

PROBLEM SET 2 – APMA 0360

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Solution 1.

a. We have

$$\begin{aligned}\hat{g}(k) &= \int_{-\infty}^{\infty} e^{ikx} g(x) dx \\ \hat{g}(k) &= \int_{-\infty}^{\infty} e^{ikx} f(x-a) dx\end{aligned}$$

Letting $u = x - a$, which implies $du = dx$, we have

$$\begin{aligned}\hat{g}(k) &= \int_{-\infty}^{\infty} e^{ik(u+a)} f(u) du \\ \hat{g}(k) &= e^{ika} \int_{-\infty}^{\infty} e^{iku} f(u) du \\ \hat{g}(k) &= e^{ika} \hat{f}(k)\end{aligned}$$

And we are done.

b. We start by taking the Fourier transform of both sides, using integration by parts to deal with the transform of u_x .

$$\begin{aligned}u_t + cu_x &= 0 \\ \int_{-\infty}^{\infty} e^{ikx} u_t dx + c \int_{-\infty}^{\infty} e^{ikx} u_x dx &= 0 \\ \frac{d}{dt} \hat{u}(k, t) - ikc \int_{-\infty}^{\infty} e^{ikx} u dx &= 0 \\ \frac{d}{dt} \hat{u}(k, t) - ikc \hat{u}(k, t) &= 0\end{aligned}$$

For the sake of simplification, define $p = \hat{u}(k, t)$. Then

$$\begin{aligned}\dot{p} - ikcp &= 0 \\ \frac{dp}{dt} &= ikcp \\ \int p^{-1} dp &= \int ikc dt \\ \ln(p) &= ikct + C(x) \\ p &= e^{ikct+C(x)} \\ \hat{u}(k, t) &= e^{ikct} B(x)\end{aligned}$$

Note that the Fourier transform of $u(x, 0) = f(x)$ is simply $\hat{u}(k, 0) = \hat{f}(k)$. So, we have

$$\begin{aligned}\hat{u}(k, 0) &= e^{ik \cdot 0} B(x) \\ \hat{u}(k, 0) &= B(x) = \hat{f}(k) \\ \hat{u}(k, t) &= e^{ikct} \hat{f}(k)\end{aligned}$$

Finally, by part (a) with $a = ct$, we get the desired result

$$u(x, t) = f(x - ct)$$

Solution 2.

- a. Given a solution of the form $u(x, t) = e^{\lambda t} e^{ikx}$, we start by plugging it into the heat equation

$$\begin{aligned}u_t &= Du_{xx} \\ \lambda e^{\lambda t} e^{ikx} &= -Dk^2 e^{\lambda t} e^{ikx}\end{aligned}$$

Since e^{ikx} and $e^{\lambda t}$ are necessarily non-zero, we can divide out by them, giving

$$\lambda = -Dk^2$$

And we are done.

- b. Using our answer from (a), we write

$$\begin{aligned}u(x, t) &= \int_{-\infty}^{\infty} a(k) e^{\lambda t} e^{ikx} dk \\ u(x, t) &= \int_{-\infty}^{\infty} a(k) e^{-Dk^2 t} e^{ikx} dk\end{aligned}$$

Now, we use the heat equation to find u_t and u_{xx}

$$\begin{aligned}u_t &= \int_{-\infty}^{\infty} -Dk^2 a(k) e^{-Dk^2 t} e^{ikx} dk \\ u_{xx} &= \int_{-\infty}^{\infty} -k^2 a(k) e^{-Dk^2 t} e^{ikx} dk\end{aligned}$$

Plugging this into the heat equation gives

$$\begin{aligned} u_t &= Du_{xx} \\ \int_{-\infty}^{\infty} -Dk^2 a(k) e^{-Dk^2 t} e^{ikx} dk &= D \int_{-\infty}^{\infty} -k^2 a(k) e^{-Dk^2 t} e^{ikx} dk \\ \int_{-\infty}^{\infty} -Dk^2 a(k) e^{-Dk^2 t} e^{ikx} dk &= \int_{-\infty}^{\infty} -Dk^2 a(k) e^{-Dk^2 t} e^{ikx} dk \end{aligned}$$

And since the LHS and the RHS are the exact same expression, $u(x, t)$ is a solution. Note that $a(k)$ being a part of the Schwartz class is important to ensure that the integral is finite.

Solution 3.

We have

$$\begin{aligned} u_t &= u_{xxxx} \\ \int_{-\infty}^{\infty} e^{ikx} u_t dx &= \int_{-\infty}^{\infty} e^{ikx} u_{xxxx} dx \end{aligned}$$

We use integration by parts four times on the RHS to simplify the integral as follows

$$\begin{aligned} \int_{-\infty}^{\infty} e^{ikx} u_t dx &= \int_{-\infty}^{\infty} (ik)^4 e^{ikx} u dx \\ \frac{\partial}{\partial t} \hat{u}(k, t) &= k^4 \hat{u}(k, t) \end{aligned}$$

This is an ODE of the form $y' = ky$, which we can solve easily as follows

$$\hat{u}(k, t) = C(k) \cdot e^{k^4 t}$$

Note that

$$\hat{u}(k, 0) = \int_{-\infty}^{\infty} e^{ikx} u(x, 0) dx = \int_{-\infty}^{\infty} e^{ikx} f(x) dx = \hat{f}(k)$$

So

$$\hat{u}(k, 0) = C(k) \cdot e^{k^4 \cdot 0} = C(k) = \hat{f}(k)$$

Giving us a final solution of

$$\hat{u}(k, t) = \hat{f}(k) \cdot e^{k^4 t}$$

Solution 4.

a. We start by finding v_t and v_{yy} . We have

$$\begin{aligned} v_t &= \frac{\partial}{\partial t} u(y - ct, t) e^{at} \\ &= a \cdot u(y - ct, t) e^{at} + (-c \cdot u_x(y - ct, t) + u_t(y - ct, t)) e^{at} \\ v_{yy} &= u_{xx}(y - ct, t) e^{at} \end{aligned}$$

We can plug this into the equation $v_t = Dv_{yy}$ to get

$$e^{at}(au(y-ct, t) - cu_x(y-ct, t) + u_t(y-ct, t)) = De^{at}(u_{xx}(y-ct, t))$$

Since e^{at} is necessarily non-zero, we can divide out by it, giving

$$\begin{aligned} au(y-ct, t) - cu_x(y-ct, t) + u_t(y-ct, t) &= Du_{xx}(y-ct, t) \\ u_t(y-ct, t) &= Du_{xx}(y-ct, t) + cu_x(y-ct, t) - au(y-ct, t) \\ u_t &= Du_{xx} + cu_x - au \end{aligned}$$

Which is given as true, and we are done.

- b. Based on the equation $u_t = u_{xx} + 2u_x - 0.5u$, we know $D = 1$, $c = 2$, and $a = 0.5$. Then, we use the fact that $x = y - ct \rightarrow y = x + ct$ and substitute to get

$$\begin{aligned} v(y, t) &= u(y - ct, t)e^{at} \\ &= u(y - 2t, t)e^{0.5t} \\ \frac{1}{\sqrt{4\pi t}}e^{-\frac{y^2}{4t}} &= u(y - 2t, t)e^{0.5t} \\ \frac{1}{\sqrt{4\pi t}}e^{-\frac{y^2}{4t}}e^{-0.5t} &= u(y - 2t, t) \\ x = y - 2t &\rightarrow y = x + 2t \\ u(x, t) &= \exp\left(-0.5t - \frac{(x + 2t)^2}{4t}\right) \cdot \frac{1}{\sqrt{4\pi t}} \end{aligned}$$

Which is our $u(x, t)$.