

## 1 Structure of the integrals

The general structure of the quantities to be computed in APFEL is the combination of terms  $I$  that have the form of Mellin convolutions between an operator  $O$  and distribution function  $d$ , that is:

$$I(x) = x \int_0^1 dz \int_0^1 dy O(z) d(y) \delta(x - yz) = x \int_x^1 \frac{dz}{z} O\left(\frac{x}{z}\right) d(z) = x \int_x^1 \frac{dz}{z} O(z) d\left(\frac{x}{z}\right). \quad (1.1)$$

However, very often, typically in the presence of mass effects, the integration phase space is modified and the convolution in eq. (1.1), limiting ourselves to the leftmost identity, is generalised as:

$$I(x, \eta) = x \int_{x/\eta}^1 \frac{dz}{z} O(z, \eta) d\left(\frac{x}{\eta z}\right). \quad (1.2)$$

where  $\eta \leq 1$ , with  $\eta = 1$  reproducing eq. (1.1). However, for purposes that will become clear later, we want to write the integral in eq. (1.2) in the form of eq. (1.1), that is in such a way that lower bound of the integral is not the rescaled variable  $x/\eta$  but the physical Bjorken  $x$ . To do so, one needs to perform the change of variable  $y = \eta z$ , so that:

$$I(x, \eta) = \int_x^\eta dy O\left(\frac{y}{\eta}, \eta\right) \frac{x}{y} d\left(\frac{x}{y}\right). \quad (1.3)$$

In order to precompute the expensive part of the integral in eq. (1.5), we use the standard interpolation formula to the distribution  $d$ :

$$\frac{x}{y} d\left(\frac{x}{y}\right) = \sum_{\alpha=0}^{N_x} x_\alpha d(x_\alpha) w_\alpha^{(k)}\left(\frac{x}{y}\right), \quad (1.4)$$

where  $\alpha$  runs over the node of a give grid in  $x$  space and the weights  $w_\alpha$  are typically polynomials of degree  $k$  (*i.e.* Laplace interpolants). Now we use eq. (1.4) in eq. (1.5) and at the same time we limit the computation of the integral  $I$  to the point  $x_\beta$  of the grid used in eq. (1.4). This way we get:

$$I(x_\beta, \eta) = \sum_{\alpha=0}^{N_x} \bar{d}_\alpha \int_{x_\beta}^\eta dy O\left(\frac{y}{\eta}, \eta\right) w_\alpha^{(k)}\left(\frac{x_\beta}{y}\right). \quad (1.5)$$

where we have defined  $\bar{d}_\alpha = x_\alpha d(x_\alpha)$ . In the particular case of the Laplace interpolants (in the APFEL procedure), one can show that:

$$w_\alpha^{(k)}\left(\frac{x_\beta}{y}\right) \neq 0 \quad \text{for} \quad c < y < d, \quad (1.6)$$

with:

$$c = \max(x_\beta, x_\beta/x_{\alpha+1}) \quad \text{and} \quad d = \min(\eta, x_\beta/x_{\alpha-k}), \quad (1.7)$$

and thus eq. (1.5) can be adjusted as:

$$I(x_\beta, \eta) = \sum_{\alpha=0}^{N_x} \bar{d}_\alpha \int_c^d dy O\left(\frac{y}{\eta}, \eta\right) w_\alpha^{(k)}\left(\frac{x_\beta}{y}\right). \quad (1.8)$$

Finally, we change back the integration variable  $z = y/\eta$  so that eq. (1.8) becomes:

$$I(x_\beta, \eta) = \sum_{\alpha=0}^{N_x} \bar{d}_\alpha \underbrace{\left[ \eta \int_{c/\eta}^{d/\eta} dz O(z, \eta) w_\alpha^{(k)}\left(\frac{x_\beta}{\eta z}\right) \right]}_{\Gamma_{\alpha\beta}}. \quad (1.9)$$

The quantity in squared brackets is the interesting bit.

The expressions for the operator  $O$  that we have to deal with have this general form:

$$O(z, \eta) = R(z, \eta) + \sum_{i=0}^n \left[ \frac{\ln^i(1-z)}{1-z} \right]_+ S^{(i)}(z, \eta) + L(\eta) \delta(1-z), \quad (1.10)$$

where  $R$  and  $S^{(i)}$  is a regular function in  $z = 1$ , that is:

$$R(1, \eta) = \lim_{z \rightarrow 1} R(z, \eta) = K(\eta) \quad \text{and} \quad S^{(i)}(1, \eta) = \lim_{z \rightarrow 1} S^{(i)}(z, \eta) = J^{(i)}(\eta), \quad (1.11)$$

being  $K$ ,  $J^{(i)}$ , and  $L$  a finite function of  $\eta$ . Notice that the sum over  $i$  in eq. (1.10) is between zero and  $n$  where  $n$  is typically not bigger than two as it depends on the perturbative order of the expressions. Plugging eq. (1.10) into the definition of  $\Gamma_{\alpha\beta}$  in eq. (1.9) and making use of the definition of plus-prescription, we obtain:

$$\begin{aligned} \Gamma_{\beta\alpha} = & \eta \int_{c/\eta}^{d/\eta} dz \left\{ R(z, \eta) w_{\alpha}^{(k)} \left( \frac{x_{\beta}}{\eta z} \right) + \sum_{i=0}^n \frac{\ln^i(1-z)}{1-z} \left[ S^{(i)}(z, \eta) w_{\alpha}^{(k)} \left( \frac{x_{\beta}}{\eta z} \right) - S^{(i)}(1, \eta) w_{\alpha}^{(k)} \left( \frac{x_{\beta}}{\eta} \right) \right] \right\} \\ & + \eta \left[ \sum_{i=0}^n \frac{1}{(i+1)!} S^{(i)}(1, \eta) \ln^{i+1} \left( 1 - \frac{c}{\eta} \right) + L(\eta) \right] w_{\alpha}^{(k)} \left( \frac{x_{\beta}}{\eta} \right). \end{aligned} \quad (1.12)$$

Eq. (1.12) expresses the full complexity of the task. However, there are a few remarks that apply to some particular cases and that reduce the complexity. In the case of the ( $\overline{\text{MS}}$ ) splitting functions, there are two main simplifications: the first is that  $\eta = 1$ , the second is that  $n = 0$  in the sums. Considering that:

$$w_{\alpha}^{(k)}(x_{\beta}) = \delta_{\alpha\beta}, \quad (1.13)$$

and that the expressions can be manipulated in such a way that the coefficient of the plus-prescribed term  $S$  is a constant, we have that:

$$\Gamma_{\beta\alpha} = \int_c^d dz \left\{ R(z) w_{\alpha}^{(k)} \left( \frac{x_{\beta}}{z} \right) + \frac{S}{1-z} \left[ w_{\alpha}^{(k)} \left( \frac{x_{\beta}}{z} \right) - \delta_{\alpha\beta} \right] \right\} + [S \ln(1-c) + L] \delta_{\alpha\beta}. \quad (1.14)$$

The same kind of simplifications apply to the case of the Zero-Mass (ZM) coefficient functions with the only exception that the sum over  $i$  extends to  $n$ 's larger than zero. In particular, since we will use exact expressions only up to  $\mathcal{O}(\alpha_s)$ , *i.e.* NLO, we have  $n = 1$  so that:

$$\begin{aligned} \Gamma_{\beta\alpha} = & \int_c^d dz \left\{ R(z) w_{\alpha}^{(k)} \left( \frac{x_{\beta}}{z} \right) + \frac{S^{(0)} + S^{(1)} \ln(1-z)}{1-z} \left[ w_{\alpha}^{(k)} \left( \frac{x_{\beta}}{z} \right) - \delta_{\alpha\beta} \right] \right\} \\ & + \left[ S^{(0)} \ln(1-c) + \frac{1}{2} S^{(1)} \ln^2(1-c) + L \right] \delta_{\alpha\beta}. \end{aligned} \quad (1.15)$$

As far as the massive coefficient functions up to  $\mathcal{O}(\alpha_s^2)$  things can be more complicated and we will discuss it later. In the next section we will consider the case of hadronic observables in the ZM scheme showing how to pre-compute the integral for the double-differential distributions in Drell-Yan production.

## 1.1 Double convolution

In the case of Drell-Yan or Semi-Inclusive DIS (SIDIS), cross sections are the result of double convolutions between partonic cross sections and a pair of non-perturbative distributions. In this case the kind of expressions we are dealing with has the following form:

$$\begin{aligned} J(x_{\delta}, x_{\gamma}) &= x_{\delta} x_{\gamma} \int_{x_{\delta}}^1 \frac{dy_1}{y_1} \int_{x_{\gamma}}^1 \frac{dy_2}{y_2} O(y_1, y_2) d^{(1)} \left( \frac{x_{\delta}}{y_1} \right) d^{(2)} \left( \frac{x_{\gamma}}{y_2} \right) \\ &= \int_{x_{\delta}}^1 dy_1 \int_{x_{\gamma}}^1 dy_2 O(y_1, y_2) \left[ \frac{x_{\delta}}{y_1} d^{(1)} \left( \frac{x_{\delta}}{y_1} \right) \right] \left[ \frac{x_{\gamma}}{y_2} d^{(2)} \left( \frac{x_{\gamma}}{y_2} \right) \right] \\ &= \sum_{\alpha=0}^{N_x} \sum_{\beta=0}^{N_x} \underbrace{\bar{d}_{\beta}^{(1)} \bar{d}_{\alpha}^{(2)} \left[ \int_{x_{\delta}}^1 dy_1 \int_{x_{\gamma}}^1 dy_2 O(y_1, y_2) w_{\beta}^{(k)} \left( \frac{x_{\delta}}{y_1} \right) w_{\alpha}^{(k)} \left( \frac{x_{\gamma}}{y_2} \right) \right]}_{\Theta^{\beta\delta, \alpha\gamma}}. \end{aligned} \quad (1.16)$$

In Eq. (1.16) we assume that there are no mass corrections and thus the convolutions take the simplest form. In the case of double convolutions, the partonic cross sections  $O$  are functions of two variables,  $y_1$  and  $y_2$ , and, as in the case of the single convolutions, they get three kinds of contributions in both these variables: local terms proportional to  $\delta$ -functions, singular terms proportional to plus-prescribed functions, and regular terms. The complication here is that these contributions from the two variables  $y_1$  and  $y_2$  mix and thus, for example,

terms local in  $y_1$  and singular in  $y_2$  appear. It is thus necessary to identify the general structure of the function  $O$  to see whether it is possible to decompose the double operator  $\Theta^{\beta\delta,\alpha\gamma}$  into products of single operators of the kind  $\Gamma^{\beta\delta}$  and  $\Gamma^{\alpha\gamma}$ .

In the case of the NLO corrections to SIDIS, the general structure of the function  $O$  can be inferred looking at Eqs. (C.2)-(C.7) of Ref. [1]:

$$\begin{aligned}
O(y_1, y_2) = & \text{LL} \delta(1-y_1)\delta(1-y_2) + \text{LS} \delta(1-y_1) \left[ \frac{\ln(1-y_2)}{1-y_2} \right]_+ + \delta(1-y_1) \text{LR}(y_2) \\
& + \text{SL} \left[ \frac{\ln(1-y_1)}{1-y_1} \right]_+ \delta(1-y_2) + \text{SS} \left[ \frac{1}{1-y_1} \right]_+ \left[ \frac{1}{1-y_2} \right]_+ + \left[ \frac{1}{1-y_1} \right]_+ \text{SR}(y_2) \\
& + \text{RL}(y_1) \delta(1-y_2) + \text{RS}(y_1) \left[ \frac{1}{1-y_2} \right]_+ + \sum_i K_i R_i^{(1)}(y_1) R_i^{(2)}(y_2).
\end{aligned} \tag{1.17}$$

It is clear that in Eq. (1.17) all terms factorise into a part only depending on  $y_1$  and a part only depending on  $y_2$ . This is crucial to use the same technology developed above for the single convolutions. Plugging Eq. (1.17) into Eq. (1.16), one finds that:

$$\begin{aligned}
\Theta^{\beta\delta,\alpha\gamma} = & C_{\text{LL}} \delta_{\beta\delta} \delta_{\alpha\gamma} + [\text{LS} + \text{SS} \ln(1-c_{\beta\delta})] \delta_{\beta\delta} \Gamma_{\alpha\gamma}^{\text{S1}} + \delta_{\beta\delta} [\Gamma_{\alpha\gamma}^{\text{LR}} + \ln(1-c_{\beta\delta}) \Gamma_{\alpha\gamma}^{\text{SR}}] \\
& + \Gamma_{\beta\delta}^{\text{S1}} [\text{SL} + \text{SS} \ln(1-c_{\alpha\gamma})] \delta_{\alpha\gamma} + \text{SS} \Gamma_{\beta\delta}^{\text{S0}} \Gamma_{\alpha\gamma}^{\text{S0}} + \Gamma_{\beta\delta}^{\text{S0}} \Gamma_{\alpha\gamma}^{\text{SR}} \\
& + [\Gamma_{\beta\delta}^{\text{RL}} + \ln(1-c_{\alpha\gamma}) \Gamma_{\beta\delta}^{\text{RS}}] \delta_{\alpha\gamma} + \Gamma_{\beta\delta}^{\text{RS}} \Gamma_{\alpha\gamma}^{\text{S0}} + \sum_i K_i \Gamma_{\beta\delta}^{R_i^{(1)}} \Gamma_{\alpha\gamma}^{R_i^{(2)}}
\end{aligned} \tag{1.18}$$

with:

$$C_{\text{LL}} = \text{LL} + \text{LS} \frac{1}{2} \ln^2(1-c_{\alpha\gamma}) + \text{SL} \frac{1}{2} \ln^2(1-c_{\beta\delta}) + \text{SS} \ln(1-c_{\beta\delta}) \ln(1-c_{\alpha\gamma}) \tag{1.19}$$

and:

$$\begin{aligned}
\Gamma_{\alpha\beta}^{\text{S}n} &= \int_{c_{\alpha\beta}}^{d_{\alpha\beta}} dz \frac{\ln^n(1-z)}{1-z} \left[ w_\beta^{(k)} \left( \frac{x_\alpha}{z} \right) - \delta_{\alpha\beta} \right] \\
\Gamma_{\alpha\beta}^f &= \int_{c_{\alpha\beta}}^{d_{\alpha\beta}} dz f(z) w_\beta^{(k)} \left( \frac{x_\alpha}{z} \right)
\end{aligned} \tag{1.20}$$

where  $f$  is a regular function and the integration bounds are defined as:

$$c_{\alpha\beta} = \max(x_\alpha, x_\alpha/x_{\beta+1}) \quad \text{and} \quad d_{\alpha\beta} = \min(1, x_\alpha/x_{\beta-k}). \tag{1.21}$$

Despite writing the hard cross sections as in Eq. (1.17) requires a bit of analytical work, it is extremely beneficial from the numerical point of view because it allows us to reduce complicated expressions of two variables to a combination of products of functions of one single variable that can be treated as discussed in the previous sections. In particular, the pre-computation is limited to the integrals in Eq. (1.20). We stress however that the structure in Eq. (1.17) is valid only for the NLO corrections to SIDIS and (possibly) DY. We expect that, when going beyond NLO, the structure of the hard cross sections would get more complicated essentially for the presence of more involved plus-prescribed functions. However, we also expect that the same kind of “factorisation” can be implemented.

We now apply this technology to the specific cases of SIDIS first and DY second to NLO in QCD.

## 1.2 Semi-inclusive deep inelastic scattering (SIDIS)

the structure of the SIDIS observables of the expressions for the respective hard cross sections can be found in Ref. [1]. Following this paper, the SIDIS differential cross section for the exchange of a virtual photon can be written as:

$$\frac{d^3\sigma}{dx dy dz} = \frac{2\pi\alpha^2}{xyQ^2} \left[ (1 + (1-y)^2) 2xF_1(x, z, Q^2) + 2(1-y)F_L(x, z, Q^2) \right], \tag{1.22}$$

with  $Q^2 = -q^2$  the (negative) virtuality of the exchanged photon,  $x$  and  $z$  the momentum fractions of PDFs and FFs, and  $Q^2 = xys$  the definition of the inelasticity  $y$  in terms of the squared collision energy in the center of mass  $s$ . Notice that, as compared to Ref. [1], we have absorbed into the definition of  $F_L$  a factor  $x$  as customary in the definition of the longitudinal structure function in inclusive DIS.

We now use the Callan-Gross relation:

$$F_2 = 2xF_1 + F_L \quad (1.23)$$

to replace  $2xF_1$  with  $F_2$  in Eq. (1.22):

$$\frac{d^3\sigma}{dx dy dz} = \frac{2\pi\alpha^2}{xyQ^2} [Y_+ F_2(x, z, Q^2) - y^2 F_L(x, z, Q^2)] , \quad (1.24)$$

where we have defined:

$$Y_+ = 1 + (1 - y)^2 . \quad (1.25)$$

It is also useful to write Eq. (1.26) as differential in  $x$ ,  $Q^2$ , and  $z$ :

$$\frac{d^3\sigma}{dx dQ^2 dz} = \frac{2\pi\alpha^2}{xQ^4} [Y_+ F_2(x, z, Q) - y^2 F_L(x, z, Q)] . \quad (1.26)$$

The structure functions  $F_2$  and  $F_L$  are given at NLO by:

$$\begin{aligned} F_{2,L}(x, z, Q) &= x \sum_{q,\bar{q}} e_q^2 \left[ q(x, Q) \otimes C_{qq}^{2,L}(x, z) \otimes D_q(z, Q) \right. \\ &\quad \left. + q(x, Q) \otimes C_{gq}^{2,L}(x, z) \otimes D_g(z, Q) + g(x, Q) \otimes C_{qg}^{2,L}(x, z) \otimes D_q(z, Q) \right] , \end{aligned} \quad (1.27)$$

where  $\{q, g\}$  are the quark and gluon PDFs and  $\{D_q, D_g\}$  are the quark and gluon FFs,  $e_q$  is the electric charge of the quark  $q$  and  $\{C_{qq}^{2,L}, C_{gq}^{2,L}, C_{qg}^{2,L}\}$  are the relevant partonic cross sections. The partonic cross sections allow for a perturbative expansion in power of  $\alpha_s$ :

$$C = \sum_{n=0} \left( \frac{\alpha_s}{4\pi} \right)^n C^{(n)} \quad (1.28)$$

that we truncate to NLO, *i.e.* to  $n = 1$ . At LO ( $n = 0$ ) we have the simple expression:

$$C_{qq}^{2,(0)}(x, z) = \delta(1-x)\delta(1-z) \quad \text{and all others zero.} \quad (1.29)$$

At NLO ( $n = 1$ ) we take the expressions from Appendix C of Ref. [1] being careful to take into account an addition factor two due to the difference in the expansion parameter ( $\alpha_s/4\pi$  rather than  $\alpha_s/4\pi$ ) and to combine the expressions for  $F_1$  and  $F_L$  using Eq. (1.23) to obtain the partonic cross sections for  $F_2$ . We start with the partonic cross sections for  $F_L$  that read:

$$\begin{aligned} C_{qq}^{L,(1)} &= 8C_F xz , \\ C_{gq}^{L,(1)} &= 8C_F x(1-z) , \\ C_{qg}^{L,(1)} &= 8x(1-x) , \end{aligned} \quad (1.30)$$

while those for  $F_2$  read:

$$\begin{aligned}
\frac{C_{qq}^{2,(1)}}{2C_F} &= -8\delta(1-x)\delta(1-z) + 2\delta(1-x) \left( \frac{\ln(1-z)}{1-z} \right)_+ + \delta(1-x) \left[ \frac{1+z^2}{1-z} \ln z + (1-z) - (1+z) \ln(1-z) \right] \\
&+ 2 \left( \frac{\ln(1-x)}{1-x} \right)_+ \delta(1-z) + 2 \left( \frac{1}{1-x} \right)_+ \left( \frac{1}{1-z} \right)_+ - \left( \frac{1}{1-x} \right)_+ (1+z) \\
&+ \left[ -\frac{1+x^2}{1-x} \ln x + (1-x) - (1+x) \ln(1-x) \right] \delta(1-z) - (1+x) \left( \frac{1}{1-z} \right)_+ + (2+6xz) \\
\frac{C_{gq}^{2,(1)}}{2C_F} &= \delta(1-x) \left[ \frac{1+(1-z)^2}{z} \ln[z(1-z)] + z \right] \\
&+ \left( \frac{1}{1-x} \right)_+ \frac{1+(1-z)^2}{z} \\
&+ 2(1+3x) - 6xz - (1+x) \frac{1}{z} \\
C_{qg}^{2,(1)} &= \left[ (x^2 + (1-x)^2) \ln \left( \frac{1-x}{x} \right) + 2x(1-x) \right] \delta(1-z) + (x^2 + (1-x)^2) \left( \frac{1}{1-z} \right)_+ \\
&+ 2(1+6x+2x^2) + (x^2 + (1-x)^2) \frac{1}{z}
\end{aligned} \tag{1.31}$$

By inspection of Eqs. (1.30) and (1.31) we can deduce the various coefficients of Eq. (1.17). In particular, as usual,  $F_L$  involves only regular functions so that all contributions are zero but the fully regular ones:

$$\begin{aligned}
C_{qq}^{L,(1)}(x, z) &: K_1 = 8C_F, \quad R_1^{(1)}(x) = x, \quad R_1^{(2)}(z) = z, \\
C_{gq}^{L,(1)}(x, z) &: K_1 = 8C_F, \quad R_1^{(1)}(x) = x, \quad R_1^{(2)}(z) = 1-z, \\
C_{qg}^{L,(1)}(x, z) &: K_1 = 1, \quad R_1^{(1)}(x) = x(1-x), \quad R_1^{(2)}(z) = 1.
\end{aligned} \tag{1.32}$$

For  $F_2$  the situation is more complicate but we can still identify the different contributions:

$$\begin{aligned}
C_{qq}^{2,(1)} : \quad & \text{LL} = -16C_F, \quad \text{LS} = 4C_F, \quad \text{LR}(z) = 2C_F \left[ \frac{1+z^2}{1-z} \ln z + (1-z) - (1+z) \ln(1-z) \right] \\
& \text{SL} = 4C_F, \quad \text{SS} = 4C_F, \quad \text{SR}(z) = -2C_F(1+z) \\
& \text{RL}(x) = 2C_F \left[ -\frac{1+x^2}{1-x} \ln x + (1-x) - (1+x) \ln(1-x) \right], \quad \text{RS}(x) = -2C_F(1+x), \\
& \left\{ K_1 = 4C_F, R_1^{(1)}(x) = 1, R_1^{(2)}(z) = 1 \right\}, \\
& \left\{ K_2 = 12C_F, R_2^{(1)}(x) = x, R_2^{(2)}(z) = z \right\}, \\
C_{gq}^{2,(1)} : \quad & \text{LR}(z) = 2C_F \left[ \frac{1+(1-z)^2}{z} \ln[z(1-z)] + z \right], \\
& \text{SR}(z) = 2C_F \left[ \frac{1+(1-z)^2}{z} \right], \\
& \left\{ K_1 = 4C_F, R_1^{(1)}(x) = 1+3x, R_1^{(2)}(z) = 1 \right\}, \\
& \left\{ K_2 = -12C_F, R_2^{(1)}(x) = x, R_2^{(2)}(z) = z \right\}, \\
& \left\{ K_3 = -2C_F, R_3^{(1)}(x) = 1+x, R_3^{(2)}(z) = \frac{1}{z} \right\}, \\
C_{qg}^{2,(1)} : \quad & \text{RL}(x) = [x^2 + (1-x)^2] \ln \left( \frac{1-x}{x} \right) + 2x(1-x), \quad \text{RS}(x) = x^2 + (1-x)^2, \\
& \left\{ K_1 = 2, R_1^{(1)}(x) = 1+6x+2x^2, R_1^{(2)}(z) = 1 \right\}, \\
& \left\{ K_2 = 1, R_2^{(1)}(x) = x^2 + (1-x)^2, R_2^{(2)}(z) = \frac{1}{z} \right\}.
\end{aligned} \tag{1.33}$$

Analogously, for the only LO partonic cross sections we find that:

$$C_{qq}^{2,(0)} : \quad \text{LL} = 1. \tag{1.34}$$

All the coefficients that are not mentioned are equal to zero. We can now implement explicitly Eq. (1.18). The one think that is left to sort out is the structure of the structure functions in terms of the appropriate PDF and FF combinations. Considering Eq. (1.27), we observe that none of the coefficient functions depends on the particular quark flavour (this is a feature of the ZM scheme). Therefore, simplifying the notation, we can rewrite Eq. (1.27) as:

$$F/x = C_{qq} \sum_q e_q^2 [qD_q + \bar{q}D_{\bar{q}}] + C_{gq} \sum_q e_q^2 [q + \bar{q}] D_g + C_{qg} \sum_q e_q^2 [D_q + D_{\bar{q}}], \tag{1.35}$$

where now the sums run only over the quark flavours and not the antiflavours.

## References

- [1] D. de Florian, M. Stratmann and W. Vogelsang, Phys. Rev. D **57** (1998) 5811 doi:10.1103/PhysRevD.57.5811 [hep-ph/9711387].