## 1 Quantum Shannon unitary decomposition

The algorithm for Quantum Shannon Decomposition (QSD) is a variation of the CSD algorithm presented in cosine-sine unitary decomposition. The change is made in the decomposition of the uniformly controlled operators L and R in the first identity of Figure 1 (explained in cosine-sine unitary decomposition). Being  $U = U_0 \oplus U_1$  (as well as the matrices L and R of CSD), we can find the unitaries Q and W and the diagonal D that satisfy [1]

$$U = (I \otimes Q)(D \oplus D^{\dagger})(I \otimes W). \tag{1}$$

From Equation (1), we have that  $U_0 = QDW$  and  $U_1 = QD^{\dagger}W$ , which implies  $U_0U_1^{\dagger} = QD^2Q^{\dagger}$ . We then have that D and Q can be calculated through diagonalization. The unitary W can be found using the relationships  $W = D^{\dagger}Q^{\dagger}U_0$  or  $W = DQ^{\dagger}U_1$ . Knowing that the diagonal D is composed of the eigenvalues of the diagonalization, the matrix  $D \oplus D^{\dagger}$  can be represented by a uniformly controlled  $R_z$  rotation applied to the most significant qubit (Figure 2) [1].



Figure 1: Identities used for the cosine-sine decomposition (CSD) of quantum operators.

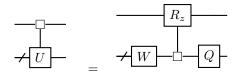


Figure 2: Identity used for quantum Shannon decomposition (QSD) of quantum operators.

Similarly to the CSD algorithm, the QSD applies the identities defined in Figure 1 and Figure 2 recursively. It starts by applying the first identity from Figure 1 to decompose the operator  $U \in \mathbb{C}^{N \times N}$  into a central matrix  $D_{cs}$  and unitary matrices  $L_i$  and  $R_i$ . Next, the identity in Figure 2 is applied to the operators L and R, decomposing them into the diagonal  $D \oplus D^{\dagger}$  and the  $N/2 \times N/2$  dimensional unitaries W and Q (Figure 3). The procedure is restarted from the first identity, acting on the operators W and Q, and continues until these operators have a dimension of  $2 \times 2$  representing one-qubit gates. The ZYZ decomposition used to represent one-qubit gates does not have CNOTs, thus only the uniformly controlled rotations contribute to the total number of CNOTs in the circuit  $\frac{3}{4}4^n - \frac{3}{2}2^n$  [1].

If the recursion terminates when the operators Q and W act on two qubits, there will be  $4^{n-2}$  such operators and  $3 \times 4^{n-2}$  fewer control-rotations by a qubit. Using the decomposition explained in two-qubit unitary decomposition, each two-qubit operator contributes with 3 CNOTs. Each control-rotation acting on two qubits contributes with 2 CNOTs [2]. Therefore, the CNOT count decreases by  $3 \times 4^{n-2}$  and the total is reduced to  $\frac{9}{16}4^n - \frac{3}{2}2^n$ . This variant of the algorithm was called QSD (l=2) by [1], where l indicates the number of qubits where the recursion terminates. The previous case, which terminates at one-qubit operators, is called QSD (l=1).

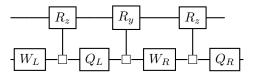


Figure 3: Circuito para a decomposição de Shannon (QSD) de um operador quântico de dois qubits.

The QSD (l=2) can still be optimized in two ways [1]. The first, when the recursion ends with V and W operators acting on two qubits, applies the optimization indicated by Figure 4 and explained in two-qubit unitary decomposition, reducing  $4^{n-2}-1$  CNOTs (one CNOT per two-qubit operator, except for the last one). The second optimization uses CZ (Controlled-Z) gates instead of CNOTs to decompose the central matrices, as the last CZ can be absorbed by the neighboring multiplexer, saving  $(4^{n-2}-1)/3$  CNOTs. With both optimizations, the total number of CNOTs is reduced to  $\frac{23}{48}4^n - \frac{3}{2}2^n + \frac{4}{3}$ . So far, this is the most efficient decomposition for general unitaries. [1] referred to this optimized decomposition as QSD (l=2, optimized).

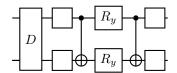


Figure 4: Diagonalization of arbitrary two-qubit unitaries.

Listing 1: Quantum Shannon decomposition (l = 2, optimized).

## References

- [1] V.V. Shende, S.S. Bullock, and I.L. Markov. Synthesis of quantum-logic circuits. <u>IEEE Transactions on Computer-Aided Design of Integrated Circuits and Systems</u>, 25(6):1000–1010, 2006.
- [2] Raban Iten, Roger Colbeck, Ivan Kukuljan, Jonathan Home, and Matthias Christandl. Quantum circuits for isometries. Physical Review A, 93(3):032318, 2016.