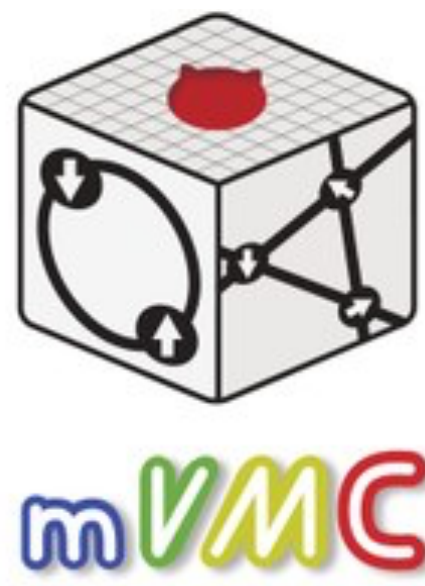


[24aPS-16]多変数変分モンテカルロ法のオープンソフトウェア  
mVMC ver1.0の新機能・利用事例の紹介

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1. What is mVMC ?

A low-energy solver for a wide range of quantum lattice models such as the multi-orbital Hubbard model, Heisenberg model, and Kondo-lattice model by using variational Monte Carlo method. Users can obtain high-accuracy wave functions for ground states of the above models. Users flexibly choose the correlation factors in wavefunctions such as Gutzwiller, Jastrow, and doublon-holon binding factors and optimize more than ten thousand variational parameters. It also allows to obtain low-energy excited states by specifying quantum numbers by using the quantum-number projections.

Open-source software distributed through GitHub (latest version: ver.1.0.2)  
Parallelization: Hybrid (MPI & OpenMP)  
License: GNU GPL version3

2. What you can do with mVMC

2-1. Summary of features

Target :

- Hubbard model (particle # conserved system)
- Heisenberg model (spin 1/2)
- Kondo lattice system (spin 1/2+ itinerant electrons)

Method: Variational wave functions

D. Tahara and M. Imada, JPSJ 77, 114701 (2008)  
T. Misawa and M. Imada, PRB, 90, 115137 (2014)

$|\psi\rangle = \mathcal{P}_G \mathcal{P}_J \mathcal{P}_{d-h}^{(2)} \mathcal{P}_{d-h}^{(4)} \mathcal{L}^S \mathcal{L}^K |\phi_{\text{pair}}\rangle$

Pair-product part

$|\phi_{\text{pair}}\rangle = \left[ \sum_{i,j} f_{ij} c_{i\uparrow}^\dagger c_{j\downarrow}^\dagger \right]^{N/2} |0\rangle$

Generalized BCS wave func.

→ correlated metal,  
AF, CO ordered states ,  
superconducting state

Correlation factors

Gutzwiller-Jastrow  $\mathcal{P}_G \mathcal{P}_J$   
doublon-holon  $\mathcal{P}_{d-h}^{(2)} \mathcal{P}_{d-h}^{(4)}$

Projections

$\mathcal{L}^S$ : Total spin, ex.  $S=0$   
 $\mathcal{L}^K$ : Total momentum, ex.  $K=0$

Update (SR method)

Minimizing  $E_{\vec{\alpha}} = \langle H \rangle_{\vec{\alpha}}$   
 $\vec{\alpha}_{\text{new}} - \vec{\alpha}_{\text{old}} = -X^{-1} \vec{g}$

$g_k = \frac{\partial E_{\vec{\alpha}}}{\partial \alpha_k}$  gradient of energy  
(MC Sampling)

$X_{\alpha\beta} = \langle \bar{\psi}_\alpha | \bar{\psi}_\beta \rangle, |\bar{\psi}_\alpha\rangle = \frac{\partial |\psi\rangle}{\partial \alpha}, |\bar{\psi}\rangle = \frac{|\psi\rangle}{\sqrt{\langle \psi | \psi \rangle}}$

2-2. Flow of calculations

Standard mode

Expert mode

ex. 1D Hubbard model  
Lx= 16  
model = “Hubbard”  
lattice = “chain”  
U = 4  
t = 1  
nelec=16  
2Sz=0

Users can treat general Hamiltonians  
 $\mathcal{H} = \mathcal{H}_T + \mathcal{H}_I,$   
 $\mathcal{H}_T = - \sum_{i,j} \sum_{\sigma_1, \sigma_2} t_{ij\sigma_1\sigma_2} c_{i\sigma_1}^\dagger c_{j\sigma_2},$   
 $\mathcal{H}_I = \sum_{i,j,k,l} \sum_{\sigma_1, \sigma_2, \sigma_3, \sigma_4} I_{ijkl\sigma_1\sigma_2\sigma_3\sigma_4} c_{i\sigma_1}^\dagger c_{j\sigma_2} c_{k\sigma_3}^\dagger c_{l\sigma_4}$

automatically  
generated

users prepare all necessary files

input files  
- Files for Hamiltonian  
- Files for wave functions

optimization by SR method  
[NVMCCalMode=0]

✓ Download

Source codes & manuals @ GitHub

✓ Prerequisite

- C compiler (Intel, Fujitsu, GNU)
- LAPACK or Scalapack
- MPI library

output  
- optimization process  
- optimized wave function

power Lanczos

[NLanczosMode=1,2]

calculating correlation  
functions [NVMCCalMode=1]

output: thermodynamic quantities,  
one/two-body Green’s functions  
 $\langle c_{i\sigma_1}^\dagger c_{j\sigma_2} \rangle \quad \langle c_{i\sigma_1}^\dagger c_{j\sigma_2} c_{k\sigma_3}^\dagger c_{l\sigma_4} \rangle$

1st-step power Lanczos  
 $|\psi_1\rangle = (1 + \alpha_1 H) |\psi\rangle$

fourier tool

output: momentum distribution,  
charge/spin structure factors  
 $n(k), N(q), S(q)$

- mVMC is pre-installed in sekirei (ISSP system B)

3. Basics of mVMC

Variational principle

$\min_{\alpha} E(\alpha) = \min_{\alpha} \frac{\langle \psi(\alpha) | \hat{H} | \psi(\alpha) \rangle}{\langle \psi(\alpha) | \psi(\alpha) \rangle}$

$\alpha$ : variational parameters

Physical properties [MC sampling]  $x$ : real-space configurations

$\frac{\langle \psi | \hat{A} | \psi \rangle}{\langle \psi | \psi \rangle} = \sum_x \frac{\langle \psi | \hat{A} | x \rangle \langle x | \psi \rangle}{\langle \psi | \psi \rangle} = \sum_x \rho(x) \frac{\langle \psi | \hat{A} | x \rangle}{\langle \psi | x \rangle}$

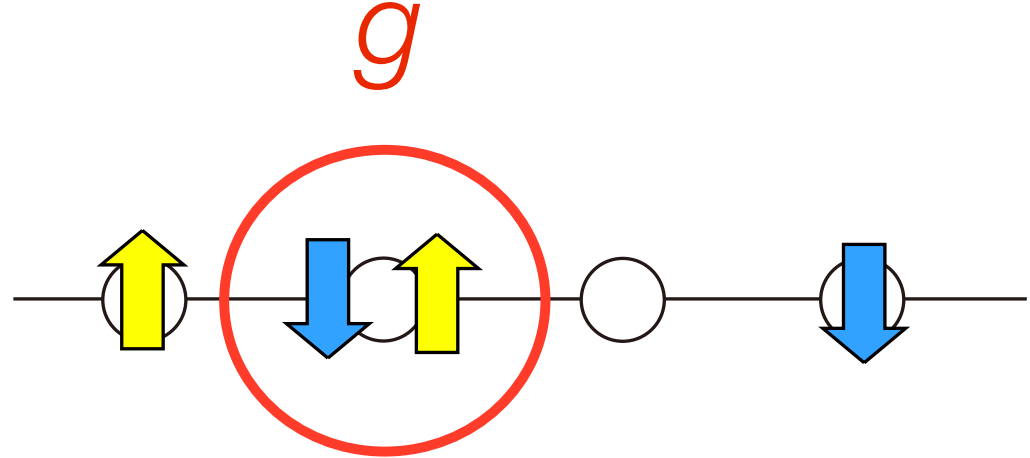
$\sim \frac{1}{N_{\text{MC}}} \sum_{\text{MC sampling}} \frac{\langle \psi | \hat{A} | x \rangle}{\langle \psi | x \rangle}$

Non-negative weight

$\rho(x) = \frac{|\langle \psi | x \rangle|^2}{\langle \psi | \psi \rangle} > 0$

Correlation factors

Ex. Gutzwiller factors  $\mathcal{P}_G = e^{-g \sum_i n_{i\uparrow} n_{i\downarrow}}$



$\mathcal{P}_G |x\rangle = |x\rangle e^{-g D(x)}$

$\langle \psi | x \rangle = \langle \phi_0 | x \rangle \mathcal{P}_G(x)$

determinant or Pfaffian

Conventional VMC: [ # of parameters~10]

Ex. Antiferromagnetic phase

Only AF order parameter  
is the variational parameter

$|\phi_{\text{AF}}\rangle = \prod_{|\mathbf{k}| < k_F, \sigma} a_{\mathbf{k}\sigma}^\dagger |0\rangle$   
 $a_{\mathbf{k}\sigma}^\dagger = u_{\mathbf{k}} c_{\mathbf{k},\sigma}^\dagger + \sigma v_{\mathbf{k}} c_{\mathbf{k}+\mathbf{Q},\sigma}^\dagger$   
 $u_{\mathbf{k}}^2 = \frac{1}{2} \left( 1 - \frac{E_{\mathbf{k}}}{\sqrt{E_{\mathbf{k}}^2 + \Delta_{\text{AF}}^2}} \right)$

$\Delta_{\text{AF}}$ :variational  
parameters

many-variable VMC (mVMC): Large-degrees of freedom in  
pair-product part [# of parameters >=10000]

$|\phi_{\text{Pf}}\rangle = \sum_{i,j,\sigma,\tau} \left( F_{ij\sigma\tau} c_{i\sigma}^\dagger c_{j\tau}^\dagger \right)^{N_e/2} |0\rangle$

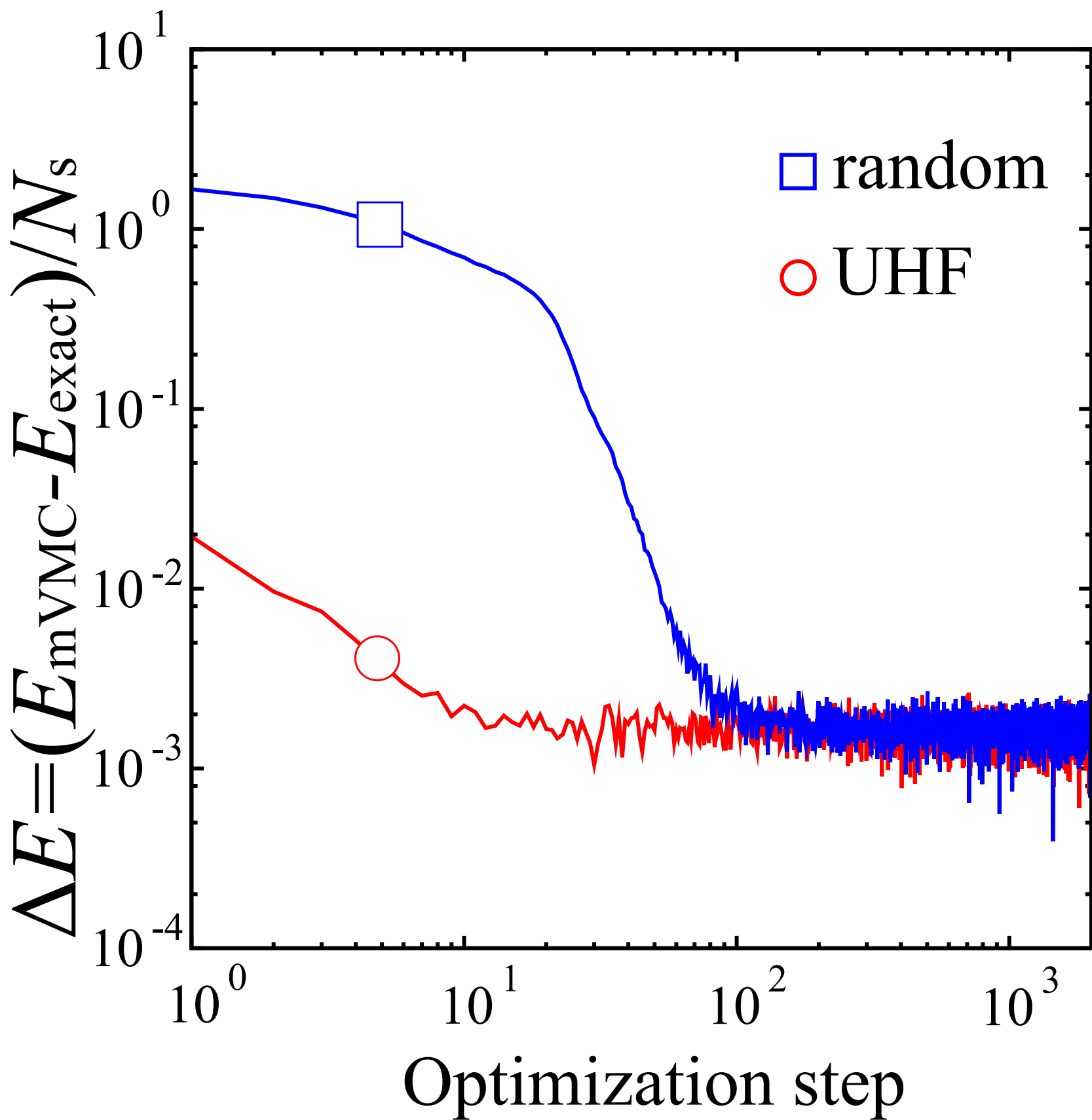
$F_{ij\sigma\tau}$ :variational parameters

$F_{ij\sigma\tau}$  [i,j real-space indices ]→ AF order, superconducting  
phases, spin-orbit couplings, non-collinear magnetism

4. Example Hubbard model on the square lattice

$H = -t \sum_{i,j,\sigma} (c_{i\sigma}^\dagger c_{j\sigma} + H.c.) + U \sum_i n_{i\uparrow} n_{i\downarrow}$

Optimization process ( $U/t=4$ , half filling)



Exact diagonalization  
by using HPhi



Spin correlations

$S(q) = \frac{1}{3N_s} \sum_{i,j} \langle \mathbf{S}_i \cdot \mathbf{S}_j \rangle e^{iq \cdot (\mathbf{r}_i - \mathbf{r}_j)}$   
 $S_{\text{nn}} = \frac{1}{4N_s} \sum_{i,\mu} \langle \mathbf{S}_{\mathbf{r}_i} \cdot \mathbf{S}_{\mathbf{r}_i + \mathbf{e}_\mu} \rangle$

Comparison with exact diagonalization

|             | $E/N_s$     | doublon    | $S_{\text{nn}}$ | $S(\pi, \pi)$ |
|-------------|-------------|------------|-----------------|---------------|
| ED          | -0.85136    | 0.11512    | -0.2063         | 0.05699       |
| mVMC(2 × 2) | -0.84982(4) | 0.11529(5) | -0.2062(1)      | 0.05773(4)    |
| mVMC(4 × 4) | -0.85068(4) | 0.1153(4)  | -0.2062(5)      | 0.0573(2)     |

mVMC well reproduces exact results !

5. Summary

- Simple & Flexible user interfaces
- *Very easy* to study standard models
- *Easy* to study general Hamiltonians

Examples and lecture notes of mVMC are available at  
mVMC-tutorial @ GitHub

