DETERMINATION OF THE ROTATIONAL CONSTANT OF $^{15}\,\mathrm{NH_3}$

A. F. Krupnov, L. I. Gershtein, V. G. Shustrov, and V. V. Polyakov

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The rotational spectrum of ^{15}N ammonia has not previously been observed. We have examined the lower rotational transition $J=0 \rightarrow 1$, K=0 for $^{15}NH_3$ in the 0.525-mm range, which has given the rotational constant for this species of ammonia.

We used a gas spectroscope with video detection [1]. The cell contained $^{14}\mathrm{NH_3}$ and $^{15}\mathrm{NH_3}$ in the ratio of about 8:1. On sweeping the oscillator frequency (backward-wave tube [2]), the lines for $J=0\to1$, K=0 for the two forms were clearly seen (Fig. 1). The distance between the lines was deduced from the frequency modulation of the source [3], which was recorded by a ChZ-9 frequency meter. The difference between the two transition frequencies was 384.8 MHz with an error of \pm 0.2 MHz. The transition frequency for $^{14}\mathrm{NH_3}$ has been measured [1] as $\nu_{0-1}(N^{14}\mathrm{H_3})=572,499$,4 MHz \pm 3 MHz, so for $^{15}\mathrm{NH_3}$ we have

$$v_{0-1}(N^{15}H_3) = 572115 \text{ MHz},$$

with the accuracy as for $\nu_{0-1}(N^{14}H_3)$.

The rotational constant B₀ was deduced as for ¹⁴NH₃ [1] and is

$$B_0(N^{15}H_3) = \frac{1}{2} \left[v_{0-1}(N^{15}H_3) + 4D_J + \frac{v_{00} + v_{01}}{2} \right],$$

where D_J is the centrifugal perturbation constant and $(\nu_{0\ 0} + \nu_{0\ 1})/2$ is half the sum of the frequencies for the inversion splitting of the levels J=K=0 and J=1, K=0. No value has been published for D_J for $^{15}{\rm NH_3}$, but the quantity is determined by an expression of the form ${\rm B}^3/\omega^2$, where B is the rotational constant and ω is the vibrational frequency [3], so the value for $^{15}{\rm NH_3}$ is virtually that for $^{14}{\rm NH_3}$, which is 19 MHz [3]. It is not possible to measure $\nu_{0\ 0}$ and $\nu_{0\ 1}$ because $^{15}{\rm NH_3}$ lacks one of the inversion levels in each pair for K=0, so one should use the dependence of the inversion frequency on J and K as found from transitions with $K\neq 0$ by using the corresponding J and K.

An empirical formula [4] for the frequencies in the inversion spectrum of ¹⁵NH₃ has a standard deviation of 308 MHz and in the appropriate region (J and K small) gives values that differ from experiment by 3-10 MHz, which reduces the accuracy of the calculation. We therefore drew up the following better approximation from the measured frequencies [5]:

$$\gamma_{JK}^{(\text{MHz})} = 22714 \exp \left[-0.6489815 \cdot 10^{-2} J(J+1) + 0.902306 \cdot 10^{-2} K^2 + 0.850786 \cdot 10^{-6} J^2 (J+1)^2 - 0.178845 \cdot 10^{-5} J(J+1) K^2 + 0.52645 \cdot 10^{-6} K^4 \right],$$
(1)

which has a standard deviation of 3.6 MHz and differences of 0.2-4.5 MHz from experiment at small J and K. Use of (1) gives a more accurate B_0 for $^{15}{\rm NH_3}$ as

$$B_0(N^{15}H_3) = 297379.3$$
 MHz.

The method of [1] gave $\nu_{0-1}(N^{15}N_3) = 572,111.44$ MHz for the frequency of the $J=0 \rightarrow 1$, K=0 transition in $^{15}NH_3$, and so

$$B_0(N^{15}H_3) = 297377.5$$
 MHz.

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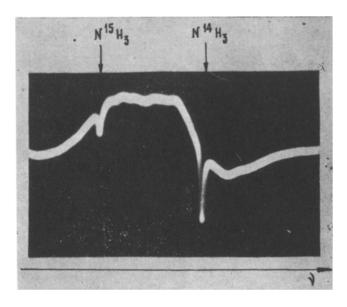


Fig. 1.

The results agree to $6 \cdot 10^{-6}$ or so, and the measured frequency for $^{14}NH_3$ is $\nu_{0-1}(N^{14}H_3) = 572,495.6$ MHz, which corresponds to $B_0(N^{14}H_3) = 298,104,2$ MHz. The result for the frequency difference is 384.1 MHz, which agrees well with the value above within the error of the present method.

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