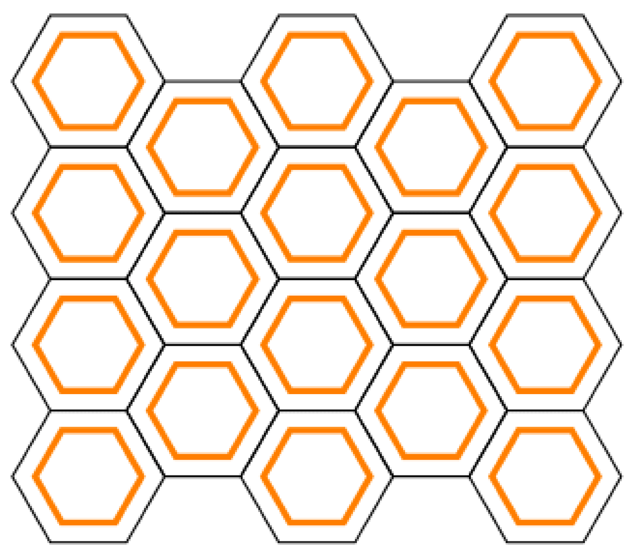
Featureless Bosonic Insulators

Brayden Ware — UCSB
Itamar Kimchi — UCB
Siddarth Parameswaran — UCI
Bela Bauer — Station Q

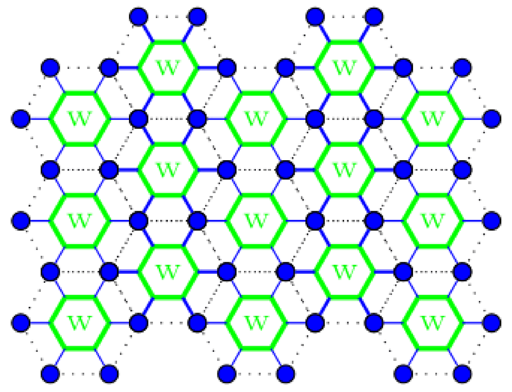
Honeycomb Lattice Proposed Wavefunction



$$|\psi\rangle = \prod_{\hexagon}\left(\sum_{i\in\hexagon} b_i^\dagger\right)|0\rangle$$

PEPS Construction of Honeycomb F.B.I.

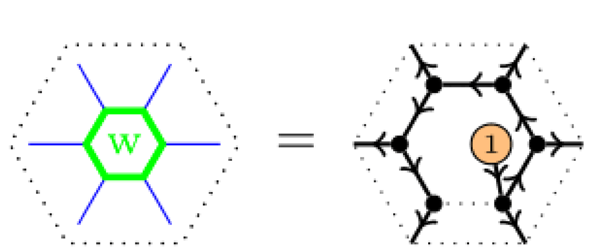
A wavefunction written as a product of local operators acting on a product state can be turned into a tensor network.



Virtual W-state on each plaquette used to synchronize the creation operators in the sum $\sum_{i\in\hexagon} b_i^\dagger$

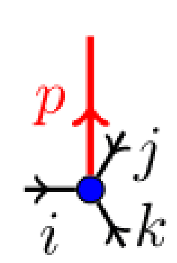
$$|W) = |000001) + |000010) + |000100) + |001000) + |010000) + |100000)$$

PEPS Construction of Honeycomb FBI



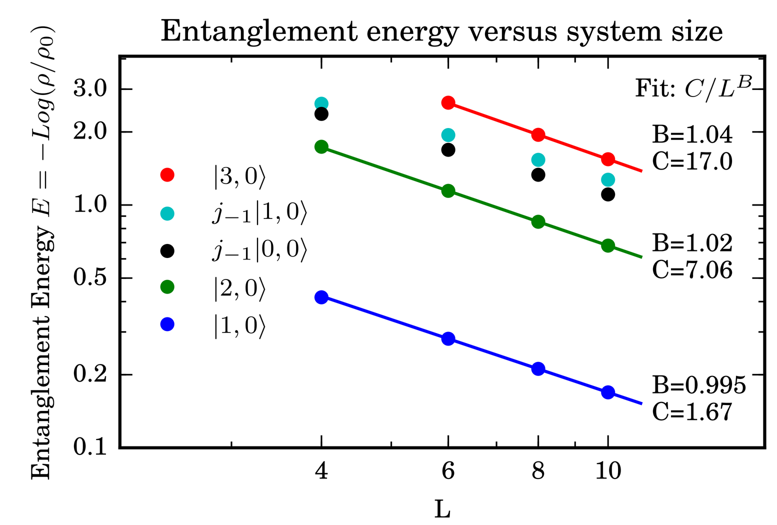
$$|W) = |100\ldots) + \ldots$$

- W-State can be factored and put on hexagon sites
- Each black directed string has either charge 0 or 1
- Charge conserved



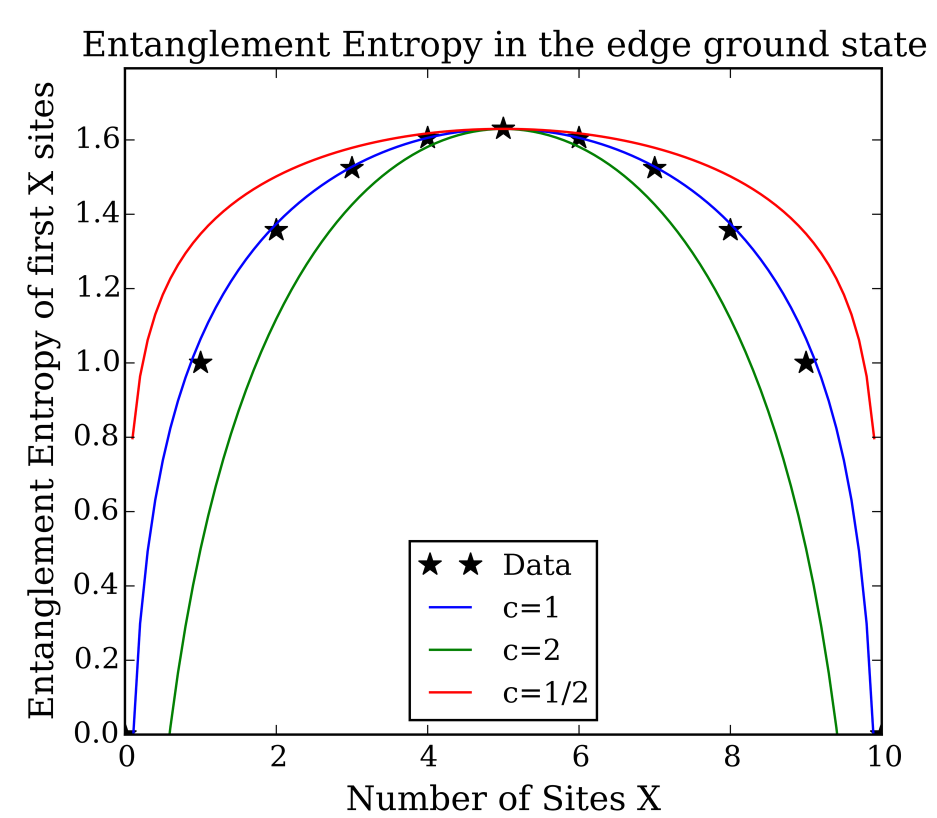
- 'Projects' the three virtual qubits coming into each vertex on to a state in physical Hilbert space
- Physical state is either $|0\rangle, |1\rangle, |2\rangle, |3\rangle$
- Charge conserved

Finite Size Analysis of Entanglement Spectra



- Fix this to show topological entanglement entropy is 0
- Low energy modes show gapless $1/L$ behavior

Conformal Charge



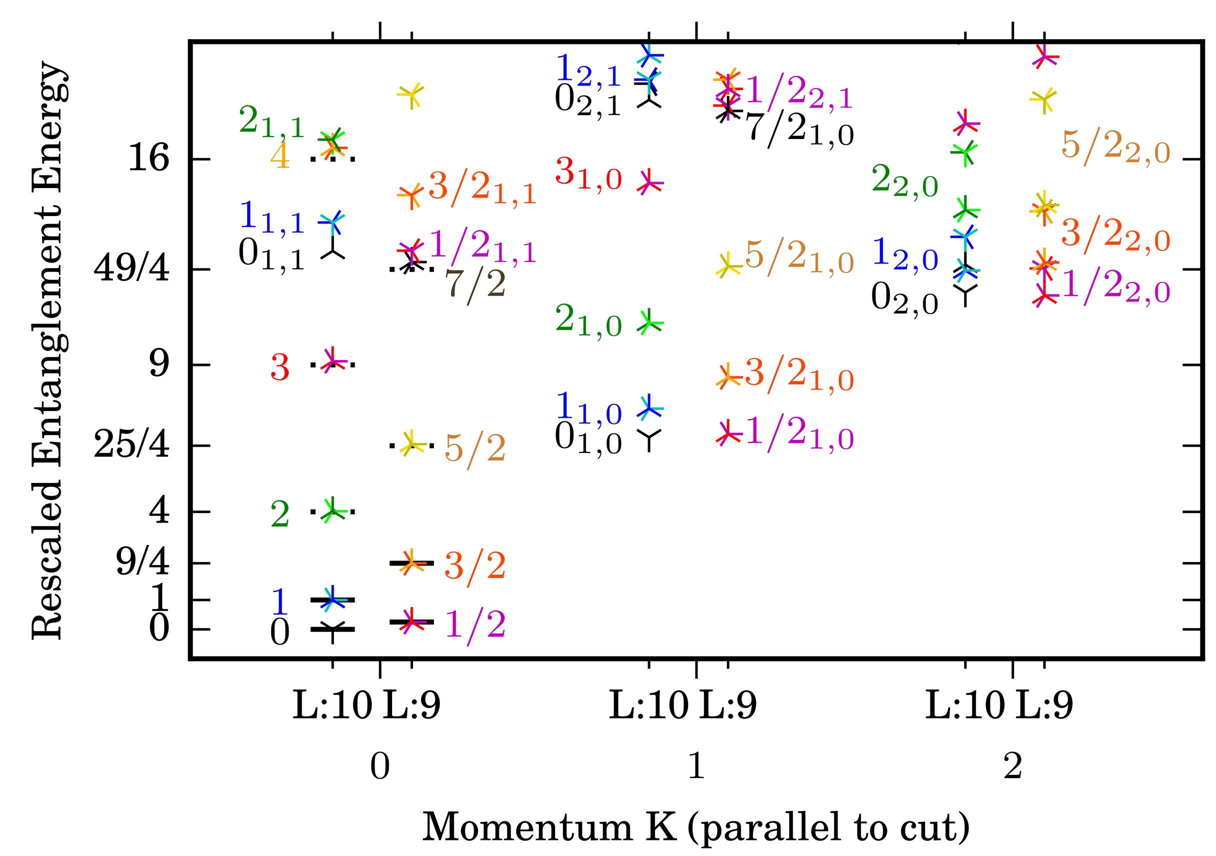
$\boxed{c = 1}$

Conformal Weights

We can match the rescaled entanglement energies to the conformal weights of a free bosonic CFT.

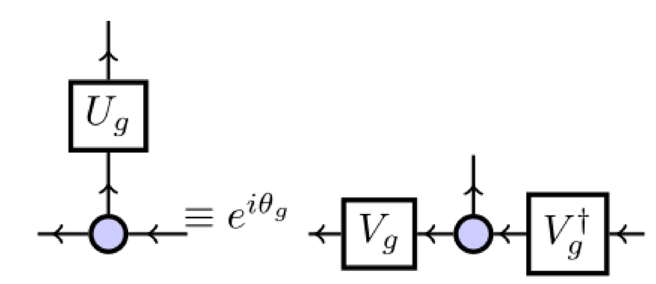
$$\begin{aligned}
\mathbf{P} &= \frac{2\pi}{L}(\mathbf{L}_0 - \bar{\mathbf{L}}_0) \\
&= \frac{2\pi}{L}(em + n - \bar{n}) \\
\widetilde{\mathbf{P}} &= em + n - \bar{n} \\
\mathbf{H} &= \frac{2\pi}{L}(\mathbf{L}_0 + \bar{\mathbf{L}}_0) \\
&= \frac{2\pi}{L}\left(\frac{\kappa e^2}{2} + \frac{m^2}{2\kappa} + \frac{n+\bar{n}}{2}\right) \\
\widetilde{\mathbf{H}} &= \frac{L}{2\pi\kappa}\mathbf{H} \\
&= e^2 + \frac{m^2}{4\kappa^2} + \frac{1}{\kappa}(n + \bar{n})
\end{aligned}$$

CFT Identification of Gapless Entanglement Edge

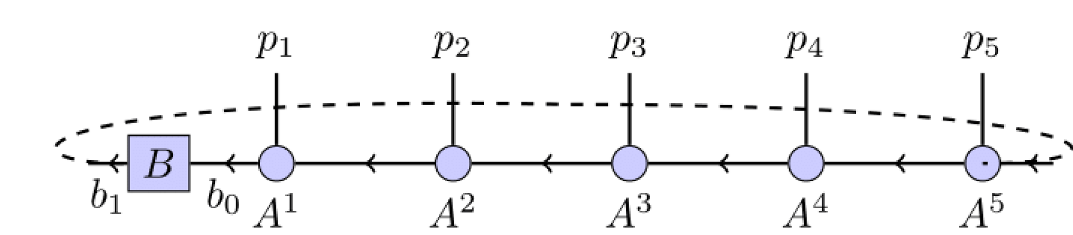


Detecting 1D SPT Order

If U_g is a global symmetry and $|\psi\rangle$ is translationally invariant, then the MPS representation satisfies:



Boundary conditions on a MPS can be represented by a matrix B which acts like:



With PBC ($B = I$), the group action leaves the state invariant.
With OBC ($B = |i\rangle\langle i|$), the group action rotates between states that differ only near the boundary; these edge states transform as $V_g \otimes V_g^\dagger$. V_g represents the group projectively. Equivalence classes of projective representations (enumerated by $H^2(G; U(1))$) classify 1D SPT phases.

Symmetry Protection of the Honeycomb FBI

For the state on a cylinder with odd circumference, and the zig-zag entanglement cut defined in the upper left picture, we have the following:

$\mathbf{G}$	$\mathbf{U_g}$	$\theta_\mathbf{g}$	$\mathbf{V_g}$	$\mathbf{V_g V_g^*}$
$U(1)$				
π				
$\mathcal{I}$				
$\pi\mathcal{I}$				

Since

$$V_{\pi\mathcal{I}}V_{\pi\mathcal{I}}^* = -I \quad \text{or} \quad V_\pi V_\mathcal{I} = -V_\mathcal{I} V_\pi,$$

the representation is in the nontrivial class of

$$\boxed{H^2(\mathbb{Z}_2 \times \mathbb{Z}_2^\mathcal{I}; U(1)) = \mathbb{Z}_2.}$$

Relation to known 1D physics

- Haldane insulator
  - Unitarily equivalent to the AKLT state
  - Distinct phase under $\mathbb{Z}_2 \times \mathbb{Z}_2^\mathcal{I}$
  - Can be connected adiabatically to $L = 1$ cylinder FBI
- Two dimensional classification is $H^3(\mathbb{Z}_2 \times \mathbb{Z}_2^\mathcal{I}; U(1)) = \mathbb{Z}_2^4$

References

- D. W. Kribs, R. Laflamme, D. Poulin, M. Lesosky, Quantum Inf. & Comp. **6** (2006), 383-399.
- P. Zanardi, M. Rasetti, Phys. Rev. Lett. **79**, 3306 (1997).