

Fast Polarization Manipulation

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Research Proposal

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1 The Proposal

1.1 Problem Definition and Solution Approach.

The presence of the splitting in the excitonic state is an unwanted effect in QDs that causes a degradation in the degree of entanglement between the two photons in the biexciton-exciton radiative cascade [1]. This is mainly due to the asymmetry of the QD which lifts the degeneracy of the levels. Several methods were investigated to control such as inducing strain and magnetic and electric fields, but the problem continues to persist. A more straightforward approach is to ignore the causes for the splitting and assume that

1.2 Theoretical formulation (exciton).

$$|\Psi(t)\rangle = \alpha(|H_{XX} \otimes H_X\rangle \cdot e^{\frac{-iE_H t}{\hbar}} + |V_{XX} \otimes V_X\rangle \cdot e^{\frac{-iE_V t}{\hbar}}) \quad (1)$$

1.3 Theoretical formulation (biexciton-exciton)

$$|\Psi(t)\rangle = \alpha(|H_{XX} \otimes H_X\rangle \cdot e^{\frac{-iE_H t}{\hbar}} + |V_{XX} \otimes V_X\rangle \cdot e^{\frac{-iE_V t}{\hbar}}) \quad (2)$$

$$|\Psi(t)\rangle = \alpha(|H_{XX} \otimes H_X\rangle + |V_{XX} \otimes V_X\rangle \cdot e^{\frac{-i(\Delta E t)}{\hbar}}) \quad (3)$$

1.4 Rotation of the exciton's polarization.

1.5 Restoring the entanglement of the photons in the biexciton-exciton radiative cascade.

If we construct two interferometers that allow us to induce phase shift to both the biexciton and exciton polarizations, then we can write the time-dependent phase

relations as:

$$\begin{aligned}
\Phi_{H_{XX}}(t, t_{prop}) &= k_{H_{XX}} \cdot (t - t_{prop}) + \Phi_{H_{XX}}^0 \\
\Phi_{V_{XX}}(t, t_{prop}) &= k_{V_{XX}} \cdot (t - t_{prop}) + \Phi_{V_{XX}}^0 \\
\Phi_{H_X}(t, t_{prop}) &= k_{H_X} \cdot (t - t_{prop}) + \Phi_{H_X}^0 \\
\Phi_{V_X}(t, t_{prop}) &= k_{V_X} \cdot (t - t_{prop}) + \Phi_{V_X}^0
\end{aligned} \tag{4}$$

Here the K 's are the different slopes that introduce the shift to the photons' polarizations, and Φ^0 's are the initial phase of the photons at the time of emission. t_{prop} is the propagation times of the photons from the quantum dot to the device. Since it's constant time we can simplify the function by including it in the constant phase Φ^0 . By taking the starting time of our system as the biexciton excitation time, we can write the state using the t_x and t_{xx} (where t_{xx} and t_x are the random emission times of the biexciton and exciton respectively), as follows:

$$\begin{aligned}
|\Psi(t)\rangle &= \alpha(|H_{XX}\rangle \cdot e^{i\Phi_{H_{XX}}(t_{XX}-t_{start}^x)} \otimes |H_X\rangle \cdot e^{i\Phi_{H_X}(t_X-t_{start}^x)} + \\
&|V_{XX}\rangle \cdot e^{i\Phi_{V_{XX}}(t_{XX}-t_{start}^x)} \otimes |V_X\rangle \cdot e^{i\Phi_{V_X}(t_X-t_{start}^x)} \cdot e^{-i(\Delta E t)/\hbar})
\end{aligned} \tag{5}$$

next, plug it in

$$\Psi(t) = |H_{xx}H_x\rangle + e^{i\Phi} |V_{xx}V_x\rangle \tag{6}$$

where $\Phi(t)$:

$$\begin{aligned}
\Phi(t) &= (K_{V_{XX}} - K_{H_{XX}} + K_{V_X} - K_{H_X}) \cdot t_{xx} + (K_{V_X} - K_{H_X} + \Delta E/\hbar) * t_x + \\
&(K_{V_{XX}} - K_{H_{XX}}) \cdot -t_{start}^{XX} + (K_{V_X} - K_{H_X}) \cdot -t_{start}^X + \\
&(\Phi_{V_{XX}}^0 - \Phi_{H_{XX}}^0 + \Phi_{V_X}^0 - \Phi_{H_X}^0).
\end{aligned} \tag{7}$$

Conditions:

$$K_{V_X} - K_{H_X} = -\Delta E/\hbar \quad (8)$$

$$(K_{V_{XX}} - K_{H_{XX}}) = -(K_{V_X} - K_{H_X}) \quad (9)$$

1.6 Additional proposed advances

1.6.1 Combining the radiative cascade into the knitting machine.

1.6.2 Using the experimental system for feed-forward operations in 1D cluster states.

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Appendices

Appendix I

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Appendix 2

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