# Hercules-Aquila and Virgo Clouds with Gaia DR2. Evidence for a common origin

Iulia T. Simion<sup>1★</sup>, Vasily Belokurov<sup>2,3</sup> and Sergey E. Koposov<sup>2,4</sup>

- 1 Key Laboratory for Research in Galaxies and Cosmology, Shanghai Astronomical Observatory, 80 Nandan Road, Shanghai 200030, China
- <sup>2</sup>Institute of Astronomy, Madingley Rd, Cambridge, CB3 0HA
- <sup>3</sup> Center for Computational Astrophysics, Flatiron Institute, 162 5th Avenue, New York, NY 10010, USA

Accepted XXX. Received YYY; in original form ZZZ

## ABSTRACT

200 words for Letters. No references should appear in the abstract.

**Key words:** keyword1 – keyword2 – keyword3

# 1 INTRODUCTION

How do you hide the evidence for a massive impact event that caused the extinction of most of the dinosaurs as well as 75% of all species on Earth? You bury it deep under the sea, covered with a layer of sediment taller than the Empire State Building (Hildebrand et al. 1991). Without the discovery of the giant Chicxulub crater, the meteorite impact hypothesis would remain a neat theory supported by striking but indirect evidence. A hypothesis of an ancient dramatic collision between the Milky Way and an unidentified massive dwarf galaxy was put forward by Deason et al. (2013) to explain a particular feature in the overall stellar halo density profile (Watkins et al. 2009; Deason et al. 2011). Most recently, through a study of a portion of the nearby stellar halo, Belokurov et al. (2018b) demonstrated that the great impactor must have collided with the young Milky Way on a nearly radial orbit, thus swamping the inner stellar halo with metal-rich material with orbital anisotropy (see Binney & Tremaine 2008) close to unity. Merger events like this tend to leave behind a battery of debris clouds and shells (see Amorisco 2015; Hendel & Johnston 2015), which - akin to the peak rings of impact craters (see e.g. Morgan et al. 2016) - if discovered could help to reconstruct the collision as well as pin down the properties of the progenitor (e.g. Sanderson & Helmi 2013; Johnston 2016).

Before the Data Release 2 (Gaia Collaboration et al. 2018) of the ESA's Gaia mission (Gaia Collaboration et al. 2016), five large and diffuse cloud-like structures had been discovered in the Galaxy's halo. These include: the Virgo Over-Density (VOD, Vivas et al. 2001; Newberg et al. 2002; Duffau et al. 2006; Jurić et al. 2008; Bonaca et al. 2012), the Hercules-Aquila Cloud (HAC, Belokurov et al. 2007; Simion et al. 2014), the Trinagulum-Andromeda structure (Tri-And, Rocha-Pinto et al. 2004; Majewski et al. 2004;

\* E-mail:isimion@shao.ac.cn

Deason et al. 2014), the Pisces Over-density (Sesar et al. 2007; Watkins et al. 2009; Nie et al. 2015) and the Eridanus-Phoenix over-density (Eri-Pho, Li et al. 2016). According to the most recent investigations, Tri-And likely comprises of Galactic disc stars kicked out of the plane in a recent interaction with a dwarf galaxy, probably the Sagittarius dSph (e.g. Price-Whelan et al. 2015; Bergemann et al. 2018; Hayes et al. 2018). Of the remaining four, the Pisces overdensity clearly stands out as it reaches much larger Galacto-centric distances. On the other hand, the VOD, HAC and Eri-Pho structures occupy a very similar range of distances, between 10 and 20 kpc from the Galactic center. This lead Li et al. (2016) to suggest that these three Clouds could all be part of one merger event, a galaxy accreted onto the Milky Way on a polar orbit (see also Jurić et al. 2008).

As demonstrated by the recent re-interpretation of the Monoceros Ring (and the associated sub-structures) and the Tri-And, deciphering the nature of halo over-densities is often impossible without either high-resolution spectroscopy (e.g. Bergemann et al. 2018) or accurate astrometry (e.g. de Boer et al. 2018; Deason et al. 2018a). In this Letter, we look for clues to the formation of the Hercules-Aquila and Virgo Clouds using proper motions provided as part of the Gaia DR2. At our disposal are highly pure samples of members of each Cloud, namely the RR Lyrae stars that i) are co-spatial with HAC and VOD in 3-D and ii) that have their line-ofsight velocities measured. By complementing the publicly available 4-D data with the GDR2 proper motions, we build a large tracer set with complete 6-D phase space information and study the make-up of each structure using the orbital properties of the constituent stars.

<sup>&</sup>lt;sup>4</sup>Department of Physics, McWilliams Center for Cosmology, Carnegie Mellon University, 5000 Forbes Avenue, Pittsburgh, PA 15213, USA

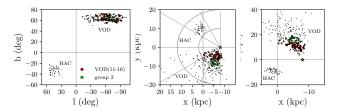


Figure 1. Spatial distribution of the RR Lyrae used in this work with full 6-D phase space measurements, in Galactic coordinates (left panel) and in the x-y (middle) and y-z (right) planes. The HAC field contains 44 RR Lyrae which likely belong to the Cloud with measured line-of-sight velocities (Simion et al. 2018) and Gaia DR2 proper motions. The VOD field contains 411 RRL which belong to several halo associations, including the Sagittarius stream and the VOD, with line-of-sight velocities provided by Vivas et al. 2016 and proper motions from Gaia DR2. In particular we mark group 2, a 'high significance' kinematical group, which contains 18 stars (green circles). The semi-circles are centred on the Sun's position and have radius of 10 and 20 kpc. The Sun (yellow star) is located at  $(x_{\odot}, y_{\odot}, z_{\odot}) = (0,-8,0)$  kpc and the Galactic centre at (0,0,0) - black circle.

## 2 DATA AND ANALYSIS

#### 2.1 4-D data

The Hercules-Aquila and Virgo Clouds are revealed unambiguously as overdensities of RR Lyrae in the halo, which peak at similar heliocentric distances, ~17 kpc (HAC: Simion et al. 2014) and ~19 kpc (VOD: Vivas & Zinn 2006; Duffau et al. 2014; Vivas et al. 2016) respectively, but in opposite quadrants (e.g. their distribution in Galactic coordinates in Fig. 1). Other tracers (e.g. BHBs, MSTO, K giants) have been used to pin down the morphology of the Clouds with slightly inconsistent results between them, caused by poor distance determination, scarcity or simply trace the structures differently; therefore, the most recent kinematic followups of the clouds have used RR Lyrae selected in the vicinity of the overdensities peak:

## 2.2 4-D RR Lyrae data

The Hercules-Aquila and Virgo Clouds are diffuse stellar over-densities in the inner stellar halo, located on the opposite sides of the Galaxy (see Figure 1). At high latitudes, these are detected as peaks in RR Lyrae number counts - curiously - at similar heliocentric distances, i.e. ~17 kpc (HAC: Watkins et al. 2009; Simion et al. 2014) and ~19 kpc (VOD: Vivas & Zinn 2006; Duffau et al. 2014; Vivas et al. 2016). Note that other tracers (e.g. BHBs, MSTO and K and M giants) have also been used to pin down the morphology of the Clouds (see e.g. Belokurov et al. 2007; Jurić et al. 2008; Sharma et al. 2010; Bonaca et al. 2012; Conroy et al. 2018). The RR Lyrae, however offer the clearest view of these halo sub-structures thanks to the associated accurate distances and minuscule Galactic foreground contamination. Therefore, in this work, we have focused on the two recently published samples of RR Lyrae towards the Clouds, where each star has a well-measured line-of-sight velocity.

• Simion et al. (2018) provide a table of 46 RRL with radial velocity measurements (45 observed with MDM and

- 1 from SDSS) with heliocentric distances between 15 and 18 kpc;
- Vivas et al. (2016) compiled a catalog of 412 RRL in the region of the sky covered by the VOD with distances between 4 and 75 kpc from the Sun with radial velocity measurements of stars from La Silla-QUEST, QUEST, CRTS and LINEAR.

#### 2.3 From 4-D to 6-D. Velocity distributions

By cross-matching to the GDR2 data with an aperture of 2". we have found Gaia counterparts to 44 HAC stars and 411 VOD. From the VOD sample, we remove 112 stars likely belonging to the Sgr stream as identified by (Vivas et al. 2016, , their group 1). The spatial distribution of the remaining stars (44 from Simion et al. 2018 and 299 from Vivas et al. 2016) with full 6-D phase space measurements is given in Figure 1, in Galactic coordinates in the left panel and in the x-y (x-z) Galactic plane in the middle (right) panel. We adopt a left-handed Galactic Cartesian coordinates with the Sun located at  $(x_{\odot}, y_{\odot}, z_{\odot}) = (-8,0,0)$  kpc, the x-axis positive in the direction of the Galactic center, y-axis oriented along the Galactic rotation and the z-axis directed towards the north Galactic pole. While Vivas et al. 2016 identify 6 significant kinematical groups in the VOD region (their Table 5), only groups 1 (Sagittarius stream) and 2 (likely members of the VOD, with  $\langle v_{GSR} \rangle = 135 \text{ km/s}$ ) contain more than 10 stars. We mark group 2 with green circles in Figure 1. We also mark with red circles the location of a group of stars selected to have galactocentric distances similar to the HAC sample, i.e.  $11 \! < \! r_{GC}/kpc \! < 16$  to facilitate a fair comparison of their velocities and orbits.

Figure 2 shows the spherical polar components of the velocities (radial  $v_r$ , azimuthal  $v_\theta$  and polar  $v_\phi$ ) of stars in the HAC (top row) and VOD (bottom row) fields. Group 2 stars (green) can be seen clustering at  $v_r=135$  km/s (by design), while the stars at intermediate  $r_{GC}$  (red) seem to have a velocity distribution very similar to those in the HAC, shown in the top row. To estimate the uncertainty in each velocity component we propagate the measured proper motion and line-of-sight velocity errors using Monte-Carlo re-sampling, where we take into account the covariances between the measurements of the Right Ascension and Declination components of proper motion as provided in GDR2. We use the standard deviations of the resulting  $\{v_r, v_\theta, v_\phi\}$  distributions as the upper limits of the velocity uncertainties; these are shown in Fig. 2.

As evident from the Figure, the velocity distributions are highly anisotropic, with the dispersion in the radial component dominating the tangential ones, especially in the HAC data. To describe the shape of the velocity distributions, we model each stellar sample as a single-component multi-variate Gaussian using the Extreme Deconvolution module implemented in astroML (Vanderplas et al. 2012). The resulting parameters and the associated uncertainties of the velocity ellipsoids are taken to be the median and the standard deviation values of 500 bootstrap re-sampling trials. The velocity ellipsoid shape can be summarized using the anisotropy parameter  $\beta = 1 - (\sigma_{\theta}^2 + \sigma_{\phi}^2)/2\sigma_r^2$  (see Binney & Tremaine 2008). We find the HAC stars have radially biased orbits with  $\beta = 0.91 \pm 0.03$  while for the whole of

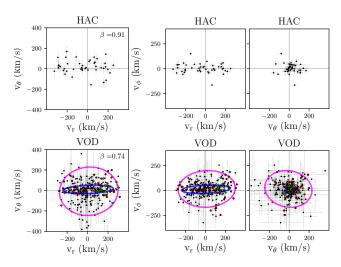
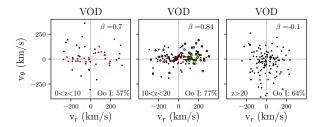


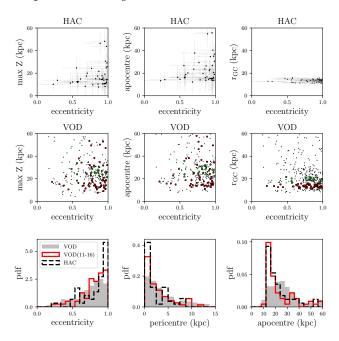
Figure 2. RRL velocity distribution in spherical polar coordinates  $(v_r, v_\theta, v_\phi)$  are the radial, azimuthal and polar components respectively) in the HAC field (top row) and the VOD field (middle and bottom rows). The error on the velocity components of each star i,  $[\sigma^i_{v_r}, \sigma^i_{v_\theta}, \sigma^i_{v_\theta}]$ , has been propagated by randomly drawing 1000 stars from a multivariate Gaussian distribution with mean the measurement  $({\rm ra}^i, {\rm dec}^i, {\rm d}^i, {\rm pmra}^i, {\rm pmdec}^i, v^i_h)$  and full covariance matrix (takes into account the covariances between ra,dec and proper motions). The orbital anisotropy, is highly radial in the HAC field  $(\beta=0.91\pm0.03)$  where the stars are most likely members of the Cloud and mildly radial in the VOD field  $(\beta=0.74\pm0.04)$  in which stars span a much wider range of distances (see Fig. 1).



**Figure 3.** Radial versus azimuthal velocity in the VOD field, in three distance ranges above the Galactic plane. The fraction of RR Lyrae of Oosterhoff type I is reported in each panel.

the VOD sample,  $\beta=0.74\pm0.04$ ). Note, however, that the two samples span very different ranges in Galactic l,b and distances. According to Belokurov et al. (2018b), the inner stellar halo can be viewed as a mixture of two debris components with distinct properties. Accordingly, we fit a model with two multivariate Gaussians to the VOD velocity data using Extreme Deconvolution (see Bovy et al. 2011). The VOD sample appears to be composed of roughly two thirds of stars with  $\beta_1=0.96^{+0.02}_{-0.44}$  (marked in magenta in Fig. 2) and a third with  $\beta_2=0.44^{+0.45}_{-0.20}$  (blue) in good agreement with the results for the local halo presented in Belokurov et al. (2018b).

Using virial theorem, Myeong et al. (2018a) concluded that the radially anisotropic component of the stellar halo is also significantly flattened vertically. To test this hypothesis, we split the VOD sample into 3 groups according to their distance from the Galactic plane and show the behaviour of the



**Figure 4.** Orbital properties of the stars in the HAC and VOD fields. 'group 2' has similar orbital properties to the HAC, however it does not display a sausage velocity distribution (see middle row figure 2) - they are concentrated at  $vr = 135 \mathrm{km/s}$  as calculated by Vivas et al. 2016.

azimuthal  $v_{\theta}$  and radial  $v_r$  velocity distributions in Figure 3. Additionally, for each z slice we have calculated the fraction of Oosterhoff type I (Oo I) RR Lyrae, using equations 1 and 2 in Belokurov et al. (2018a) to explore the changes in the RRL mixture. In the 10 < z/kpc < 20 range, where the velocity anisotropy is the highest ( $\beta = 0.84 \pm 0.03$ ) approaching the value in the HAC field, the Oo I type dominates (77%). In the same slice, 73% of the stars belong to the more squashed (or 'sausage' looking) velocity ellipsoid. The same behaviour but less accentuated can be noticed in the 0<z/kpc<10 slice where the anisotropy  $\beta = 0.7 \pm 0.1$  is lower but the fraction of Oo I stars is somewhat lower at 57%. Further from the plane, at z>20 kpc, the velocity ellipsoid changes dramatically to almost isotropic with  $\beta = -0.1 \pm 0.2$ . Please note that in this particular z bin, the  $\beta$  value may still be affected by the presence of the Sagittarius stream. Group 2 and the stars sharing the same galactocentric distance range with the HAC, are all located at z < 20 kpc. Interestingly, the HAC counterparts in the VOD (red points) are all clearly part of the anisotropic component.

# 2.4 Orbital properties of the HAC and VOD

For orbit integration we use the galpy package Bovy (2015) in the recommended MWPotential2014 model for the Galactic potential which is composed of a Miyamoto-Nagai disc, a bulge with a power-law density profile that is exponentially cut-off, and a dark matter halo described by a NFW potential. The parameters are given in Bovy (2015), table 1. The resulting orbital parameters of the HAC and VOD are given in Fig 4. To compute the errors (not shown for VOD to simply the figure) we integrated 500 orbits for each

# 4 Iulia T. Simion et al

star where the orbits were initialised on parameters resampled from data, as in the previous section. The pdfs of the eccentricities, apocentres and pericentres are also shown. If the HAC orbits are clearly highly eccentric, we find that the majority of stars in VOD are also on quite eccentric orbits: in particular, the stars with  $11 < r_{\rm GC}/\rm kpc < 16$  (in red), follow closely the HAC orbital properties.

### 3 ARE THE VOD AND THE HAC RELATED?

In this section we discuss the possible connection between the two clouds. We use the 6-D phase space measurements as initial conditions for backward orbit integration. We integrate the orbits for 8 Gyrs and show them in Fig. 5.

## 4 CONCLUSIONS

Using a sample of ~500 RR Lyrae with full 6-D phase space information, we have studied the orbital properties of the Hercules-Aquila and Virgo Clouds. Both Clouds appear dominated by stars on highly eccentric orbits. Assuming that the kinematics of each structure is well described by a single Gaussian, the orbital anisotropy of the HAC is  $\beta = 0.91$  and for the VOD,  $\beta = 0.74$ . Note, however that the original criteria applied to the RR Lyrae stars to select targets for spectroscopic follow-up differ drastically between the HAC and the VOD datasets. The HAC sample covers a very limited region in the of l, b, D space, while the VOD dataset spans a wide range of longitudes, latitudes and heliocentric distances. It is therefore likely that the VOD dataset contains a mixture of several halo sub-structures (see Vivas et al. 2016, for a detailed discussion). For the entirety of the analysis described here, we made sure to cull the probable Sgr stream members. Additionally, we explore how the VOD's make-up changes with Galactic height and demonstrate that for |z| < 20 kpc, the VOD orbital anisotropy is  $\beta\sim0.84,$  while above this threshold, it quickly changes to  $\beta \sim 0$ . We conclude therefore, that an assumption of a single Gaussian for the entire VOD sample is not appropriate. Modeling the kinematics of the VOD stars with a mixture of 2 multi-variate Gaussians, we show that two thirds of the VOD sample are the stars with  $\beta = 0.96$ , while the remainder has  $\beta = 0.44$ , in good agreement with the local measurement presented in Belokurov et al. (2018b).

As revealed by Gaia, the two structures are composed of stars on nearly radial orbits, with peaks in the eccentricity distribution at 0.9 (0.8) for the HAC (VOD). The distributions of the peri-centric and apo-centric distances also match: the stars in the Clouds turn around at 1-3and 15-30 kpc. Not only the HAC and the VOD look alike kinematically, their orbital composition is in perfect agreement with the stellar halo properties as analysed locally by Belokurov et al. (2018b) and globally (out to 40 kpc) by Deason et al. (2018b). As these authors demonstrate, the inner halo is dominated by metal-rich debris from an old and massive accretion event. In particular, Belokurov et al. (2018b) use Cosmological simulations of Milky Way halo formation, to bracket the time of the merger - between 8 and 11 Gyr ago - and its mass, which they show to be in excess of  $10^{10} M_{\odot}$ . The tell-tale sign of this dramatic head-on collision is the particular shape of the corresponding stellar velocity ellipsoid, which is stretched so much in the radial direction compared to the tangential ones, that it resembles a sausage. An alternative view of this merger can be found in Myeong et al. (2018a), where the local stellar halo is mapped out in the action space. Here, the metal-rich stars are shown have extended radial action distribution in addition to a prominent spray of material on retrograde orbits. The high mass of the progenitor is evidenced not only by the metallicity distribution of its likely member or the numerical simulations of halo formation, but also by a sizeable number of Globular Clusters that could be attributed to the same event (see Myeong et al. 2018b).

Johnston et al. (2012) attempted to connect various clouds together.

#### ACKNOWLEDGEMENTS

The research leading to these results has received funding from the European Research Council under the European Union's Seventh Framework Programme (FP/2007-2013) / ERC Grant Agreement n. 308024.

#### REFERENCES

Amorisco N. C., 2015, MNRAS, 450, 575

Belokurov V., Deason A. J., Koposov S. E., Catelan M., Erkal D., Drake A. J., Evans N. W., 2018a, MNRAS, 477, 1472

Belokurov V., Erkal D., Evans N. W., Koposov S. E., Deason A. J., 2018b, MNRAS, 478, 611

Belokurov V., Evans N. W., Bell E. F., Irwin M. J. et al., 2007, ApJ, 657, L89

Bergemann M., Sesar B., Cohen J. G., Serenelli A. M. et al., 2018, Nature, 555, 334

Binney J., Tremaine S., 2008, Galactic Dynamics: Second Edition. Princeton University Press

Bonaca A., Jurić M., Ivezić Ž., Bizyaev D. et al., 2012, AJ, 143, 105

Bovy J., 2015, ApJS, 216, 29

Bovy J., Hogg D. W., Roweis S. T., 2011, Annals of Applied Statistics. 5

Conroy C., Bonaca A., Naidu R. P., Eisenstein D. J., Johnson B. D., Dotter A., Finkbeiner D. P., 2018, ArXiv e-prints

de Boer T. J. L., Belokurov V., Koposov S. E., 2018, MNRAS, 473, 647

Deason A. J., Belokurov V., Evans N. W., 2011, MNRAS, 416, 2903

Deason A. J., Belokurov V., Evans N. W., Johnston K. V., 2013, ApJ, 763, 113

Deason A. J., Belokurov V., Hamren K. M., Koposov S. E. et al., 2014, MNRAS, 444, 3975

Deason A. J., Belokurov V., Koposov S. E., 2018a, MNRAS, 473,

Deason A. J., Belokurov V., Koposov S. E., Lancaster L., 2018b, ArXiv e-prints

Duffau S., Vivas A. K., Zinn R., Méndez R. A., Ruiz M. T., 2014, A&A, 566, A118

Duffau S., Zinn R., Vivas A. K., Carraro G., Méndez R. A., Winnick R., Gallart C., 2006, ApJ, 636, L97

Gaia Collaboration, Brown A. G. A., Vallenari A., Prusti T., de Bruijne J. H. J., Babusiaux C., Bailer-Jones C. A. L., 2018, ArXiv e-prints

Gaia Collaboration, Prusti T., de Bruijne J. H. J., Brown A. G. A. et al., 2016, A&A, 595, A1

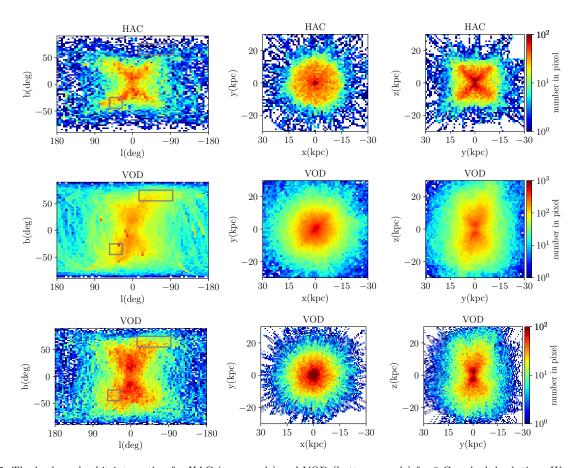
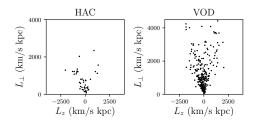


Figure 5. The backward orbit integration for HAC (top panels) and VOD (bottom panels) for 8 Gyrs look back time. We use  $M_{vir} = 0.8 \times 10^{12} M_{\odot}$ , the default galpy value . log(N) shown, notice the change in colour scale between top and bottom rows. The present day loci of HAC and VOD are marked with gray rectangles. The initial conditions of 44 stars with heliocentric distances between 15 and 18 kpc were used for the HAC backward orbit integration and of 299 stars with heliocentric distances between 4 and 75 kpc for the VOD orbit integration.



**Figure 6.** EXTRA: angular momentum - I can combine these two and also add another panel with energy vs total L when I figure out the energy.

Hayes C. R., Majewski S. R., Hasselquist S., Beaton R. L. et al., 2018, ApJ, 859, L8

Hendel D., Johnston K. V., 2015, MNRAS, 454, 2472

Hildebrand A. R., Penfield G. T., Kring D. A., Pilkington M., Camargo Z. A., Jacobsen S. B., Boynton W. V., 1991, Geology, 19, 867

Johnston K. V., 2016, in Astrophysics and Space Science Library, Vol. 420, Tidal Streams in the Local Group and Beyond, Newberg H. J., Carlin J. L., eds., p. 141

Johnston K. V., Sheffield A. A., Majewski S. R., Sharma S., Rocha-Pinto H. J., 2012, ApJ, 760, 95

Jurić M., Ivezić Ž., Brooks A., Lupton R. H. et al., 2008, ApJ, 673, 864 Li T. S., Balbinot E., Mondrik N., Marshall J. L. et al., 2016, ApJ, 817, 135

Majewski S. R., Ostheimer J. C., Rocha-Pinto H. J., Patterson R. J., Guhathakurta P., Reitzel D., 2004, ApJ, 615, 738

Morgan J. V., Gulick S. P. S., Bralower T., Chenot E. et al., 2016, Science, 354, 878

Myeong G. C., Evans N. W., Belokurov V., Sanders J. L., Koposov S. E., 2018a, ApJ, 856, L26

Myeong G. C., Evans N. W., Belokurov V., Sanders J. L., Koposov S. E., 2018b, ArXiv e-prints

Newberg H. J., Yanny B., Rockosi C., Grebel E. K. et al., 2002, ApJ, 569, 245

Nie J. D., Smith M. C., Belokurov V., Fan X. H. et al., 2015, ApJ,

Price-Whelan A. M., Johnston K. V., Sheffield A. A., Laporte C. F. P., Sesar B., 2015, MNRAS, 452, 676

Rocha-Pinto H. J., Majewski S. R., Skrutskie M. F., Crane J. D., Patterson R. J., 2004, ApJ, 615, 732

Sanderson R. E., Helmi A., 2013, MNRAS, 435, 378

Sesar B., Ivezić Ž., Lupton R. H., Jurić M. et al., 2007, AJ, 134, 2236

Sharma S., Johnston K. V., Majewski S. R., Muñoz R. R., Carlberg J. K., Bullock J., 2010, ApJ, 722, 750

Simion I. T., Belokurov V., Irwin M., Koposov S. E., 2014, MN-RAS, 440, 161

Simion I. T., Belokurov V., Koposov S. E., Sheffield A., Johnston K. V., 2018, MNRAS, 476, 3913

# 6 Iulia T. Simion et al

- Vanderplas J., Connolly A., Ivezić Ž., Gray A., 2012, in Conference on Intelligent Data Understanding (CIDU), pp. 47 –54
- Vivas A. K., Zinn R., 2006, AJ, 132, 714
- Vivas A. K., Zinn R., Andrews P., Bailyn C. et al., 2001, ApJ,  $554,\,\mathrm{L}33$
- Vivas A. K., Zinn R., Farmer J., Duffau S., Ping Y., 2016, ApJ, 831, 165
- Watkins L. L., Evans N. W., Belokurov V., Smith M. C. et al., 2009, MNRAS, 398, 1757

This paper has been typeset from a TEX/IATEX file prepared by