

Line Broadening and Shifting Studies of the $J = 5 \leftarrow 4$ Transition of Carbon Monoxide Perturbed by CO, N₂, and O₂

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The collisional relaxation of the $J=5\leftarrow 4$ rotational transition of CO induced by carbon monoxide, nitrogen, and oxygen has been studied at room temperature. Pressure-broadening parameters were determined as 3.29(2), 2.61(2), and 2.30(2) MHz/Torr for CO, N₂, and O₂ buffer gases, respectively. Experimental deviations from the Voigt line shape profile have been observed which are mostly the effect of a narrowing in the spectral line core. The difference between the model profile and the experimental profile is less than 0.5% of the maximum line amplitude in the investigated pressure range 0.2–5 Torr. In addition, a small positive collision-induced shift of the line center frequency was observed for the pure gas, corresponding to a pressure self-shift parameter of 6(3) kHz/Torr. © 2002 Elsevier Science (USA)

I. INTRODUCTION

Studies of the effect of collisional relaxation phenomena of gas-phase molecular spectral lines are usually done by observation of the pressure dependence of absorption line profiles. An accurate modeling of such experimental line shapes is particularly important with respect to the interpretation of spectra from Earth's and other planetary atmospheres in order to better understand their composition and their dynamical behavior (1, 2). Precise lineshape parameters for description of collision-induced effects are required not only to interpret atmospheric absorption, but also to better understand intermolecular interactions and to develop more advanced theories describing their effect on the lineshapes of rotational or rovibrational transitions (3).

For these reasons, the CO molecule is of special interest. Carbon monoxide is a trace gas molecule in Earth's as well as other planetary atmospheres and is also rather abundant in the interstellar medium (4, 5). Detection of emission lines of CO from some galaxies could also help to clarify the star formation processes, see for example (6). Thus, there is a necessity to investigate rotational lineshapes of CO in many respects, as particularly pointed out in a recent investigation of the $J=3\leftarrow 2$ transition in the submillimeter wavelength range (7). Moreover, carbon monoxide is an often-studied diatomic molecule with well-known parameters of intermolecular interaction, and therefore it is very interesting to compare the experimental results with a prediction of the theoretical models (8).

The frequency of the CO $J = 5 \leftarrow 4$ transition in the ground vibrational state was determined previously (9-12) but, to our knowledge, collisional relaxation effects for this rotational

line have not been investigated before. A few papers on lineshape investigations were concerned with the study of the R(4) (m = 5) line in the fundamental vibrational band of CO (13–19). In this paper we report on the results of an experimental investigation of the collisional relaxation of the $J = 5 \leftarrow 4$ rotational line of CO induced by carbon monoxide, nitrogen, and oxygen. The description of the experimental setup is also presented.

II. EXPERIMENT AND ANALYSIS

The study was carried out by using a RAD spectrometer (20) with acoustic detection based on a broadband submm Backward Wave Oscillator (BWO) (21), stabilized against the reference signal of a mm-wave synthesizer, "Kvarz" (Nizhii Novgorod, Russia) (22).

A block diagram of the RAD spectrometer with the BWO phase stabilization circuit is shown in Fig. 1. A tube (526-714 GHz) placed in an electromagnet with a magnetic field of about 1.0 T and driven by a high-voltage power supply (voltage ripples <10 mV on high-voltage output ≈ 4000 V, voltage drift <0.1 V/h, overcurrent protection circuit with switch off time $<10 \,\mu s$) with a highly stable filament current output is used as radiation source. High spectral purity and low phase noise (below -25 dBc at 10 kHz in the 600-GHz range) was obtained by means of a phase-lock loop (PLL) circuit.

The PLL-circuit consists of a home-made PLL synchronizer, a quasioptical harmonic mixer with a planar Schottkybarrier diode, a beam splitter (mica plate, thickness $50 \,\mu\text{m}$), and the "Kvarz" mm-wave synthesizer (78-118 GHz, output



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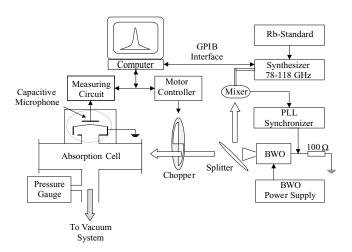


FIG. 1. Block diagram of the RAD spectrometer.

power \approx 10 mW) as the local reference oscillator. A 5-MHz Rb-standard (frequency accuracy $\Delta f/f \approx 10^{-11}$) served as a primary clock. Similar PLL systems of BWOs were used in (23–25).

For precise measurement of lineshape parameters, amplitude modulation at a frequency of 170 Hz was used by employing a synchronous mechanical chopper. A step-by-step frequency scan over a frequency range of about seven linewidths was used for all pressures, controlling the mm-wave synthesizer via an NI GPIB interface. The signal at each frequency step point (integration time about 1 s) during a frequency scan was acquired by a homemade digital synchronous detector (DSD). Both forward and backward frequency scans over the line were performed in order to avoid any amplitude drift instability. Two quadrature channels of the DSD board were used to control the phase between the detected and the reference signal to avoid signal distortion. The electromagnet with the BWO's jack system, the high-voltage power supply, the PLL synchronizer, the harmonic mixer, and the DSD board were made at the Institute of Applied Physics.

The procedure to investigate the pressure dependence of the lineshape was basically the same as described in (26). The studies of self-broadened linewidths were carried out at pressures of the pure gas ranging within 0.2–5 Torr. For investigation of foreign gas pressure broadening, the RAD cell was first filled with pure carbon monoxide to a fixed pressure of about 1.0 Torr. After that, a controlled amount of foreign gas was added and the corresponding lineshape parameters were determined. Further adding the foreign gas up to a partial pressure of about 4 Torr continued this procedure. The total pressure in the cell was measured by a MKS Baratron absolute pressure gauge (Type 122AAX-00010ADS) with a stated accuracy of 0.15%. All measurements were done at room temperature $(295 \pm 2 \text{ K})$.

The experimental data set for the RAD cell signal was fitted by using the Nelder–Meed algorithm (27) for the following

model function related to a Voigt lineshape with a fixed Doppler halfwidth,

$$I(\nu_n) = a_0 \cdot (1 + a_1 \cdot (\nu_n - \nu_0)) \cdot \left(1 - e^{-a_2 \cdot F(\nu_n - \nu_0, \Delta \nu_L)}\right) + b_0 + b_1 \cdot (\nu_n - \nu_0),$$
[1]

where a_0 is an amplitude factor proportional to the incident radiation power, the coefficient a_1 corresponds to the linear trend of the BWO output power, ν_n is the frequency, ν_0 is the line center frequency, a_2 is proportional to the optical depth of the sample gas, $F(\nu_n - \nu_0, \Delta \nu_L)$ is the Voigt profile function with a Doppler half width $\Delta \nu_D$ at a given temperature ($\Delta \nu_D = 0.670$ MHz for the $J = 5 \leftarrow 4$ CO line at 295 K), $\Delta \nu_L$ is the half width at half maximum (HWHM) of the Lorentzian lineshape which accounts for pressure broadening, and b_0 and b_1 are first coefficients of polynomial series describing the baseline fringes. As an example, a recording of the self-broadened $J = 5 \leftarrow 4$ line of CO at 0.38 Torr pressure with the residuum of the fit to the Voigt model function (Eq. [1]) is presented in Fig. 2.

It is worth to point out that because of the short length of the cell (l = 10 cm) in comparison to conventional absorption cells and the distinctive feature of "volume" detection of the photoacoustic cell (the cell itself is a selective power detector),

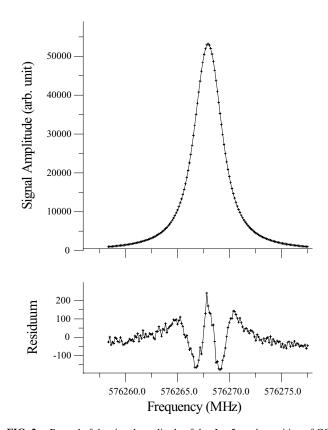


FIG. 2. Record of the signal amplitude of the $J=5 \leftarrow 4$ transition of CO line at a pressure of 0.38 Torr (crosses) together with the theoretical Voigt profile (straight line). The residuum resulting from a fit of 150 data points to the Voigt profile model is also shown.

the slope and amplitude of the baseline are strongly reduced what is extremely important for lineshape studies.

III. RESULTS AND DISCUSSION

The results of the present investigation on pressure broadening and shift of line center of the carbon monoxide $J=5 \leftarrow 4$ transition in the ground vibrational state at room temperature are presented in Table 1. The broadening and shift parameters given in the table are the coefficients which describe the linear dependence of $\Delta \nu_L$ (HWHM) and ν_0 (see Eq. [1]) on the pure gas pressure or the foreign gas partial pressures, respectively. The pressure dependence of $\Delta \nu_L$ is depicted in Fig. 3 for broadening by CO, N₂, and O₂ together with linear regression lines from least squares analysis.

The experimental results demonstrate the quite good sensitivity of our spectrometer. The residues which were found from the fitting to a Voigt profile line function (for example, see Fig. 2) show a pattern similar to the ones observed in a recent investigation on the $J=3\leftarrow 2$ of CO perturbed by N_2 and O_2 (7), but have somewhat less residuum in our case (about 0.5% of the line peak amplitude for all pressures).

In (7) these significant deviations (typically about 2% at a total gas pressure of 0.18 Torr) were explained by a narrowing of the line due to speed dependence effects of relaxation rate (28) as well as from velocity-changing collisions (collisional Dicke narrowing) (29). In the present study collisional broadening is, for the higher pressures, much more important than Doppler broadening. The effect of Dicke narrowing may be then be neglected, but speed dependence of relaxation rates and effects of velocity averaging may still contribute significantly to distortions on the Voigt (30) or Lorentzian (31) lineshape profiles. However, in view of the fact that pressure-broadening parameters which are derived from a speed-independent model are normally not much affected by an analysis based on more sophisticated lineshape models, no such further analysis was done in the present paper. We did also not measure and consider the behavior in the far wings of absorption lines where non-Lorentzian profiles prevail.

Additionally, it should be pointed out that the observed deviations from Voigt profiles may also result from distortions such as

TABLE 1 Pressure-Broadening and Frequency-Shift Parameters of the $J=5\leftarrow4$ Rotational Line of Carbon Monoxide Due to Collisionally Induced Relaxation by CO, N_2 , and O_2 at Room Temperature

| CO and Perturber | Broadening parameter* MHz/Torr | Shift parameter* kHz/Torr | |
|---------------------|-----------------------------------|------------------------------|--|
| СО | 3.29(2) | 6(3) | |
| N_2 | 2.61(2) | 5(5) | |
| O ₂ | 2.30(2) | 6(6) | |

^{*}Error given in parentheses is in units of the last digit quoted.

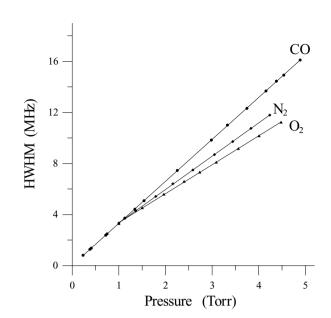


FIG. 3. Pressure dependence of the HWHM of the $J=5 \leftarrow 4$ rotational line of carbon monoxide, broadened by CO, N₂, and O₂. The error bars retrieved from Voigt profiles for the individual points are for one standard deviation less than the size of the data point symbols.

transit time effects, e.g., collisions with the walls of the sample cell (34).

Apart from pressure broadening of the investigated line, we also observed a small, but significant pressure dependence of the line center frequency of the transition for the pure CO gas. Figure 4 shows the results for this positive selfshift of the

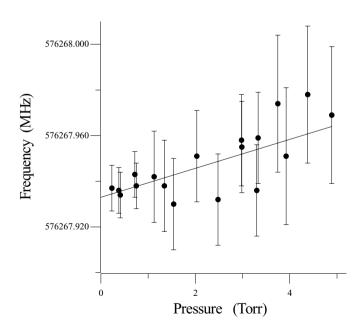


FIG. 4. Pressure dependence of center frequency of the $J=5 \leftarrow 4$ line of pure carbon monoxide. The fitted straight line with zero pressure intercept $v_0=576267.933(5)$ MHz and slope (selfshift parameter = 6(3) kHz/Torr) is also shown.

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TABLE 2
Zero-Pressure CO Line Center
Frequency of $J = 5 \leftarrow 4$

| Frequency (MHz) | Reference |
|-----------------|-----------|
| 576267.75(10) | (9) |
| 576267.934 | (10) |
| 576267.910(72) | (11) |
| 576267.9305(5) | (12) |
| 576267.933(5) | this work |

line center. The corresponding slope value (shift parameter) is also given in Table 1 together with our results for the mixtures with the foreign gases nitrogen and oxygen, respectively, which also indicate a small positive but poorly determined pressure shift.

IV. CONCLUSION

The zero-pressure line center frequency of the $J = 5 \leftarrow 4$ transition of carbon monoxide, $v_0 = 576267.933(5)$ MHz, determined in this work demonstrates good consistency with values of previous measurements (9–12), especially with the results of Gordy's group (10) and a sub-Doppler investigation by Winnewisser *et al.* (12); see Table 2.

Our results from the pressure-broadening investigation may be compared with broadening coefficients at room temperature reported from IR studies on the R(4) (m=5) line of the fundamental vibrational band (13-19), which involves the same lower (J=4) and upper (J=5) rotational state as for the line reported here. As shown in Table 3, the results from the two spectroscopic ranges are in fair agreement for self-broadening, and very consistent for foreign gas broadening with N_2 and O_2 . The smaller value reported in (I5) for the self-broadening parameter may

TABLE 3
Comparison of the Pressure-Broadening Parameter for the $J=5\leftarrow 4$ Line in the Ground Vibrational State of Carbon Monoxide Perturbed by CO, N₂, and O₂, with Literature Data for the *R*-branch, m=5 Transition of the Fundamental Vibrational Band

| CO and Perturber | Broadening parameter MHz/Torr | Temperature K | References | |
|---------------------|----------------------------------|------------------|------------|--|
| СО | 3.29(2) | 295 | this work | |
| | 2.84(1) | 300 | (15) | |
| | 2.95(1) | 296 | (17) | |
| N_2 | 2.61(2) | 295 | this work | |
| - | 2.68 | 298 | (13) | |
| | 2.61 | 300 | (14) | |
| | 2.67(5) | 300 | (15) | |
| | 2.61 | 298 | (16) | |
| | 2.58(3) | 301 | (18) | |
| | 2.57(3) | 301 | (19) | |
| O_2 | 2.30(2) | 295 | this work | |
| | 2.22(5) | 300 | (15) | |

TABLE 4

Comparison of the Pressure-Broadening Parameter for the $J=5\leftarrow 4$ Line of Carbon Monoxide at Room Temperature, Perturbed by CO, N_2 , and O_2 , with Literature Data for Other Rotational Transition of CO in the Ground Vibrational State

| Transition $J' \leftarrow J$ | CO MHz/Torr | N ₂ MHz/Torr | O ₂ MHz/Torr | References |
|------------------------------|----------------|----------------------------|----------------------------|------------|
| 1 ← 0 | 3.43(1) | 3.22(2) | 2.83(2) | (26) |
| $2 \leftarrow 1$ | 3.36(7) | 3.01(5) | 2.64(5) | (35) |
| $3 \leftarrow 2$ | a | 2.80(2) | 2.46(2) | (7) |
| 5 ← 4 | 3.29(2) | 2.61(2) | 2.30(2) | this work |

^a 3.10(6) from IR measurement on the R(2) (m = 3) line in the fundamental vibrational band [15].

indicate some systematic deviation to smaller values for the results from the measurement in the IR region; see also Table 4 below, which compares with pressure-broadening parameters of other rotational transitions of CO. The observed consistency of linewidth data for pure rotational and rovibrational transitions allows the conclusion that line broadening of the rovibrational line is predominantly governed by rotational relaxation phenomena involving rates due to collision-induced transitions between rotational levels which are not much dependent on the vibrational state. In contrast to this behavior, the lineshift data for the observed pure rotational transition (see Table 1) differ significantly from the corresponding values for the R-branch m = 5 transition of the fundamental vibrational band, determined to -82(10) kHz/Torr for CO-CO (17), -72.8 kHz/ Torr for CO-N₂ (32), and -84.9 kHz/Torr for CO-O₂ (33), respectively. Such differences are related to the fact that the rovibrational lines are predominantly shifted due to the isotropic part of the shifting cross section which does not contribute to the shifts of pure rotational transitions (36).

Finally, it may be interesting to compare our results with the experimentally determined broadening coefficients for the $J=1\leftarrow 0$ (26), $J=2\leftarrow 1$ (35), and $J=3\leftarrow 2$ (7) pure rotational lines of CO; see Table 4. For the pure gas, as well as for the mixtures with the foreign gases N_2 and O_2 , our results fit very well to the row of J-quantum numbers which show a slight decrease of the broadening parameters for a $J+1\leftarrow J$ transition with increasing J.

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