

Entanglement in Featureless Mott Insulators

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Featureless Insulators

Definition of 'Featureless Insulator'

- Gapped
- Symmetric
- No topological order

- Unique ground state:

$$E_1 - E_0 \geq \text{const.}$$

Alternate Definition

- Unique ground state on all boundary-less systems
- Possibly with 'features' localized to edge of system

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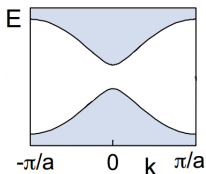
- Gapless modes:
 $E_1 - E_0 \sim \frac{1}{L^z}$

- Spontaneous symmetry breaking:
 $E_1 - E_0 = 0$

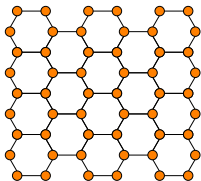
- Topological order:
 $E_1 - E_0 \sim e^{-L/\xi}$

Examples of Featureless Insulators

Classical Insulators

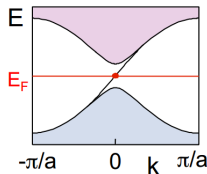


Free fermion band insulator

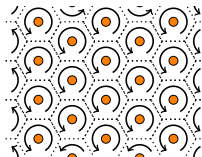


Atomic picture

Topological Insulators



Band insulator with chiral edge



Atomic picture breaks down

Obstructions to Featurelessness

Fundamental Result

A featureless insulator must have an integer charge per unit cell

- (Lieb, Schultz, Mattis 1961)
- (Hastings 2004)

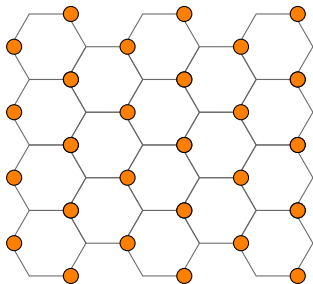
For certain lattices, not all integers are possible

- (Parameswaran 2013)

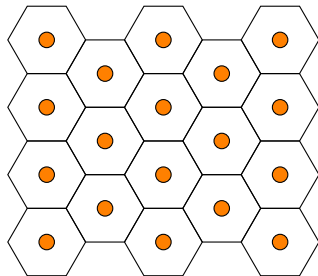
For this talk, we will look at a proposed honeycomb lattice featureless insulator with charge 1 per unit cell.

Honeycomb Bosonic Mott Insulators

Does there exist a featureless bosonic insulator with charge 1 per unit cell on the honeycomb lattice?



Breaks rotational symmetry

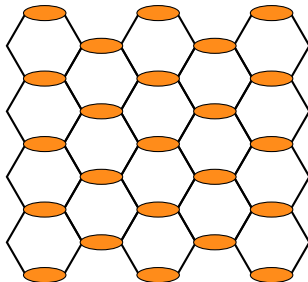


Leaves honeycomb lattice

'Classical cartoons and usual tricks' lead to symmetry breaking, as noticed by Parameswaran et al. (2013a)

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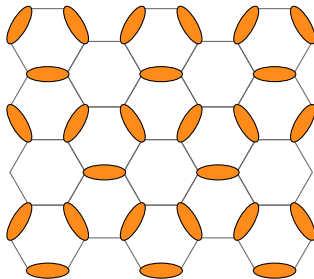


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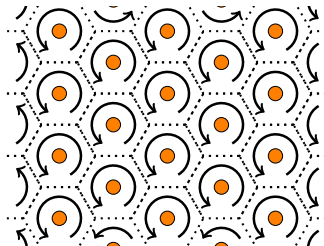
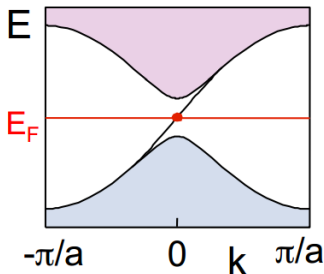


Breaks translationally symmetry, unit cell is 3 times larger

'Classical cartoons and usual tricks' lead to symmetry breaking, as noticed by Parameswaran et al. (2013a)

Honeycomb Bosonic Mott Insulators

Does there exist a featureless bosonic insulator with charge 1 per unit cell on the honeycomb lattice?



Band insulator with chiral edge ¹

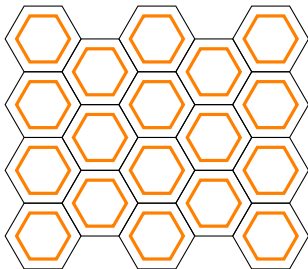
The Haldane Chern insulator is NOT an example. D_6 explicitly broken.

'Classical cartoons and usual tricks' lead to symmetry breaking, as noticed by Parameswaran et al. (2013a)

¹(Hasan and Kane, 2010)

Honeycomb Bosonic Mott Insulators

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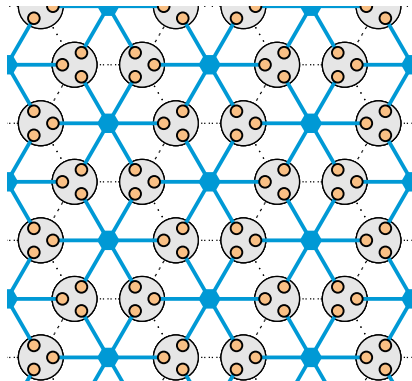


$$|\psi\rangle = \prod_{\hexagon} \left(\sum_{i \in \hexagon} b_i^\dagger \right) |0\rangle$$

Proposed Solution by Kimchi et al. (2013)

Bosons filled into non-orthogonal, plaquette centered orbitals works. Numerics confirm the expected wavefunction properties, but no known parent Hamiltonian.

Construction of Honeycomb FBI



$$\text{3 orange dots} = 3 \sqrt{3!}$$

$$\text{2 orange dots} = 2 \sqrt{2!}$$

$$\text{1 orange dot} = 1$$

$$\text{0 orange dots} = 0$$

$$\text{Site with 3 dots} = \text{Sum of 6 hexagonal clusters (each with 1 dot in a different position)}$$

$$|\psi\rangle = \prod_{\text{hex}} \left(\sum_{i \in \text{hex}} b_i^\dagger \right) |0\rangle$$

Wavefunction proposed by
Kimchi et al. (2013)

Known Results for Honeycomb FBI

Correlations

$$\langle b_i^\dagger b_j \rangle$$

- Looks rotationally symmetric
- Decays exponentially
- Correlation length
 $\xi/a \sim 3.6$
- $\langle n_i n_j \rangle$
- Looks rotationally symmetric
- Decays exponentially
- Correlation length
 $\xi/a \sim 1.6$

Hamiltonian Construction

Try filling plaquette orbitals

$$b_{\hexagon} = \sum_{i \in \hexagon} \frac{1}{\sqrt{6}} b_i^\dagger$$

$$H = \sum_{\hexagon} -\frac{t}{6} b_{\hexagon}^\dagger b_{\hexagon} + V n_{\hexagon} n_{\hexagon}$$

$$= \left(\sum_{\hexagon} \sum_{i,j \in \hexagon} -t b_i^\dagger b_j \right) - \frac{3t}{6} N + V \dots$$

- Fails, gapless modes
- Parent Hamiltonian not known

Known Results for Honeycomb FBI

Correlations

$$\langle b_i^\dagger b_j \rangle$$

- Looks rotationally symmetric
- Decays exponentially
- Correlation length $\xi/a \sim 3.6$

$$\langle n_i n_j \rangle$$

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Hamiltonian Construction

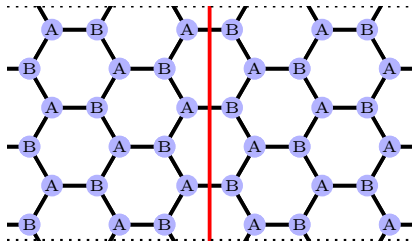
To get a parent Hamiltonian:

- Need symmetric, exponentially localized, orthogonal orbitals
- Such as the Wannier orbitals of a classical band insulator
- Parameswaran et al. (2013a)

Other lattices:

- Need a fixed point of all lattice symmetries
- Fails on nonsymmorphic lattices
- Extension of LSM theorem
- Parameswaran et al. (2013b)

Edge Geometry

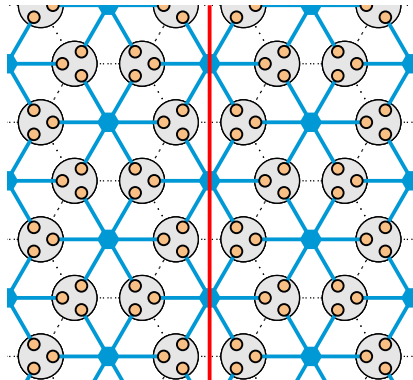


Generic honeycomb lattice PEPS on zig-zag cylinder with $L=3$

In cylindrical geometry:

- Treat state as 1D
- Use MPS techniques
- On-site translational symmetry parallel to cut
- Physical site dimension 4^{2L}

Edge Geometry

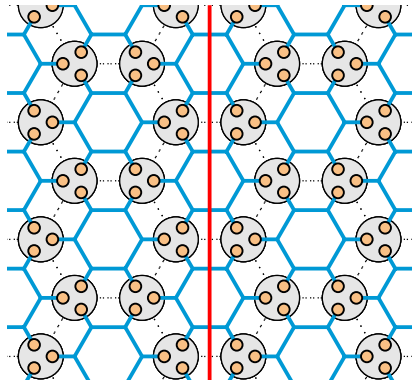


In cylindrical geometry:

- Treat state as 1D
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Honeycomb lattice PEPS on zig-zag cylinder with $L=3$, achieved by factoring W-state of plaquette bosons

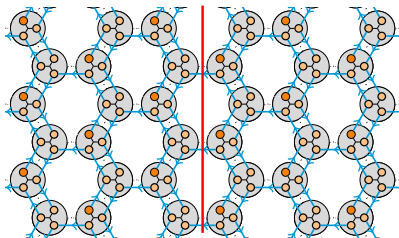
Edge Geometry



Honeycomb lattice PEPS on zig-zag cylinder with $L=3$, achieved by factoring W-state of plaquette bosons

- In cylindrical geometry:
- Treat state as 1D
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- On-site translational symmetry parallel to cut
- Physical site dimension 4^{2L}
- MPS bond dimension = Rank of $\rho_r = 2^L$
- Entanglement spectrum $\{\epsilon_i\}$ defined from eigenvalues $\{\rho_i\}$ of ρ_r via $\epsilon_i = e^{-\rho_i}$
- Charge and Translation represented linearly on edge

Edge Geometry

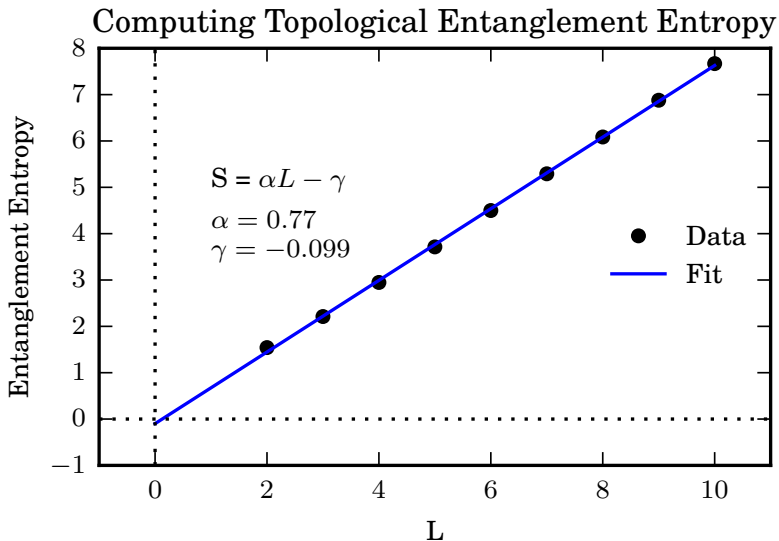


Honeycomb lattice PEPS on zig-zag cylinder with $L=3$, achieved by factoring W-state of plaquette bosons

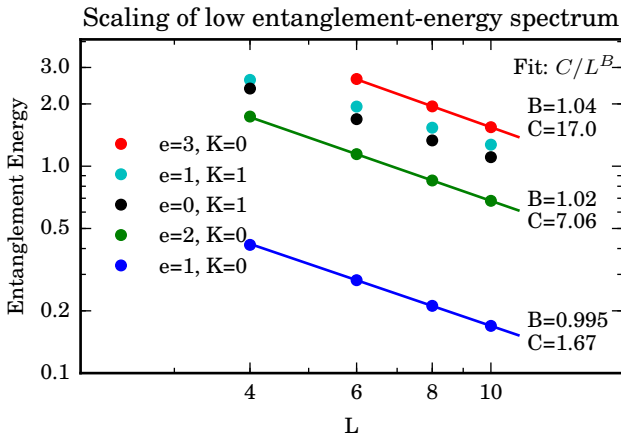
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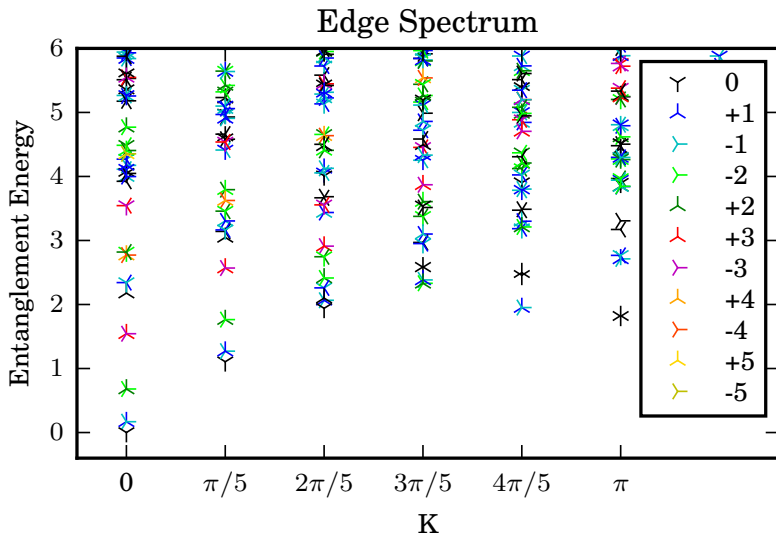
Finite Size Analysis



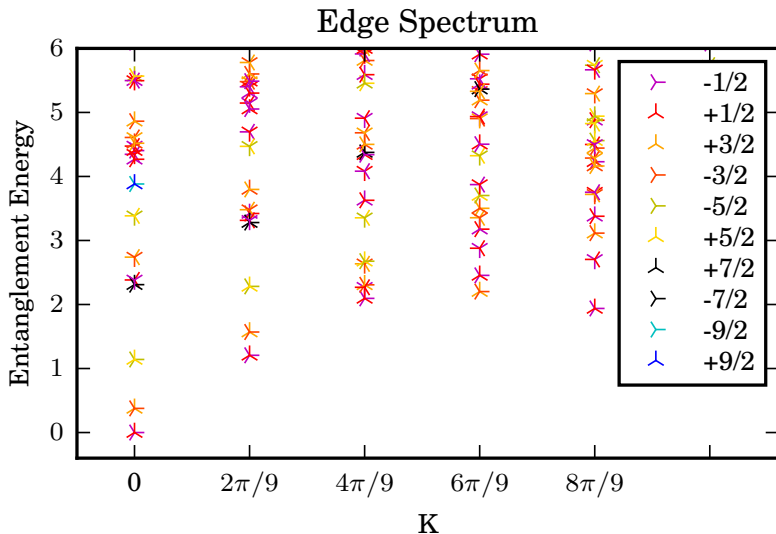
Finite Size Analysis



Entanglement Spectrum

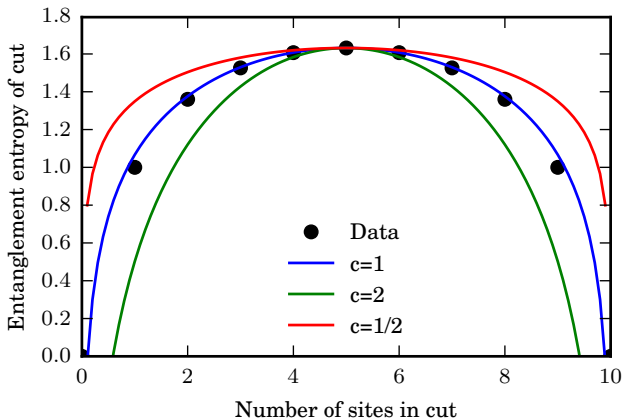


Entanglement Spectrum



Identification of Edge CFT

Conformal Charge



$$c = 1$$

Identification of Edge CFT

Conformal Weights

We can match the rescaled entanglement energies to the conformal weights of a free bosonic CFT.

$$\begin{aligned}\mathbf{P} &= \frac{2\pi}{L}(\mathbf{L}_0 - \bar{\mathbf{L}}_0) = \frac{2\pi}{L}(em + n - \bar{n}) \\ \mathbf{H} &= \frac{2\pi}{L}(\mathbf{L}_0 + \bar{\mathbf{L}}_0) = \frac{2\pi}{L}\left(\frac{\kappa e^2}{2} + \frac{m^2}{2\kappa} + \frac{n + \bar{n}}{2}\right)\end{aligned}$$

$$\mathbf{H} \propto e^2 + \frac{m^2}{\kappa^2} + \frac{1}{\kappa}(n + \bar{n})$$

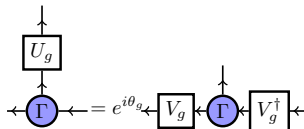
Conformal primary identification in entanglement spectra

Symmetry Protection of Degenerate Edge

1D Symmetry Protection

On-site symmetries g come with projective representation V_g

- V_g acts on sets of degenerate Schmidt states
- Charge and translation represented linearly on edge

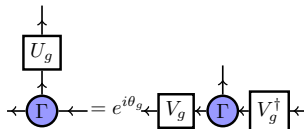


Symmetry Protection of Degenerate Edge

1D Symmetry Protection

Time reversal symmetry τ
represented by antiunitary $V_\tau K$ on
the edge

- $\tau^2 = +1$ on this edge



Symmetry Protection of Degenerate Edge

1D Symmetry Protection

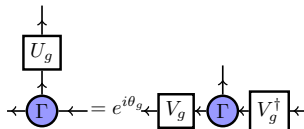
Inversion \mathcal{I}

- \mathcal{I} in combination with swapping Schmidt states represented by antiunitary operation $V_{\mathcal{I}}K$ on the edge

- $\mathcal{I}^2 = V_{\mathcal{I}}V_{\mathcal{I}}^* = 1$

Inversion \mathcal{I} combined with $\pi = e^{i\pi N}$

- $\pi\mathcal{I}$ represented antiunitarily on the edge by $V_{\pi\mathcal{I}}K$
- $(\pi\mathcal{I})^2 = 1$ but $V_{\pi\mathcal{I}}V_{\pi\mathcal{I}}^* = -1$



Future Work

Entanglement properties with different geometries

- Armchair cylinder edge
- Finite size clusters
- Explain results for arbitrary geometries with tensor network properties, e.g. 'MPO injectivity'

Find a 2D local Hamiltonian and confirm with numerics

$$H_{EBH} = \left(\sum_{\hexagon} \sum_{i,j \in \hexagon} -tb_i^\dagger b_j + V n_i n_j \right) + \mu N?$$

Physical properties of the phase

Can we construct an SU(2) symmetric FI?

Resources

- Hasan, M. Z. and Kane, C. L. (2010). *Colloquium: Topological insulators. Reviews of modern physics*, 82(4):3045–3067.
- Kimchi, I., Parameswaran, S. A., Turner, A. M., Wang, F., and Vishwanath, A. (2013). Featureless and nonfractionalized mott insulators on the honeycomb lattice at $1/2$ site filling. *Proceedings of the National Academy of Sciences*, 110(41):16378–16383.
- Parameswaran, S. A., Kimchi, I., Turner, A. M., Stamper-Kurn, D. M., and Vishwanath, A. (2013a). Wannier permanent wave functions for featureless bosonic mott insulators on the $1/3$ -filled kagome lattice. *Phys. Rev. Lett.*, 110:125301.
- Parameswaran, S. A., Turner, A. M., Arovas, D. P., and Vishwanath, A. (2013b). Topological order and absence of band insulators at integer filling in non-symmorphic crystals. *Nature Physics*, 9(5):–303?
- Pollmann, F., Turner, A. M., Berg, E., and Oshikawa, M. (2010). Entanglement spectrum of a topological phase in one

Questions?

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Bonus slides

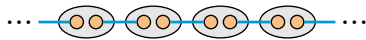
Construction of 1D Featureless Insulators

Classical Insulators



1D Trivial Chain

Topological Insulators



1D Topological Chain

$$\begin{aligned}
 \text{orange dot} - \text{blue line} - \text{orange dot} &= \text{white circle} + \text{orange dot} + \text{orange dot} + \text{white circle} \\
 \text{gray oval with two white circles} &= 0 \\
 \text{gray oval with one white circle and one orange dot} &= 1 \\
 \text{gray oval with two orange dots} &= 2
 \end{aligned}$$

Projectors and entangled pairs (PEPS) used in state construction

Construction of 1D Featureless Insulators

Classical Insulators



1D Trivial Chain

Product state with one boson per site

Topological Insulators



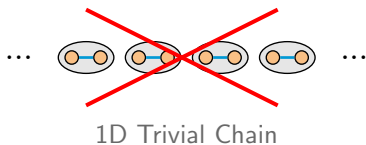
1D Topological Chain

Haldane Insulator Phase
Pollmann et al. (2010)

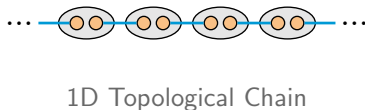
- Unitarily related to AKLT
- No $SU(2)$ symmetry
- Symmetry protected 2-fold edge degeneracy

Construction of 1D Featureless Insulators

Classical Insulators



Topological Insulators



$$\begin{aligned}
 \bullet\bullet &= \circ - \bullet\bullet\circ \\
 \bullet\bullet &= -\sqrt{2} \\
 \circ\bullet &= 0 \\
 \bullet\bullet &= +\sqrt{2}
 \end{aligned}$$

Projectors and entangled pairs (PEPS) for $SU(2)$ symmetric state