Featureless Bosonic Insulators

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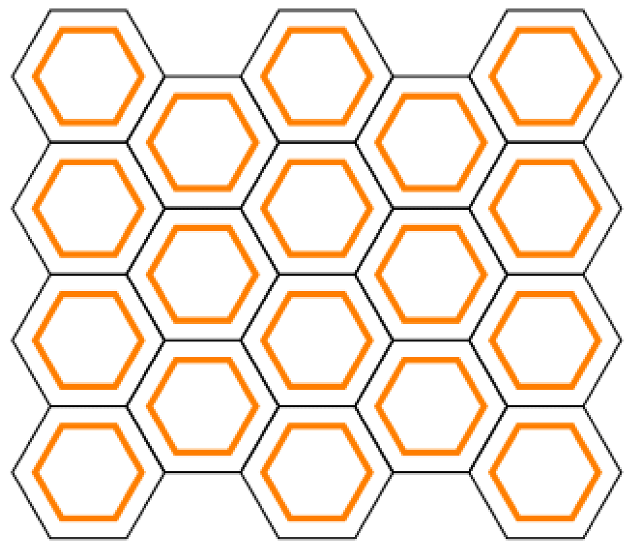
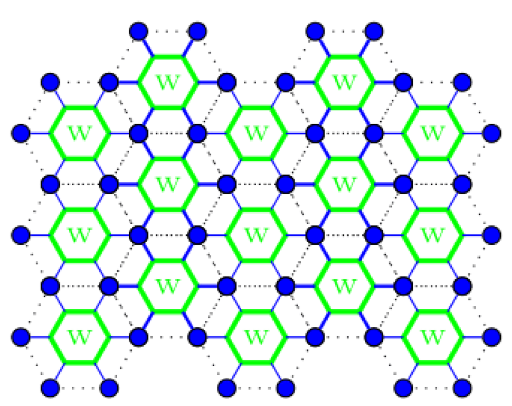
Honeycomb Lattice Proposed Wavefunction



$$|\psi\rangle = \prod_{\hexagon}\left(\sum_{i\in\hexagon} b_i^\dagger\right)|0\rangle$$

PEPS Construction of Honeycomb F.B.I.

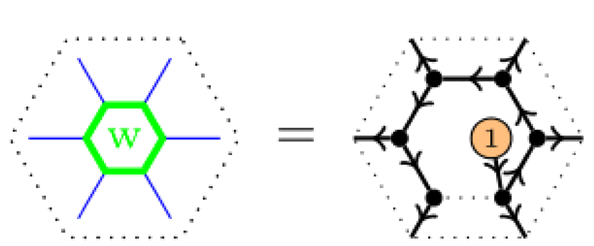
A wavefunction written as a product of local operators acting on a product state can be turned into a tensor network.

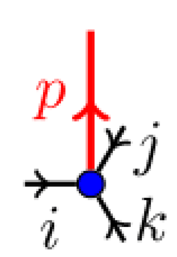


Virtual W-state on each plaquette used to synchronize the creation operators in the sum $\sum_{i\in\hexagon} b_i^\dagger$

$|W) = |000001) + |000010) + |000100) + |001000) + |010000) + |100000)$

PEPS Construction of Honeycomb FBI



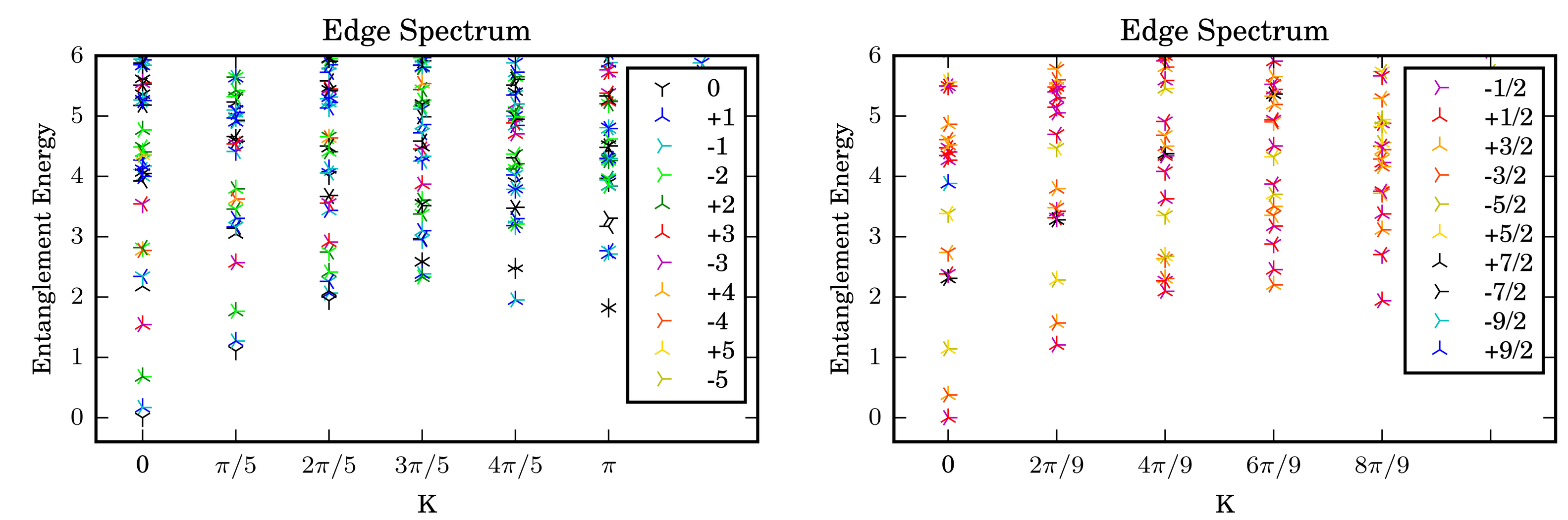


$|W) = |100\ldots) + \ldots$

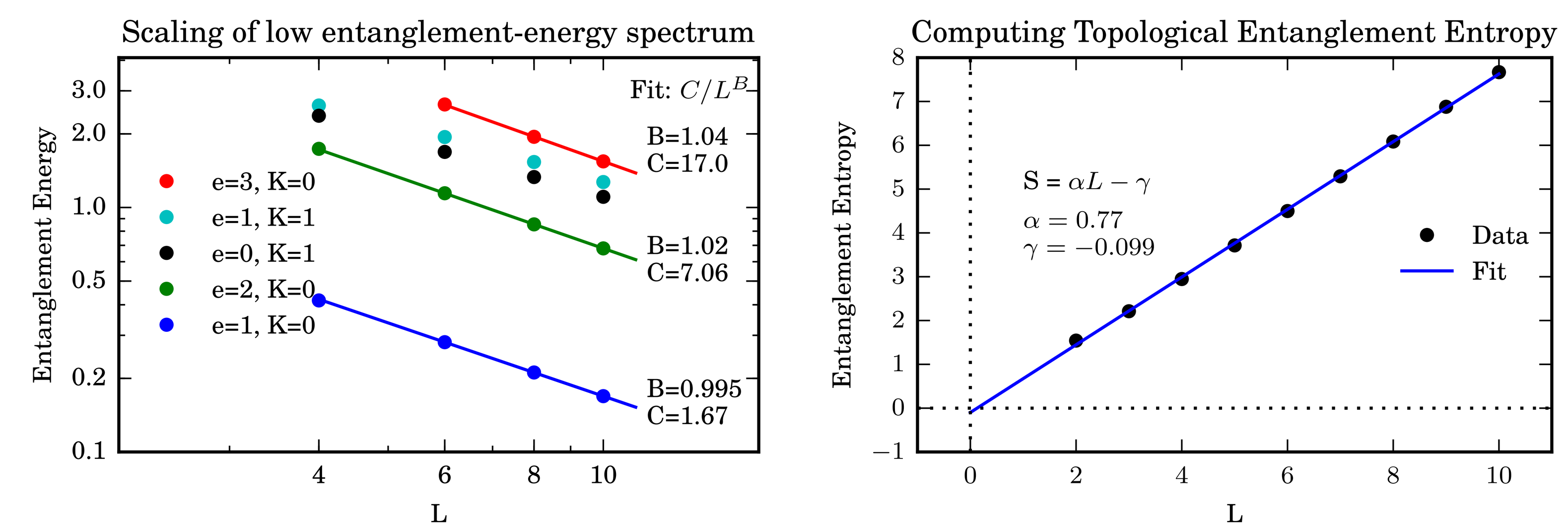
- W-State can be factored and put on hexagon sites
- Each black directed string has either charge 0 or 1
- Charge conserved

- 'Projects' the three virtual qubits coming into each vertex on to a state in physical Hilbert space
- Physical state is either $|0\rangle, |1\rangle, |2\rangle, |3\rangle$
- Charge conserved

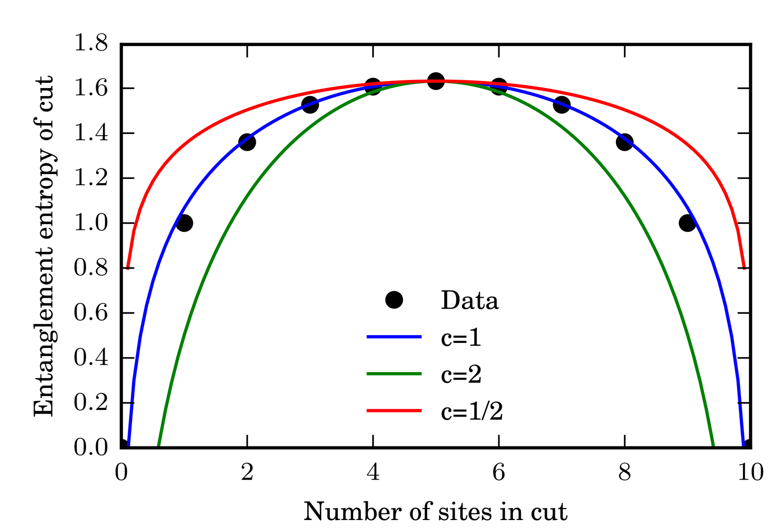
Entanglement Spectrum



Finite Size Analysis of Entanglement Spectra



Conformal Charge



$\boxed{c = 1}$

Conformal Weights

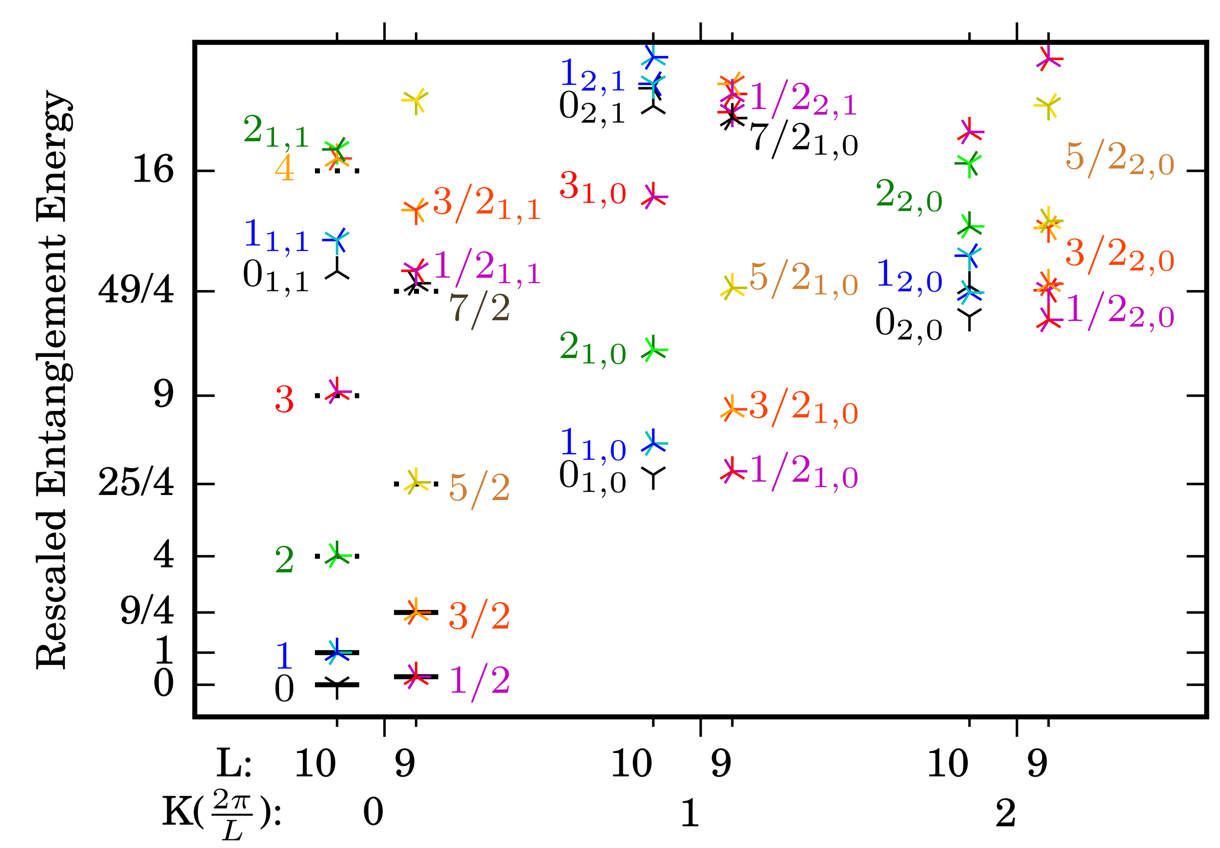
We can match the rescaled entanglement energies to the conformal weights of a free bosonic CFT.

$$\mathbf{P} = \frac{2\pi}{L}(\mathbf{L}_0 - \bar{\mathbf{L}}_0) = \frac{2\pi}{L}(em + n - \bar{n})$$

$$\mathbf{H} = \frac{2\pi}{L}(\mathbf{L}_0 + \bar{\mathbf{L}}_0) = \frac{2\pi}{L}\left(\frac{\kappa e^2}{2} + \frac{m^2}{2\kappa} + \frac{n+\bar{n}}{2}\right)$$

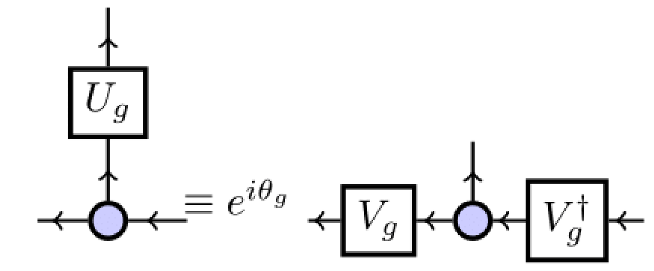
$$\boxed{\mathbf{H} \propto e^2 + \frac{m^2}{\kappa^2} + \frac{1}{\kappa}(n + \bar{n})}$$

CFT Identification of Gapless Entanglement Edge

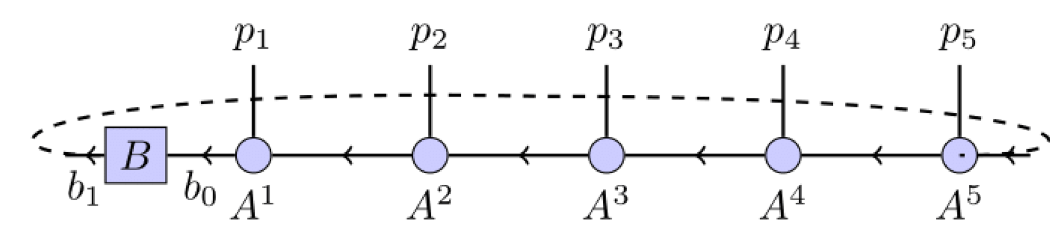


Detecting 1D SPT Order

If U_g is a global symmetry and $|\psi\rangle$ is translationally invariant, then the MPS representation satisfies:



Boundary conditions on a MPS can be represented by a matrix B which acts like:



With PBC $(B = I)$, the group action leaves the state invariant.
With OBC $(B = |i\rangle\langle i|)$, the group action rotates between states that differ only near the boundary; these edge states transform as $V_g \otimes V_g^\dagger$. V_g represents the group projectively. Equivalence classes of projective representations (enumerated by $H^2(G; U(1))$) classify 1D SPT phases.

Symmetry Protection of the Honeycomb FBI

For the state on a cylinder with odd circumference, and the zig-zag entanglement cut defined in the upper left picture, we have the following:

$\mathbf{G}$	$\mathbf{U_g}$	$\theta_{\mathbf{g}}$	$\mathbf{V_g}$	$\mathbf{V_g V_g^*}$
$U(1)$				
π				
$\mathcal{I}$				
$\pi\mathcal{I}$				

Since

$$V_{\pi\mathcal{I}}V_{\pi\mathcal{I}}^* = -I \quad \text{or} \quad V_\pi V_\mathcal{I} = -V_\mathcal{I} V_\pi,$$

the representation is in the nontrivial class of

$$\boxed{H^2(\mathbb{Z}_2 \times \mathbb{Z}_2^\mathcal{I}; U(1)) = \mathbb{Z}_2.}$$

Relation to known 1D physics

- Haldane insulator
  - Unitarily equivalent to the AKLT state
  - Distinct phase under $\mathbb{Z}_2 \times \mathbb{Z}_2^\mathcal{I}$
  - Can be connected adiabatically to $L = 1$ cylinder FBI
- Two dimensional classification is $H^3(\mathbb{Z}_2 \times \mathbb{Z}_2^\mathcal{I}; U(1)) = \mathbb{Z}_2^4$

References

- I. Kimchi, S. A. Parameswaran, A. M. Turner, F. Wang, and A. Vishwanath, Featureless and non-fractionalized mott insulators on the honeycomb lattice at 1/2 site filling,(2012), arXiv:1207.0498 [cond-mat.str-el].
- S. A. Parameswaran, A. M. Turner, D. P. Arovas, and A. Vishwanath, Topological order and absence of band insulators at integer fillng in Non-Symmorphic crystals, (2012), arXiv:1212.0557 [cond-mat.str-el]