

USER MANUAL

NEMEC OUTPUT

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NEMEC EQUILIBRIUM OUTPUT

The program **readnemec.f90** reads the equilibrium output of the NEMEC code, which is stored in the output file **wout.xxx**. Both, ASCII and binary data files are possible. **Please note**, the equilibrium output is provided in double precision: **real-size 64**. Therefore, on a LINUX system the program has to be compiled with

ifort -real-size 64 or **mpiifort -real-size 64**

Please also read the comments in program **readnemec.f90**.

The NEMEC code [1, 2, 3] uses the left-handed flux coordinates (s, θ, ζ) with the radial coordinate s being the normalized toroidal flux, the poloidal coordinate, θ , ($0 \leq \theta < 2\pi$), and the toroidal coordinate, ζ , ($0 \leq \zeta < 2\pi$).

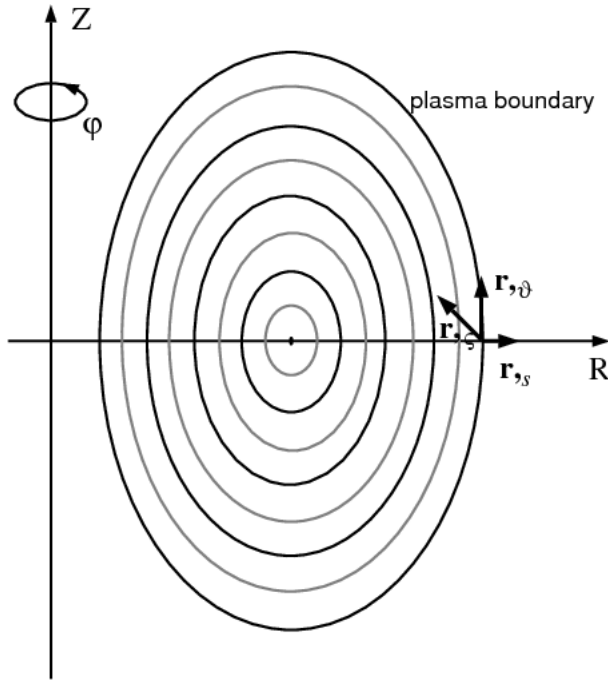


Fig. 1: Flux surfaces from the magnetic axis to the plasma boundary for a poloidal cross-section in the cylindrical coordinate system (R, φ, Z) . The solid lines mark the flux surfaces (full mesh) on which the Fourier coefficients $\{\hat{r}_{m,n}^c(s_i), \hat{r}_{m,n}^s(s_i), \hat{z}_{m,n}^c(s_i), \hat{z}_{m,n}^s(s_i)\}$ are given, while the dashed lines belong to surfaces (half mesh) on which the Fourier coefficients of the magnetic field are defined. The arrows represent the curvilinear basis vectors $(\mathbf{r}_s, \mathbf{r}_\theta, \mathbf{r}_\zeta)$ used in the NEMEC code.

The flux coordinates are related with the cylindrical coordinates (R, φ, Z) :

$$R = \sum_{m=0, n=-n_b}^{m_b, n_b} \hat{r}_{m,n}^c(s) \cos(m\theta - n\zeta N_p) + \hat{r}_{m,n}^s(s) \sin(m\theta - n\zeta N_p),$$

$$\varphi = \zeta, \tag{1}$$

$$Z = \sum_{m=0, n=-n_b}^{m_b, n_b} \hat{z}_{m,n}^c(s) \cos(m\theta - n\zeta N_p) + \hat{z}_{m,n}^s(s) \sin(m\theta - n\zeta N_p),$$

with $\{\hat{r}_{m,n}^c(s), \hat{r}_{m,n}^s(s), \hat{z}_{m,n}^c(s), \hat{z}_{m,n}^s(s)\}$ being the Fourier coefficients of a flux surface with normalized toroidal flux s . N_p is the number of periods, $0 \leq m \leq m_b$ are the poloidal harmonics, and $-n_b \leq n \leq n_b$ are the toroidal harmonics.

The covariant form of the magnetic field reads

$$\mathbf{B} = B^\theta \mathbf{r}_{,\theta} + B^\zeta \mathbf{r}_{,\zeta}, \quad (2)$$

with the contravariant components

$$B^s = 0,$$

$$B^\theta = \sum_{m=0, n=-n_b}^{m_b, n_b} \hat{b}_{m,n}^{\theta,c}(s) \cos(m\theta - n\zeta N_p) + \hat{b}_{m,n}^{\theta,s}(s) \sin(m\theta - n\zeta N_p), \quad (3)$$

$$B^\zeta = \sum_{m=0, n=-n_b}^{m_b, n_b} \hat{b}_{m,n}^{\zeta,c}(s) \cos(m\theta - n\zeta N_p) + \hat{b}_{m,n}^{\zeta,s}(s) \sin(m\theta - n\zeta N_p),$$

and the Fourier coefficients $\{\hat{b}_{m,n}^{\theta,c}(s), \hat{b}_{m,n}^{\theta,s}(s), \hat{b}_{m,n}^{\zeta,c}(s), \hat{b}_{m,n}^{\zeta,s}(s)\}$.

Its contravariant form is given by

$$\mathbf{B} = B_s \nabla s + B_\theta \nabla \theta + B_\zeta \nabla \zeta \quad (4)$$

with the covariant components

$$B_s = \sum_{m=0, n=-n_b}^{m_b, n_b} \hat{b}_{s,c}^{m,n}(s) \cos(m\theta - n\zeta N_p) + \hat{b}_{s,s}^{m,n}(s) \sin(m\theta - n\zeta N_p)$$

$$B_\theta = \sum_{m=0, n=-n_b}^{m_b, n_b} \hat{b}_{\theta,c}^{m,n}(s) \cos(m\theta - n\zeta N_p) + \hat{b}_{\theta,s}^{m,n}(s) \sin(m\theta - n\zeta N_p) \quad (5)$$

$$B_\zeta = \sum_{m=0, n=-n_b}^{m_b, n_b} \hat{b}_{\zeta,c}^{m,n}(s) \cos(m\theta - n\zeta N_p) + \hat{b}_{\zeta,s}^{m,n}(s) \sin(m\theta - n\zeta N_p)$$

and the Fourier coefficients $\{\hat{b}_{s,c}^{m,n}(s), \hat{b}_{s,s}^{m,n}(s), \hat{b}_{\theta,c}^{m,n}(s), \hat{b}_{\theta,s}^{m,n}(s), \hat{b}_{\zeta,c}^{m,n}(s), \hat{b}_{\zeta,s}^{m,n}(s)\}$.

Furthermore, the NEMEC code uses the following Fourier representation of the single-valued function λ

$$\lambda = \sum_{m=0, n=-n_b}^{m_b, n_b} \hat{\lambda}_{m,n}^c(s) \cos(m\theta - n\zeta N_p) + \hat{\lambda}_{m,n}^s(s) \sin(m\theta - n\zeta N_p) \quad (6)$$

with the Fourier coefficients $\{\hat{\lambda}_{m,n}^c(s), \hat{\lambda}_{m,n}^s(s)\}$.

The NEMEC code provides the Fourier coefficients

$$\{\hat{r}_{m,n}^c(s_i), \hat{r}_{m,n}^s(s_i), \hat{z}_{m,n}^c(s_i), \hat{z}_{m,n}^s(s_i), \hat{\lambda}_{m,n}^c(s_i), \hat{\lambda}_{m,n}^s(s_i)\},$$

for a discrete number, N_s of nested flux surfaces i with $0 \leq i \leq N_s - 1$ (full mesh), while the Fourier coefficients

$$\{\hat{b}_{m,n}^{s,c}(s_j), \hat{b}_{m,n}^{s,c}(s_j), \hat{b}_{m,n}^{\theta,c}(s_j), \hat{b}_{m,n}^{\theta,s}(s_j), \hat{b}_{m,n}^{\phi,c}(s_j), \hat{b}_{m,n}^{\phi,s}(s_j), \hat{b}_{\theta,c}^{m,n}(s_j), \hat{b}_{\theta,s}^{m,n}, \hat{b}_{\phi,c}^{m,n}(s_j), \hat{b}_{\phi,s}^{m,n}(s_j)\}$$

are defined on flux surfaces in between (half mesh), that is, $s_j = (s_i + s_{i+1})/2$ (for details see Fig. 1).

The equilibrium output (file **wout.xxx**) of the NEMEC code, which is written in subroutine **wrout.f90**, has the following form:

List of variables:

Dimensions of the fields

gamma	real	adiabatic constant
enfp	real	N_p = number of periods
enrho	real	N_s = number of flux surfaces ($0 \leq i \leq N_s - 1$)
empol	real	$m_b + 1$ = number of poloidal harmonics ($0 \leq m \leq m_b$)
entor	real	n_b = maximum toroidal harmonic ($-n_b \leq n \leq n_b$)
empmnt	real	total number of harmonics $m_b(2n_b + 1) + (n_b + 1)$
eiasym	real	symmetry of the equilibrium
	= 0	stellarator-symmetric equilibrium
	= 1	asymmetric equilibrium
phiedge	real	total toroidal flux

Please note, these variables are transformed into integers after reading.

Fourier coefficients of R , Z , B^θ , B^ζ , λ , B_θ , B_ζ , and B_s

frmnc(m,n,i)	real	$\hat{r}_{m,n}^c(s_i)$
fzmns(m,n,i)	real	$\hat{z}_{m,n}^s(s_i)$
frmns(m,n,i)	real	$\hat{r}_{m,n}^s(s_i)$
fzmnc(m,n,i)	real	$\hat{z}_{m,n}^c(s_i)$
hbumnc_up(m,n,j)	real	$\hat{b}_{m,n}^{\theta,c}(s_j)$
hbvmnc_up(m,n,j)	real	$\hat{b}_{m,n}^{\zeta,c}(s_j)$
hbumns_up(m,n,j)	real	$\hat{b}_{m,n}^{\theta,s}(s_j)$
hbvmns_up(m,n,j)	real	$\hat{b}_{m,n}^{\zeta,s}(s_j)$

flmns(m,n,i)	real	$\hat{\lambda}_{m,n}^s(s_i)$
flmnc(m,n,i)	real	$\hat{\lambda}_{m,n}^c(s_i)$
hbumnc_dw(m,n,j)	real	$\hat{b}_{\theta,c}^{m,n}(s_j)$
hbvmnc_dw(m,n,j)	real	$\hat{b}_{\zeta,c}^{m,n}(s_j)$
hbsmns_dw(m,n,j)	real	$\hat{b}_{s,s}^{m,n}(s_j)$
hbumns_dw(m,n,j)	real	$\hat{b}_{\theta,s}^{m,n}(s_j)$
hbvmns_dw(m,n,j)	real	$\hat{b}_{\zeta,s}^{m,n}(s_j)$
hbsmnc_dw(m,n,j)	real	$\hat{b}_{s,c}^{m,n}(s_j)$

Radial profiles: The definitions below relate the NEMEC output quantities to the corresponding ASDEX Upgrade quantities in SI units.

hiota(j)	real	rotational transform ι
hmass(j)	real	mass function
hpres(j)	real	$\mu_0 p$ with p = pressure in [Pa]
hhiph(j)	real	$\frac{1}{2\pi} \frac{\partial \Phi(s)}{\partial s}$ with Φ = toroidal flux in [Tm ²]
hbuc(j)	real	$-\frac{\mu_0}{2\pi} I$ with I = toroidal current in [A]
hbvco(j)	real	$-\frac{\mu_0}{2\pi N_p} J$ with J = poloidal current in [A]
hphi(j)	real	$\Phi(s)$ = toroidal flux in [Tm ²]
hvp(j)	real	$\frac{1}{4\pi^2 N_p} \frac{\partial V(s)}{\partial s}$ with $V(s)$ = volume in [m ³]
hoverr(j)	real	overr(j) in NEMEC
fjcuru(i)	real	$-\frac{1}{2\pi N_p} \frac{\partial J}{\partial s}$
fjcurv(i)	real	$\frac{1}{2\pi} \frac{\partial I}{\partial s}$
hspecw(j)	real	specw(j) in NEMEC

The prefixes **f** and **h** denote full mesh and half mesh quantities, respectively. The suffix **up** marks the contravariant vector components, while **dw** denotes the covariant ones.

Below ASDEX Upgrade and NEMEC algebraic signs of the toroidal current, I , the poloidal current, J , the toroidal magnetic field, B_{tor} , the poloidal magnetic field, B_{pol} , the toroidal flux, Φ , and its radial derivative, Φ' , the poloidal flux, Ψ , and its radial derivative, Ψ' , and the safety factor $q = \Phi'/\Psi'$ are defined.

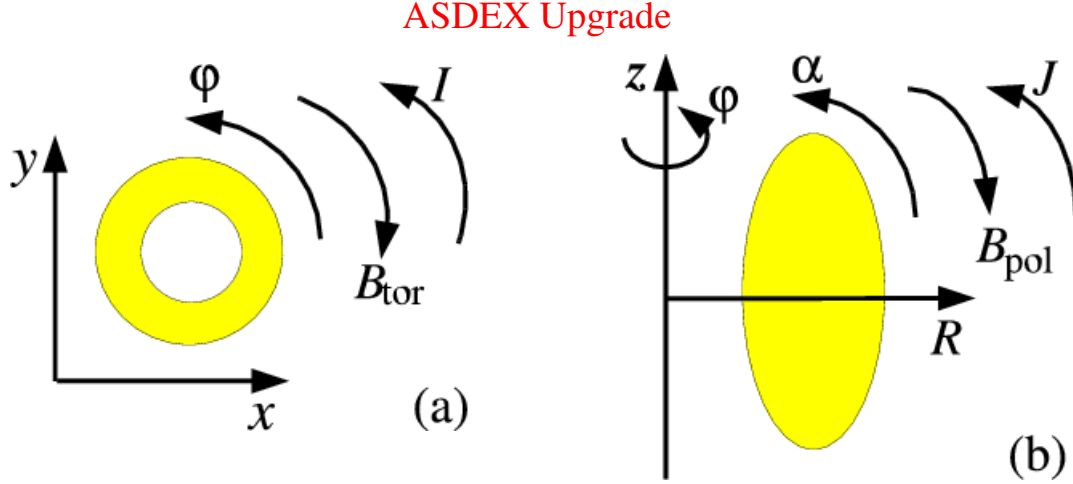


Fig. 2: (a) Toroidal and (b) poloidal field and current directions of ASDEX Upgrade.

Algebraic signs of AUG and NEMEC quantities									
	I	J	B_{tor}	B_{pol}	Φ	Φ'	Ψ	Ψ'	q
AUG	+	+	-	-	-	-	+	-	+
NEMEC	+	+	-	-	-	-	-	-	+

Please note, $\Psi_{\text{NEMEC}}(s) = \Psi_{\text{AUG}}(s) - \Psi_{\text{AUG}}(s = 0)$

References

- [1] Hirshman S P and Whitson J C 1983 ‘Steepest-decent moment method for three-dimensional magnetohydrodynamic equilibria.’ *Phys. Fluids* **26** 3553. doi:10.1063/1.864116.
- [2] Hirshman S P and Lee D K 1986 ‘MOMCON: a spectral code for obtaining three-dimensional magnetohydrodynamic equilibria.’ *Comput. Phys. Commun.* **39** 161.
- [3] Hirshman S P, van Rij W I, and Merkel P 1986 ‘Three-dimensional free boundary calculations using a spectral Greens’s function method.’ *Comput. Phys. Commun.* **43** 143.