

Reminders for NAPDE

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Reminders on calculus

$$\begin{aligned}\int_{\Omega} -\Delta uv &= \int_{\Omega} \nabla u \cdot \nabla v - \underbrace{\int_{\Gamma_D} \nabla u \cdot \mathbf{n} v}_{=0 \text{ if } v|_{\Gamma_D}=0} \\ \int_{\Omega} \operatorname{div} u &= \int_{\partial\Omega} u \cdot \mathbf{n} \\ - \int_{\Omega} \mathbf{v} \cdot \nabla p &= \int_{\Omega} p \operatorname{div} \mathbf{v} + \int_{\partial\Omega} \mathbf{v} \cdot \mathbf{n} p \\ \int_{\Omega} \frac{\partial}{\partial x} u^2 &= \frac{1}{2} \int_{\partial\Omega} u \cdot \mathbf{n}\end{aligned}$$

Lifting Operators

If u on $\partial\Omega$ is non-null, we need to solve this problem, otherwise we cannot use test functions that vanish at the boundary. To do so, given $u = g$ on $\partial\Omega$ we use a lifting operator $Rg \in H^1(\Omega) : Rg|_{\partial\Omega} = g$, and modify our solution such that $u = \overset{\circ}{u} + Rg$, so the function $\overset{\circ}{u}$ has the properties we need. We then look for bilinear formulation such as $a(\overset{\circ}{u}, v)$ and add to the right-hand side $-a(Rg, v)$.

Weak Formulations

Elliptic equations

$$\begin{aligned}\begin{cases} -\operatorname{div}(\mu \nabla u) + \mathbf{b} \cdot \nabla u + \sigma u = f & \text{in } \Omega & g \in L^2(\Gamma_N) \\ u = 0 & \text{on } \Gamma_D & \partial\Omega = \Gamma_D \cup \Gamma_N \\ \mu \nabla u \cdot \mathbf{n} = g & \text{on } \Gamma_N & \Gamma_D^{\circ} \cap \Gamma_N^{\circ} = \emptyset \end{cases} \\ \Downarrow \\ \underbrace{\int_{\Omega} \mu \nabla u \cdot \nabla v + \int_{\Omega} \mathbf{b} \cdot \nabla uv + \int_{\Omega} \sigma uv}_{=: a(u, v)} = \int_{\Omega} f v + \underbrace{\int_{\Gamma_D} \mu \nabla u \cdot \mathbf{n} v}_{=0 \text{ if } v|_{\Gamma_D}=0} + \underbrace{\int_{\Gamma_N} \mu \nabla u \cdot \mathbf{n} v}_{=: g} \\ \Downarrow \\ \begin{cases} \text{find } u \in V & V = \{v \in H^1(\Omega), v|_{\Gamma_D} = 0\} =: H_{\Gamma_D}^1(\Omega) \\ a(u, v) = \langle F, v \rangle & \forall v \in V \end{cases}\end{aligned}$$

Parabolic equations

$$\begin{cases} \frac{\partial u}{\partial t} - \nu \frac{\partial^2 u}{\partial x^2} = f & 0 < x < d, t > 0 \\ u(x, 0) = u_0(x) & 0 < x < d \\ u(0, t) = u(d, t) = 0 & t > 0 \end{cases}$$

\Downarrow

$$\int_{\Omega} \frac{\partial u(t)}{\partial t} v \, d\Omega + a(u(t), v) = \int_{\Omega} f(t) v \, d\Omega \quad \forall v \in V$$

\Downarrow

for each $t > 0$, we need to find $u_h(t) \in V_h$ s.t.

$$\int_{\Omega} \frac{\partial u_h(t)}{\partial t} v_h \, d\Omega + a(u_h(t), v_h) = \int_{\Omega} f(t) v_h \, d\Omega \quad \forall v_h \in V_h$$

Discontinuous Galerkin

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases}$$

\Downarrow

$$\sum_{\mathcal{K} \in \mathcal{T}_h} \int_{\mathcal{K}} \nabla u \cdot \nabla v - \sum_{\mathcal{K} \in \mathcal{T}_h} \int_{\partial\Omega \mathcal{K}} \nabla u \cdot \mathbf{n}_{\mathcal{K}} v = \int_{\Omega} f v$$

\Downarrow

$$\sum_{\mathcal{K} \in \mathcal{T}_h} \int_{\partial\Omega \mathcal{K}} u \cdot \mathbf{n}_{\mathcal{K}} v = \sum_{F \in \mathcal{F}} \int_F \llbracket u \rrbracket \cdot \llbracket v \rrbracket + \sum_{F \in \mathcal{F}'_h} \int_F \llbracket u \rrbracket \llbracket v \rrbracket \quad \text{Magic formula}$$

\Downarrow

$$\sum_{\mathcal{K} \in \mathcal{T}_h} \int_{\mathcal{K}} \nabla u \nabla v - \sum_{F \in \mathcal{F}_h} \int_F \llbracket \nabla u \rrbracket \cdot \llbracket v \rrbracket - \theta \sum_{F \in \mathcal{F}_h} \int_F \llbracket u \rrbracket \cdot \llbracket \nabla_h v \rrbracket + \sum_{F \in \mathcal{F}_h} \int_F \gamma \llbracket u \rrbracket \cdot \llbracket v \rrbracket = \int_{\Omega} f v$$

Numerical formulation

Continuous Galerkin

Space $V_h = \{v_h \in \mathcal{C}^0(\overline{\Omega}) : v_h|_{\mathcal{K}} \in \mathbb{P}^r(\mathcal{K}) \, \forall \mathcal{K} \in \mathcal{T}_h, v_h|_{\Gamma_D} = 0\}$

$$\text{find } u_h \in V_h : a(u_h, v_h) = F(v_h) \quad \forall v_h \in V_h$$

Discontinuous Galerkin

Space $V_h^p = \{v_h \in L^2(\Omega) : v_h|_{\mathcal{K}} \in \mathbb{P}^{p_{\mathcal{K}}}(\mathcal{K}) \, \forall \mathcal{K} \in \mathcal{T}_h\} \not\subset H_0^1(\Omega)$

$$\text{find } u_h \in V_h^p : \mathcal{A}(u_h, v_h) = \int_{\Omega} f v_h \quad \forall v_h \in V_h^p$$

where

$$\begin{aligned} \mathcal{A}(u, v) = & \sum_{\mathcal{K} \in \mathcal{T}_h} \int_{\mathcal{K}} \nabla u \nabla v - \sum_{F \in \mathcal{F}_h} \int_F \llbracket \nabla_h u \rrbracket \cdot \llbracket v \rrbracket \\ & - \theta \sum_{F \in \mathcal{F}_h} \int_F \llbracket u \rrbracket \cdot \llbracket \nabla_h v \rrbracket + \sum_{F \in \mathcal{F}_h} \int_F \gamma \llbracket u \rrbracket \cdot \llbracket v \rrbracket \end{aligned}$$

- $\theta = 1$ Symmetric interior penalty

- $\theta = -1$ Non-symmetric interior penalty
- $\theta = 0$ Incomplete interior penalty
- $\gamma = \alpha \frac{p^2}{h}$

In the case of Dirichlet B.C. $u = g$ on $\partial\Omega$ modify r.h.s. and introduce the set of boundary faces \mathcal{F}_h^B

$$\int_{\Omega} f v - \theta \sum_{F \in \mathcal{F}_h^B} \int_F g \nabla_h v \cdot n + \sum_{F \in \mathcal{F}_h^B} \int_F \gamma g v$$

In case of Neumann B.C. $\nabla u \cdot n = g$ on $\partial\Omega$ we introduce the set of interior faces \mathcal{F}_h'

$$\begin{aligned} \mathcal{A}(u, v) = & \sum_{\mathcal{K} \in \mathcal{T}_h} \int_{\mathcal{K}} \nabla u \nabla v - \sum_{F \in \mathcal{F}_h'} \int_F \{\nabla_h u\} \cdot [v] \\ & - \theta \sum_{F \in \mathcal{F}_h'} \int_F [u] \cdot \{\nabla_h v\} + \sum_{F \in \mathcal{F}_h'} \int_F \gamma [u] \cdot [v] \end{aligned}$$

and

$$\int_{\Omega} f v + \sum_{F \in \mathcal{F}_h^B} \int_F g v$$

0.1 Galerkin-NI

Space $V_N = \{v_N \in \mathbb{P}_N(\Omega) : v|_{\Gamma_D} = g\}$

$$\text{find } u_N \in V_N : a_N(u_N, v_N) = f_N(v_N) \quad \forall v_N \in V_N$$

SEM-NI

Space $X_{\delta} = \left\{ v_{\delta} \in \mathcal{C}^0(\Omega) : v|_{\mathcal{K}} = \hat{v}_{\delta} \circ \mathbf{F}_{\mathcal{K}}^{-1}, \text{ with } \hat{v}_{\delta} \in \mathbb{Q}_p(\hat{\mathcal{K}}) \quad \forall \mathcal{K} \in \mathcal{T}_h \right\}$

$$\text{find } u_{\delta} \in X_{\delta} : a_{\delta}(u_{\delta}, v_{\delta}) = F_{\delta}(v_{\delta}) \quad \forall v_{\delta} \in X_{\delta}$$

Stability

Discontinuous Galerkin

Introduce broken Sobolev space

$$H^s(\mathcal{T}_h) = \{v \in L^2(\Omega) : v|_{\mathcal{K}} \in H^s(\mathcal{K}) \quad \forall \mathcal{K} \in \mathcal{T}_h\}$$

and the norms

$$\begin{aligned} \|v\|_{H^s(\mathcal{T}_h)}^2 &= \sum_{\mathcal{K} \in \mathcal{T}_h} \|v\|_{H^s(\mathcal{K})}^2 & \|v\|_{L^2(\mathcal{F}_h)}^2 &= \sum_{F \in \mathcal{F}} \|v\|_{L^2(F)}^2 \\ \|v\|_{DG}^2 &= \|\nabla_h v\|_{L^2(\Omega)}^2 + \|\sqrt{\gamma} [v]\|_{L^2(\mathcal{F}_h)}^2 & \|v\|_{DG}^2 &= \|v\|_{DG}^2 + \left\| \frac{1}{\sqrt{\gamma}} \{\nabla_h v\} \right\|_{L^2(\mathcal{F}_h)}^2 \end{aligned}$$

And some key properties to stability

- Continuity on $H^2(\mathcal{T}_h) \times V_h^p$

$$|\mathcal{A}(u, v_h)| \lesssim \|u\|_{DG} \|v_h\|_{DG} \quad \forall u \in H^2(\mathcal{T}_h), \forall v_h \in V_h^p$$

- Coercivity on $V_h^p \times V_h^p$

$$\mathcal{A}(v_h, v_h) \gtrsim \|v\|_{DG}^2 \quad \forall v_h \in V_h^p$$

- Strong-consistency (Galerkin orthogonality):

$$\mathcal{A}(u, v_h) = \int_{\Omega} f v_h \quad \forall v_h \in V_h^p \Rightarrow \mathcal{A}(u - u_h, v_h) = 0 \quad \forall v_h \in V_h^p$$

- Approximation. Let $\Pi_h^p u \in V_h^p$ be a suitable approximation of u , then

$$\|u - \Pi_h^p u\|_{DG} \lesssim \frac{h^{\min(p,s)}}{p^{s-\frac{1}{2}}} \|u\|_{H^{s+1}(\mathcal{T}_h)}$$

If $p \geq s$

$$\|u - \Pi_h^p u\|_{DG} \lesssim \left(\frac{h}{p}\right)^s p^{\frac{1}{2}} \|u\|_{H^{s+1}(\mathcal{T}_h)}$$

Convergence rates

Galerkin-NI

$$\|u - u_N^{\text{GNI}}\|_{H^1(\Omega)} \leq C(s) \left(\frac{1}{N}\right)^s \left(\|u\|_{H^{s+1}(\Omega)} + \|f\|_{H^s(\Omega)}\right)$$

Discontinuous Galerkin

General (if α large enough for SIP and NIP):

$$\|u - u_h\| \lesssim \frac{h^{\min(p,s)}}{p^{s-\frac{1}{2}}} \|u\|_{H^{s+1}(\mathcal{T}_h)}$$

L^2 norm (if Ω is a convex domain):

- SIP $\theta = 1$

$$\|u - u_h\|_{L^2(\Omega)} \lesssim \frac{h^{\min(p,s)+1}}{p^{s+\frac{1}{2}}} \|u\|_{H^{s+1}(\mathcal{T}_h)}$$

- NIP $\theta = -1$ and IIP $\theta = 0$

$$\|u - u_h\|_{L^2(\Omega)} \lesssim \frac{h^{\min(p,s)}}{p^{s-\frac{1}{2}}} \|u\|_{H^{s+1}(\mathcal{T}_h)}$$

Navier-Stokes

In case of inf-sup (LBB) condition satisfied by V and Q

$$\|\mathbf{u} - \mathbf{u}_h\|_V + \|p - p_h\|_Q \leq C(\alpha_h, \beta_h, \gamma, \delta) \left\{ \inf_{\mathbf{v}_h \in V_h} \|\mathbf{u} - \mathbf{v}_h\|_V + \inf_{q_h \in Q_h} \|p - q_h\|_Q \right\}$$

where

- α_h is the coercivity constant on the subspace V_h of divergence free velocities
- β_h is the LBB constant
- γ is the continuity constant of $a(\cdot, \cdot)$
- δ is the continuity constant of $b(\cdot, \cdot)$

In case of Taylor-Hoods elements

$$\|\mathbf{u} - \mathbf{u}_h\|_V + \|p - p_h\|_Q \leq Ch(\|\mathbf{u}\|_{H^{k+1}} + \|p\|_{H^k})$$

Code implementation

CG-FEM

- Matrix A ;

$$A_{ij} = \int_{\Omega} \nabla \varphi_j \nabla \varphi_i$$

Loop on all the elements and compute locally (elements with $\hat{\cdot}$ are computed on the reference element):

$$A_{locij} = \det(\mathbf{B}_{\mathcal{K}}) \int_{\hat{\mathcal{K}}} \hat{\nabla} \hat{\varphi}_j^T \mathbf{B}_{\mathcal{K}}^{-1} \mathbf{B}_{\mathcal{K}}^{-1} \hat{\nabla} \hat{\varphi}_i = \frac{\det(\mathbf{B})}{2} \hat{\nabla} \hat{\varphi}_j^T \mathbf{B}_{\mathcal{K}}^{-1} \mathbf{B}_{\mathcal{K}}^{-T} \hat{\nabla} \hat{\varphi}_i$$

Can be implemented as

```
function [K_loc]=C_lap_loc(Grad,w_2D,nln,BJ)
K_loc=zeros(nln,nln);
for i=1:nln
    for j=1:nln
        for k=1:length(w_2D)
            Binv = inv(BJ(:, :, k)); % inverse
            Jdet = det(BJ(:, :, k)); % determinant
            K_loc(i,j) = K_loc(i,j) + (Jdet.*w_2D(k)) .* ( (Grad(k, :, i)
                * Binv) * (Grad(k, :, j) * Binv )');
        end
    end
end
```

- Mass matrix M :

$$M_{ij} = \int_{\Omega} \varphi_j, \varphi_i$$

Loop on all the elements and calculate the local mass matrix

$$M_{locij} = \det(\mathbf{B}_{\mathcal{K}}) \int_{\hat{\mathcal{K}}} \hat{\varphi}_j^T \hat{\varphi}_i$$

Can be implemented as

```
function [M_loc]=C_mass_loc(dphiq,w_2D,nln,BJ)
M_loc=zeros(nln,nln);
for i=1:nln
    for j=1:nln
        for k=1:length(w_2D)
            Jdet = det(BJ(:, :, k)); % determinant
            M_loc(i,j) = M_loc(i,j) + (Jdet.*w_2D(k))
                .* dphiq(1,k,i) .* dphiq(1,k,j);
        end
    end
end
```

- Transport matrix T

Can be implemented as

```

function [ADV_loc]=C_adv_loc(Grad,dphiq,beta,w_2D,nln,BJ)
ADV_loc=sparse(nln,nln);
for i=1:nln
    for j=1:nln
        for k=1:length(w_2D)
            Binv=inv(BJ(:,:,k));    % inverse
            Jdet=det(BJ(:,:,k));    % determinant
            ADV_loc(i,j) = ADV_loc(i,j)+(Jdet.*w_2D(k)).* dphiq(1,k,i)
                                *( (beta)*(Grad(k,:,j) * Binv )');
        end
    end
end
end

```

- Right-hand side **b**:

$$b_i = \int_{\Omega} f \varphi_i$$

which is computed

```

function [f]=C_loc_rhs2D(force,dphiq,BJ,w_2D,pphys_2D,nln,mu)
f = zeros(nln,1);
x = pphys_2D(:,1);
y = pphys_2D(:,2);
F = eval(force);
for s = 1:nln
    for k = 1:length(w_2D)
        Jdet = det(BJ(:,:,k)); % determinant
        f(s) = f(s) + w_2D(k)*Jdet*F(k)*dphiq(1,k,s);
    end
end
end

```