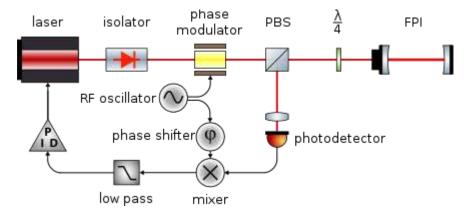
### Laser stabilization and control

### **PDH** method

对激光器的频率稳定,我们主要采用 PDH 稳频方案,该方案的 PID,low pass filter,mixer,phase shifter,RF oscillator 这些伺服电子设备都由 Toptica DL pro 激光控制器自带。



如果是加低频±1级边带,那么激光内置 RF oscillator 就可以搞定。若要加 MHz 及以上级别调制需要外部的 EOM、信号发生器和驱动 EOM 的放大器。

#### PDH 方案的要点:

1. error signal: 
$$P_{\varepsilon} = \sqrt{P_c P_{s,1}} \operatorname{Im} \left( F(\omega_c) F^*(\omega_c + \omega_m) - F^*(\omega_c) F(\omega_c - \omega_m) \right) \cos(\varphi)$$

 $P_c$  is the optical power of the carrier wave,  $P_{s,1}$  is the power of the first side-band wave,  $\varphi$  is the phase difference between the mixed signals.

#### Error signal:

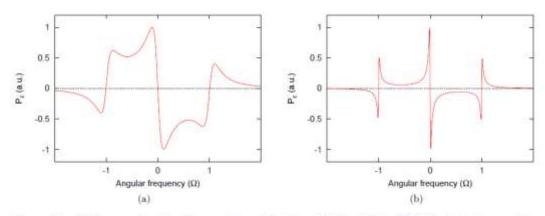


Figure 2.5: PDH error signal for finesse values of (a)  $\mathcal{F} \approx 30\,000$ , (b)  $\mathcal{F} \approx 290\,000$ . In both cases the parameters  $\Omega = 3$  MHz, d = 47.635 mm,  $\beta = 1.08$  and  $\phi = \pi$  were fixed. Clearly high finesse values entail a sharper fringe structure, a steeper center slope and a weaker signal between -1 and 0 as well as 0 and 1.

- 2. The best modulation index is  $\beta=1.08$  (  $\frac{P_{s,1}}{P_c}\approx 0.42$  ). In this condition, the slope of the error signal will be most steep.
- 3. The capture range of the PID is the side-band frequency from the carrier frequency, i.e. the modulation frequency.
- 4. The locking point is integer multiples of the free spectral range of the reference cavity. The integer is very large, typically around 10<sup>5</sup>.

PDH 误差信号的推导:

入射光经过相位调制:

$$E_{inc} = E_0 e^{i(\omega t + \beta \sin(\Omega t))}$$

将上式做傅里叶级数展开:

$$E_{inc} = E_0 e^{i\omega t} \left( \sum_{n=-\infty}^{\infty} a_n e^{in\Omega t} \right)$$

其中
$$a_n = \frac{\Omega}{2\pi} \int_{-\pi/\Omega}^{\pi/\Omega} e^{i\beta \sin\Omega t - in\Omega t} dt$$
,即:

$$a_n = \frac{1}{2\pi} \int_{-\tau}^{\pi} e^{i\beta \sin \tau - in\tau} d\tau = J_n(\beta)$$

 $J_n(\beta)$ 为 n 阶 Bessel 函数。故:

$$E_{inc} = E_0 \left( J_0(\beta) e^{i\omega t} + J_1(\beta) e^{i(\omega + \Omega)t} - J_1(\beta) e^{i(\omega - \Omega)t} + \sum_{|n| \ge 2} J_n(\beta) e^{i(\omega + n\Omega)t} \right)$$

当光场从参考腔反射后,会附加上参考腔的共振信息,数学形式就是乘上反射系数 $F(\omega)$ :

$$\begin{split} E_{ref} &= E_0 \Big( J_0(\beta) F(\omega) e^{i\omega t} + J_1(\beta) F(\omega + \Omega) e^{i(\omega + \Omega)t} - J_1(\beta) F(\omega - \Omega) e^{i(\omega - \Omega)t} \Big) \\ &+ \sum_{|n| > 2} J_n(\beta) E_0 F(\omega + n\Omega) e^{i(\omega + n\Omega)t} \end{split}$$

当调制深度  $\beta$  约为 1 时,  $|n| \ge 2$  的 Bessel 函数值都很小故可以只保留到 1 阶项,反射光强信号为:

$$\begin{split} P_{ref} &= P_0 \big| F(\omega) \big|^2 + P_1 \Big\| F(\omega + \Omega) \big|^2 + \big| F(\omega - \Omega) \big|^2 \Big) \\ &+ 2 \sqrt{P_0 P_1} \big\{ \text{Re} \Big[ F(\omega) F^*(\omega + \Omega) - F^*(\omega) F(\omega + \Omega) \Big] \cos \Omega t \\ &+ \text{Im} \Big[ F(\omega) F^*(\omega + \Omega) - F^*(\omega) F(\omega - \Omega) \Big] \sin \Omega t \big\} + \big( 2\Omega \_terms \big) \end{split}$$

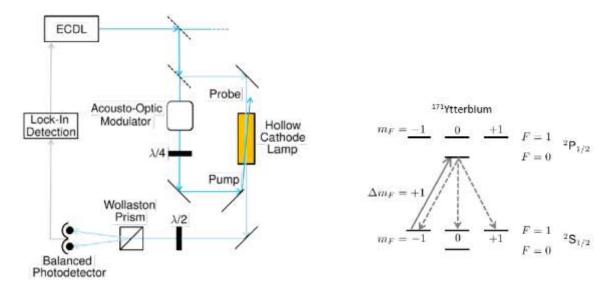
其中  $P_i = E_0^2 J_i^2(\beta)$ , 经过 mixer 和低通滤波的解调过程后得到误差信号为:

$$\varepsilon = P_{ref} \sin(\Omega t + \varphi) \xrightarrow{low\_pass} \sqrt{P_0 P_1} \operatorname{Im} \left[ F(\omega) F^*(\omega + \Omega) - F^*(\omega) F(\omega - \Omega) \right] \cos \varphi$$

## Polarization spectroscopy method

采用 PDH 方案有一个难以避免的问题,那就是参考腔腔长的单向漂移会导致激光锁定的中心频率漂移。为了防止这一现象,一个有效方案是采用阴极灯的 Doppler free spectroscopy 和这里的 Polarization spectroscopy 方法。

下面左图是用 Polarization spectroscopy 方法锁定 Yb 离子 369.5nm 的装置图:



Polarization spectroscopy 方法的要点:

1. There are two beams going through the hollow cathode lamp, one is pump beam(right-handed polarization) which is used to deplete the population of F=1, m<sub>F</sub>=-1. The other one is probe beam(linear porization= $\sigma^+ + \sigma^-$ ). Because of the depletion of population in the F=1, m<sub>F</sub>=-1, the absorption of  $\sigma^+$  component of the probe beam will be different from the  $\sigma^-$  component because the absorption is saturated for the  $\sigma^+$  component but not for the  $\sigma^-$  component.

2. Then there will be intensity difference between the vertical polarization beam and horizontal polarization beam:

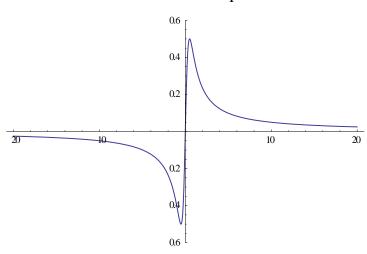
$$I = I_V - I_H = I_0 e^{-2\Delta \alpha L} \cos \left( 2\phi + \frac{\omega L \Delta n}{c} \right)$$

The Wollaston prism will split the vertical polarization beam and horizontal polarization beam, then the balanced photon-detector can detect the intensity difference.

3. The error signal:

$$\phi = \frac{\pi}{4} \text{ , then } \quad I \approx I_0 e^{-2\Delta \alpha L} \Delta \alpha_0 L \frac{\left(\omega - \omega_0\right) \frac{2}{\Gamma}}{1 + \left(\omega - \omega_0\right)^2 \frac{4}{\Gamma^2}} \text{ , so the error signal is:}$$

$$\varepsilon \propto \frac{(\omega - \omega_0)^{\frac{2}{\Gamma}}}{1 + (\omega - \omega_0)^2 \frac{4}{\Gamma^2}}$$

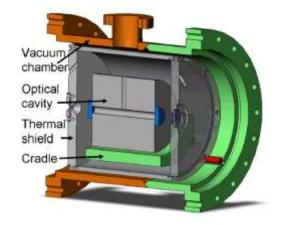


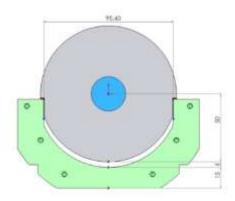
# PDH method with ULE reference cavity

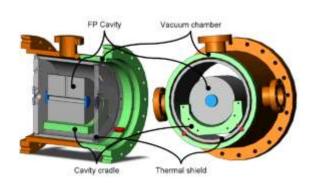
另一种抑制中心频率漂移的方法是在 PDH 方案中把参考腔换成特制的 ULE(ultra low expansion)腔,即腔长极为稳定的 F-P 腔。

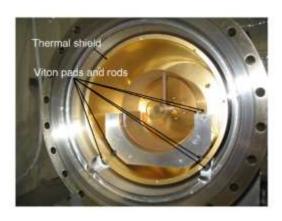
ULE 腔制作参考:

1. Development of ultra-stable laser sources and long-distance optical link via telecommunication

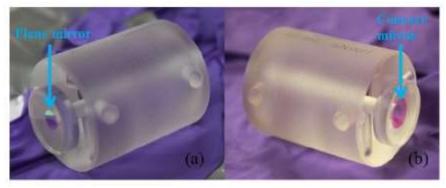




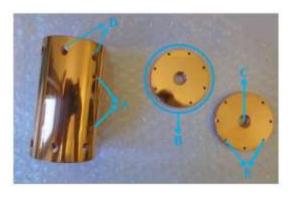


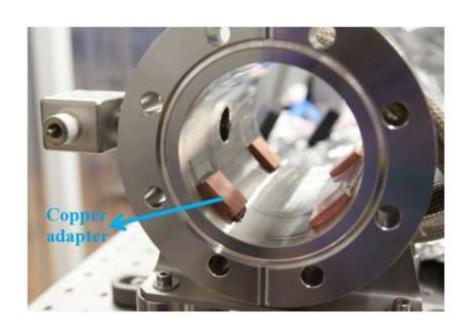


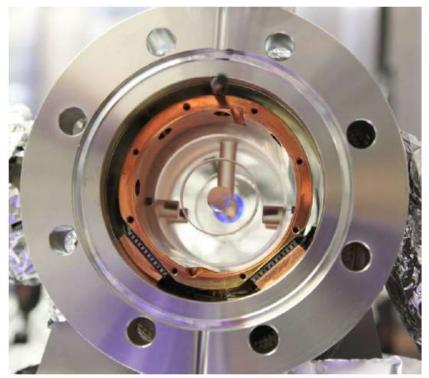
## 2. Diode laser frequency stabilization onto an optical cavity











制作:初步打算先定制一个微晶玻璃 spacer 或者 ULE 玻璃 spacer,它们的线性热膨胀系数 -般为 $\alpha$ < $10^{-7}$   $K^{-1}$ ,然后采用光学粘合法(自己可能不能做)或直接用低热膨胀系数的胶 把反射镜粘到 spacer 上。我们希望胶的热膨胀系数控制在 $10^3\alpha$  以内,这样假设胶厚度为 0.1mm,spacer 有 100mm,则最终整体的线性热膨胀系数为:

$$\alpha = \frac{\Delta L}{L} \frac{1}{\Delta T} = \frac{\Delta L_{glue} + \Delta L_{spacer}}{L_{spacer}} \frac{1}{\Delta T} = \frac{L_{glue}}{L_{spacer}} \alpha_{glue} + \alpha_{spacer} \approx 2\alpha_{spacer}$$

微晶玻璃 spacer 参考公司: 肖特

粘胶参考公司: Permabond、Bacon adhesive、United adhesive

腔自由光谱程漂移:

$$\frac{\Delta v}{v} = \sqrt{\left(\frac{\Delta L}{L}\right)^2 + \left(\frac{\Delta n}{n}\right)^2 + 2\frac{\Delta L}{L}\frac{\Delta n}{n}}$$

其中 L 为腔长, n 为腔中空气柱的折射率。假设空气折射率和空气压成正比,则:

$$\frac{\Delta v}{v} = \sqrt{\left(\frac{\Delta L}{L}\right)^2 + \left(\frac{\Delta p}{p}\right)^2 + 2\frac{\Delta L}{L}\frac{\Delta p}{p}}$$

 $\frac{\Delta L}{L} = \alpha \Delta T$ 。假设一个小时内气压变化有 1 毫巴量级,温度是 0.1K 量级,则主要是气压变化占主导,故 $\frac{\Delta \nu}{\nu} = 10^{-3}$ 。假设腔的自由光谱程为 1.5GHz,则自由光谱程漂移为 1.5MHz。假设我们锁定 369.5nm,则锁点频率约是自由光谱程 50 万倍。假设我们忽略锁频反馈时间,则激光频率将会被锁点带着漂移 0.75THz。所以**气压变化是占主导**的,在参考腔外罩真空腔是绝对必要的,不然参考腔热膨胀系数再低都没有意义。

## Demodulation and servo control system for frequency stabilization

- 1、给 370nm 锁频的 PDH 方案: demodulation 基本考虑从 mini-circuits 买 mixer、low-pass filter (DC-10MHz) 以及 phase shifter 自己搭; servo control 为 Monroe 组高速 PID,但缺少慢速 PID。(注: 370 腔镜的精细度大约 1500,故加>40MHz 会看到比较好的 error signal)
- 2、给 370nm 稳频的阴极灯方案: NIST 组的数字 lock-in amplifier 和 PID; TOPTICA dlc pro 自带的 lock-in 和 PID。
- 3、给 871nm 稳频的 PDH 方案: 可以考虑 TOPTICA mFALC 110 模块+信号源或者 FALC 110+PDD 110/F 模块(注: mFALC 具有混频器和高低速 PID 功能; FALC 是高低速 PID; PDD110/F 具有内置信号源和 demodulation 功能。)

#### Laser control

由于对离子进行特定的幺正操控涉及激光的频率和相位参数的设定(<u>Cold trapped ions</u> as quantum information processors),故我们需要通过 AOM 间接操控这些参数,这等价于设定 AOM 信号源的频率和相位。

- 1、Phase Continuous switching: 上一段波形和下一段波形是连续的。这种切换用单个 AD9910 dds(direct digital synthesizer)即能实现,因为 AD9910 内置相位寄存器,能记录当前的相位,只要在事先对其内置的 8 个 profile 寄存器写好信号参数然后做 profile 切换即可实现 phase continuous switching。
- 2、Phase coherent switching: 上一段波形和下一段波形是断开的。这种情况下,如果用单个dds 实现可能需要事先算好切换点的相位然后提前写入到 profile 寄存器;或者在切换时转换成 PARALLEL DATA PORT MODULATION MODE,然后修改信号参数。如果用多个 dds 则可以比较方便,比如用一个 AD9959/AD9958 芯片,让内置的多个 dds 芯片以不同频率同时进行相位累加,在需要切换时让一个 dds 的 DAC 输出幅度降为零(即关闭该通道),让另一个的输出幅度设为全幅度(即打开该通道)即可实现 phase coherent switching。或者用一个外置的 phase accumulator 和 dds 内置的 phase accumulator 同时运转,当需要切换时将外置的相位寄存器数据写给 dds,这样也可以做到切换。

