

1. 如图所示,滑块与杆 AB 铰接,静止放在光滑的水平面上。小球 D 以垂直于杆的速度 v 与杆 AB 的端点 B 发生碰撞。恢复因数为 e=0.5。 滑块与杆 AB 的质量均为 m,小球 D 的质量为 2m。

求

- (1) 碰撞后滑块 A 的速度和杆 AB 的角速度。
- (2) A处的约束冲量

解

$$m(v_{1\tau} - 0) = I_1 \quad (1)$$

$$m(v_{2\tau} - 0) = I_2 - I_1 \quad (2)$$

$$\frac{1}{12}ml^2(\omega_{2\tau} - 0) = \frac{l}{2}(I_2 + I_1) \quad (3)$$

$$2m(v_{3\tau} - (-v)) = I_2 \quad (4)$$

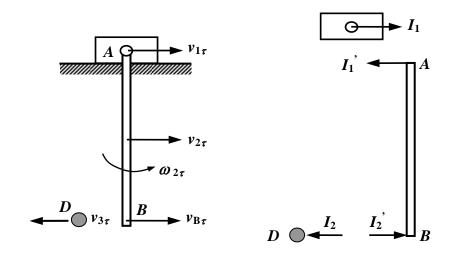
$$\frac{v_{B\tau} - (-v_{3\tau})}{v} = e = \frac{1}{2} \quad (5)$$

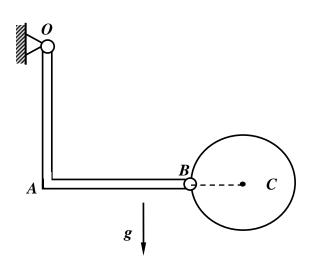
$$v_{2\tau} = v_{1\tau} + \frac{l}{2}\omega_{2\tau} \quad (6)$$

$$v_{B\tau} = v_{1\tau} + l\omega_{2\tau} \quad (7)$$

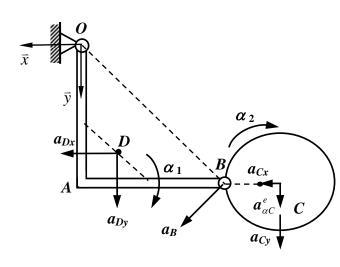
解得
$$v_{1\tau} = -\frac{6}{37}v$$
, $\omega_{2\tau} = \frac{54}{37} \cdot \frac{v}{l}$

$$I_1 = m(v_{1\tau} - 0) = -\frac{6}{37}mv$$





2. 如图所示,半径为 r 的圆盘与匀质折杆 OAB 在 B 处铰接,OA = AB = 2r 。设圆盘的质量为 m, 折杆 OAB 的质量为 2m, 图示位置 AB 水平, BC 水平。用达朗贝尔原理求系统在图示位置无初速 开始运动时折杆 OAB 的角加速度和 B 端作用于圆盘的约束力。



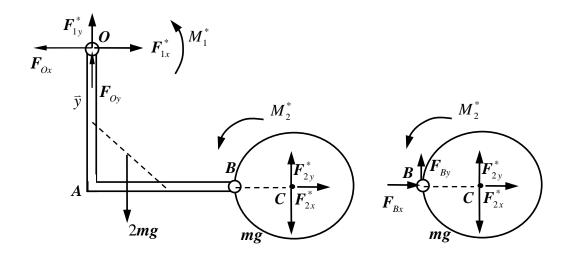
$$a_{Dx} = \frac{3}{4}l\alpha_1 = \frac{3}{2}r\alpha_1$$
, $a_{Dy} = \frac{1}{4}l\alpha_1 = \frac{1}{2}r\alpha_1$ $a_{CX} + a_{CY} = a_B + a_{\alpha C}^e$, $a_{\alpha C}^e = r\alpha_2$, $a_B = 2r\sqrt{2}\alpha_1$ 得到

$$a_{CX} = 2r\alpha_1, \quad a_{CY} = 2r\alpha_1 + r\alpha_2$$

$$F_{1x}^* = 2ma_{Dx} = 3mr\alpha_1, \quad F_{1y}^* = 2ma_{Dy} = mr\alpha_1$$

$$F_{2x}^* = ma_{Cx} = 2mr\alpha_1, \quad F_{2y}^* = m(2r\alpha_1 + r\alpha_2)$$

$$\begin{split} J_O &= \frac{ml^2}{3} + \frac{ml^2}{12} + m \left[\left(\frac{l}{2} \right)^2 + l^2 \right] = \frac{5ml^2}{12} + \frac{5ml^2}{4} = \frac{20ml^2}{12} = \frac{5ml^2}{3} = \frac{20}{3}mr^2 \\ M_1^* &= J_O \alpha_1 = \frac{20}{3}mr^2 \alpha_1, \ M_2^* = J_C \alpha_2 = \frac{1}{2}mr^2 \alpha_2 \end{split}$$



取系统为对象,对 O 点取矩

$$M_1^* + M_2^* + 2rF_{2x}^* + 3rF_{2y}^* - 2mg \cdot \frac{r}{2} - mg \cdot 3r = 0$$

$$\frac{20}{3}mr^2\alpha_1 + \frac{1}{2}mr^2\alpha_2 + 2r\cdot 2mr\alpha_1 + 3mr\left(2r\alpha_1 + r\alpha_2\right) - 4mgr = 0$$

即

$$\frac{50}{3}mr^{2}\alpha_{1} + \frac{7}{2}mr^{2}\alpha_{2} - 4mgr = 0 \qquad (1)$$

取圆盘为对象,对B点取矩

$$M_2^* + r F_{2y}^* - mgr = 0$$

$$\frac{1}{2}mr^2\alpha_2 + mr(2r\alpha_1 + r\alpha_2) - mgr = 0$$

即

$$2mr^{2}\alpha_{1} + \frac{3}{2}mr^{2}\alpha_{2} - mgr = 0 \qquad (2)$$

得到:
$$\alpha_2 = \frac{2g}{3r} - \frac{4}{3}\alpha_1$$
 (3)

将(3)代入(1),解得:

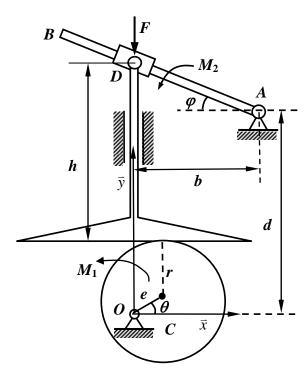
$$\alpha_1 = \frac{5g}{36r}, \quad \alpha_2 = \frac{13g}{27r}$$

取圆盘为对象

$$F_{Bx} = -F_{2x}^* = -2mr\alpha_1 = -\frac{5}{18}mg$$

$$F_{2y}^* = m(2r\alpha_1 + r\alpha_2) = m\left(\frac{5}{18}g + \frac{13}{27}g\right) = \frac{41}{54}mg$$

$$F_{By} = mg - F_{2y}^* = mg - \frac{41}{54}mg = \frac{13}{54}mg$$

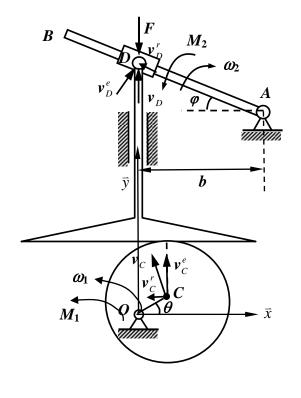


3. 凸轮机构在图示位置处于平衡。设凸轮的半径为r,轮心偏离转轴的距离为e=r/2。铰点A到y轴的距离为 $b=\sqrt{3}r$,铰点A到x轴的距离为d=3r,挺杆的高度为 $h=\frac{11}{4}r$ 。图示位置 $\theta=30^\circ$,

 $\varphi=30^{\circ}$ 。设各构件的重量和摩擦力不计, 用虚位移原理分析

- (1)系统平衡时 M_1 , M_2 与F之间的关系
- (2) 铰点 0 垂直方向的约束力

解[1]: (1) $v_D = v_C^e = \cos \theta v_C = \cos \theta e \omega_1$ (1)



$$v_D^e = \cos \varphi v_D = \frac{b}{\cos \varphi} \omega_2$$

得到

$$\frac{b}{\cos^2 \varphi} \omega_2 = \cos \theta e \omega_1 \quad (2)$$

(1)和(2)两边乘以 dt,得到

$$dy_D = \cos\theta e d\theta$$

 $bd\varphi = \cos\theta\cos^2\varphi ed\theta$

对于定常约束,有

$$\delta y_D = \cos \theta e \delta \theta \qquad (3)$$

$$\delta\varphi = \frac{e}{b}\cos\theta\cos^2\varphi\,\delta\theta \qquad (4)$$

由虚位移原理

$$M_1 \delta \theta - F \delta y_D - M_2 \delta \varphi = 0$$
,

将(3)和(4)代入:

$$\left(M_{1} - F\cos\theta e - M_{2}\frac{e}{b}\cos\theta\cos^{2}\varphi\right)\delta\theta = 0 \ \text{#} \ \theta = \varphi = 30^{\circ} \text{#}\lambda:$$

$$M_1 - \frac{\sqrt{3}}{4}rF - \frac{3}{16}M_2 = 0$$

(2)释放O点Y方向的约束

为两个自由度问题,广义坐标为 θ, φ

S为凸轮的速度顺心

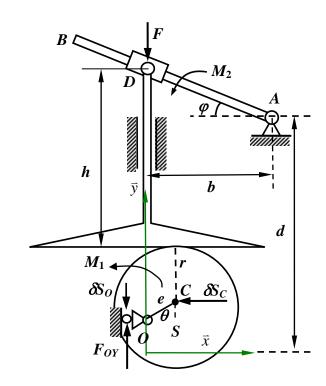
$$\delta S_o = \delta \theta e \cos \theta \quad (5)$$

由虚位移原理 $M_1\delta\theta - F_{OV}\delta S_O = 0$

将(5)代入:

$$(M_1 - F_{OY}e\cos\theta)\delta\theta = 0,$$

得到
$$F_{OY} = \frac{M_1}{e\cos\theta} = \frac{4\sqrt{3}}{3} \frac{M_1}{r}$$



或者取广义坐标为 y_o , θ , $\vartheta \delta \theta = 0$, 根据关系式 $y_D = y_o + e \sin \theta + r + h$

$$\delta y_D = \delta y_O$$
, 又根据关系式 $y_D = y_O + e \sin \theta + r + h = d + b \operatorname{tg} \varphi$

$$\delta y_o = b \sec^2 \varphi \delta \varphi$$
, $\forall \varphi = \frac{\cos^2 \varphi}{b} \delta y_o$

由虚位移原理 $F_{OY}\delta y_O - F\delta y_D - M_2\delta \varphi = 0$

$$F_{OY}\delta y_O - F\delta y_O - M_2 \frac{\cos^2 \varphi}{b} \delta y_O = 0$$

$$F_{OY} = F + M_2 \frac{\cos^2 \varphi}{b} = F + \frac{\sqrt{3}}{4} \cdot \frac{M_2}{r}$$

(由于
$$M_1 - \frac{\sqrt{3}}{4}rF - \frac{3}{16}M_2 = 0$$
, $F + \frac{\sqrt{3}}{4} \cdot \frac{M}{r} = \frac{4\sqrt{3}}{3} \frac{M_1}{r}$), 说明取 θ , φ 和 y_o , θ 得到的

结果相同。前面取法简单。

解[2]: (1) 根据几何关系, $y_D = h + r + e \sin \theta = d + b \operatorname{tg} \varphi$

两边求变分: $\delta y_D = e \cos \theta \delta \theta = b \sec^2 \varphi \delta \varphi$

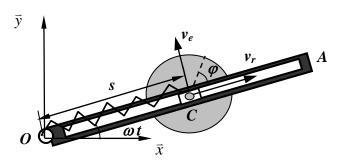
由虚位移原理 $M_1\delta\theta - F\delta y_D - M_2\delta\varphi = 0$

代入得到 $M_1\delta\theta - Fe\cos\theta\delta\theta - M_2\frac{e}{h}\cos\theta\cos^2\varphi\delta\theta = 0$

将 $\theta = 30^{\circ}$, $\varphi = 30^{\circ}$ 代入, 化简得到:

$$M_1 - \frac{\sqrt{3}}{4}rF - \frac{3}{16}M_2 = 0$$
 (2) 同前

4. 如图所示,<u>在水平面内</u>,杆 OA 以匀角速度 ω 绕 O 点转动。滑块可沿 OA 内的滑槽无摩擦地滑动,滑块与圆盘在圆盘中心 C 处铰接。设 s 为圆盘中心 C 相对杆 OA 的位移, φ 为圆盘相对杆 OA 的转角。杆 OA 和圆盘的质量均为 m,滑块的质量不计。杆 OA 关于 OZ 的转动惯量为 $m\rho^2$ 。圆盘中心 C 与 O 点通过线弹簧连接,弹簧刚度为 k,原长为 s_0 。以 φ 和 s 为独立坐标,写出系统的拉格朗日函数和初积分。



盤.

$$v_r = \dot{s}$$
, $v_e = s\omega$, $\omega_C = \omega + \dot{\varphi}$

$$T = \frac{1}{2}m\rho^{2}\omega^{2} + \frac{1}{2}m(v_{r}^{2} + v_{e}^{2}) + \frac{1}{2}\cdot\frac{1}{2}mr^{2}(\omega + \dot{\varphi})^{2}$$

$$= \frac{1}{2}m\rho^{2}\omega^{2} + \frac{1}{2}m(\dot{s}^{2} + s^{2}\omega^{2}) + \frac{1}{4}mr^{2}(\omega + \dot{\varphi})^{2}$$

$$= T_{0} + T_{1} + T_{2}$$

其中

$$\begin{split} T_0 &= \frac{1}{2} m \rho^2 \omega^2 + \frac{1}{2} m s^2 \omega^2 + \frac{1}{4} m r^2 \omega^2 \\ T_1 &= \frac{1}{2} m r^2 \omega \dot{\varphi} \;, \quad T_2 &= \frac{1}{2} m \dot{s}^2 + \frac{1}{4} m r^2 \dot{\varphi}^2 \\ V &= \frac{1}{2} k \left(s - s_0 \right)^2 \end{split}$$

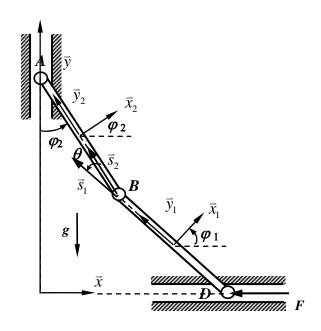
$$L = T - V = \frac{1}{2}m\rho^{2}\omega^{2} + \frac{1}{2}ms^{2}\omega^{2} + \frac{1}{4}mr^{2}\omega^{2} + \frac{1}{2}mr^{2}\omega\dot{\phi}$$
$$+ \frac{1}{2}m\dot{s}^{2} + \frac{1}{4}mr^{2}\dot{\phi}^{2} - \frac{1}{2}k(s - s_{0})^{2}$$

因为L不显含 φ ,得到循环积分式:

$$\frac{\partial L}{\partial \dot{\varphi}} = \frac{1}{2} mr^2 \omega + \frac{1}{2} mr^2 \dot{\varphi} = C_1$$

因为L不显含t,且是非定常约束,得到广义能量积分式:

$$\begin{split} &T_2 - T_0 + V \\ &= \frac{1}{2}m\dot{s}^2 + \frac{1}{4}mr^2\dot{\phi}^2 - \frac{1}{2}m\rho^2\omega^2 - \frac{1}{2}ms^2\omega^2 - \frac{1}{4}mr^2\omega^2 + \frac{1}{2}k(s - s_0)^2 = C_2 \end{split}$$



5. 如图动力学系统由长均为 2l 的杆 AB (B₂)和杆 BD (B₁)组成,杆 AB 和杆 BD 在 B 处铰接,销子 A 和销子 D 分别可以在铅垂滑槽和水平滑槽内滑动。水平力 F 作用于 D 点。不计摩擦。设杆 AB 和杆 BD 的质量均为 m。 B 处在两杆间有一卷簧,弹簧系数为 k, B₁ 相对 B₂ 的转角为θ。当两杆在一直线时,卷簧无力偶。以系统位形坐标写出封闭的带拉格朗日乘子的第一类拉格朗日方程。

解: 系统位形坐标为 $\mathbf{q} = \begin{bmatrix} x_1 & y_1 & \varphi_1 & x_2 & y_2 & \varphi_2 \end{bmatrix}^T$ 约束方程为

$$\boldsymbol{\Phi} = \begin{bmatrix} \begin{bmatrix} 1 & 0 \end{bmatrix} \begin{pmatrix} \begin{bmatrix} x_2 \\ y_2 \end{bmatrix} + A_2 \begin{bmatrix} 0 \\ l \end{bmatrix} \end{pmatrix} \\ \begin{bmatrix} 0 & 1 \end{bmatrix} \begin{pmatrix} \begin{bmatrix} x_1 \\ y_1 \end{bmatrix} + A_1 \begin{bmatrix} 0 \\ -l \end{bmatrix} \end{pmatrix} \\ \begin{bmatrix} x_1 \\ y_1 \end{bmatrix} + A_1 \begin{bmatrix} 0 \\ l \end{bmatrix} - \begin{bmatrix} x_2 \\ y_2 \end{bmatrix} - A_2 \begin{bmatrix} 0 \\ -l \end{bmatrix} \end{bmatrix} = \mathbf{0}$$

$$\boldsymbol{\Phi} = \begin{bmatrix} x_2 - l \sin \varphi_2 \\ y_1 - l \cos \varphi_1 \\ x_1 - x_2 - l \sin \varphi_1 - l \sin \varphi_2 \\ y_1 - y_2 + l \cos \varphi_1 + l \cos \varphi_2 \end{bmatrix} = \mathbf{0}$$

加速度约束方程为 $\boldsymbol{\Phi}_{\!\!\!\!q}\,\ddot{q}=\boldsymbol{\gamma}$, 其中

$$\boldsymbol{\varPhi}_{q} = \begin{bmatrix} 0 & 0 & 0 & 1 & 0 & -l\cos\varphi_{2} \\ 0 & 1 & l\sin\varphi_{1} & 0 & 0 & 0 \\ 1 & 0 & -l\cos\varphi_{1} & -1 & 0 & -l\cos\varphi_{2} \\ 0 & 1 & -l\sin\varphi_{1} & 0 & -1 & -l\sin\varphi_{2} \end{bmatrix}, \quad \boldsymbol{\gamma} = \begin{bmatrix} -l\sin\varphi_{2}\dot{\varphi}_{2}^{2} \\ -l\cos\varphi_{1}\dot{\varphi}_{1}^{2} \\ -l\sin\varphi_{1}\dot{\varphi}_{1}^{2} - l\sin\varphi_{2}\dot{\varphi}_{2}^{2} \\ l\cos\varphi_{1}\dot{\varphi}_{1}^{2} + l\cos\varphi_{2}\dot{\varphi}_{2}^{2} \end{bmatrix}$$

带拉格朗日乘子的第一类拉格朗日方程为

$$\mathbf{Z}\ddot{\mathbf{q}} + \mathbf{\Phi}_{\mathbf{q}}^{\mathrm{T}} \boldsymbol{\lambda} = \hat{\mathbf{F}}^{a}$$

其中,
$$\mathbf{Z} = diag\left(m, m, \frac{1}{12}m(2l)^2, m, m, \frac{1}{12}m(2l)^2\right)$$

$$\hat{\mathbf{F}}^{a} = \begin{bmatrix} -F & -mg & -Fl\cos\varphi_1 + M & 0 & -mg & -M \end{bmatrix}^{\mathrm{T}}$$

其中,
$$M = -k(\theta - \theta_0)$$
, $\theta = \varphi_1 - \varphi_2$, $\theta_0 = 0$

即

$$\hat{\mathbf{F}}^{a} = \begin{bmatrix} -F & -mg & -Fl\cos\varphi_{1} - k(\varphi_{1} - \varphi_{2}) & 0 & -mg & k(\varphi_{1} - \varphi_{2}) \end{bmatrix}^{T}$$

$$\lambda = \begin{bmatrix} \lambda_1 & \lambda_2 & \lambda_3 & \lambda_4 \end{bmatrix}^T$$

封闭的带拉格朗日乘子的第一类拉格朗日方程为

$$\begin{bmatrix} \boldsymbol{Z} & \boldsymbol{\mathcal{\Phi}}_{q}^{\mathrm{T}} \\ \boldsymbol{\mathcal{\Phi}}_{q} & \mathbf{0} \end{bmatrix} \begin{bmatrix} \ddot{\boldsymbol{q}} \\ \boldsymbol{\lambda} \end{bmatrix} = \begin{bmatrix} \hat{\boldsymbol{F}}^{a} \\ \boldsymbol{\gamma} \end{bmatrix}$$