

# 今後の将来展望: HΦ

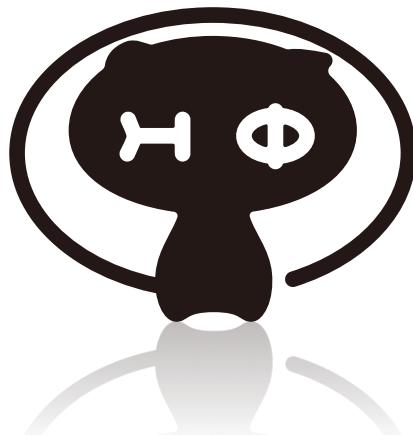
山地 洋平

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1.  $K\omega$ の応用: Finite- $T$  linear response
2. 今後の開発目標 Perspective



Computational  
Science  
Alliance  
The University of Tokyo

# New function will be implemented: Finite- $T$ linear response Combination of TPQ and $K\omega$

Y. Yamaji, T. Suzuki, & M. Kawamura, arXiv:1802.02854.



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# Finite-Temperature Spectra

$$\mathcal{G}_\beta^{AB}(\omega) = \sum_{n,m} \frac{e^{-\beta E_n}}{Z(\beta)} \frac{\langle n | \hat{A}^\dagger | m \rangle \langle m | \hat{B} | n \rangle}{\omega + i\delta + E_n - E_m}$$

$$Z(\beta) = \sum_n e^{-\beta E_n}$$

$$\mathcal{G}_\beta^{AB}(\omega) = \sum_n \frac{e^{-\beta E_n}}{Z(\beta)} \langle n | \hat{A}^\dagger \frac{1}{\omega + i\delta + E_n - \hat{H}} \hat{B} | n \rangle$$

Complexity  $\mathcal{O}(N_H^3)$

Memory  $\mathcal{O}(N_H^2)$

Is it necessary? Answer is No

# Finite-Temperature Spectra by Real-Time Evolution of Wave Functions

- T. litaka and T. Ebisuzaki, Phys. Rev. Lett. 90, 047203 (2003).  
R. Steinigeweg, J. Gemmer, and W. Brenig, Phys. Rev. Lett. 112, 120601 (2014).  
T. Monnai and A. Sugita, J. Phys. Soc. Jpn. 83, 094001 (2014).  
C. Karrasch, D. M. Kennes, and J. E. Moore, Phys. Rev. B 90, 155104 (2014).  
F. Jin, R. Steinigeweg, F. Heidrich-Meisner, K. Michielsen, and H. De Raedt,  
Phys. Rev. B 92, 205103 (2015).

# Finite-Temperature Spectra by Micorocanonical Ensemble

- M. W. Long, P. Prelovsek, S. El Shawish, J. Karadamoglou, and X. Zotos,  
Phys. Rev. B 68, 235106 (2003).  
X. Zotos, Phys. Rev. Lett. 92, 067202 (2004).

# An Intuitive Description of TPQ States and Green's Function at Finite Temperature

A normalized TPQ state

$$|\psi_\beta\rangle \equiv \frac{|\phi_\beta\rangle}{\sqrt{\langle\phi_\beta|\phi_\beta\rangle}} \sim \sum_n e^{i\varphi_n} \frac{e^{-\frac{\beta}{2}E_n}}{\sqrt{Z(\beta)}} |n\rangle$$

Spectral projector  $\hat{P}_n = |n\rangle\langle n|$

Green's function rewritten by using a TPQ state

$$\mathcal{G}_\beta^{AB}(\zeta) \sim \sum_n \langle\psi_\beta|\hat{P}_n \hat{A}^\dagger \frac{1}{\zeta + E_n - \hat{H}} \hat{B} \hat{P}_n |\psi_\beta\rangle$$

# An Alternative to Spectral Projection

T. Kato, Progress of Theoretical Physics 4, 514 (1949).

$$\hat{P}_{\gamma,\rho} = \frac{1}{2\pi i} \oint_{C_{\gamma,\rho}} \frac{dz}{z - \hat{H}} \quad z = \rho e^{i\theta} + \gamma$$

$$|\phi\rangle = \sum_n d_n |n\rangle$$
$$\hat{P}_{\gamma,\rho} |\phi\rangle = \sum_{E_n \in (\gamma-\rho, \gamma+\rho)} d_n |n\rangle$$

Discretized by Riemann sum

T. Sakurai and H. Sugiura,  
J. Comput. Appl. Math. 159, 119 (2003).  
T. Ikegami, T. Sakurai, and U. Nagashima,  
J. Comput. Appl. Math. 233, 1927 (2010).

$$\hat{P}_{\gamma,\rho,M} = \frac{1}{M} \sum_{j=1}^M \frac{\rho e^{i\theta_j}}{\rho e^{i\theta_j} + \gamma - \hat{H}}$$

$$\theta_j = 2\pi(j - 1/2)/M$$

# Shifted Krylov Subspace Method

$$\vec{x} = \frac{1}{\rho e^{i\theta_j} + \gamma - \hat{H}} \vec{b}$$

Liner equations

$$(z\mathbf{1} - H)\vec{x} = \vec{b} \quad \vec{b} \doteq \hat{O}|\psi\rangle$$
$$\Rightarrow G_{\hat{O}}(z) = \vec{b}^\dagger \vec{x} \quad \vec{x} \doteq (z\mathbf{1} - \hat{H})^{-1} \hat{O}|\psi\rangle$$

← Solvable by Shifted Krylov subspace method

A. Frommer (1995, 2003)

T. Sogabe, T. Hoshi, S. L. Zhang, and T. Fujiwara, *A numerical method for calculating the Green's function arising from electronic structure theory*, In Frontiers of Computational Science. pp.189-195, 2007.

# Shifted CG: Algorithm

Initial  $\vec{r}_0 = \vec{b}$ ,  $\alpha_{-1} = 1$ ,  $\rho_{-1} = +\infty$ ,  
 $\pi_0^\sigma = \pi_{-1}^\sigma = 1$ ,  $\vec{p}_{-1}^\sigma = \vec{x}_{-1}^\sigma = \vec{0}$

For  $k = 0, 1, \dots, m$

-Seed equations

$$\rho_k = \vec{r}_k^T \vec{r}_k$$

$$\beta_{k-1} = \frac{\rho_k}{\rho_{k-1}}$$

$$\alpha_k = \frac{\rho_k}{\vec{r}_k^T A \vec{r}_k - \beta_{k-1} \frac{\rho_k}{\alpha_{k-1}}}$$

$$\vec{r}_{k+1} = \left(1 + \frac{\alpha_k \beta_{k-1}}{\alpha_{k-1}}\right) \vec{r}_k - \alpha_k A \vec{r}_k - \frac{\alpha_k \beta_{k-1}}{\alpha_{k-1}} \vec{r}_{k-1}$$

-Shifted equations

$$\pi_{k+1}^\sigma = (1 + \alpha_k \sigma) \pi_k^\sigma - \frac{\alpha_k \beta_{k-1}}{\alpha_{k-1}} (\pi_{k-1}^\sigma - \pi_k^\sigma)$$

$$\vec{p}_k^\sigma = \frac{1}{\pi_k^\sigma} \vec{r}_k + \beta_{k-1} \left( \frac{\pi_{k-1}^\sigma}{\pi_k^\sigma} \right)^2 \vec{p}_{k-1}^\sigma$$

$$\vec{x}_k^\sigma = \vec{x}_{k-1}^\sigma + \frac{\pi_k^\sigma}{\pi_{k+1}^\sigma} \alpha_k \vec{p}_k^\sigma$$

Seed switch

S. Yamamoto, *et al.*,  
J. Phys. Soc. Jpn. 77, 114713 (2008).

Library  $Kw$

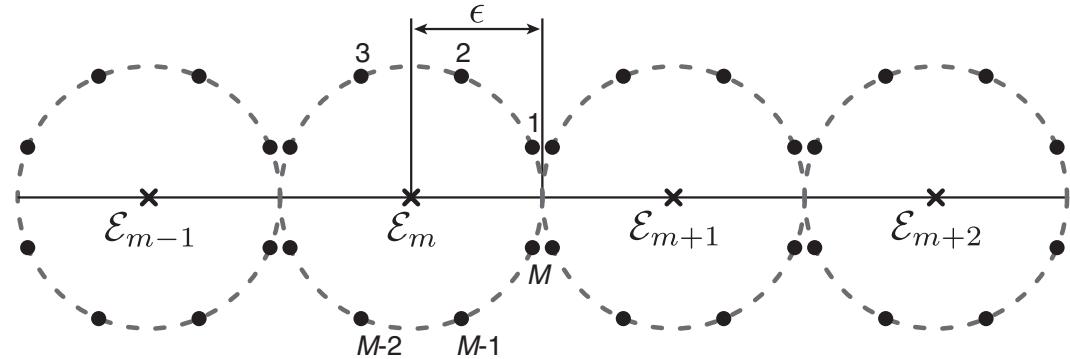
by Dr. Kawamura (ISSP)



# Finite-Temperature Green's Function by Typical Pure States

$$|\psi_{\beta,\delta}^m\rangle = \hat{P}_{\mathcal{E}_m, \epsilon, M} |\psi_\beta\rangle$$

$$\delta = (E_0, \epsilon, M)$$



$$\mathcal{E}_m = E_0 + (2m - 1)\epsilon$$

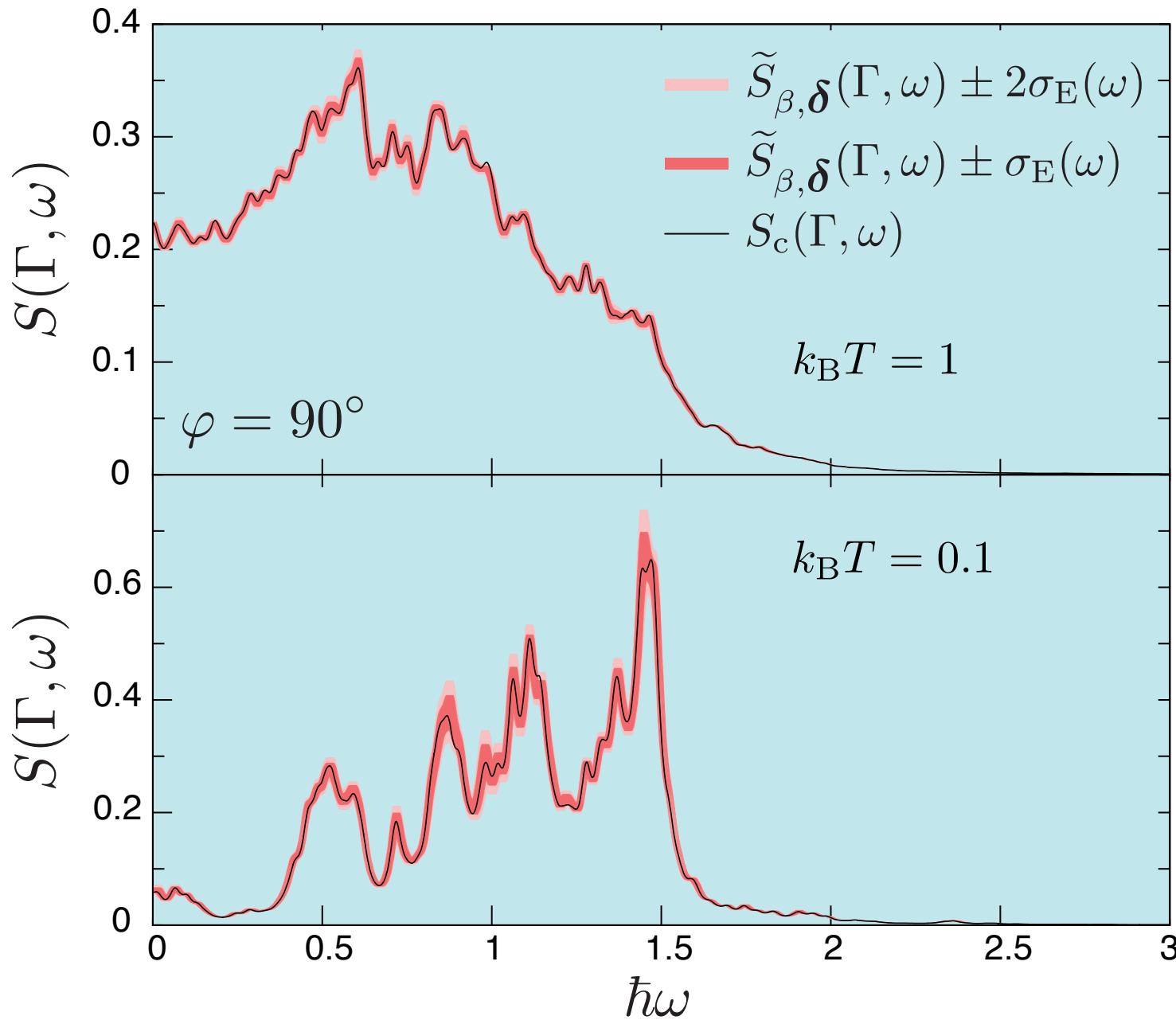
Green's function

$$\tilde{\mathcal{G}}_{\beta,\delta}^{AB}(\zeta) = \sum_{m \geq 0} \langle \psi_{\beta,\delta}^m | \hat{A}^\dagger \frac{1}{\zeta + \mathcal{E}_m - \hat{H}} \hat{B} | \psi_{\beta,\delta}^m \rangle$$

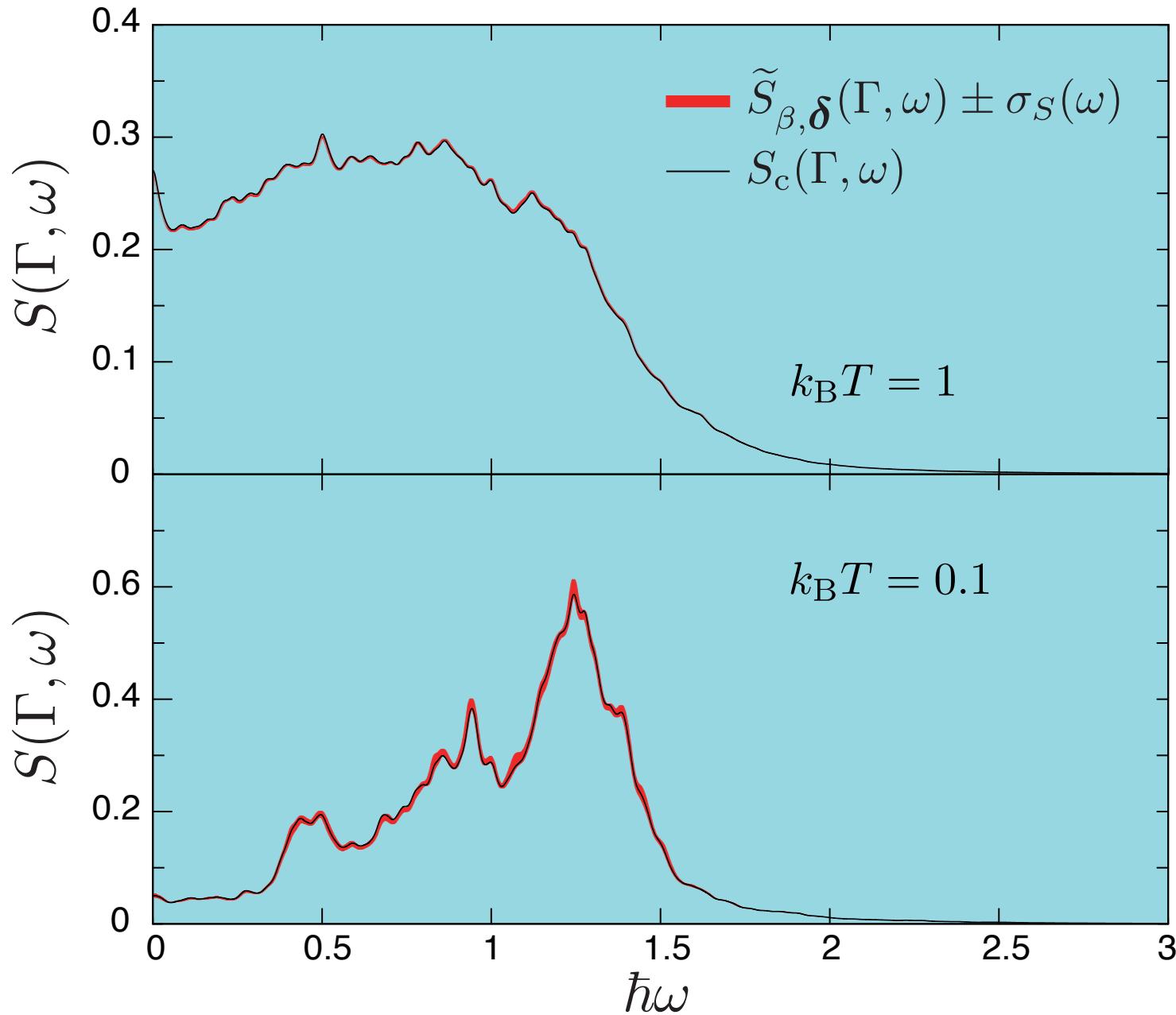
$$\mathcal{G}_{\beta}^{AB}(\zeta) = \lim_{\epsilon \rightarrow +0} \lim_{M \rightarrow +\infty} \mathbb{E} \left[ \tilde{\mathcal{G}}_{\beta,\delta}^{AB}(\zeta) \right]$$

Probability distribution

$$\tilde{P}_{\delta}(\mathcal{E}_m) = \langle \psi_{\beta,\delta}^m | \psi_{\beta,\delta}^m \rangle$$

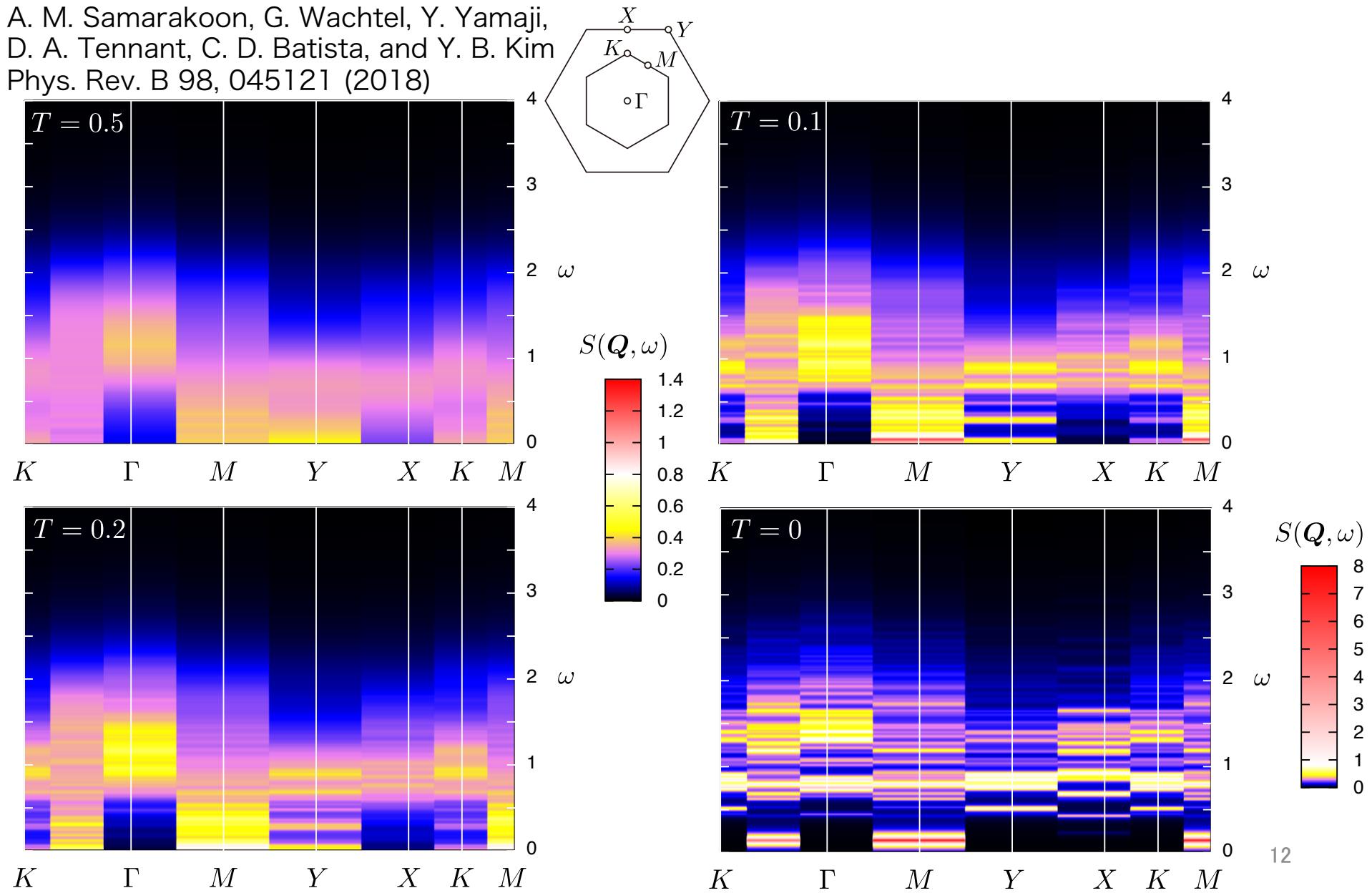


# 18 site AF Kitaev Standard deviation



# Finite- $T$ $S(Q,\omega)$ of a Frustrated Magnets: $\Gamma$ model

A. M. Samarakoon, G. Wachtel, Y. Yamaji,  
D. A. Tennant, C. D. Batista, and Y. B. Kim  
Phys. Rev. B 98, 045121 (2018)



# Future Plan

New functions will be implemented

1. Finite- $T$  linear response:  
Combination of TPQ and  $K\omega$
2.  $N$  spin/body interaction and Green's function
3. Tool for optimizing model parameters to  
fit experimental measurements
  - Example: Find an effective spin Hamiltonian that  
reproduces an observed magnetization process
4. Symmetry
  - Reduction of dimension of Hilbert space
  - Analysis of wave functions