

Stability and instability in inverse problems

Mikhail I. Isaev

supervisor: Roman G. Novikov

Moscow Institute of Physics and Technology (the state university)

Centre de Mathématiques Appliquées, École Polytechnique

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Plan of the presentation

- The Gel'fand inverse problem with boundary measurements represented as a Dirichlet-to-Neumann map
- The Gel'fand inverse problem with boundary measurements represented as an impedance boundary map (Robin-to-Robin map)
- Inverse scattering problems

Basic assumptions

Consider the Schrödinger equation

$$-\Delta\psi + v(x)\psi = E\psi \quad \text{for } x \in D, \quad (1)$$

where

- D is an open bounded domain in \mathbb{R}^d ,
- $d \geq 2$,
- $\partial D \in C^2$,
- $v \in L^\infty(D)$.

Statement of the problem

Let

$$\mathcal{C}_v(E) = \left\{ \left(\psi|_{\partial D}, \frac{\partial \psi}{\partial \nu} |_{\partial D} \right) : \begin{array}{l} \text{for all sufficiently regular solutions } \psi \text{ of} \\ \text{equation (1) in } \bar{D} = D \cup \partial D \end{array} \right\}.$$

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Problem 1.

- Given $\mathcal{C}_v(E)$.
- Find v .

Standart representation of the Cauchy data

The Dirichlet-to-Neumann map $\hat{\Phi}_v(E)$ is defined by

$$\hat{\Phi}_v(E)(\psi|_{\partial D}) = \frac{\partial \psi}{\partial \nu}|_{\partial D}.$$

Here we assume also that

E is not a Dirichlet eigenvalue for operator $-\Delta + v$ in D .

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Problem 1a.

- Given $\hat{\Phi}_v(E)$.
- Find v .

Questions

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- Reconstruction.

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- Reconstruction.
- Stability: there is some function ϕ such that

$$\|v_2 - v_1\|_{L^\infty(D)} \leq \phi(\|\hat{\Phi}_{v_2}(E) - \hat{\Phi}_{v_1}(E)\|),$$
$$\phi(t) \rightarrow 0 \text{ as } t \rightarrow +0.$$

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Problem 1 was formulated for the first time by Gel'fand (1954).

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<i>Uniqueness:</i>	Novikov (1988)	Bukhgeim (2008)
<i>Reconstruction:</i>	Novikov (1988)	Bukhgeim (2008)
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The Calderón inverse problem (of the electrical impedance tomography):

Calderón (1980), Slichter (1933), Tikhonov (1949), Druskin (1982), Sylvester-Uhlmann (1987), Nachman (1996), Liu (1997).

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- Mandache (2001), optimality of logarithmic stability results in the case of zero energy (the Gel'fand-Calderón inverse problem) up to the value of some exponent.
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- **Lipschitz stability in the case of piecewise constant potentials:**

- Alessandrini-Vessella (2005), the Calderón inverse problem.
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- **Regularity and/or energy dependent stability estimates:**

- Novikov (2011), effectivization of the result of Alessandrini (1988).
- Novikov (1998, 2005, 2008), Isakov (2011), Santacesaria (2013), the phenomena of increasing stability for the high-energy case.

Logarithmic and Hölder-logarithmic stability estimates

Theorem 1 (Isaev, Novikov [IN1]).

Let basic assumptions of Problem 1a hold and

- $d \geq 3$, $m > d$, $N > 0$ and $\text{supp } v_j \subset D$,
- $v_j \in W^{m,1}(\mathbb{R}^d)$ and $\|v_j\|_{m,1} \leq N$, $j = 1, 2$,

Then, for any $s \geq 0$, $s \leq (m - d)/d$,

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq C \left(\ln \left(3 + \delta^{-1} \right) \right)^{-s}, \quad (2)$$

where constant C depends only on N , D , m , s , E ,

$$\delta = \|\hat{\Phi}_{v_1}(E) - \hat{\Phi}_{v_2}(E)\|_{\mathbb{L}^\infty(\partial D) \rightarrow \mathbb{L}^\infty(\partial D)}.$$

In addition, for $E \geq 0$, $\tau \in (0, 1)$ and any $\alpha, \beta \geq 0$, $\alpha + \beta \leq (m - d)/d$,

$$\|v_1 - v_2\|_{L^\infty(D)} \leq A(1 + \sqrt{E})\delta^\tau + B(1 + \sqrt{E})^{-\alpha} \left(\ln \left(3 + \delta^{-1} \right) \right)^{-\beta}, \quad (3)$$

where constants $A, B > 0$ depend only on N , D , m , α , β , τ .

Relations to known results

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- Estimate (2) with $s = s_0$ is a variation of the result of Alessandrini (1988). This stability result was improved by Novikov (2011) for $\mathbf{E} = \mathbf{0}$ and $\mathbf{d} = \mathbf{3}$: estimate (2) holds for $s = s_2$ also. A principal advantage is that

$$s_1 \rightarrow +\infty \text{ and } s_2 \rightarrow +\infty \quad \text{as} \quad m \rightarrow +\infty.$$

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- Mandache (2001), estimate (2) **can not hold** for $E = 0$ and dimension $d \geq 2$
 - when $s > 2m - \frac{m}{d}$ for real-valued potentials,
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 - when $s > 2m - \frac{m}{d}$ for real-valued potentials,
 - when $s > m$ for complex-valued potentials.
- If we put $\alpha = s_1$, $\beta = 0$ in estimate (3), we get that

$$\|v_1 - v_2\|_{L^\infty(D)} \leq A(1 + \sqrt{E})\delta^\tau + B(1 + \sqrt{E})^{-\frac{m-d}{d}}.$$

Similar estimates (but with modified exponent) follows from the approximate reconstruction algorithms of Novikov (1999, 2005).

Instability results

Consider the union of the energy intervals $S = \bigcup_{j=1}^K I_j$ such that DtN maps $\hat{\Phi}_{v_1}(E)$, $\hat{\Phi}_{v_2}(E)$ are correctly defined for any $E \in S$.

It was shown in [Isaev1] that estimate

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq C \sup_{E \in S} (\ln(3 + \delta(E)^{-1}))^{-s},$$

where $C = C(N, D, m, s, S)$, can not hold with $s > 2m$ for real-valued potentials and with $s > m$ for complex potentials.

Instability results

Let $A, B, \alpha, \beta, \kappa, \tau \geq 0$. We consider class of estimates of the type

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq A(1 + \sqrt{E})^\kappa \delta^\tau + B(1 + \sqrt{E})^{-\alpha} \left(\ln(3 + \delta^{-1}) \right)^{-\beta}.$$

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- for $\alpha + 2\beta > 2m$ **can not hold**

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$$\bullet \text{ for } \alpha + \beta \leq \frac{m-d}{d} \quad \text{hold}$$

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$$\bullet \text{ for } \alpha + 2\beta > 2m \quad \text{can not hold}$$

In particular, results of [Isaev2] show optimality of the estimate

$$\|v_1 - v_2\|_{L^\infty(D)} \leq A(1 + \sqrt{E})\delta^\tau + B(1 + \sqrt{E})^{-\frac{m-d}{d}}.$$

The weakness

Bad news: stability estimates given earlier make no sense if

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or too weak if energy E is close to Dirichlet spectrum.

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Idea: let us consider another operator representation of the Cauchy data set

$$\mathcal{C}_v(E) = \left\{ \left(\psi|_{\partial D}, \frac{\partial \psi}{\partial \nu} |_{\partial D} \right) : \begin{array}{l} \text{for all sufficiently regular solutions } \psi \text{ of} \\ \text{equation (1) in } \bar{D} = D \cup \partial D \end{array} \right\} :$$

$$\hat{M}_{c_1, c_2, c_3, c_4} \left(c_1 \psi|_{\partial D} + c_2 \frac{\partial \psi}{\partial \nu} |_{\partial D} \right) = \left(c_3 \psi|_{\partial D} + c_4 \frac{\partial \psi}{\partial \nu} |_{\partial D} \right).$$

Impedance boundary map (Robin-to-Robin map)

Let consider the map $\hat{M}_{\alpha,v}(E)$ defined by

$$\hat{M}_{\alpha,v}[\psi]_{\alpha} = [\psi]_{\alpha-\pi/2}$$

for all suffuciently regular solutions ψ of equation (1) in $\bar{D} = D \cup \partial D$, where

$$[\psi]_{\alpha} = [\psi(x)]_{\alpha} = \cos \alpha \psi(x) - \sin \alpha \frac{\partial \psi}{\partial \nu}|_{\partial D}(x), \quad x \in \partial D.$$

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Problem 1b.

- Given $\hat{M}_{\alpha,v}(E)$ for some fixed E and α .
- Find v .

Impedance boundary map (Robin-to-Robin map)

We have that

- there is not more than a countable number of α such that E is an eigenvalue for the operator $-\Delta + v$ in D with the boundary condition

$$\cos \alpha \psi|_{\partial D} - \sin \alpha \frac{\partial \psi}{\partial \nu}|_{\partial D} = 0,$$

- the map \hat{M}_α is reduced to the Dirichlet-to-Neumann map if $\alpha = 0$ is reduced to the Neumann-to-Dirichlet map if $\alpha = \pi/2$.

Stability estimates for $d \geq 3$

Theorem 2 (Isaev, Novikov [IN2]).

Let basic assumptions of Problem 1b hold and

- $d \geq 3$, $m > d$, $N > 0$ and $\text{supp } v_j \subset D$,
- $v_j \in W^{m,1}(\mathbb{R}^d)$ and $\|v_i\|_{m,1} \leq N$, $j = 1, 2$,

Then, for any $s \geq 0$, $s \leq (m - d)/m$,

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq C_\alpha \left(\ln \left(3 + \delta_\alpha^{-1} \right) \right)^{-s}, \quad (4)$$

where constant $C_\alpha = C_\alpha(N, D, m, s, E)$,

$$\delta_\alpha = \|\hat{M}_{\alpha, v_1}(E) - \hat{M}_{\alpha, v_2}(E)\|_{\mathbb{L}^\infty(\partial D) \rightarrow \mathbb{L}^\infty(\partial D)}.$$

Estimate (4) with $\alpha = 0$ is a variation of the result of Alessandrini (1988).

Stability estimates for $d = 2$

Theorem 3 (Isaev, Novikov [IN2]).

Let basic assumptions of Problem 1b hold and

- $d = 2$, $N > 0$ and $\text{supp } v_j \subset D$,
- $v_j \in C^2(\bar{D})$ and $\|v_j\|_{C^2(\bar{D})} \leq N$, $j = 1, 2$,

Then, for any $0 < s \leq 3/4$,

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq C_\alpha \left(\ln \left(3 + \delta_\alpha^{-1} \right) \right)^{-s} \left(\ln \left(3 \ln \left(3 + \delta_\alpha^{-1} \right) \right) \right)^2,$$

where constant $C_\alpha = C_\alpha(N, D, s, E)$,

$$\delta_\alpha = \|\hat{M}_{\alpha, v_1}(E) - \hat{M}_{\alpha, v_2}(E)\|_{\mathbb{L}^\infty(\partial D) \rightarrow \mathbb{L}^\infty(\partial D)}.$$

Theorem 3 for $\alpha = 0$ was given by Novikov-Santacesaria (2010) with $s = 1/2$ and by Santacesaria (2012) with $s = 3/4$.

Stability of determining a potential from its Cauchy data

Theorems 2 and 3 imply, in particular, that

- For $d \geq 3$ and $0 < s \leq (m - d)/m$

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq \min_{\alpha \in \mathbb{R}} C_\alpha \left(\ln \left(3 + \delta_\alpha^{-1} \right) \right)^{-s}.$$

- For $d = 2$ and $0 < s \leq 3/4$,

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(D)} \leq \min_{\alpha \in \mathbb{R}} C_\alpha \left(\ln \left(3 + \delta_\alpha^{-1} \right) \right)^{-s} \left(\ln \left(3 \ln \left(3 + \delta_\alpha^{-1} \right) \right) \right)^2.$$

Idea of the proofs

For any sufficiently regular solutions ψ_1 and ψ_2 of equation (1) in $\bar{D} = D \cup \partial D$ with $v = v_1$ and $v = v_2$, respectively, the following identity holds (see [IN2]):

$$\int_D (v_1 - v_2) \psi_1 \psi_2 dx = \int_{\partial D} [\psi_1]_\alpha \left(\hat{M}_{\alpha, v_1}(E) - \hat{M}_{\alpha, v_2}(E) \right) [\psi_2]_\alpha dx. \quad (5)$$

Identity (5) for $\alpha = \mathbf{0}$ is reduced to Alessandrini's identity.

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Corollary (Isaev, Novikov [IN2]).

Under basic assumptions real-valued potential v is uniquely determined by its Cauchy data $\mathcal{C}_v(E)$ at fixed real energy E in dimension $d \geq 2$.

To our knowledge the result of this corollary for $d \geq 3$ was not yet completely proved in the literature.

Schema of reconstruction of a potential v from $\hat{M}_{v,\alpha}(E)$

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- 3 $S_E \rightarrow v$ as described by Grinevich (1988, 2000), Henkin-Novikov (1987), Novikov (1992 – 2009), Novikov-Santacesaria (2013).

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- ③ $S_E \rightarrow v$ as described by Grinevich (1988, 2000), Henkin-Novikov (1987), Novikov (1992 – 2009), Novikov-Santacesaria (2013).

In addition, numerical efficiency of related inverse scattering techniques was shown by the research group headed by V.A. Burov (2000, 2008, 2009, 2012), see also Bikowski-Knudsen-Mueller (2011).

Basic assumptions

Consider the three-dimensional stationary acoustic equation at frequency ω in an inhomogeneous medium with refractive index n

$$\Delta\psi + \omega^2 n(x)\psi = 0, \quad x \in \mathbb{R}^3, \quad \omega > 0, \quad (6)$$

where

- $(1 - n) \in W^{m,1}(\mathbb{R}^3)$ for some $m > 3$,
- $\operatorname{Im} n(x) \geq 0, \quad x \in \mathbb{R}^3,$
- $\operatorname{supp} (1 - n) \subset B_{r_1}$ for some $r_1 > 0$,

where $W^{m,1}(\mathbb{R}^3)$ denotes the Sobolev space of m -times smooth functions in \mathbb{L}^1 and B_r is the open ball of radius r centered at 0 .

The Green function

Let $G^+(x, y, \omega)$ denote the Green function for the operator $\Delta + \omega^2 n(x)$ with the Sommerfeld radiation condition:

$$\begin{aligned}
 (\Delta + \omega^2 n(x)) G^+(x, y, \omega) &= \delta(x - y), \\
 \lim_{|x| \rightarrow \infty} |x| \left(\frac{\partial G^+}{\partial |x|}(x, y, \omega) - i\omega G^+(x, y, \omega) \right) &= 0, \\
 &\text{uniformly for all directions } \hat{x} = x/|x|, \\
 x, y &\in \mathbb{R}^3, \quad \omega > 0.
 \end{aligned}$$

It is known that, under basic assumptions, the function G^+ is uniquely specified, see, for example, Colton-Kress (1998), Hähner-Hohage (2001).

Near-field inverse scattering problem

We consider, in particular, the following near-field inverse scattering problem for equation (6):

Problem 2.

- Given G^+ on $\partial B_r \times \partial B_r$ for fixed $\omega > 0$ and $r > r_1$.
- Find n on B_{r_1} .

Scattering amplitude

Consider also the solutions $\psi^+(x, k)$, $x \in \mathbb{R}^3$, $k \in \mathbb{R}^3$, $k^2 = \omega^2$, of equation (6) specified by the following asymptotic condition:

$$\psi^+(x, k) = e^{ikx} - 2\pi^2 \frac{e^{i|k||x|}}{|x|} f\left(k, |k| \frac{x}{|x|}\right) + o\left(\frac{1}{|x|}\right) \quad (7)$$

$$\text{as } |x| \rightarrow \infty \left(\text{uniformly in } \frac{x}{|x|} \right),$$

with some a priory unknown f .

The function f on $\mathcal{M}_\omega = \{k \in \mathbb{R}^3, l \in \mathbb{R}^3 : k^2 = l^2 = \omega^2\}$ arising in (7) is the classical scattering amplitude for equation (6).

Far-field inverse scattering problem

In addition to Problem 2, we consider also the following far-field inverse scattering problem for equation (6):

Problem 3.

- Given f on \mathcal{M}_ω for some fixed $\omega > 0$.
- Find n on B_{r_1} .

Far-field inverse scattering problem

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- It was shown by Berezanskii (1958) that the near-field scattering data of Problem 2 are uniquely determined by the far-field scattering data of Problem 3 and vice versa.
 - Global uniqueness for Problems 2 and 3 was proved for the first time in Novikov (1988); in addition, this proof is constructive.
 - Stability estimates were given for the first time by Stefanov(1990).

Stability estimate

Theorem 4 (Isaev, Novikov [IN4]).

Let $N > 0$ and $r > r_1$ be fixed constants. Then there exists a positive constant C (depending only on m, ω, r_1, r and N) such that for all refractive indices n_1, n_2 satisfying

- $\|1 - n_1\|_{m,1}, \|1 - n_2\|_{m,1} < N,$
- $\text{supp}(1 - n_1), \text{supp}(1 - n_2) \subset B_{r_1},$

the following estimate holds:

$$\|n_1 - n_2\|_{\mathbb{L}^\infty(\mathbb{R}^3)} \leq C \left(\ln \left(3 + \delta^{-1} \right) \right)^{-s}, \quad s = \frac{m-3}{3}, \quad (8)$$

where $\delta = \|G_1^+ - G_2^+\|_{\mathbb{L}^2(\partial B_r \times \partial B_r)}$ and G_1^+, G_2^+ are the near-field scattering data for the refractive indices n_1, n_2 , respectively, at fixed frequency ω .

For some regularity dependent s but always smaller than 1 the stability estimate of Theorems 4 was proved by Hähner-Hohage (2001).

Stability estimate

Theorem 5 (Isaev, Novikov [IN4]).

Let $N > 0$ and $0 < \epsilon < \frac{m-3}{3}$ be fixed constants. Then there exists a positive constant C (depending only on m, ϵ, ω, r_1 and N) such that for all refractive indices n_1, n_2 satisfying

- $\|1 - n_1\|_{m,1}, \|1 - n_2\|_{m,1} < N,$
- $\text{supp}(1 - n_1), \text{supp}(1 - n_2) \subset B_{r_1},$

the following estimate holds:

$$\|n_1 - n_2\|_{L^\infty(\mathbb{R}^3)} \leq C \left(\ln \left(3 + \delta^{-1} \right) \right)^{-s+\epsilon}, \quad s = \frac{m-3}{3}, \quad (9)$$

where $\delta = \|f_1 - f_2\|_{L^2(\mathcal{M}_\omega)}$ and f_1, f_2 denote the scattering amplitudes for the refractive indices n_1, n_2 , respectively, at fixed frequency ω .

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Solution of the open problem

Possibility of estimates (8), (9) with $\mathbf{s} > \mathbf{1}$ was formulated by Hähner-Hohage (2001) as an open problem.

Solution of the open problem

Possibility of estimates (8), (9) with $s > 1$ was formulated by Hähner-Hohage (2001) as an open problem.

Our estimates (8), (9) with $s = \frac{m-3}{3}$ give a solution of this problem. Indeed,

$$s = \frac{m-3}{3} \rightarrow +\infty \quad \text{as} \quad m \rightarrow +\infty.$$

Instability result

Result of Stefanov (1990): for some s always smaller than 1

$$\|n_1 - n_2\|_{\mathbb{L}^\infty(\mathbb{R}^3)} \leq C \left(\ln \left(3 + \|f_1 - f_2\|_S^{-1} \right) \right)^{-s},$$

where some special norm $\|f_1 - f_2\|_S$ is used and

$$\|f_1 - f_2\|_{\mathbb{L}^2(\mathcal{M}_\omega)} \leq c \|f_1 - f_2\|_S.$$

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It was shown in [Isaev3] that for any interval $I = [\omega_1, \omega_2]$, $\omega_1 > 0$, estimate

$$\|n_1 - n_2\|_{\mathbb{L}^\infty(D)} \leq C \sup_{\omega \in I} \left(\ln(3 + \|f_1 - f_2\|_S^{-1}) \right)^{-s}$$

where $C = C(N, D, m, I)$, **can not hold** with $s > 2m$ in the case of the scattering amplitude given on the interval of frequencies and with $s > 5m/3$ in the case of fixed frequency.

Basic assumptions

Now we focus on inverse scattering for the Schrödinger equation

$$L\psi = E\psi, \quad L = -\Delta + v(x), \quad x \in \mathbb{R}^d, \quad d \geq 2, \quad (10)$$

where

- v is real-valued, $v \in L^\infty(\mathbb{R}^d)$
- $v(x) = O(|x|^{-d-\varepsilon})$, $|x| \rightarrow \infty$, for some $\varepsilon > 0$.

The Green function

Consider the resolvent $\mathbf{R}(\mathbf{E})$ of the Schrödinger operator \mathbf{L} in $\mathbb{L}^2(\mathbb{R}^d)$:

$$\mathbf{R}(\mathbf{E}) = (\mathbf{L} - \mathbf{E})^{-1}, \quad \mathbf{E} \in \mathbb{C} \setminus \sigma(\mathbf{L}).$$

Let $\mathbf{R}(\mathbf{x}, \mathbf{y}, \mathbf{E})$ denote the Schwartz kernel of $\mathbf{R}(\mathbf{E})$ as an integral operator. Consider also

$$\mathbf{R}^+(\mathbf{x}, \mathbf{y}, \mathbf{E}) = \mathbf{R}(\mathbf{x}, \mathbf{y}, \mathbf{E} + i0), \quad \mathbf{x}, \mathbf{y} \in \mathbb{R}^d, \quad \mathbf{E} \in \mathbb{R}_+.$$

We recall that in the framework of equation (10) the function $\mathbf{R}^+(\mathbf{x}, \mathbf{y}, \mathbf{E})$ describes scattering of the spherical waves

$$\mathbf{R}_0^+(\mathbf{x}, \mathbf{y}, \mathbf{E}) = -\frac{i}{4} \left(\frac{\sqrt{\mathbf{E}}}{2\pi|\mathbf{x} - \mathbf{y}|} \right)^{\frac{d-2}{2}} H_{\frac{d-2}{2}}^{(1)}(\sqrt{\mathbf{E}}|\mathbf{x} - \mathbf{y}|),$$

generated by a source at \mathbf{y} (where $H_{\mu}^{(1)}$ is the Hankel function of the first kind of order μ). We recall also that $\mathbf{R}^+(\mathbf{x}, \mathbf{y}, \mathbf{E})$ is the Green function for $\mathbf{L} - \mathbf{E}$, $\mathbf{E} \in \mathbb{R}_+$, with the Sommerfeld radiation condition at infinity.

Near-field inverse scattering problem

In addition, the function

$$S^+(x, y, E) = R^+(x, y, E) - R_0^+(x, y, E), \\ x, y \in \partial B_r, E \in \mathbb{R}_+, r \in \mathbb{R}_+,$$

is considered as near-field scattering data for equation (10).

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We consider, in particular, the following near-field inverse scattering problem for equation (10):

Problem 4.

- Given S^+ on $\partial B_r \times \partial B_r$ for some fixed $r, E \in \mathbb{R}_+$.
- Find v on B_r .

This problem can be considered under the assumption that v is a priori known on $\mathbb{R}^d \setminus B_r$. We consider Problem 4 under the assumption that $v \equiv 0$ on $\mathbb{R}^d \setminus B_r$ for some fixed $r \in \mathbb{R}_+$.

Approaches to the problem

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- It is well-known that the near-field scattering data of Problem 4 uniquely and efficiently determine the scattering amplitude f for equation (10) at fixed energy E , see Berezanskii (1958).

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- It is also known that the near-field data of Problem 4 uniquely determine the Dirichlet-to-Neumann map in the case when E is not a Dirichlet eigenvalue for operator L in B_r , see Nachman (1988), Novikov (1988).

Hölder-logarithmic stability estimate for $d \geq 3$

Theorem 6 ([Isaev4]).

Let $E > 0$ and $r > r_1 > 0$ be given constants. Let dimension $d \geq 3$ and potentials v_1, v_2 be real-valued such that

- $v_j \in W^{m,1}(\mathbb{R}^d)$, $m > d$, $\text{supp } v_j \subset B_{r_1}$,
- $\|v_j\|_{m,1} \leq N$ for some $N > 0$, $j = 1, 2$.

Let $S_1^+(E)$ and $S_2^+(E)$ denote the near-field scattering data for v_1 and v_2 , respectively. Then for $\tau \in (0, 1)$ and any $s \in [0, s_1]$ the following estimate holds:

$$\|v_2 - v_1\|_{\mathbb{L}^\infty(B_r)} \leq A(1 + E)^{\frac{5}{2}} \delta^\tau + B(1 + E)^{\frac{s-s_1}{2}} \left(\ln \left(3 + \delta^{-1} \right) \right)^{-s},$$

where $s_1 = \frac{m-d}{d}$, $\delta = \|S_1^+(E) - S_2^+(E)\|_{\mathbb{L}^2(\partial B_r \times \partial B_r)}$, and constants $A, B > 0$ depend only on N, m, d, r, τ .

Logarithmic stability estimate for $d = 2$

Theorem 7 ([Isaev4]).

Let $E > 0$ and $r > r_1 > 0$ be given constants. Let dimension $d = 2$ and potentials v_1, v_2 be real-valued such that

- $v_j \in C^2(\mathbb{R}^d)$, $\text{supp } v_j \subset B_{r_1}$,
- $\|v_j\|_{m,1} \leq N$ for some $N > 0$, $j = 1, 2$.

Let $S_1^+(E)$ and $S_2^+(E)$ denote the near-field scattering data for v_1 and v_2 , respectively. Then

$$\|v_1 - v_2\|_{\mathbb{L}^\infty(B_r)} \leq C \left(\ln \left(3 + \delta^{-1} \right) \right)^{-3/4} \left(\ln \left(3 \ln \left(3 + \delta^{-1} \right) \right) \right)^2,$$

where $\delta = \|S_1^+(E) - S_2^+(E)\|_{\mathbb{L}^2(\partial B_r \times \partial B_r)}$ and $C > 0$ depends only on N, m, r .

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The end

Thank you for your attention!