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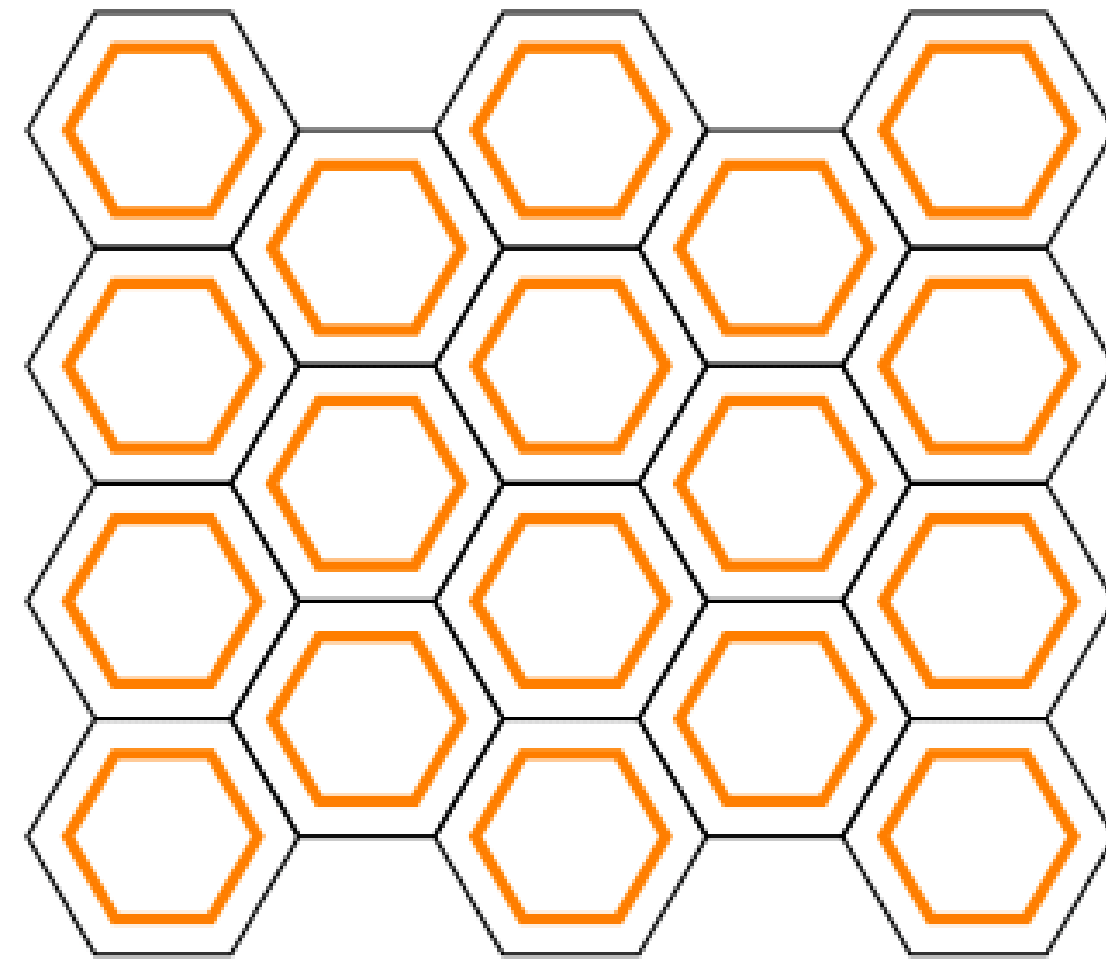
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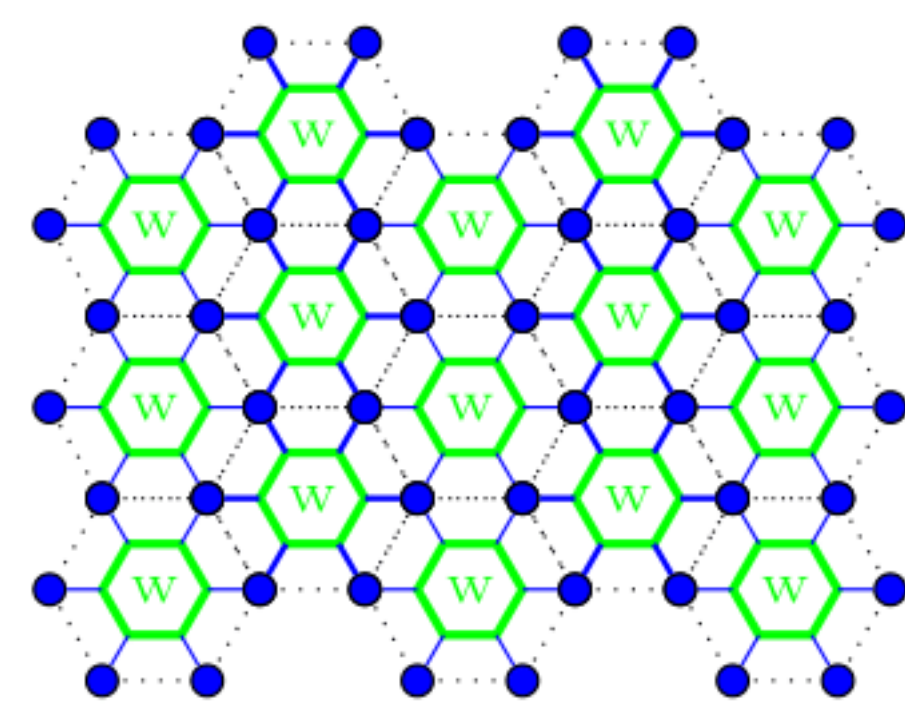
Honeycomb Lattice Proposed Wavefunction



$$|\psi\rangle = \prod_{\mathbf{O}} \left(\sum_{i \in \mathcal{O}} b_i^\dagger \right) |0\rangle$$

PEPS Construction of Honeycomb F.B.I.

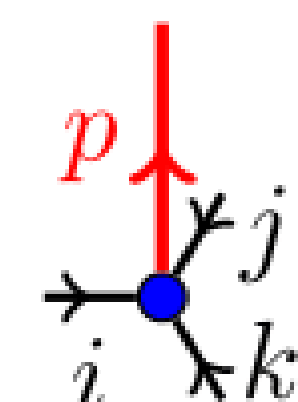
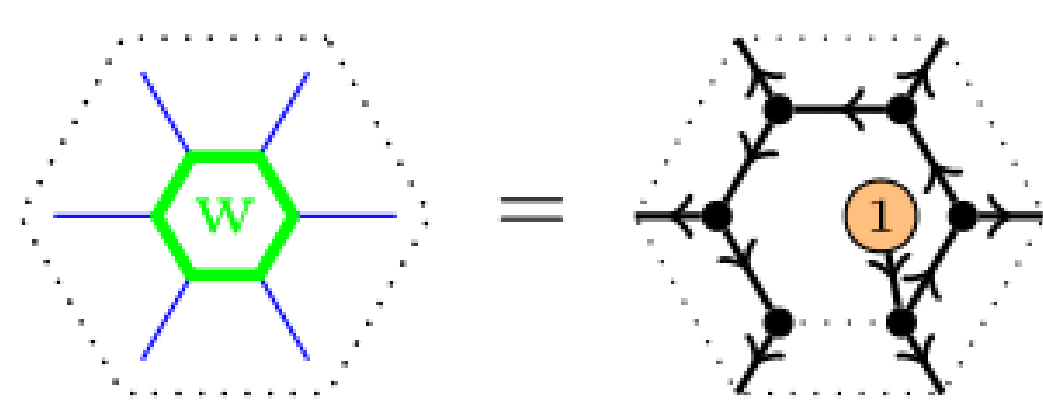
A wavefunction written as a product of local operators acting on a product state can be turned into a tensor network.



Virtual W-state on each plaquette used to synchronize the creation operators in the sum $\sum_{i \in \mathcal{O}} b_i^\dagger$

$$|W\rangle = |000001\rangle + |000010\rangle + |000100\rangle + |001000\rangle + |010000\rangle + |100000\rangle$$

PEPS Construction of Honeycomb FBI



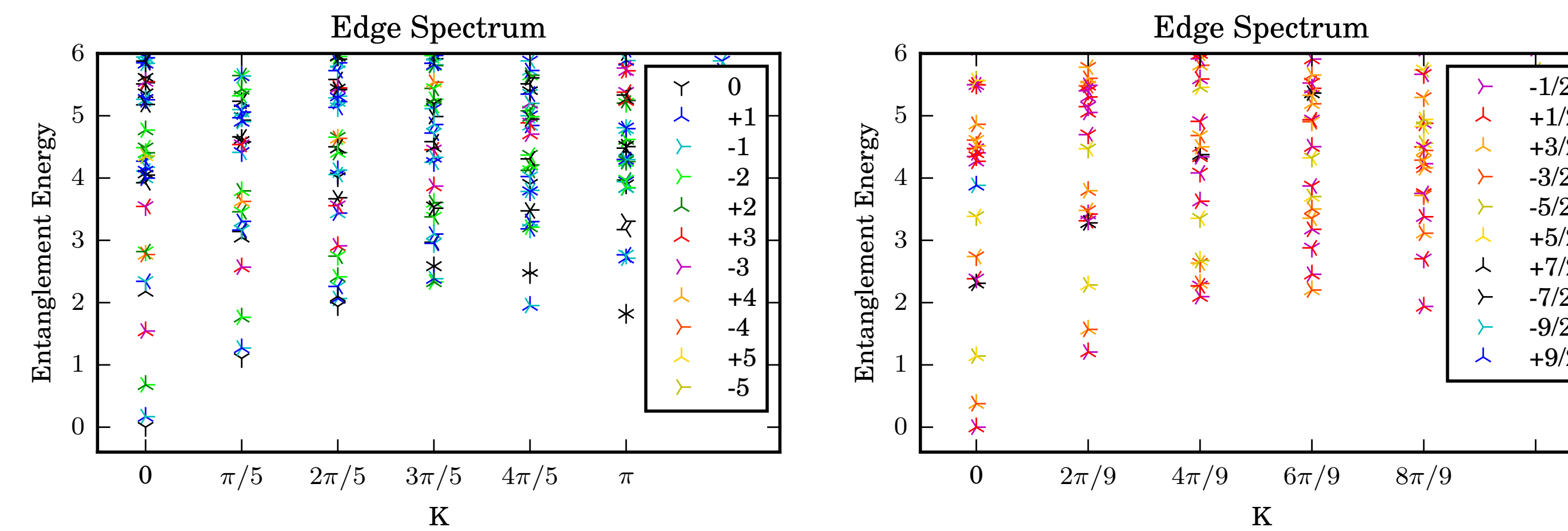
$$|W\rangle = |100\dots\rangle + \dots$$

- W-State can be factored and put on hexagon sites
- Each black directed string has either charge 0 or 1
- Charge conserved

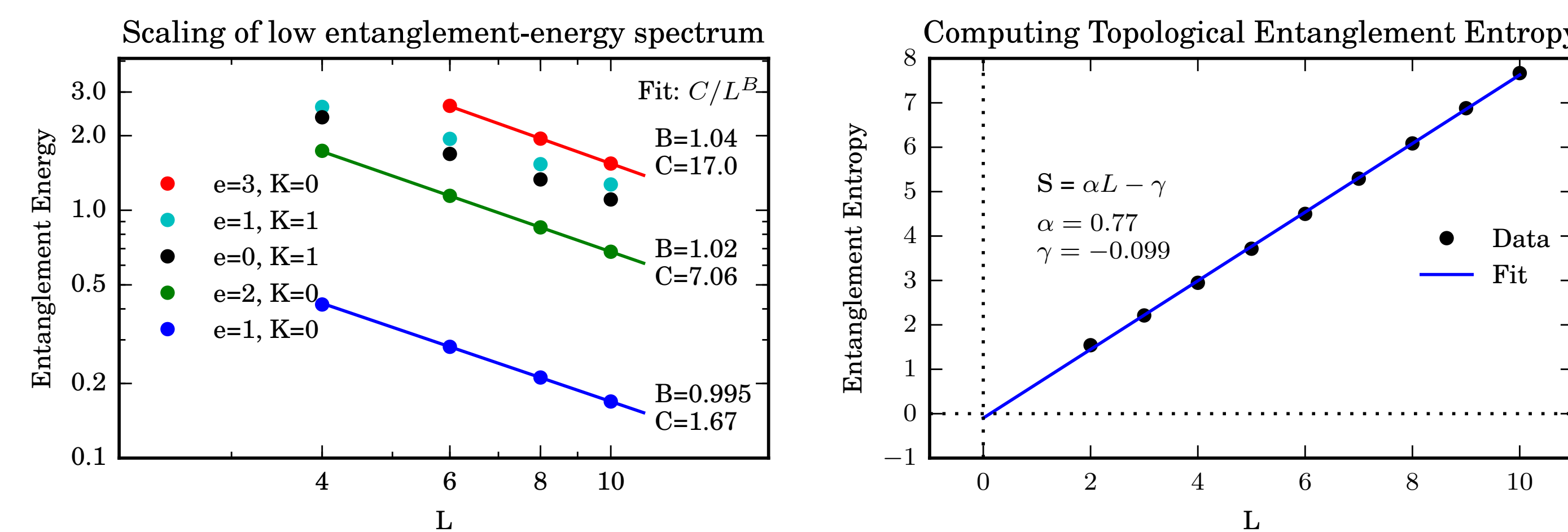
- 'Projects' the three virtual qubits coming into each vertex on to a state in physical Hilbert space
- Physical state is either $|0\rangle, |1\rangle, |2\rangle, |3\rangle$
- Charge conserved

Featureless Bosonic Insulators

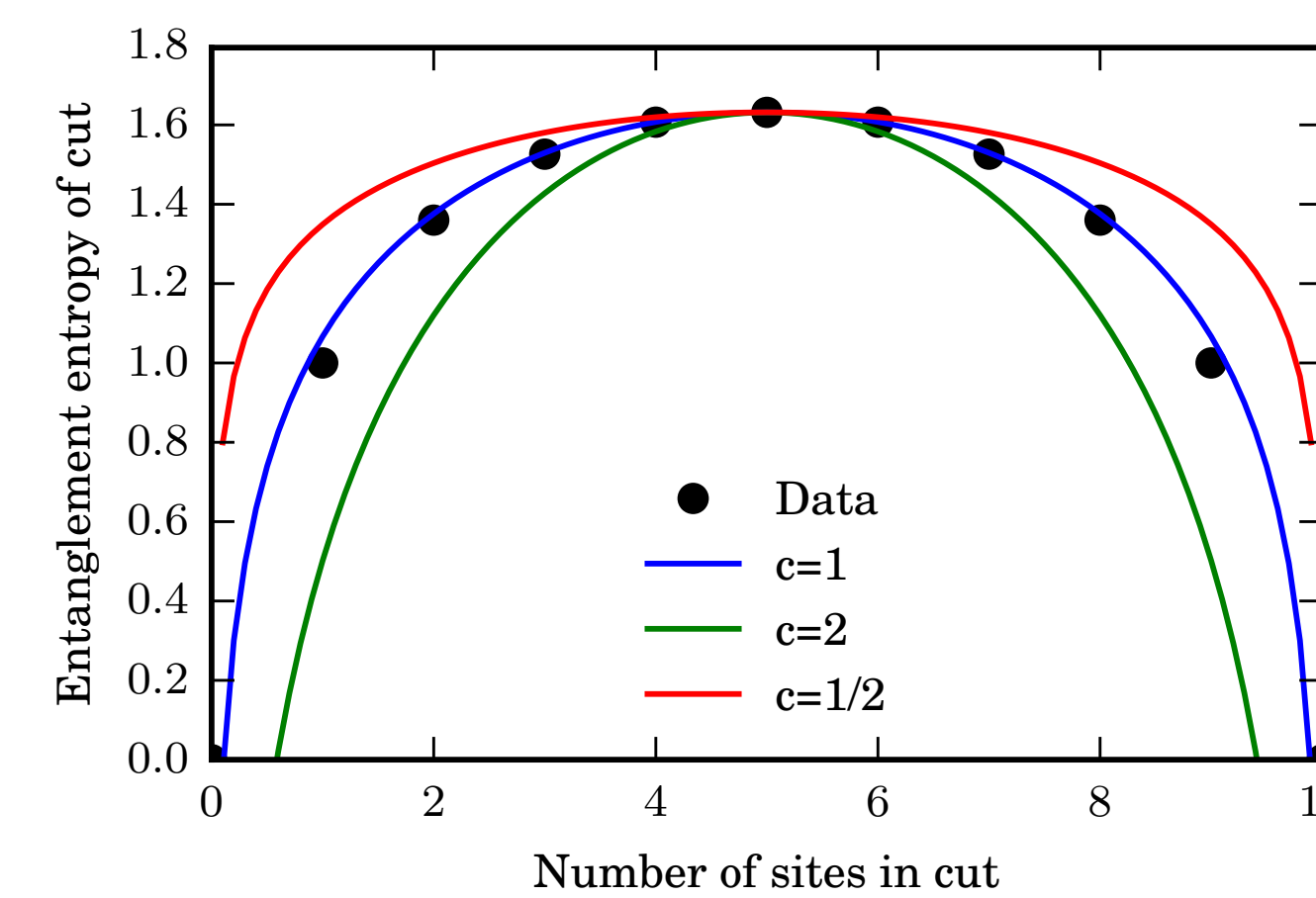
Entanglement Spectrum



Finite Size Analysis of Entanglement Spectra



Conformal Charge



$$c = 1$$

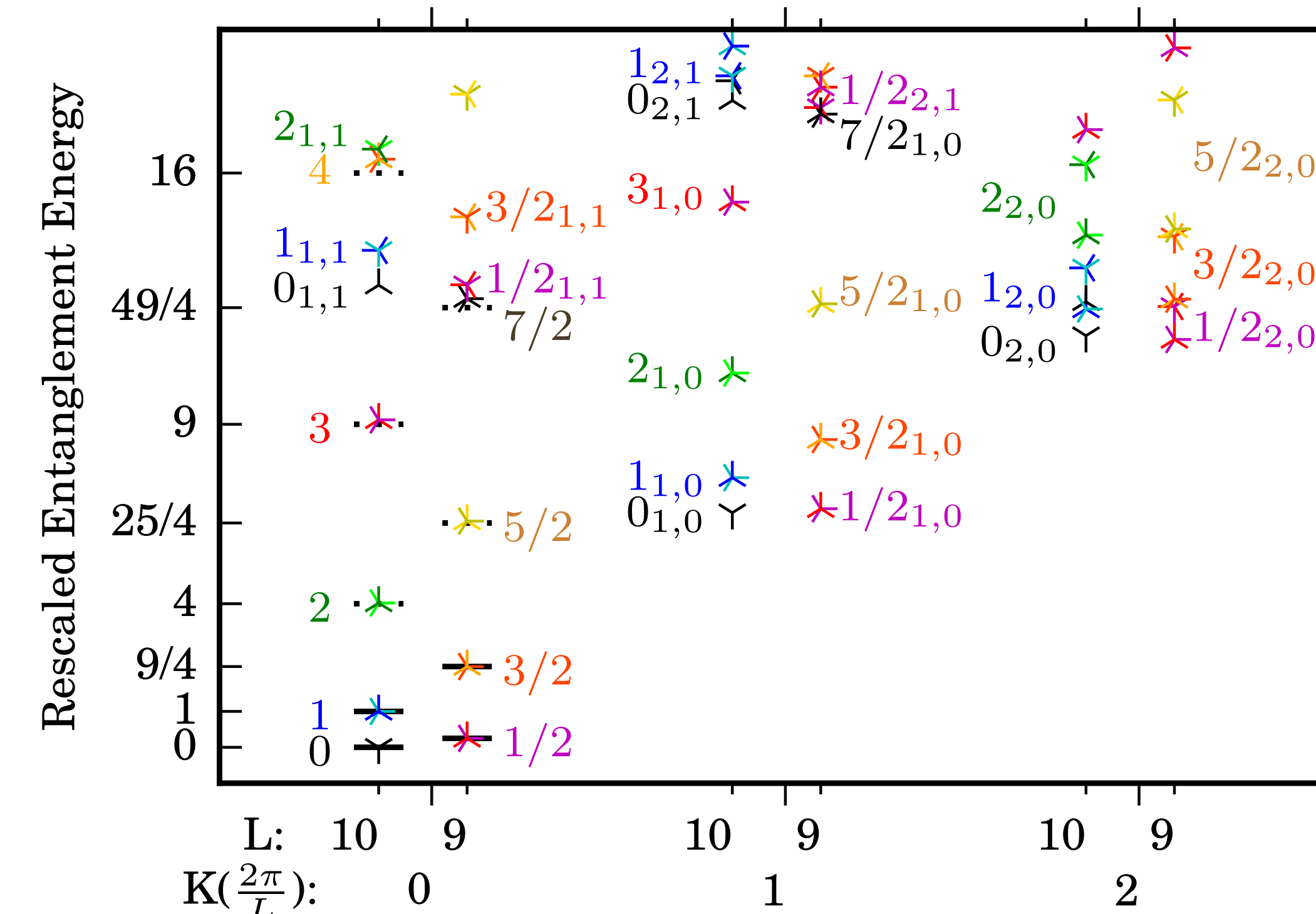
Conformal Weights

We can match the rescaled entanglement energies to the conformal weights of a free bosonic CFT.

$$\begin{aligned} \mathbf{P} &= \frac{2\pi}{L}(\mathbf{L}_0 - \bar{\mathbf{L}}_0) = \frac{2\pi}{L}(em + n - \bar{n}) \\ \mathbf{H} &= \frac{2\pi}{L}(\mathbf{L}_0 + \bar{\mathbf{L}}_0) = \frac{2\pi}{L}\left(\frac{\kappa e^2}{2} + \frac{m^2}{2\kappa} + \frac{n + \bar{n}}{2}\right) \end{aligned}$$

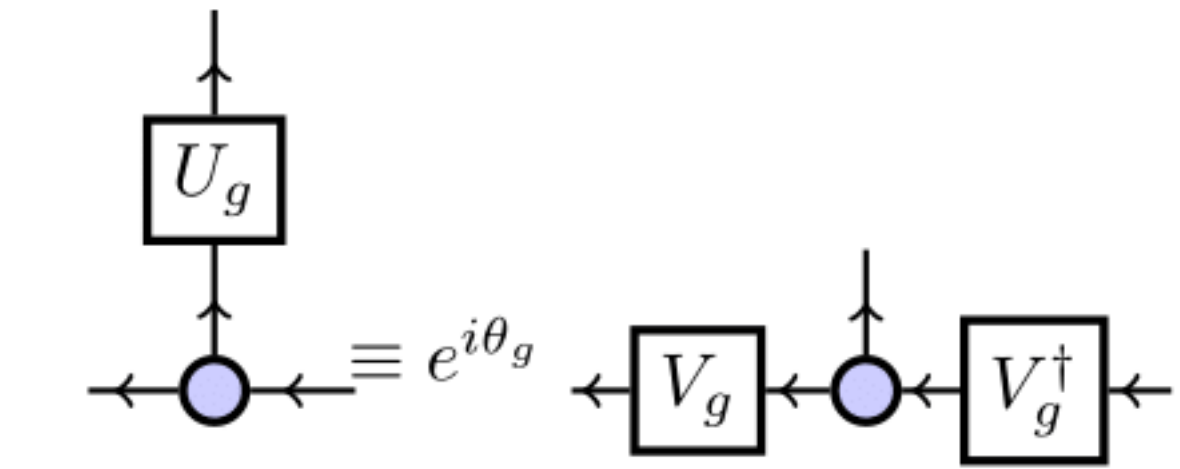
$$\mathbf{H} \propto e^2 + \frac{m^2}{\kappa^2} + \frac{1}{\kappa}(n + \bar{n})$$

CFT Identification of Gapless Entanglement Edge

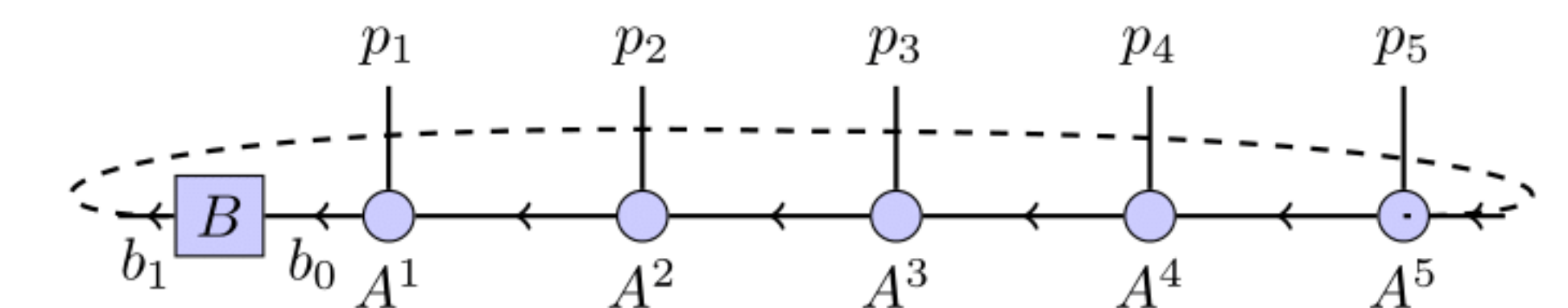


Detecting 1D SPT Order

If U_g is a global symmetry and $|\psi\rangle$ is translationally invariant, then the MPS representation satisfies:



Boundary conditions on a MPS can be represented by a matrix B which acts like:



With PBC ($B = I$), the group action leaves the state invariant. With OBC ($B = |i\rangle\langle i|$), the group action rotates between states that differ only near the boundary; these edge states transform as $V_g \otimes V_g^\dagger$. V_g represents the group projectively. Equivalence classes of projective representations (enumerated by $H^2(G; U(1))$) classify 1D SPT phases.

Symmetry Protection of the Honeycomb FBI

For the state on a cylinder with odd circumference, and the zig-zag entanglement cut defined in the upper left picture, we have the following:

G	U_g	θ_g	V_g	$V_g V_g^*$
$U(1)$				
π				
\mathcal{I}				
$\pi\mathcal{I}$				

Since

$$V_{\pi\mathcal{I}} V_{\pi\mathcal{I}}^* = -I \quad \text{or} \quad V_{\pi} V_{\mathcal{I}} = -V_{\mathcal{I}} V_{\pi},$$

the representation is in the nontrivial class of

$$H^2(\mathbb{Z}_2 \times \mathbb{Z}_2^{\mathcal{I}}; U(1)) = \mathbb{Z}_2.$$

Relation to known 1D physics

- Haldane insulator
 - Unitarily equivalent to the AKLT state
 - Distinct phase under $\mathbb{Z}_2 \times \mathbb{Z}_2^{\mathcal{I}}$
 - Can be connected adiabatically to $L = 1$ cylinder FBI
- Two dimensional classification is $H^3(\mathbb{Z}_2 \times \mathbb{Z}_2^{\mathcal{I}}; U(1)) = \mathbb{Z}_2^4$

References

1. Kimchi, S. A. Parameswaran, A. M. Turner, F. Wang, and A. Vishwanath, Featureless and non-fractionalized mott insulators on the honeycomb lattice at 1/2 site filling,(2012), arXiv:1207.0498 [cond-mat.str-el].
2. S. A. Parameswaran, A. M. Turner, D. P. Arovas, and A. Vishwanath, Topological order and absence of band insulators at integer filling in Non-Symmorphic crystals, (2012), arXiv:1212.0557 [cond-mat.str-el]