = e^{-\int_t^T \alpha(s) ds} u(\eta, \tilde{t})$$

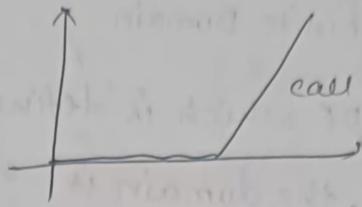
$$\frac{\partial V}{\partial t} + \frac{\sigma^2}{2} S^2 \frac{\partial^2 V}{\partial S^2} + (\mu - r) S \frac{\partial V}{\partial S} - rV = 0$$

$$V(S, T) = (S - K)^+$$



$$V_p(S, t) = 0, S \rightarrow \infty$$

$$V_p(S, t) = \quad S \rightarrow 0$$



$$V_c(S, t) = 0, S = 0$$

$$V_c(S, t) = \quad S \rightarrow \infty$$

$$V_p + V_c = e^{-rt} \left[S - \frac{V_0(2(r-\lambda))}{2\sigma} + \frac{V_0(2(r-\lambda))}{2\sigma} e^{-r(T-t)} \right] C - P = S - Ke^{-rT}$$

$$V_c(S, t) = S - Ke^{-r(T-t)} \quad S \rightarrow \infty$$

$$V_p(S, t) = Ke^{-r(T-t)} - S \quad S \approx 0, V = (7, 2) V$$

$$S = Ke^n, \quad t = T - 2\tau$$

$$q = \frac{2\tau}{6^2} (3-1) 2 q_8 = 12 \frac{(r-\lambda)}{6^2}$$

$$V(S, t) = V\left(Ke^n, T - \frac{2\tau}{6^2}\right) = v(n, \tau)$$

$$v(n, \tau) = K e^{\eta \tau} \left\{ \exp \left\{ -\frac{1}{2} (q_8 - 1) n - \frac{1}{2} (q_8 - 1)^2 \tau \right\} \right\} y(n, \tau)$$

$$V(S, T) = (S - K)^+ = \max(S - K, 0)$$

$$V(S, T) = V_T(S)$$

$$y(n, 0) = y(n, 0) = \begin{cases} \text{Call} \\ \text{Put} \end{cases}$$

$$\frac{\partial^2 y}{\partial x^2} = \frac{\partial y}{\partial x}$$

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \hookrightarrow \text{get Black Scholes}$$

$V(S, t)$
 $S \in (-\infty, \infty).$

Transform - II Finite Domain

Black Scholes PDE which is defined in semi-infinite strip & transformed to, the domain is $S \in (-\infty, \infty)$ so, it will be difficult to deal with unbounded domain, therefore we transform Black Scholes PDE in domain $[0, 1]$ with appropriate transformation.

$$\frac{\partial V}{\partial t} + \frac{\sigma^2 S^2}{2} \frac{\partial^2 V}{\partial S^2} + (\mu - \delta) S \frac{\partial V}{\partial S} - \alpha V = 0$$

$0 \leq S < \infty, t \leq T$

$$\textcircled{1} \quad V(S, T) = V_T(S) = (S - K)^+$$

$$= (K - S)^+$$

Consider following transformation

$$\textcircled{2} \quad \begin{aligned} \xi &= \frac{S}{S+P} & P \neq 0, P > 0, (S+P)\xi &= S \\ P &= T-t & P\xi &= S(1-\xi) \\ V(S, t) &\leftarrow V(\xi, P). \end{aligned}$$

$$S = \frac{P\xi}{1-\xi}$$

$$\begin{aligned} S+P &= \frac{P\xi + P\xi + P}{1-\xi} \\ &= \frac{P}{1-\xi} \end{aligned}$$

$$\begin{aligned} \frac{\partial V}{\partial S} &= \frac{\partial \xi}{\partial S} = \frac{S+P-S}{(S+P)^2} = \frac{P}{(S+P)^2} \\ &\approx \frac{P(1-\xi)}{P^2} \\ &\approx \frac{(1-\xi)^2}{P} \end{aligned}$$

$$V(S)$$

$$\frac{\partial V}{\partial t} = \frac{\partial}{\partial t} [(S+P) \otimes \bar{V}(\varepsilon_e, \tau)] = -(S+P) \frac{\partial \bar{V}}{\partial \tau}$$

$$= -\left(\frac{P}{1-\varepsilon_e}\right) \frac{\partial \bar{V}}{\partial \tau}$$

$$\frac{\partial V}{\partial S} = \frac{\partial}{\partial S} [(S+P) \otimes \bar{V}(\varepsilon_e, \tau)]$$

$$= \bar{V}(\varepsilon_e, \tau) + (S+P) \frac{\partial \bar{V}}{\partial \varepsilon_e} \cdot \frac{\partial \varepsilon_e}{\partial S}$$

$$= \bar{V}(\varepsilon_e, \tau) + (S+P) \cdot \frac{\partial \bar{V}}{\partial \varepsilon_e} \cdot \frac{P}{(S+P)^2}$$

New domain:

$$0 \leq S < \infty$$

$$\varepsilon_e = \frac{S}{S+P}$$

$$0 \leq \varepsilon_e \leq 1.$$

$$t \leq \tau$$

$$\tau \geq 0$$

$$\frac{\partial V}{\partial S} = \bar{V}(\varepsilon_e, \tau) + \frac{\partial \bar{V}}{\partial \varepsilon_e} \cdot (1-\varepsilon_e)$$

$$\frac{\partial^2 V}{\partial S^2} = \cancel{\frac{\partial \bar{V}}{\partial \varepsilon_e} \cdot \frac{\partial \varepsilon_e}{\partial S}} + \frac{\partial^2 \bar{V}}{\partial \varepsilon_e^2} \cdot \frac{\partial \varepsilon_e}{\partial S} (1-\varepsilon_e) + \frac{\partial \bar{V}}{\partial \varepsilon_e} \cdot \cancel{\left(-\frac{\partial \varepsilon_e}{\partial S}\right)}$$

$$= \cancel{\frac{\partial \bar{V}}{\partial \varepsilon_e} \times \left(\frac{(1-\varepsilon_e)^2}{P}\right)} + \frac{\partial^2 \bar{V}}{\partial \varepsilon_e^2} \cdot \frac{(1-\varepsilon_e)^3}{P} +$$

$$\frac{\partial^2 V}{\partial S^2} = \frac{\partial^2 \bar{V}}{\partial \varepsilon_e^2} \left(\frac{(1-\varepsilon_e)^3}{P}\right).$$

$$\left(\frac{P}{1-\varepsilon_e}\right) \frac{\partial \bar{V}}{\partial \tau} + \cancel{\sigma \left(\frac{P\varepsilon_e}{(1-\varepsilon_e)}\right)} \frac{P\varepsilon_e^2}{(1-\varepsilon_e)^2} \cdot \frac{\partial^2 \bar{V}}{\partial \varepsilon_e^2} \times \left(\frac{(1-\varepsilon_e)^2}{P}\right)$$

$$+ (\lambda - \sigma) \left(\frac{P\varepsilon_e}{1-\varepsilon_e}\right) \left(\bar{V} + \frac{\partial \bar{V}}{\partial \varepsilon_e} (1-\varepsilon_e)\right)$$

$$- \lambda \left(\frac{P}{1-\varepsilon_e}\right) \bar{V} = 0$$

$$-\left(\frac{P}{1-\varepsilon_e}\right) \frac{\partial \bar{V}}{\partial \tau} + \frac{\sigma^2 (\varepsilon_e) P \varepsilon_e^2 (1-\varepsilon_e) \frac{\partial^2 \bar{V}}{\partial \varepsilon_e^2}}{2} + (\lambda - \sigma) P \varepsilon_e \frac{\partial \bar{V}}{\partial \varepsilon_e} + \frac{P \bar{V}}{(1-\varepsilon_e)} ((\lambda - \sigma) \varepsilon_e - \lambda) = 0$$

$$\frac{\partial \bar{V}}{\partial \bar{z}} = \frac{\sigma^2(\varepsilon_e)}{2} \frac{\varepsilon_e^2(1-\varepsilon_e)^2}{\partial \varepsilon_e^2} \frac{\partial^2 \bar{V}}{\partial \varepsilon_e^2} + (4-8) \varepsilon_e(1-\varepsilon_e) \frac{\partial \bar{V}}{\partial \varepsilon_e} \quad \rightarrow (4)$$

$$- [4(1-\varepsilon_e) + 8\varepsilon_e] \bar{V} \quad 0 \leq \varepsilon_e \leq 1$$

$$V(S, t) = (S+P) \bar{V}(\varepsilon_e, \bar{t})$$

$$\bar{t} = T$$

$$V(S, T) = (S+P) \bar{V}(\varepsilon_e, 0) \quad P=0$$

$$V_T(S) = \frac{P}{(1-\varepsilon_e)} \bar{V}(\varepsilon_e, 0)$$

$$\bar{V}(\varepsilon_e, 0) = \frac{(1-\varepsilon_e)}{P} V_T(S)$$

$$\boxed{\bar{V}(\varepsilon_e, 0) = \frac{P(1-\varepsilon_e)}{P} V_T \left(\frac{P\varepsilon_e}{1-\varepsilon_e} \right)} \quad \text{Initial}$$

$$\varepsilon_e = 0$$

$$\varepsilon_e = 1$$

$$\frac{\partial \bar{V}}{\partial \bar{z}} = -8\bar{V}$$

$$\frac{\partial \bar{V}}{\partial \bar{z}} = -\frac{8\bar{V}}{9} \frac{\bar{V}''}{\bar{V}'} \quad \frac{\bar{V}''}{\bar{V}'} = \frac{1}{9}$$

In order to solve eq.(4) we require the boundary conditions.
at $\varepsilon_e=0$ & $\varepsilon_e=1$. when $\varepsilon_e=0$ & $\varepsilon_e=1$ given in (4).

degenerates to ODE:

when
 $\varepsilon_e=0$

$$\frac{\partial \bar{V}}{\partial \bar{z}} = -8\bar{V}$$

$$\frac{\partial \bar{V}(0, \bar{t})}{\partial \bar{z}} = -8\bar{V}(0, \bar{t})$$

when $\varepsilon_e=1$:

$$\frac{\partial \bar{V}}{\partial \bar{z}} = -8\bar{V}$$

$$\boxed{\bar{V}(1, \bar{t}) = \bar{V}(1, 0) e^{-8\bar{t}}}$$

Boundary

$$\boxed{\bar{V}(0, \bar{t}) = \bar{V}(0, 0) e^{-8\bar{t}}}$$

$$\bar{V}(\varepsilon_e, 0) = \frac{(1-\varepsilon_e)}{P} V_T \left(\frac{P\varepsilon_e}{1-\varepsilon_e} \right)$$

call

$$V(S, T) = \max(S - k, 0) = \max\{S - k, 0\} = V_T(S)$$

Put

$$V(S, T) = (k - S)^+ = \max\{k - S, 0\} = V_T(S)$$

consider call.

$$\bar{V}(\varepsilon_e, 0) = \frac{(1-\varepsilon_e)}{P} \max \left\{ \frac{P\varepsilon_e - k + k\varepsilon_e}{(1-\varepsilon_e)}, 0 \right\}$$

$$= \frac{(1-\varepsilon_e)}{P} \max \left\{ \frac{P\varepsilon_e - k}{1-\varepsilon_e}, 0 \right\} = \max \left\{ \varepsilon_e - \frac{(1-\varepsilon_e)k}{P}, 0 \right\}.$$

consider put

$$\bar{V}(\varepsilon_e, 0) = \frac{(1-\varepsilon_e)}{P} \max \left\{ k - \frac{P\varepsilon_e}{1-\varepsilon_e}, 0 \right\} = \max \left\{ \frac{k(1-\varepsilon_e)}{P} - \varepsilon_e, 0 \right\}$$

$$P > 0, P = k.$$

call

$$\bar{V}(\varepsilon_e, 0) = \max \left\{ \varepsilon_e - (1-\varepsilon_e), 0 \right\}$$

$$= \max \left\{ 2\varepsilon_e - 1, 0 \right\}.$$

Put:

$$\bar{V}(\varepsilon_e, 0) = \max \left\{ 1 - 2\varepsilon_e, 0 \right\}$$

Remark:

- ① we assume the volatility σ is function of S (i.e. asset & price) and δ & γ (interest rate & dividend) are constant, the same transformation will be applicable if $r(s, t)$, $r(s, t)$, $\sigma s(s, t)$
- ② When $\varepsilon_e = 1$, we obtain the asymptotic expression of Black-Scholes formula. (i.e. solution of Black-Scholes PDE).

$$V(S, t) \xrightarrow[S \rightarrow \infty]{} \bar{V}(1, \varepsilon_e)$$

$$V(S, t) = (S + P) \bar{V}(1, \varepsilon_e)$$

$$= (S + P) \bar{V}(1, 1) = (S + P) \bar{V}(1, 0) e^{-\delta T} \approx V(S, T) e^{-\delta(T-t)}$$

$$\bar{V}(1, 0)(S + P) \approx V(S, T)$$

Obtain the Black Scholes formula from 1D heat conduction eqn.

$$\textcircled{1} \quad \begin{cases} \frac{\partial u}{\partial \bar{x}} = \frac{\partial^2 u}{\partial n^2}, & -\infty < n < \infty \\ & \bar{x} \geq 0 \\ u(n, 0) = U_0(n), & \\ u(n, \bar{x}) = V(S, t) & \end{cases}$$

Cauchy Pbm
Pure IVP.
No Boundary conditions
consider $n \in (0, \infty)$

$$u(0, \bar{x}) = \phi(\bar{x})$$

$$u(d, \bar{x}) = \psi(\bar{x})$$

$$u(n, \bar{x}) = X(n)T(\bar{x})$$

Since the domain for n is $(-\infty, \infty)$, we don't have Boundary condition.
 \therefore Method of separation of Variable is not applicable to pbm. ①
 We will try to look for a special solution of this form

$$\boxed{u(n, \bar{x}) = \frac{1}{\sqrt{\bar{x}}} v(\bar{x})}$$

$$\eta = \frac{n - \epsilon \bar{x}}{\sqrt{\bar{x}}}$$

$\bar{x}^{-1/2} v(\bar{x})$
 ϵ - parameter

$$\frac{\partial u}{\partial \bar{x}} = \cancel{\frac{1}{\sqrt{\bar{x}}}} v_{\bar{x}} + \frac{1}{\sqrt{\bar{x}}} \frac{\partial v(n)}{\partial n} \cdot \frac{1}{\sqrt{\bar{x}}} = \frac{1}{\bar{x}} \frac{\partial v(n)}{\partial n}$$

$$\frac{\partial u}{\partial n^2} = \frac{1}{\bar{x}} \frac{\partial^2 v(n)}{\partial n^2} + \frac{1}{\sqrt{\bar{x}}} \frac{\partial^2 v(n)}{\partial n \partial \bar{x}} \cdot \frac{1}{\sqrt{\bar{x}}} = \frac{1}{\bar{x}^2} \frac{\partial^2 v(n)}{\partial n^2}$$

$$\begin{aligned} \frac{\partial u}{\partial \bar{x}^2} &= \frac{-1}{2} \bar{x}^{-3/2} v(n) + \frac{\partial v}{\partial n} \frac{1}{\sqrt{\bar{x}}} (n - \epsilon \bar{x}) \cdot \left(-\frac{1}{2} \bar{x}^{-3/2}\right) \\ &= -\frac{\bar{x}^{-3/2}}{2} \left(v + \eta \frac{\partial v}{\partial n}\right) \end{aligned}$$

$$\frac{\partial^2 u}{\partial n^2} = \frac{\partial u}{\partial \hat{r}}$$

$$\hat{e}^{-\hat{r}/2} \frac{\partial^2 U}{\partial \eta^2} = -\frac{\hat{r}-3/2}{2} \frac{\partial(\eta U)}{\partial \eta}$$

$$\Rightarrow \frac{\partial^2 U}{\partial n^2} + \frac{1}{2} \frac{\partial(nU)}{\partial \eta} = 0$$

$$\Rightarrow \frac{\partial}{\partial \eta} \left[\frac{\partial U}{\partial \eta} + \frac{1}{2} n U \right] = 0$$

$$\therefore \boxed{\frac{\partial U}{\partial \eta} + \frac{1}{2} n U = C}$$

$$\frac{\partial U}{\partial \eta} = C - \frac{1}{2} n U$$

for simplicity
C = 0

$$\frac{\partial U}{C - \frac{1}{2} n U} = d\eta$$

$$\frac{\partial U}{U} = -\frac{n}{2} d\eta$$

$$\log U = -\frac{n^2}{4} \eta + C \quad \text{or} \quad U = C_1 e^{-\frac{n^2}{4} \eta}$$

$$\boxed{U(\eta) = C_1 e^{-\frac{(\eta - \xi_e)^2}{4\hat{r}}}}$$

We require that

$$u(n, \hat{r}) = \frac{1}{\sqrt{\hat{r}}} U(n) = C \hat{r}^{-1/2} e^{-\frac{(n - \xi_e)^2}{4\hat{r}}}$$

We require that $\int_{-\infty}^{\infty} C \hat{r}^{-1/2} e^{-\frac{(n - \xi_e)^2}{4\hat{r}}} d\xi_e = 1$.

$$C = \frac{1}{\int_{-\infty}^{\infty} \hat{r}^{-1/2} e^{-\frac{(n - \xi_e)^2}{4\hat{r}}} d\xi_e}$$

$$\boxed{\frac{n - \xi_e}{\sqrt{\hat{r}}} = \eta}$$

$$= \frac{1}{\sqrt{2} \int_{-\infty}^{\infty} e^{-\frac{\eta^2}{2}} d\eta} = \frac{1}{2\sqrt{\pi}}$$

$$u(n, \hat{t}) = \frac{1}{2\sqrt{\pi\hat{t}}} e^{-(n-\varepsilon_e)^2/4\hat{t}}$$

\hookrightarrow fundamental soln. of Green's fn.

$$u(n, 0) = u_0(n)$$

$$\frac{\partial u}{\partial \hat{t}} = \frac{\partial^2 u}{\partial n^2}$$

For any function $G_1(\varepsilon_e, n, \hat{t})$ where ε_e is a parameter. S.t.

$$\frac{\partial G_1(\varepsilon_e, n, \hat{t})}{\partial \hat{t}} = \frac{\partial^2 G_1(\varepsilon_e, n, \hat{t})}{\partial n^2}$$

$$\int_{-\infty}^{\infty} u_0(\varepsilon_e) \frac{\partial}{\partial \hat{t}} G_1(\varepsilon_e, n, \hat{t}) d\varepsilon_e = \int_{-\infty}^{\infty} u_0(\varepsilon_e) \frac{\partial^2 G_1(\varepsilon_e, n, \hat{t})}{\partial n^2} d\varepsilon_e$$

$$\frac{\partial}{\partial \hat{t}} \int_{-\infty}^{\infty} u_0(\varepsilon_e) G_1(\varepsilon_e, n, \hat{t}) d\varepsilon_e = \frac{\partial^2}{\partial n^2} \int_{-\infty}^{\infty} u_0(\varepsilon_e) G_1(\varepsilon_e, n, \hat{t}) d\varepsilon_e.$$

$$u(n, \hat{t}) = \int_{-\infty}^{\infty} u_0(\varepsilon_e) \frac{1}{2\sqrt{\pi\hat{t}}} e^{-(n-\varepsilon_e)^2/4\hat{t}} d\varepsilon_e.$$

$$\lim_{\hat{t} \rightarrow 0} u(n, \hat{t}) = u_0(n)$$

$$\lim_{\hat{t} \rightarrow 0} \frac{1}{2\sqrt{\pi\hat{t}}} e^{-(n-\varepsilon_e)^2/4\hat{t}} = \begin{cases} 0 & n \neq \varepsilon_e \\ \infty & n = \varepsilon_e \end{cases} = \delta(n - \varepsilon_e)$$

Further we have $\int_{-\infty}^{\infty} \frac{1}{2\sqrt{\pi\hat{t}}} e^{-(n-\varepsilon_e)^2/4\hat{t}} d\varepsilon_e = 1$.

$$\lim_{\hat{t} \rightarrow 0} \frac{1}{2\sqrt{\pi\hat{t}}} e^{-(n-\varepsilon_e)^2/4\hat{t}} = \delta(n - \varepsilon_e).$$

$$u_0(n) := \lim_{\hat{\tau} \rightarrow 0} \int_{-\infty}^{\infty} u_0(\varepsilon_e) \frac{1}{2\sqrt{\pi}\hat{\tau}} e^{-(n-\varepsilon_e)^2/4\hat{\tau}} d\varepsilon_e.$$

(eq.)

$$u(n, \hat{\tau}) = \frac{1}{2\sqrt{\pi}\hat{\tau}} e^{-(n-\varepsilon_e)^2/4\hat{\tau}}$$

$u(n, \hat{\tau})$ along with $u_0(n)$ satisfies heat conduction condition with initial conditions. we have to retrieve Black Scholes formula from (eq.).

$$V(S, t) = e^{-r(T-t)} u(n, \hat{\tau})$$

$$= e^{-r(T-t)} \int_{-\infty}^{\infty} u_0(\varepsilon_e) \frac{1}{2\sqrt{\pi}\hat{\tau}} e^{-(n-\varepsilon_e)^2/4\hat{\tau}} d\varepsilon_e.$$

$$V(S, t) = N_T(S)$$

$$n = \ln S + \left(r - \delta - \frac{\sigma^2}{2}\right) T$$

$$n - \varepsilon_e = \left(\ln S + \left(r - \delta - \frac{\sigma^2}{2}\right)\right)$$

$$\hat{\tau} = \frac{1}{2} \sigma^2 T$$

$$4\hat{\tau} = 2\sigma^2 T$$

$$V(S, t) = e^{-r(T-t)} \frac{1}{\sigma\sqrt{2\pi}(T-t)} \int_{-\infty}^{\infty} u_0(\varepsilon_e) e^{-\frac{(n - \varepsilon_e)^2}{2\sigma^2(T-t)}} d\varepsilon_e$$

$$V(S, t) = e^{-r(T-t)} \frac{1}{\sigma\sqrt{2\pi}(T-t)} \int_0^{\infty} V_T(S') \times \exp\left(-\ln S - \left(\ln S + \left(r - \delta - \frac{\sigma^2}{2}\right)(T-t)\right)\right) / 2\sigma^2 \times dS'$$

Black Scholes formula can be written.

$$V(S, t) = e^{-rt} \int_0^{\infty} V_T(\tilde{S}) G_1(\tilde{S}, T, S, t) d\tilde{S}$$

$$G_1(\cdot) = \frac{1}{\tilde{S}} \frac{1}{\sigma \sqrt{2\pi(T-t)}} \exp\left(-\frac{\ln \tilde{S} - \left(\ln S + \left(r - \frac{\sigma^2}{2}\right)(T-t)\right)}{\frac{\sigma^2}{2}(T-t)}\right)^2$$

G_1 is Green's function for Black Scholes PDE.

The PDF of a log normal distribution with $m = \left(r - \frac{\sigma^2}{2}\right)$ and

$$E(\tilde{S}) = S \exp((r - \frac{\sigma^2}{2})(T-t))$$

$$G_1(\cdot) = \frac{1}{\tilde{S}} \left[\exp\left(-\frac{(\ln(\tilde{S}/S) + \frac{\sigma^2}{2})^2}{2\sigma^2}\right) \right] \cdot \frac{1}{\sqrt{2\pi} \sigma \tilde{S}}$$

$$a = S e^{(r-\frac{\sigma^2}{2})(T-t)}$$

$$b = \sigma \sqrt{T-t}$$

$$\text{Call} \rightarrow N(d_1) S - N(d_2) k e^{-rt}$$

In order to obtain Black & Scholes formula we have to prove
following identity:

$$\int_c^\infty G_1(\tilde{S}, T, S, t) d\tilde{S} = N\left(\frac{\ln(a/c) - b^2/2}{b}\right)$$

$$\int_c^\infty \tilde{S} G_1(\tilde{S}, T, S, t) d\tilde{S} = a N\left(\frac{\ln(a/c) + b^2/2}{b}\right).$$

$$N(\phi) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\phi} e^{-\frac{e^2}{2}} de$$

$$\text{let } \eta(\tilde{s}) = \frac{\ln(s/a) + b^2/2}{b}$$

$$\tilde{s} = ae^{b\eta} - b^2/2$$

$$d\tilde{s} = ae^{-b^2/2} b \cdot e^{b\eta} d\eta$$

$$= b \tilde{s} d\eta$$

$$C_1 = e^{-\frac{\eta^2}{2}} \cdot \frac{1}{\sqrt{2\pi} b \tilde{s}}$$

B-S model.

- ① FDM
- ② FEM
- ③ Spline Method
- ④ NNS

$$\int_C^\infty C_1 d\tilde{s} = \int_C^\infty \frac{e^{-\eta^2/2}}{\sqrt{2\pi} b \tilde{s}} d\tilde{s} = \int_C^\infty \frac{e^{-\eta^2/2}}{\sqrt{2\pi}} d\eta$$

$$\int_C^\infty \frac{e^{-\eta^2/2}}{\sqrt{2\pi}} d\eta$$

$$\int_C^\infty e^{-\eta^2/2} d\eta$$

$$\frac{\ln(\frac{c}{a}) + \frac{b^2}{2}}{b}$$

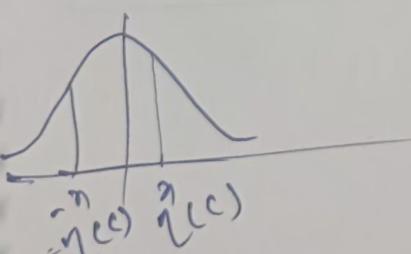
$$\int_C^\infty \frac{e^{-\eta^2/2}}{\sqrt{2\pi}} d\eta$$

$$\frac{\ln(\frac{c}{a}) + \frac{b^2}{2}}{b}$$

$$\frac{(b^2 - b^2)/b}{(b^2 - b^2)/b}$$

$$\int_{-\infty}^0 \frac{e^{-\eta^2/2}}{\sqrt{2\pi}} d\eta$$

$$N \left(\frac{\ln(\frac{c}{a}) + \frac{b^2}{2}}{b} \right)$$



$$(z^0 - z^-)^{-1} \int_{-\infty}^0$$

$$\frac{(z^0 - z^-)^{-1}}{b} dz$$

$$\left(\frac{(z^0 - z^-)^{-1}}{b} \right)^2 dz$$

$$\left(\frac{(z^0 - z^-)^{-1}}{b} \right)^2 dz$$

European PDE

1. Generalized one

2. Non linear

3. Jump difference.

American

free BVP

Exotic option

$$\int_c^{\infty} \tilde{G}_1 d\tilde{s} = \int_{\eta(c)}^{\infty} \frac{e^{-n^2/2}}{\sqrt{2\pi b\tilde{s}}} b\tilde{s} d\eta$$

~~Integrate by parts~~

$$\int_{\eta(c)}^{\infty} \frac{e^{-n^2/2}}{\sqrt{2\pi b\tilde{s}}} \cdot \tilde{s} \psi d\eta$$

$$\tilde{s} = ae^{bn - \frac{b^2}{2}}$$

$$\int_{\eta(c)}^{\infty} \frac{e^{-n^2/2} - \frac{b^2}{2} + bn}{\sqrt{2\pi}} d\eta$$

$$= \int_{\eta(c)}^{\infty} \frac{e^{-\frac{(n-b)^2}{2}}}{\sqrt{2\pi}} d\eta$$

$$= \frac{a}{\sqrt{2\pi}} \int_{\eta(c)}^{\infty} e^{-\frac{(n-b)^2}{2}} d\eta$$

$$= \frac{a}{\sqrt{2\pi}} \int_{\eta(c)-b}^{\infty} e^{-n^2/2} dn$$

$$\therefore a N(b - \eta(c))$$

$$\therefore a N\left(b + \frac{\ln(a/c) - b^2/2}{b}\right)$$

$$\therefore a N\left(\frac{\ln(a/c) + b^2/2}{b}\right)$$

$$V(S, t) = \int_0^{\infty} e^{-r(T-t)} \int_0^{\infty} \max(S - k, 0) G_1 dS \quad \left| \begin{array}{l} V_T(S) = \max(S - K, 0) \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right.$$

$$= e^{-r(T-t)k} \int_0^{\infty} (S - k) G_1 dS + \int_k^{\infty} S G_1 dS - \int_k^{\infty} k G_1 dS$$

$$= \left[aN\left(\frac{\ln(a/k) + b^2/2}{b}\right) - k N\left(\frac{\ln(a/k) - b^2/2}{b}\right) \right] e^{-r(T-t)}$$

$$a = S e^{(R-\delta)(T-t)},$$

$$b = \sigma \sqrt{T-t},$$

$$\frac{b^2}{2} = \sigma^2 \frac{(T-t)}{2}.$$

$$\bullet \quad \ln\left(\frac{a}{k}\right) = \ln\left(\frac{S}{K}\right) + (R-\delta)(T-t)$$

$$\ln\left(\frac{a}{k}\right) + \frac{b^2}{2} = \frac{\ln\left(\frac{S}{K}\right) + \left(\frac{\sigma^2}{2} + (R-\delta)\right)(T-t)}{\sigma \sqrt{T-t}}$$

$$d_1 = \frac{\sigma \sqrt{T-t}}{b}$$

$$\ln\left(\frac{a}{k}\right) - \frac{b^2}{2} = \frac{\ln\left(\frac{S}{K}\right) + \left((R-\delta) - \frac{\sigma^2}{2}\right)}{\sigma \sqrt{T-t}}$$

consider

$$d_2 = \frac{\sigma \sqrt{T-t}}{b}$$

$$\boxed{V(S, t) = S e^{-\delta(T-t)} N(d_1) + k e^{-r(T-t)} N(d_2)}$$

Consider the payoff for a forward contract where.

$$V(S, T) = S - K.$$

$$V(S, t) = e^{-\delta(T-t)} \int_0^T (\tilde{S} - K) \alpha(\tilde{S}, \tilde{t}, S, t) d\tilde{S}.$$

$$\begin{aligned} &= e^{-\delta(T-t)} (S e^{(\bar{\alpha} - \delta)(T-t)} - K) \\ &= S e^{-\delta(T-t)} - K e^{-\delta(T-t)}. \end{aligned}$$

Since for a forward contract the buyer does not need to pay any premium at initial time. ∴

$$\therefore t=0$$

$$V(S, 0) = S e^{-\delta T} - K e^{-\delta T} = 0.$$

$$K = e^{(\bar{\alpha} - \delta)T} S_0$$

Greeks i.e. delta, gamma, theta, vega, rho → $\frac{\partial V}{\partial S}, \frac{\partial V}{\partial t}, \frac{\partial V}{\partial \delta}, \frac{\partial V}{\partial \bar{\alpha}}, \frac{\partial V}{\partial K}$

$\frac{\partial V}{\partial S}$: Derivative of option w.r.t. parameter, $S, t, \delta, \bar{\alpha}, K$

Calculate greeks, corresponding errors and corresponding plot

$$A\eta = b$$

1) Jacobi

2) Gauss-Seidel

3) SOR (Relaxation method)

$$\left| \begin{array}{l} \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, \quad (x, t) \in (a, b) \times [0, T] \\ u(a, t) = u(b, t) = 0 \\ u(x, 0) = \phi(x) \quad x \in [a, b] \end{array} \right.$$

In FTCS, BTCS, CN we discretize domain in both x and t .

Semi-discretization \rightarrow yield more accuracy than above
discretize either x or t \Rightarrow get system of ODEs
 \downarrow
get IVP/BVPs

Semi-discrete schemes (Method of lines) (MOL)

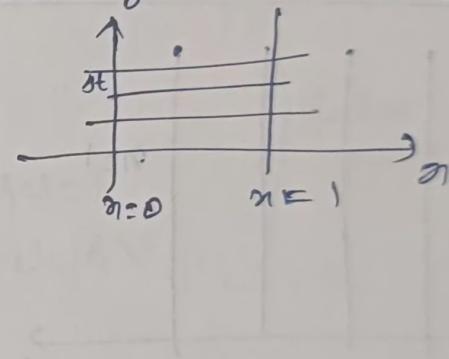
In order to solve parabolic IBVP numerically, one can use Method of lines which is known as Semi discrete scheme. In this method we discretize either time domain and preserve x domain continuous (or) vice versa. This method is Semi discrete. Time discretization \rightarrow horizontal

Consider this model.

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + f(x, t)$$

$$u(a, t) = u(b, t) = 0$$

$$u(x, 0) = \phi(x) \quad x \in [a, b]$$



Let discretize t and use Euler. (Method of horizontal lines)

$$\frac{U^{k+1}(n) - U^k(n)}{\Delta t} = \frac{\partial^2 U^k(n)}{\partial x^2} + f(n, t^k).$$

$$U^{k+1}(n) = \Delta t \left(\frac{\partial^2 U^k(n)}{\partial x^2} + f(n, t^k) \right) + U^k(n)$$

This is explicit euler
for

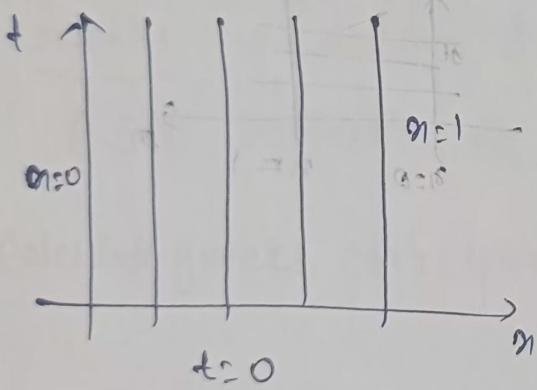
For stability condition we use implicit euler.

$$\textcircled{2} \quad \frac{U^{k+1}(n) - U^k(n)}{\Delta t} = \frac{\partial^2 U^{k+1}(n)}{\partial n^2} + f(n, t^{k+1}) \quad U^{k+1}(n) = U(n, t^{k+1})$$

$$\textcircled{3} \quad \left\{ \begin{array}{l} U^k(a) = U^k(b) = 0 \\ U^{k+1} - \Delta t \frac{\partial^2 U^{k+1}(n)}{\partial n^2} = U^k(n) + f(n, t^{k+1}) \Delta t. \end{array} \right.$$

Eq. Given in $\textcircled{3}$ are system of 2nd order ODEs with boundary conditions i.e. system of BVP which can be solved by 2nd order methods or any higher order methods.

Discretize in first and preserve t continuous (Method of vertical lines)



late discrete special domain with vertical lines.

We write the eq. $\textcircled{1}$

$$\textcircled{4} \quad \left\{ \begin{array}{l} \frac{du_i(t)}{dt} = \frac{U_{i+1}(t) - 2U_i(t) + U_{i-1}(t)}{\Delta n^2} + f_i(t) \\ U_0(t) = U_{N+1}(t) = 0 \\ U_i(0) = \phi(n_i) \end{array} \right. \quad i = 0, 1, \dots, N+1$$

Equations given in ④ are system of 1st order ODEs with initial conditions i.e. we obtain system of IVP for 1st order ODEs.

Explicit Euler \rightarrow FTCS

Implicit Euler \rightarrow BTCS

One can use either explicit or Implicit Euler Scheme, the resultant scheme will be FTCS or BTCS. In order to obtain higher order scheme one can use RK of 4th order, or multistep scheme like Adam's Bashforth or Adams

Moulton

$$\frac{dy}{dt} = f(y, t)$$

$$\left. \begin{array}{l} \frac{dy_1}{dt} = f_1(y_1, y_2, t) \\ \frac{dy_2}{dt} = f_2(y_1, y_2, t) \end{array} \right\}$$

Finite Element Method (FEM) Ch. 5

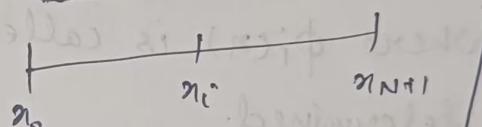
$$u'' = f(x), \quad x \in (0, 1)$$

$$u(0) = 0 \equiv u(1)$$

Adhoc approach

$$\frac{U_{i+1} - 2U_i + U_{i-1}}{h^2}, \quad f_i, \quad 1 \leq i \leq N$$

$$U_0 = U_{N+1} = 0 \quad (B.C.)$$



$$AU = F$$

$$U = A^{-1}F$$

Adhoc approach

In FEM we use systematic approach.

$$L_E = \|u'' - f\| \rightarrow 0$$

Residue or loss

(00)

$$\int_0^1 (u'' - f) \psi(n) dn = 0$$

$$\int_0^1 u'' \psi(n) dn = \int_0^1 f \psi(n) dn$$

weak formulation

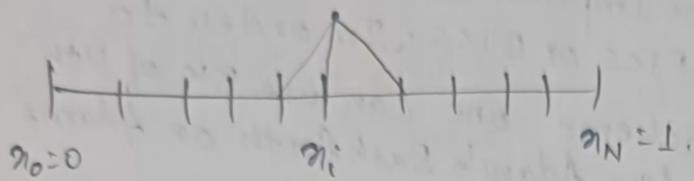
$$\begin{cases} Lu = f, \quad \Omega \\ Lu = -u'' \quad \Omega = (0,1) \\ u(0) = u(1) = 0 \end{cases}$$

FEM

Step 1: Weak formulation

Step 2: Domain discretization
finite dimensional pbm

Step 3: Numerical Integration/
System of linear algebraic



$$u(x_i) = \sum c_i \phi_i(x) \rightarrow \text{Basis function}$$

\downarrow hat function

We discretize domain $\Omega = \bigcup_{k=1}^N \Omega_k$, $\Omega_k = [x_k, x_{k+1}]$.

$$\text{Approx } w(x) = \sum_{i=1}^N c_i \phi_i(x) \quad (*)$$

where $\phi_i(x)$ is called Basis functions and constants

c_i 's has to be determined.

In order to determine the constants c_i 's we have to form n. no. of equations.

One way of calculation of c_i 's by using residual function

Residual fn. $R = Lw - f$

Idea is such that to minimise the residual in such a way that approximate soln. given in that eq. in such that a way that error is minimum.

$$\text{Error} = |u(x_i) - w(x_i)|$$

To minimise the residue we will consider N weight function (test function). ψ_1, \dots, ψ_N s.t. $\int_{\Omega} R \psi_j(x) dx = 0$
 $j=1, \dots, N$

$$Lw = - \sum c_i \phi_i''(n)$$

On substituting w in R we obtain:

$$\int_2 - \left(\sum_{i=1}^N c_i \phi_i''(n) \right) \psi_j(n) dn = 0$$

$$\left\{ \sum_{i=1}^N c_i \int_2 \phi_i''(n) \psi_j(n) dn = - \int_2 f \psi_j(n) dn \right\} \rightarrow \text{(*)}$$

$j = 1, \dots, N$

$$\int Lw \psi_j(n) dn = \int f \psi_j(n) dn \quad j = 1, \dots, N \rightarrow (3)$$

$(Lw, \psi_j) = (f, \psi_j) \rightarrow \text{Inner product.}$

If we choose

(i) $\phi_j(n) = \psi_j(n)$ i.e. test and trial function is same.

But now Galerkin.

(ii) Collocation Scheme.

$\psi_j(n) = \delta(n - n_j)$ Dirac function

$$\delta(n - n_j) = \begin{cases} \infty & n = n_j \\ 0 & n \neq n_j \end{cases}$$

Substitute, $\psi_j(n) = \delta(n - n_j)$ in (3).

$$\int f \psi_j(n) dn = \int \cancel{f(n_j)} dn \cdot f(n_j) \therefore$$

$$\int Lw \psi_j(n) dn = Lw(n_j).$$

Least Square method

$$\psi_j(n) = \frac{\partial R}{\partial c_j}$$

$$\int_{\Omega} Lw \psi_j(n) dn = \int_{\Omega} f \psi_j(n) dn.$$

$$\int_{\Omega} Lw \frac{\partial R}{\partial c_j} dn = \int_{\Omega} f \frac{\partial R}{\partial c_j} dn. \quad \frac{\partial R}{\partial c_j} = \frac{\partial}{\partial c_j} (Lw - f)$$

$$\int_{\Omega} ((Lw - f)) \frac{\partial R}{\partial c_j} dn = 0$$

$$\int_{\Omega} R \frac{\partial R}{\partial c_j} dn = 0$$

Model Pbm.

1D BVP

$$① \quad \begin{cases} -(p(n)u'(n))' + q(n)u(n) = r(n) & n \in (0, 1) \\ u(0) = 0 = u(1); \end{cases}$$

Solve using FEM.

Step 1: Weak formulation

$$v \in \mathcal{C}_0^\infty(\Omega)$$

continuous many times differentiable

$$① \times v \Rightarrow -(p(n)u')'v + q(n)uv = r(n)v$$

Integration by parts:

$$\int_0^1 -(p(n)u')'v dn + \int_0^1 q(n)uv dn = \int_0^1 r(n)v dn.$$

$$\int_0^1 uv' = (uv)_0^1 - \int_0^1 u'v du$$

$$(-(\rho(n)u')v)^{\perp} + \int_0^1 \rho(n)u'v'dn + \int_0^1 q(n)uv dn = \int_0^1 r(n)v(n) dn.$$

$\xrightarrow{?}$

find $u \in H_0(\Omega) = \mathbb{V}$

$$\textcircled{1} \quad \begin{cases} \text{s.t.} \\ \int_0^1 \rho(n)u'v'dn + \int_0^1 q(n)uv dn = \int_0^1 r(n)v(n) dn \quad \forall v \in V \end{cases}$$

$\therefore q(n)$ is continuous in $\textcircled{1}$

$\textcircled{2} \quad \therefore u \in C^2(\Omega) \text{ & } p \in C^1(\Omega)$

$$\begin{cases} u \in L^p(\Omega) \\ \text{i.e. } \left(\int |u|^p dn \right)^{1/p} < \infty \end{cases}$$

$\textcircled{3} \quad \textcircled{2} \text{ we don't need } \textcircled{1}$

bonded measurable fn.

Step ② Discretise the domain

$$\Omega = \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \quad \eta_0 = 0 \quad \eta_1 \quad \eta_N = 1.$$

V_h is a subspace of V which is of finite dimension.

$$V_h = \text{Span} \{ \phi_1, \dots, \phi_{N-1} \}$$

ϕ_i 's are piecewise basis function.
finite element problem.

Find $u_h \in V_h$ s.t.

$$\int_0^1 \rho(n)u_h v_h dn + \int_0^1 q(n)u_h v_h dn = \int_0^1 r(n)v_h(n) dn \quad \forall v_h \in V_h.$$

$\textcircled{3} \rightarrow$ Finite dimensional space

$\textcircled{2} \rightarrow$ In infinite

$$\text{Express } u_h(n) = \sum_{i=1}^{N-1} v_i \phi_i(n) \rightarrow \textcircled{4}$$

v_i 's are unknown to determine

In order to determine $N-1$ unknowns we need to have that many equations i.e. System of linear algebraic Equations

Using ④ in ③

$$\sum_{i=1}^{N-1} \left(\int_{n_i}^{n_{i+1}} p(n) \phi_i'(n) \phi_j'(n) + q(n) \phi_i(n) \phi_j(n) \right) U_i = \int_0^1 r(n) \phi_j(n) dn$$

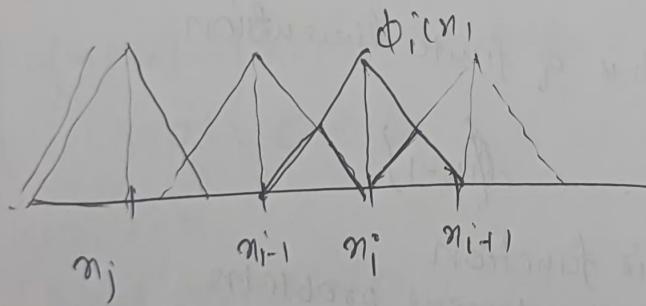
$\forall j = 1, \dots, N-1$

$AU = F \Leftarrow$ System of linear equations.

$$U = [U_1, \dots, U_{N-1}]^T$$

Assume the basis elements ϕ_i 's are piecewise linear polynomial i.e. $\phi_i(n) = a_i^n + b_i n$ $\Omega_i = [n_i, n_{i+1}]$

$$\phi_i(n_j) = \begin{cases} 1 & i=j \\ 0 & \text{o.w.} \end{cases}$$

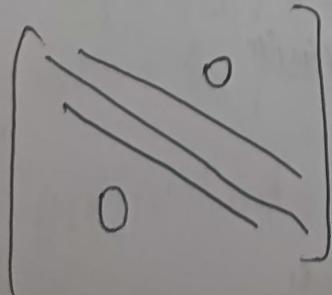


$$\phi_i \phi_j \neq 0 \text{ when } \begin{array}{l} j = i-1 \\ j = i+1 \\ j = i \end{array}$$

Matrix A is called stiffness matrix which is 3-diagonal matrix

i.e. $a_{ij} \neq 0 \quad \begin{cases} |i-j| \leq 1 \end{cases}$

$$a_{ij} = \begin{cases} 0 & |i-j| > 1 \\ \neq 0 & |i-j| \leq 1 \end{cases}$$



$$Q_{ij} = \int p(n) \phi_i^{\dagger} \phi_j^{\dagger} + \int q(n) \phi_i \phi_j dn \quad p_j = \int q(n) \phi_j(n) dn.$$

$$J_1 = \int_{n_{i-1}}^{n_i} + \int_{n_i}^{n_{i+1}}$$

Use Trapezoidal rule to evaluate Integrals

$$a_{ii} = \int_{n_{i-1}}^{n_i} p(n) \phi_i^{\dagger} \phi_i^{\dagger} + \int_{n_i}^{n_{i+1}} p(n) \phi_i^{\dagger} \phi_i^{\dagger} dn.$$

$$= p\left(\frac{1}{h}\right)\left(\frac{1}{h}\right) \int_{n_{i-1}}^{n_i} dn + p\left(\frac{-1}{h}\right)\left(\frac{-1}{h}\right) \int_{n_i}^{n_{i+1}} dn.$$

$$= \frac{2p}{h}$$

$$j = i-1$$

$$a_{i,i-1} = \int_{n_{i-1}}^{n_i} p \phi_i^{\dagger}(n) \phi_{i-1}^{\dagger}(n) dn = p\left(\frac{1}{h}\right)\left(-\frac{1}{h}\right) h = -p/h$$

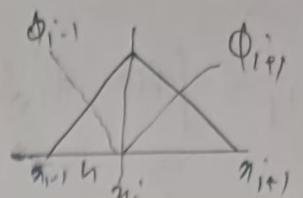
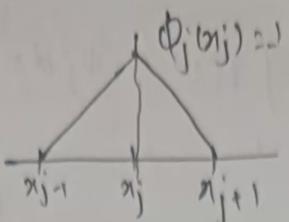
$$j = i+1$$

$$a_{i,i+1} = \int_{n_i}^{n_{i+1}} p \phi_i^{\dagger}(n) \phi_{i+1}^{\dagger}(n) dn = p\left(-\frac{1}{h}\right)\left(\frac{1}{h}\right) h = -p/h$$

$$a_{ij} = \sum_{j=1}^{N-1} \int_{n_j}^{n_i} q(n) \phi_i(n) \phi_j(n) dn. \quad i = 1, \dots, N-1.$$

$$a_{ii} = \int_{n_{i-1}}^{n_i} q \phi_i^{\dagger}(n) \phi_i^{\dagger}(n) dn + \int_{n_i}^{n_{i+1}} q \phi_i^{\dagger}(n) \phi_i^{\dagger}(n) dn.$$

$$\phi_i(x_j) = \begin{cases} 1, & j = i \\ 0, & \text{o.w.} \end{cases}$$



$$a_{ii} = \frac{qh}{2} [\phi_i(x_{i-1}) + \phi_i(x_i)] + \frac{qh}{2} [\phi_i(x_i) + \phi_i(x_{i+1})] \\ = qh.$$

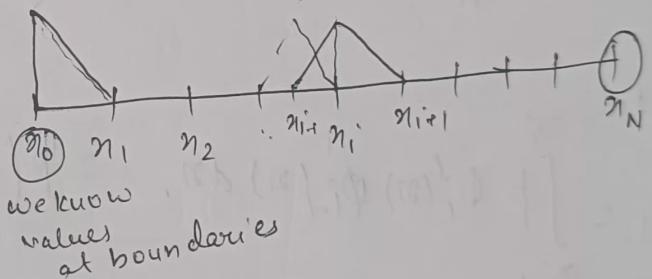
$j = i - 1$

$$a_{ii-1} = \int_{x_{i-1}}^{x_i} q \phi_{i-1} \phi_i dx = \frac{qh}{2} [\phi_{i-1}(x_{i-1}) \phi_i(x_{i-1}) + \phi_{i-1}(x_i) \phi_i(x_i)] \\ = 0$$

$j = i + 1$

using trapezoidal rule.

$$\textcircled{2} a_{ii+1} = \int_{x_i}^{x_{i+1}} q \phi_i(x) \phi_{i+1} dx = 0$$



Now we get $AU=F$ matrix:

$$a_{ii} = \cancel{\frac{2p}{h}} +$$

$$a_{ij} = \begin{cases} -p/h & j = i-1 \\ 2p/h + qh & j = i \\ -p/h & j = i+1 \end{cases}$$

$$\begin{aligned}
 \text{RHS} &= \int_0^{\infty} \varphi(n) \phi_j(n) dn \\
 &= \sum_{j=1}^{N+1} \int_{n_{j-1}}^{n_j} \varphi(n) \phi_j(n) dn \\
 &= \int_{n_{j-1}}^{n_j} \varphi(n) \phi_j(n) dn + \int_{n_j}^{n_{j+1}} \varphi(n) \phi_j(n) dn \\
 &= \frac{h}{2} \left[(\varphi(n_{j-1}) \phi_j(n_{j-1}) + \varphi(n_j) \phi_j(n_j)) \right] \\
 &\quad + \frac{h}{2} \left[\varphi(n_j) \phi_j(n_j) + \varphi(n_{j+1}) \phi_j(n_{j+1}) \right] \\
 &= \frac{h}{2} \varphi(n_j) \times 2 = \boxed{\varphi(n_j) \cdot h}
 \end{aligned}$$

$$f_j = \varphi(n_j) \cdot h$$

$$A \cup F \Rightarrow U = A \setminus F$$

After this

$$U_n(n) = \sum_{i=1}^{N+1} U_i \phi_i(n)$$

$$\textcircled{a} \quad U(n_j) \approx U_i$$

Exercise if $\phi_j(n)$ is quadratic.

Using Simpson's rule evaluate the entries

$$a_{ji} = \int_0^1 p(n) \phi_i'(n) \phi_j(n) dn + \int_0^1 q(n) \phi_i(n) \phi_j(n) dn$$

$\underbrace{\hspace{10em}}$ $\underbrace{\hspace{10em}}$

I_1 I_2

$$\begin{aligned}
 \text{when } j=1 \quad I_1 &= \int_{n_{i-1}}^{n_i} p(n) \phi_i'(n) \phi_1(n) dn + \int_{n_i}^{n_{i+1}} q(n) \phi_i'(n) \phi_1(n) dn \\
 I_1 &= \int_{n_{i-1}}^{n_i} p(n) \phi_i'(n) \phi_1(n) dn + \int_{n_i}^{n_{i+1}} q(n) \phi_i'(n) \phi_1(n) dn
 \end{aligned}$$

$$\phi_i^1 = \frac{1}{h}$$

$$I_1 = \int_{n_{i-1}}^{n_i} p(n) \phi_{i-1}^1(n) \phi_i^1(n) dn + \int_{n_i}^{n_{i+1}} p(n) \phi_i^1(n) \phi_i^1(n) dn.$$

$$\frac{1}{h^2} \int_{n_{i-1}}^{n_i} p(n) dn + \frac{1}{h^2} \int_{n_i}^{n_{i+1}} p(n) dn.$$

$$\frac{1}{h^2} \times \frac{1}{6} \left[p(n_{i+1}) + 4p\left(\frac{n_{i+1} + n_{i-1}}{2}\right) + p(n_i) \right]$$

$$+ \frac{1}{h^2} \times \frac{1}{6} \left[p(n_i) + 4p\left(\frac{n_i + n_{i+1}}{2}\right) + p(n_{i+1}) \right]$$

$$\frac{1}{6h} \left[2p(n_i) + p(n_{i-1}) + p(n_{i+1}) + 4p\left(\frac{n_i + n_{i-1}}{2}\right) + 4p\left(\frac{n_i + n_{i+1}}{2}\right) \right]$$

$$= \frac{1}{6h} \left[2p(n_i) + p(n_{i-1}) + p(n_{i+1}) + 2p(n_i) + 2p(n_{i-1}) + 2p(n_i) + 2p(n_{i+1}) \right]$$

$$\frac{2p}{6h} = \frac{1}{6h} \left[6p(n_i) + 3p(n_{i-1}) + 3p(n_{i+1}) \right] \\ = \left[\frac{p_{i-1}}{2} + p_i + \frac{p_{i+1}}{2} \right] \frac{1}{h}$$

$$a_{i(i-1)} = \int_{n_{i-1}}^{n_i} p(n) \phi_{i-1}^1(n) \phi_i^1(n) dn =$$

$$= \left(-\frac{1}{h} \right) \left(\frac{1}{h} \right) \int_{n_{i-1}}^{n_i} p(n) dn$$

$$= -\frac{1}{h^2} \times \frac{h}{6} \left[p_{i-1} + 2p_i + 2p_{i-1} + p_i \right]$$

$$\frac{2p}{6h} \left(\frac{p_{i-1} + p_i}{2} \right) \times \frac{1}{h}$$

$$\begin{aligned}
 a_{ii+1} &= \int_{n_i}^{n_{i+1}} p(n) \phi_i'(n) \phi_{i+1}'(n) dn \\
 &= \left(\frac{1}{h}\right) \left(\frac{1}{h}\right) \int_{n_i}^{n_{i+1}} p(n) dn \\
 &= -\frac{1}{h^2} \times \frac{h}{6} \left[p_{i+1} + 2p_i + 2p_{i+1} + p_{i+1} \right] \\
 &\approx -\frac{h}{2} [p_i + p_{i+1}]
 \end{aligned}$$

when $j = i$

$$\begin{aligned}
 I_2 &= \int_{n_{i-1}}^{n_i} q(n) (\phi_i(n))^2 dn + \int_{n_i}^{n_{i+1}} q(n) (\phi_i(n))^2 dn \\
 &= \frac{h}{6} \left[q_{i-1} (\phi_i(n_{i-1}))^2 + 4q\left(\frac{n_{i-1} + n_i}{2}\right) \left(\phi_i\left(\frac{n_{i-1} + n_i}{2}\right)\right)^2 + q_{i+1} (\phi_i(n_{i+1}))^2 \right. \\
 &\quad \left. + \frac{h}{6} \left[q_i (\phi_i(n_i))^2 + 4q\left(\frac{n_i + n_{i+1}}{2}\right) \left(\phi_i\left(\frac{n_i + n_{i+1}}{2}\right)\right)^2 + q(n_{i+1}) (\phi_i(n_{i+1}))^2 \right] \right] \\
 &= \frac{h}{6} \left[\cancel{4q(n_{i-1})} \frac{4q\left(\frac{n_{i-1} + n_i}{2}\right)}{2} \left(\frac{\phi_i(n_i) + \phi_i(n_{i-1})}{2}\right)^2 + q_i^2 \right] \\
 &\quad + \frac{h}{6} \left[q\left(\frac{n_{i-1} + n_i}{2}\right) + q_i \right] + \frac{h}{6} \left[q_i + q\left(\frac{n_i + n_{i+1}}{2}\right) \right] \\
 &\quad + \frac{h}{3} q_i + \frac{h}{6} \left[q\left(\frac{n_{i-1} + n_i}{2}\right) + q\left(\frac{n_i + n_{i+1}}{2}\right) \right]
 \end{aligned}$$

$$a_{ii-1} = \frac{h}{6} \left[q\left(\frac{n_{i-1} + n_i}{2}\right) + q_i \right]$$

$$a_{ii+1} = \frac{h}{6} \left[q_i + q\left(\frac{n_i + n_{i+1}}{2}\right) \right]$$

Assume $\phi_i(n)$ is quadratic then find A and F

$$\begin{aligned} -u'' &= f \\ -(p(x)u')' + q(x)u &= r(x) \quad \text{Span } \{\phi_1, \dots, \phi_{N-1}\} = V_h \\ u(0) &= u(1) = 0 \quad A U = F \end{aligned}$$

$$\int_0^1 p(x) \phi_i' \phi_j' dx + \int_0^1 q(x) \phi_i \phi_j dx = \int_0^1 f \phi_j dx$$

$$\int \phi_i' \phi_j' dx \in \int \phi_i \phi_j dx.$$

Right hand side Vector can be calculated directly or it can be calculated from $f_j B$.

$$A = \frac{1}{h} \begin{bmatrix} 1 & -1 & & & \\ -1 & 2 & -1 & & \\ & -1 & 2 & \ddots & \\ & & & \ddots & 1 \end{bmatrix}$$

$$B = h \begin{bmatrix} 2 & 1 & & & \\ 1 & 4 & 1 & & \\ & & & \ddots & \\ & & & & 1 \end{bmatrix}$$

$$\frac{\partial u}{\partial \xi} = \frac{\partial^2 u}{\partial \eta^2} \quad (\eta, \xi) \in (-\infty, \infty) \times [0, T]$$

(η_{min}, η_{max}) × [0, T]

B.C.
I.C.

$u(\eta, t) \longleftarrow$ Basis $\phi(\eta, t)$

By using Method of Horizontal lines we can convert PDE to System of ODE's which can be solved by Finite Element method.

Rather one can follow discretizing both time and space by finite element.

The soln. $u(\eta, t) = \sum_{i=1}^N w_i(t) \phi_i(\eta) + \phi_0(\eta, t) \rightarrow ②$

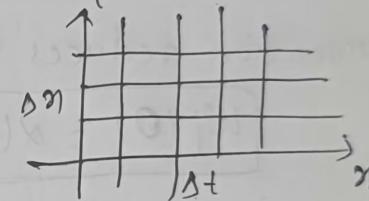
ϕ_0 Boundary conditions

where ϕ_0 satisfies the boundary and initial conditions.

In earlier case of ODE the soln. of $u_h(\eta) = \sum w_i \phi_i(\eta)$.

whereas in the case of parabolic PDEs. ~~it will be of~~ it will be of function of times also.

Using ② in eq. ①. We obtain



$$\int_{\eta_0}^{\eta_m} \left[\sum_{i=1}^{m-1} \overset{\circ}{w_i} \phi_i + \overset{\circ}{\phi_0} \right] \phi_j d\eta = \int_{\eta_0}^{\eta_m} \left[\sum_{i=1}^{m-1} \overset{\circ}{w_i} \phi_i'' + \overset{\circ}{\phi_0}'' \right] \phi_j d\eta \quad j=1, \dots, m-1. \rightarrow ③$$

$\overset{\circ}{}$ denotes w.r.t. ξ .

$\overset{\circ}{w_i}$ denotes derivative w.r.t. ξ and $\overset{\circ}{\phi_i}''$ denotes derivative w.r.t. η .

Eq. ③ can be written in terms of matrix vector form.

$$④ \boxed{B\overset{\circ}{w} + b = -Aw - a}$$

$$b(\xi) = \begin{pmatrix} \int \overset{\circ}{\phi_0} \phi_1 d\eta \\ \int \overset{\circ}{\phi_0} \phi_2 d\eta \\ \vdots \\ \int \overset{\circ}{\phi_0} \phi_{n-1} d\eta \end{pmatrix}$$

$$a(\xi) = \begin{pmatrix} \int \overset{\circ}{\phi_0}'' \phi_1 d\eta \\ \int \overset{\circ}{\phi_0}'' \phi_2 d\eta \\ \vdots \\ \int \overset{\circ}{\phi_0}'' \phi_{n-1} d\eta \end{pmatrix}$$

$$\begin{pmatrix} \int \overset{\circ}{\phi_0}'' \phi_1 d\eta \\ \int \overset{\circ}{\phi_0}'' \phi_2 d\eta \\ \vdots \\ \int \overset{\circ}{\phi_0}'' \phi_{n-1} d\eta \end{pmatrix}$$

④ is System of first order ODEs which is method of vertical lines

$$w = [w_1(t), \dots, w_{n-1}(t)]^T$$

We have to determine w which is soln. of eq. ④

At $t=0$ $\boxed{u(\pi, 0) = \alpha(\pi)} \rightarrow ⑤$

From ② and ⑤

$$\sum_{i=1}^N w_i(0) \phi_i(\pi) + \phi_0(\pi, 0) = \alpha(\pi)$$

At $\pi = \pi_j$

Since $\phi_i(\pi_j) = \delta_{ij} = \begin{cases} 1 & i=j \\ 0 & i \neq j \end{cases}$

Summation reduces to 1 term.

$$\boxed{w_j(0) = \alpha(\pi_j) - \phi_0(\pi_j, 0)}$$

Initial condition.