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Study of HD¹⁷O spectrum. Theory and experiment

Leonid N. Sinitsa^b, Nikolai F. Zobov^c, Mikhail A. Rogov^c, Jonathan Tennyson^{a,*}, Oleg L. Polyansky^{c,a}

- ^a Department of Physics and Astronomy, University College London, Gower Street, London WC1E 6BT, UK
- b V.E Zuev Optics of Atmosphere Institute, Tomsk, Russia
- ^c Institute of Applied Physics, Russian Academy of Sciences, Ulyanov Street 46, Nizhny Novgorod, 603950, Russia

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ABSTRACT

A spectrum of $\rm HD^{17}O$ in the region between 5 480 and 6 400 cm⁻¹ recorded in Tomsk is reported. In this region 4292 lines are observed belonging to 7 water isotopologues: 536 lines – $\rm H_2^{16}O$ 1019 – $\rm HD^{16}O$, 122 – $\rm D_2^{16}O$, 447 – $\rm H_2^{17}O$ 458 – $\rm HD^{17}O$, 151 – $\rm H_2^{18}O$, 181 – $\rm HD^{18}O$. A new potential energy surface (PES) of $\rm HD^{16}O$ is obtained by fitting to empirical energy levels. This PES, with a diagonal Born–Oppenheimer correction (DBOC), is used to compute the $\rm HD^{17}O$ spectrum. Pseudo-experimental isotopologue-extrapolation energy levels of $\rm HD^{17}O$ are constructed using the method of Polyansky et al (MNRAS 466, 1363 (2017)). Assignment of the measured spectrum is conducted: 68 % of the lines can be assigned using assignments from previously published work. The remaining lines are assigned using the pseudo-experimental energy levels procedure. We compare the calculated pseudo-experimental values of energy levels with both existing 1285 experimental levels of $\rm HD^{17}O$ and the 152 newly determined in this work energy levels. The standard deviation of levels with low J (up to J=10) is about 0.007 cm⁻¹ in both cases. Energy levels from both pseudo-experimental and variationally calculated are also compared with newly measured $\rm HD^{17}O$ lines above 10 000 cm⁻¹.

1. Introduction

Water molecule is arguably the most important molecule in the entire Universe. Its spectrum is very well studied. Much less well studied are its isotopologues spectra, though probably much better, than the major isotopologues of other molecules.

The energy levels of the water molecule, including its isotopologues, are important for a variety of applications. A decade ago some of us collaborated on a IUPAC (International Union of Pure and Applied Chemistry) project whose aim was to provide empirical energy levels for all isotopologues of water [1]. The project developed the MARVEL (measured active rotation–vibration energy levels) procedure [2] and used it to provides a set of recommended energy levels for all water isotopologues including $\rm HD^{16}O$, $\rm HD^{17}O$ and $\rm HD^{18}O$ [3], which we use here. While both the MARVEL procedure [4,5] and energy levels for $\rm H_2^{16}O$, and $\rm H_2^{17}O$ and $\rm H_2^{18}O$ have subsequently been regularly updated [5–8], there is as yet no systematic update of the IUPAC data for the deuterated isotopologues.

We report measurements and analysis of the vibrational–rotational absorption spectrum of the $\mathrm{HD^{17}O}$ molecule in the IR range from 5480 to 6400 cm⁻¹. The absorption spectrum of the $\mathrm{H_2^{17}O}$ molecule in this region was recently investigated by us [9] as well as the higher

frequency spectrum of $\mathrm{HD^{17}O}$ [10–12]; these studies are considered further below.

2. Measurements and data processing

The absorption spectrum of the HD¹⁷O molecule was recorded using an IFS-125M Fourier spectrometer in the range from 5480 cm⁻¹ to 6400 cm⁻¹, at a pressure of 29 mbar and an optical path length of 24 m. The Doppler FWHM in this region is 0.0166 cm⁻¹; measurements were performed with a spectral resolution of 0.02 cm⁻¹. A multi-pass optical White cell with a base length of 60 cm was used for measurements. Measurements were carried out at a temperature of 25 ± 1 °C. A halogen lamp was used as the radiation source. We used a Mertzphase correction with a 1 cm⁻¹ phase resolution, and no apodization were applied to the averaged interferograms. A background spectrum was first recorded at a resolution of 0.05 cm⁻¹ while the cell was being continuously evacuated. In order to prepare the HD¹⁷O vapor, liquid water enriched to 86% with H₂¹⁷O produced by the Sigma-Aldrich Chemical Company was mixed with an equal quantity of the D₂¹⁶O. It gives us approximate water isotopologues abundances: H₂¹⁶O -12%, H₂¹⁷O -12%, HD¹⁶O -25%, HD¹⁷O -25%, D₂¹⁶O -12%.

E-mail addresses: j.tennyson@ucl.ac.uk (J. Tennyson), o.polyansky@ucl.ac.uk (O.L. Polyansky).

^{*} Corresponding author.

Table 1
Experimental conditions.

Spectral resolution	0.02 cm ⁻¹
Optical path length	2400 cm
Pressure of H ₂ O/HDO/D ₂ O mixture	29 mbar
Temperature	$25 \pm 1 ^{\circ}\mathrm{C}$
Signal-to-noise ratio	About 10 000
Diaphragm	0.9 mm
The total measurement time	7 days

Measurements were performed during several days; each spectrum was obtained by averaging over 37 584 scans during 168 h of recording. This procedure allowed us to obtain a signal-to-noise ratio of about 1000 and to detect weak absorption lines with intensities of the order of 10^{-26} cm/molecule.

The volume of the cell is 2.765 L. The absorption spectra were recorded at a pressure of 29 mbar. The pressure measurements were performed using a pressure transducer AIR-20M. Considering the uncertainty arising from small variations of the pressure during the recording, the measurement uncertainty on the pressure was estimated as 0.5%. The signal-to-noise ratio on the strong lines is about 1000 despite the fact that the optical path was only 24 m; Fig. 1 gives and overview of the spectrum.

The FT-Spectrometer Bruker IFS-125M was not evacuated. Variations in the baseline due to changes in humidity and air pressure in the laboratory during measurements do not exceed 1.5%. Spectral noise is the most important source of error for the line strength measurements. For strong lines with an intensity $6\times 10^{-25}-2\times 10^{-25}$ cm/molecule, they are 0.5%, while for weak lines with an intensity of less than 2×10^{-25} cm/molecule they are 5%. The errors that occur when fitting a spectral line contour vary from 0.1% for strong lines to 1.5% for weak lines. This gives a total measurement uncertainty for the line broadening coefficients of 4% and 9% for strong and weak lines, respectively. The uncertainty in the determination of the line position does not exceed 0.003 cm $^{-1}$. The measurement conditions are given in Table 1.

The measured spectrum was processed using the Wxspe software package [13] that uses pattern recognition algorithms. The software searches for lines in automatic mode. The centers, intensities, and half-widths were determined by least-squares fitting Voigt profile parameters to the experimental data taking into account an appropriate instrumental function.

32 $\rm{H_2}^{16}O$ lines from the HITRAN data [14] were used to calibrate the measured line centers. In fact, the $\rm{H_2}^{16}O$ spectroscopic data were used as secondary frequency standards. As a result of the calibration performed, the measurement error of the line centers can be estimated to be better than 0.003 cm⁻¹ for strong lines and not exceeding 0.02 cm⁻¹ for weak lines with intensities less than 4×10^{-26} cm/molecule.

The partial pressure of the component mixture was determined by comparing the HITRAN 2020 [14] line intensities, which for the deuterated isotopologues are actually derived from calculations [15], with the measured ones. We found the partial pressure of $\rm H_2^{16}O$ to be 13 mbar; $\rm H_2^{17}O$ – 7 mbar; $\rm HD^{16}O$ – 5.5 mbar and that of $\rm HD^{17}O$ isotopologue 3.5 mbar. The partial pressure of the ^{18}O isotopologues was very small and was neglected. The uncertainty in determining the line intensities can be estimated as 7%–10%. The recorded spectrum was compared with water isotopologues spectrum containing 86% $\rm H_2^{17}O$ and 14% $\rm H_2^{16}O$. The comparison allowed to unambiguously tie the recorded lines with concrete water isotopologue. Fig. 2 shows a part of the recorded spectrum in the range 5524–5530 cm $^{-1}$ with water isotopologue lines identified.

3. Variational calculations for HDO

In order to perform a fit of the HDO PES, the rotation-vibration Schrödinger equation has to be solved. For that purpose we used the

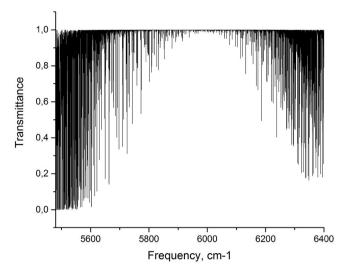


Fig. 1. Overview of the recorded spectrum.

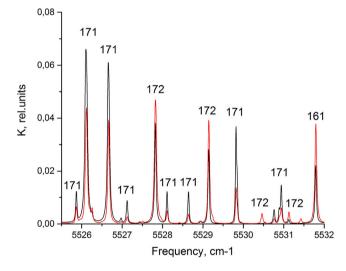


Fig. 2. A short portion of the observed spectrum. Numbers denote the isotopologue associated with each line: 161 is $\rm H_2^{16}O$, 171 is $\rm H_2^{17}O$, and 172 is $\rm HD^{17}O$. The red and black lines represent different partial concentrations of water isotopologues in the sample.

DVR3D program suite [16]. DVR3D uses an exact representation of nuclear motion kinetic energy operator so that, within constraints of the Born–Oppenheimer approximation, any uncertainty in the wavefunctions is due to the PES and the failure of the Born–Oppenheimer approximation. Because of the lower symmetry of HDO and the higher density of states, achieving accuracy in HDO calculations comparable to that routinely achieved for $\rm H_2^{16}O$ and its 17–18 oxygen isotopologues [17,18], requires the diagonalization of significantly larger matrices. The PES is represented as a polynomial in valence coordinates of the HDO molecule with 107 adjustable parameters. A file with fixed and adjusted parameters is presented in the SM.

Nuclear motion calculations were performed in Radau coordinates with the z-axis embedded close to the OH bond direction. The discrete variable representation (DVR) basis comprised 30 Gauss–Laguerre points in each radial coordinate and 48 (associated) Gauss–Legendre points in the angular coordinate. Basis set parameters for the radial coordinates were taken from Yurchenko et al. [19]. The final vibrational Hamiltonian matrix was of dimension 7000. For each rotational state J, two Hamiltonian matrices were diagonalized with dimensions $1000 \times (J+1)$ and $1000 \times J$, depending on their overall parity.

4. Fitted PES

To get accurate values for the energy levels for the ¹⁷O isotopologues of HDO one has to construct the PES in a special form which separately identifies the mass-dependent part of the PES, namely the Diagonal Born–Oppenheimer Correction (DBOC). As a result the PES consists of a mass-independent term and a mass dependent-terms which varies for the different isotopologues.

$$V^{XX} = V^{all} + DBOC^{XX}$$
, where XX = 16,17 or 18. (1)

The mass-independent term can be obtained by fitting $\mathrm{HD^{16}O}$ energy levels to the form which includes the $^{16}\mathrm{O}$ DBOC. We choose the most abundant isotopologue of HDO because it has the biggest set of experimental levels. In practice we took the HDO PES of Yurchenko et al. [19], added to it an $\mathrm{HD^{16}O}$ DBOC and then performed a refit to energy levels up to 15 000 cm⁻¹ (637 energy levels with J values of 0, 2, 5); in this work we do not need levels with higher energies. Fitting to $\mathrm{HD^{16}O}$ levels with J values of 0, 2, 5 leads to HDO PESs which gives standard deviations, σ , equal to 0.019 cm⁻¹ for $^{16}\mathrm{O}$, 0.020 cm⁻¹ for $^{17}\mathrm{O}$ and 0.027 cm⁻¹ for $^{18}\mathrm{O}$ isotopologues; we call these PES16, PES17 and PES18, respectively. PES17 and PES18 are sub-products of the PES16 fitting procedure, using $\mathrm{HD^{16}O}$ energy levels. The form (1) is chosen due to its usefulness for obtaining pseudo-experimental levels of 17 and 18 HDO isotopologues.

5. Pseudo-experimental energy levels predictions for HD¹⁷O

Polyansky et al. [20] developed a method for calculating H₂17O and H₂ ¹⁸O energy levels, based on experimental energy levels for H₂¹⁶O and the calculated differences between the parent isotopologue levels and levels of the corresponding state for the ¹⁷O or ¹⁸O isotopologues. The energy levels calculated this way gave predictions with an accuracy of about 0.005 cm⁻¹, which is close to the accuracy of the experimentally-determined energy levels of H₂¹⁶O. The high accuracy achieved and the use of experimental values for the parent isotopologue led us to label the resulting energy levels "pseudoexperimental" originally, though this technique subsequently was also called isotopologue extrapolation (IE) by McKemmish et al. [21]. Subsequently [6] we were able to compare newly measured ¹⁷O and ¹⁸O water energy levels with the previously calculated pseudo-experimental energy levels. The excellent accuracy achieved with $\sigma = 0.0039 \text{ cm}^{-1}$ for 1031 experimental energy levels confirmed the usefulness of our methodology. In this paper we use the same procedure to determine HD¹⁷O energy levels starting from the HD¹⁶O experimentally known energy levels. These pseudo-experimental levels are compared with the energy levels obtained here from our newly measured experimental HD¹⁷O line centers, see Table 2.

To obtain values for the pseudo experimental $\mathrm{HD^{17}O}$ levels one has to have experimental levels for $\mathrm{HD^{16}O}$. The 2010 IUPAC compilation of HDO levels [3] provides 7200 experimental levels of $\mathrm{HD^{16}O}$. Recent measurements [22–28] have increased this number almost to 8000. Exactly this set of 'MARVEL extended' levels are called MARVEL levels in this work. We managed to obtain 5800 pseudo experimental $\mathrm{HD^{17}O}$ energy levels. This number is lower than the number of available $\mathrm{HD^{16}O}$ energy levels due to the difficulties of correct transferring of $\mathrm{HD^{16}O}$ labeling to $\mathrm{HD^{17}O}$ isotopologue. The pseudo experimental $\mathrm{HD^{17}O}$ energy levels are in supplementary materials.

6. Analysis of HD¹⁷O spectra

HD¹⁷O line positions have been measured in the IR range 5500–6500 cm⁻¹ previously in papers [10–12,22–39]. Comparing measured in this work lines with published spectra gives us the following assignment: 536 lines – $\rm H_2^{16}O$ 1019 – $\rm HD^{16}O$, 122 – $\rm D_2^{16}O$, 447 – $\rm H_2^{17}O$ 458 – $\rm HD^{17}O$, 151 – $\rm H_2^{18}O$, 181 – $\rm HD^{18}O$. This meant that 68% of

Table 2 Standard deviations, σ , of calculated with PES17 HD¹⁷O energy levels for different J values (left) and σ for pseudo experimental levels (right). N is the number of experimental levels.

J	N	σ (cm ⁻¹)	N	σ (cm ⁻¹)
0	9	0.0169	11	0.0042
1	31	0.0166	41	0.0058
2	55	0.0159	67	0.0055
3	82	0.0144	98	0.0055
4	90	0.0113	113	0.0071
5	109	0.0095	128	0.0096
6	102	0.0113	131	0.0161
7	105	0.0162	134	0.0218
8	86	0.0206	104	0.0294
9	69	0.0258	71	0.0333
10	53	0.0295	43	0.0374
11	42	0.0339	43	0.0427
12	21	0.0399	15	0.0424
13	16	0.0386	13	0.0393
14	10	0.0436	4	0.0486
15	5	0.0305	4	0.0384

Table 3Statistics for the assigned HD¹⁷O lines.

Type	N lines	σ (cm ⁻¹)
Pseudo and MARVEL	564	0.0110
MARVEL	290	0.0044
Pseudo	274	0.0173

the observed lines could be assigned from published papers, leaving us 1368 lines to assign.

29 newly determined upper energy levels are involved in two or more transitions which allow us to check the correctness of the assignment using ground state combination differences which showed a standard deviation of $0.0071~\rm cm^{-1}$. 301 lines were assigned to $\rm HD^{17}O$ mostly using pseudo experimental levels. From these transitions we obtained values of 153 new energy $\rm HD^{17}O$ levels. 94 newly assigned lines do not directly lead to the determination of new levels as their lower levels remains unknown experimentally. These lines generally involve levels with higher values of J and K_a . The 2010 IUPAC compilation of HDO levels [3] contains just 162 experimental levels of $\rm HD^{17}O$. Recent experimental studies [22–28] have increased this number to 1285. This set of experimental levels are called MARVEL levels below.

The standard deviation between experimental frequencies of HD¹⁷O lines and our calculated ones is 0.019 cm⁻¹ (see Tables 3 and 4).

Subsequent to our main work on recording and analyzing the spectra around 6000 cm $^{-1}$, three studies analyzing HD ^{17}O spectra in $10\,000-13\,000$ cm $^{-1}$ range [10–12] became available. Table 5 shows the quality of our prediction using PES17 and pseudo experimental levels in $10\,000-13\,000$ cm $^{-1}$ range. For experimental levels around $10\,000$ cm $^{-1}$ $\sigma=0.11$ cm $^{-1}$, while for levels around $13\,000$ cm $^{-1}$ the $\sigma=0.02$ cm $^{-1}$. In this region, the quality of the levels calculated with PES17 levels is about the same as for pseudo experimental levels. For the (400) band the deviation is two times bigger — 0.202 cm $^{-1}$ then for other bands in the region; this suggests that some levels of the (400) band may not be correctly assigned.

7. Conclusions

This paper presents the results of the observation and analysis of the enriched spectrum of HDO molecule by the ^{17}O isotopologue. In order to analyze the experimental spectrum we improved the fitted PES of HD ^{17}O and calculated the 6000 pseudo-experimental values of energy levels. We compared the calculated pseudo-experimental values of energy levels with both existing 1300 levels of HD ^{17}O and the 152 newly determined in this work energy levels. The standard deviation of low J (up to J = 10) is about 0.007 cm $^{-1}$ in both cases. Note, that the newly determined number of 152 levels is higher than the entire 130

Table 4 Newly determined energy levels of $\mathrm{HD^{17}O}$. Calc.1 – energy levels calculated with the use of new $\mathrm{HD^{17}O}$ PES (PES17), Calc.2 – pseudo experimental levels. 'cd' denotes levels determined with combinations differences, 's' and denotes levels of unresolved doublet split using the values from the theoretical calculations.

$v_1v_2v_3$ state	$J K_a K_c$	Obs. (cm ⁻¹)	ObsCalc.1 (cm ⁻¹)	ObsCalc.2 (cm ⁻¹)
011	6 5 2	5799.7529	-0.006 cd	0.000
011	6 5 1	5801.6658		
011	9 5 4	6174.8925	6174.8925 0.019	
011	11 3 9	6245.5243	-0.060	-0.054
021	4 1 4	6594.1379	-0.001	-0.004
021	4 1 3	6625.0098	0.011	0.007
021	4 2 3	6664.9406	-0.002	-0.006
021	4 3 1 5 2 3	6757.0341	-0.001	-0.005
021 021	5 3 3	6753.1835 6834.9870	-0.003 -0.003 cd	-0.007 -0.007
021	5 3 2	6835.8466	0.003	-0.001
021	6 2 5	6831.9743	-0.003	-0.007
021	6 4 3	7052.8349	-0.013	-0.012
021	6 4 2	7052.9119	-0.030	-0.028
021	707	6834.8102	0.003 cd	-0.001
021	7 1 7	6836.3386	-0.001	-0.005
021	7 2 6	6936.8442	-0.010	-0.014
021	7 2 5	6971.2352	-0.006	-0.010
021	7 4 3	7162.6270	-0.038	-0.042
021	753	7317.6459	-0.049 s	-0.053
021 021	7 5 2 8 0 8	7317.6357 6941.4673	-0.049 s 0.000	-0.004 -0.013
021	827	7055.4489	-0.009	-0.013 -0.004
021	836	7162.8415	0.000	0.000
021	8 3 5	7174.4787	-0.002	0.001
021	919	7060.6397	0.000	-0.001
021	9 2 8	7187.4072	-0.004	-0.055
021	10 0 10	7191.2730	0.000	0.017
021	10 2 9	7332.3733	-0.004	-0.013
031	2 2 1	7864.1804	-0.017	-0.001
031	4 0 4	7887.1525	-0.005	0.005
040	3 3 1	5722.0495	0.036	0.005
040	4 4 1	6016.2464	0.011	
040	5 3 3	5861.1482	0.003	
040	5 4 2	6094.8725	0.000 s	-0.048
040	5 5 0	6308.6763	0.013 cd	-0.001
040	7 4 3	6299.7717	-0.011	0.070
040	8 4 4	6326.4738	0.001	-0.054
040 040	937 945	6324.3469 6568.8747	-0.007 -0.033	
040	10 2 8	6417.9730	0.031 cd	-0.024
040	11 3 9	6600.0811	-0.029	0.024
				0.001
101 101	2 1 2 3 1 2	6458.9143 6516.4022	-0.001 0.002	0.001 0.004
101	3 2 2	6552.8612	0.002	0.004
101	4 1 3	6581.4430	0.009	0.011
101	4 2 2	6618.0910	0.006	0.008
101	5 1 4	6661.7172	0.013	0.015
101	5 1 5	6621.5808	0.003	0.005
101	5 2 3	6698.5359	0.007	0.009
101	5 3 2	6763.6248	0.014	0.016
101	5 3 3	6762.7628	0.012	0.014
101	606	6699.7704 6756.4345	0.006	0.008
101 101	6 1 5 6 1 6	6756.4345 6701.7289	0.017 0.002	0.019 0.004
101	707	6793.4389	0.002 0.004 cd	0.004
101	716	6864.6429	0.019	0.021
101	7 3 5	6960.7533	0.023 cd	0.025
101	7 4 4	7059.6053	-0.002	0.000
101	808	6899.2032	0.012	0.014
101	8 1 7	6985.4310	0.027	0.029
101	8 4 5	7181.6963	-0.004	-0.002
101	909	7017.1089	0.020	0.022
101	10 1 9	7262.5573	0.047	0.049
101	10 2 9	7265.9631	0.030	0.032
111	5 3 3	8137.3415	0.005	0.006

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Table 4	(continued).			
120	1 0 1	5504.9309	-0.016	
120	202	5535.3927	0.009	0.000
120	2 1 1	5561.5630	0.003	-0.006
120 120	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	5552.3240 5622.3352	0.005 0.008	-0.004 -0.001
120	2 2 1	5622.0044	0.003 0.002 cd	-0.007
120	3 0 3	5580.1993	0.008	-0.001
120	3 1 2	5612.1608	0.007	-0.002
120	3 2 1	5669.9801	0.006	-0.003
120	3 2 2	5668.3875	0.008 cd	-0.001
120	3 3 0	5784.1382	0.006 cd	-0.003
120	3 3 1 4 0 4	5784.1207	0.008	-0.001
120 120	414	5638.5416 5648.5367	0.009 0.006	0.000 -0.003
120	413	5679.1390	0.008	-0.002
120	4 2 3	5729.9461	0.005 cd	-0.004
120	4 2 2	5734.5567	0.004 cd	-0.005
120	4 3 1	5847.0491	0.005	-0.004
120	4 3 2	5846.9152	0.003 cd	-0.006
120	5 0 5	5709.7007	-0.012	-0.021 -0.004
120 120	5 1 5 5 1 4	5716.5562 5762.0258	0.005 0.003	-0.004 -0.006
120	5 2 4	5806.4627	0.002	-0.007
120	5 2 3	5816.5965	0.002	-0.007
120	5 3 3	5925.5206	-0.001cd	-0.010
120	5 3 2	5926.0377	0.000 cd	-0.009
120	6 0 6	5793.0204	0.002	-0.007
120	616	5797.5343	0.003	-0.006
120 120	6 1 5 6 2 5	5860.1850 5897.6689	-0.003 0.001	-0.012 -0.008
120	6 2 4	5916.3080	-0.004	-0.003
120	6 3 4	6019.9244	-0.006	-0.015
120	6 4 3	6094.6072	0.003	-0.006
120	6 5 2	6263.3536	0.013	0.004
120	7 0 7	5888.4318	0.000	-0.009
120	717	5891.2613	-0.001 cd	-0.010
120 120	7 1 6 7 2 6	5972.8361 6003.2586	−0.003 −0.001 cd	-0.012 -0.010
120	7 2 5	6033.5260	-0.010	-0.019
120	7 3 5	6130.0461	-0.014	-0.023
120	7 3 4	6133.6285	-0.015	-0.024
120	7 5 3	6370.4172	-0.003 cd	-0.012
120	7 5 2	6370.4121	-0.003 s	-0.011
120	8 0 8 8 1 8	5995.8442 5997.5680	-0.002 -0.005	-0.014
120 120	817	6099.1346	-0.003 -0.006 cd	-0.015 -0.021
120	8 2 7	6122.9061	-0.012	-0.033
120	8 2 6	6167.7802	-0.024	-0.010
120	8 4 5	6325.8110	-0.001	-0.024
120	8 5 4	6255.7432	-0.015	-0.003
120	863	6690.0765	0.006 s	-0.015
120	862	6690.0762 6116.3269	0.006 s -0.006 cd	-0.015 -0.024
120 120	9 0 9 9 1 9	6115.2898	-0.006 cd -0.006	-0.024 -0.047
120	9 2 8	6256.3178	-0.015 cd	-0.019
120	936	6410.3893	-0.038	-0.049
120	10 0 10	6246.8155	-0.010	-0.054
120	10 2 8	6484.9096	-0.040	-0.059
120	10 3 8	6552.8619	-0.045	-0.021
120	10 3 7	6575.5978	-0.050	-0.018 -0.045
120 120	11 0 11 11 1 11	6390.4535 6390.8217	−0.012 cd −0.009 cd	-0.043
120	11 2 10	6563.2591	-0.036	-0.027
120	11 3 9	6723.6744	-0.058	-0.023
120	12 0 12	6546.1874	-0.018 cd	-0.015
120	12 1 12	6546.4127	-0.014	-0.024
200	4 3 2	5634.6648	-0.022	-0.001
200	4 3 1	5634.8289	-0.026	-0.005
200	4 4 1	5738.8813	-0.032	-0.011
200	5 2 3	5641.4130	-0.009	0.012
200	5 3 3	5709.6410 5813.6144	0.004 _0.033 cd	0.025
200	5 4 2	5813.6144	-0.033 cd	0.002

5946.1305

5946.1306

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-0.004

-0.007

−0.025 s

 $-0.025 \ s$

200

200

5 5 1

5 5 0

Table 4	(continued).	

Table 4 (continueu).			
200	6 4 3	5903.4151	-0.028	-0.001
200	6 4 2	5903.4895	-0.022	0.022
200	7 3 5	5904.7571	0.001	0.019
200	7 4 4	6008.3897	-0.002	0.027
200	7 4 3	6008.6431	0.006 cd	0.034
200	8 3 6	6024.5106	0.013	0.025
200	8 4 5	6128.5097	0.004	0.021
200	8 5 3	6259.2000	0.000	-0.012
200	10 2 8	6275.7877	-0.033	
200	11 3 8	6494.9810	-0.039	-0.010
200	11 4 7	6586.9720	-0.031	-0.016
200	12 1 11	6523.0623	-0.037	-0.030
200	12 3 9	6698.4137	-0.051	-0.012

Table 5 Comparison of recently measured $\mathrm{HD^{17}O}$ levels with calculations made using PES17 and pseudo experimental (isotopologue extrapolation, IE) levels. E_{\min} denotes the energy of the lowest observed level for each band.

Band	N levels	σ	N levels	σ	E_{min}
	PES17	(cm ⁻¹)	IE	(cm ⁻¹)	cm ⁻¹
(003)	138	0.127	70	0.114	10615.4
(022)	55	0.095	35	0.061	10161.4
(400)	44	0.192	40	0.202	10366.1
(102)	64	0.068	59	0.059	10069.5
(112)	25	0.035	24	0.030	11 292.5
(004)	114	0.029	50	0.027	13832.8
(023)	89	0.019	46	0.023	13 255.6
(103)	75	0.027	34	0.032	13 307.8

levels known up to 2010. The newly measured lines of $\mathrm{HD^{17}O}$ and the newly determined energy levels as well as the 6000 predicted pseudo-experimental levels increase significantly our knowledge of this water isotopologue and will provide a sound basis for the future attempts to solve completely the water isotopologue problem via solution of the non-adiabatic problem (see for example [40–42]).

CRediT authorship contribution statement

Leonid N. Sinitsa: Project administration, Methodology, Data curation, Conceptualization. Nikolai F. Zobov: Writing – review & editing, Software, Investigation, Formal analysis, Conceptualization. Mikhail A. Rogov: Investigation, Formal analysis. Jonathan Tennyson: Writing – review & editing, Methodology, Funding acquisition. Oleg L. Polyansky: Writing – original draft, Project administration, Methodology, Investigation.

Declaration of competing interest

The authors declare no conflict of interest

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Appendix A. Supporting material

The following are given as supporting material: 1. A full list of the assigned experimental spectrum; 2. The $\mathrm{HD^{16}O}$ and $\mathrm{HD^{17}O}$ potential energy surfaces; 3. The lists of $\mathrm{HD^{16}O}$ and $\mathrm{HD^{17}O}$ energy levels.

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Data availability

All data is included in the article or the supplementary material.

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