

PASSIVE HIGHER-HARMONIC RF CAVITY SIMULATIONS FOR BUNCH LENGTHENING IN THE NSLSII-U

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Abstract

Higher-harmonic cavities (HHCs) are commonly employed in modern synchrotron light sources to lengthen electron bunches without increasing the natural energy spread. This approach effectively reduces intrabeam scattering and enhances Touschek lifetime. Accurate simulations are essential for predicting beam dynamics, particularly the impact of cavity detuning on beam stability and performance. Simulations are performed with the particle-tracking code ELEGANT, and key uniform-fill results are verified with the Vlasov–Fokker–Planck solver SPACE.

INTRODUCTION

Achieving extremely low electron beam emittance is critical for advancing the capabilities of modern synchrotron light sources. However, ultra-low emittance operation often results in high particle density, which implies shorter bunch lengths and enhances collective effects such as intra-beam scattering (IBS) and impedance-induced effects [1, 2]. This leads to emittance growth and reduced beam lifetime, particularly at higher beam currents. To mitigate these effects, passive higher-harmonic cavities (HHCs) are employed to lengthen the bunch while preserving the natural energy spread. This reduces the strength of IBS and provides Landau damping for suppressing longitudinal instabilities [3].

The NSLSII-U is a proposed upgrade to NSLS-II, designed to operate at 4 GeV with a complex-bend achromat lattice [4]. As part of its intensity and stability enhancement strategy, the design includes a passive superconducting third-harmonic cavity (3HC). Table 1 summarizes the primary design parameters. In passive mode, the 3HC is detuned from the third harmonic of the main RF frequency, enabling beam-induced voltage to flatten the total RF waveform and lengthen the bunch. Although a multi-harmonic RF system combining multiple cavity orders is under consideration for NSLSII-U to further increase bunch lengthening and improve beam stability [5], this study focuses on the impact of a single 3HC. Uniform-fill simulations were carried out with the particle-tracking code ELEGANT; key results were cross-validated against the Vlasov–Fokker–Planck solver SPACE to confirm consistency.

THEORETICAL MODEL

For passive HHC operation, the total RF voltage experienced by a particle at longitudinal coordinate τ is the combination of the main RF cavity voltage and the beam-induced

Table 1: NSLSII-U High-Brightness Storage Ring Parameters for Complex-Bend Lattice

Parameters	Symbol	Value	Unit
Circumference	C	791.72	m
Beam energy	E	4	GeV
Average current	I_{av}	400	mA
Momentum compaction	α	6.67	10^{-5}
Natural energy spread	σ_δ	1.10	10^{-3}
RF Voltage	V_{rf}	4.8	MV
RF frequency	f_{rf}	499.8	MHz
Harmonic number	h	1320	
Bunch length	$\sigma_{\tau 0}$	7.87	ps
Energy loss per turn	U_0	1940	keV
Main RF shunt impedance	$R_{s,main}$	33375	$M\Omega$
Main RF quality factor	Q_{main}	7.5×10^8	—
HHC shunt impedance	R_H	8800	$M\Omega$
HHC quality factor	Q_H	1.0×10^8	—
Detuning scan range	Δf	31–35	kHz

voltage in the harmonic cavity [3]:

$$V(\tau) = V_{rf} \sin(\omega_{rf}\tau + \phi_s) - i_m R_H \cos \psi_m \cos(m\omega_{rf}\tau + \psi_m), \quad (1)$$

where V_{rf} and ω_{rf} are the amplitude and angular frequency of the main RF cavity, respectively, ϕ_s is the synchronous phase, m is the harmonic number of the HHC, R_H is the shunt impedance of the HHC, and ψ_m is the cavity detuning angle. The beam-induced current component at the harmonic frequency $m\omega_{rf}$ for an approx. Gaussian bunch of rms length σ_τ is given by $i_m = 2I_0 \exp\left[-\frac{1}{2}(m\omega_{rf}\sigma_\tau)^2\right]$, where I_0 is the average stored beam current.

The detuning angle ψ_m is determined by the cavity's frequency offset from the exact harmonic multiple of the main RF frequency as:

$$\tan \psi_m = 2Q_H \frac{\omega_m - m\omega_{rf}}{m\omega_{rf}}, \quad (2)$$

with Q_H and ω_m being the quality factor and resonant angular frequency of the superconducting harmonic cavity, respectively.

For superconducting HHC operation in the NSLSII-U storage ring, the cavity shunt impedance R_H is significantly higher than the ideal value required for passive operation. This large impedance value forces the cavity to operate at a detuning angle ψ_m close to $\pi/2$. Under these conditions, the optimal operating points satisfy $V(0) = V'(0) = 0$, leading to the relationships:

$$\sin \phi_s = \sin \phi_{s0}, \quad \cos \psi_m = -\frac{V_{rf} \cos \phi_s}{i_m m R_H}, \quad (3)$$

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where ϕ_{s0} is the synchronous phase without the presence of the harmonic cavity.

To achieve maximum bunch lengthening while preserving beam stability, the superconducting harmonic cavity must be detuned by an optimal frequency offset $\Delta\omega_H = \omega_m - m\omega_{rf}$, given by:

$$\Delta\omega_H = -\frac{m^2\omega_{rf}i_mR_H}{2Q_HV_{rf}\cos\phi_s}. \quad (4)$$

Using the NSLSII-U storage ring parameters summarized in Table 1, this detuning frequency offset evaluates to approximately 35 kHz.

NUMERICAL SIMULATIONS

We begin by loading the NSLSII-U lattice into each code and defining the RF systems in beam-loaded mode. In **ELEGANT**, the main RF cavity is modeled with the built-in beam-loading **rfmode** element, and a second **rfmode** element is configured in passive mode to represent the 3HC. The shunt impedance of the harmonic cavity is gradually ramped up over the first 20,000 turns to suppress transient effects, after which it remains constant during detuning frequency scans. During each run, we track all bunches using a particle distribution file and periodically record the bunch length and centroid using histogram and watch elements.

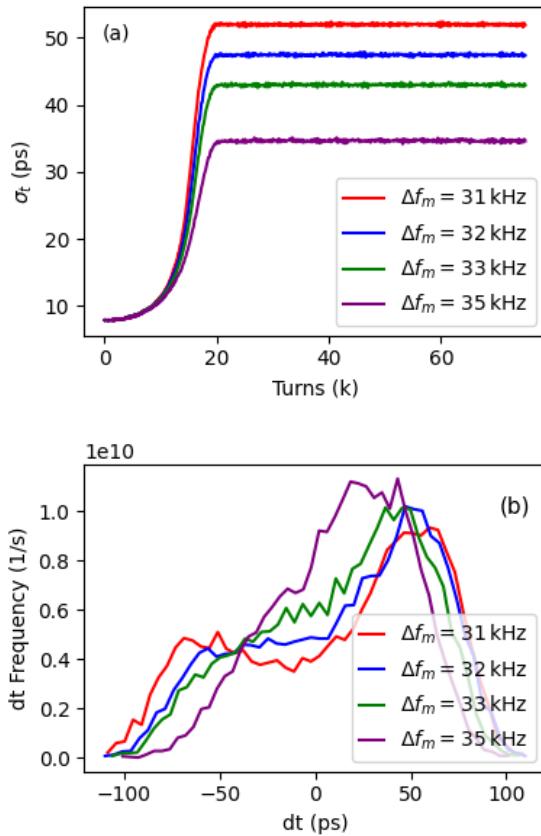


Figure 1: Simulated (a) rms bunch length and (b) bunch profiles for the uniformly filled NSLSII-U lattice with a passive 3HC, as computed by **ELEGANT**.

Figure 1 shows the simulated steady-state longitudinal beam parameters for the uniformly filled NSLSII-U lattice with a passive 3HC using **ELEGANT**. The initial rms bunch length without the harmonic cavity is $\sigma_{t,0} = 7.87$ ps. With a detuning frequency of $\Delta f = 35$ kHz, the steady-state bunch length is approximately 35 ps (bunch-lengthening factor $u! \approx 4.4$). The corresponding longitudinal density develops a broad, shallow top characteristic of a quartic RF potential, although a perfectly flat profile is not reached because the large shunt impedance of the superconducting cavity limits the achievable compensation of the sinusoidal curvature. At $\Delta f = 32$ kHz the distribution broadens further and a shallow double-bump appears, signalling an over-flattened RF potential that is still acceptable for uniform bunch lengthening.

In **SPACE**, we represent each of the 1320 bunches with 30k macroparticles and solve the Vlasov–Fokker–Planck equation self-consistently, including radiation damping, quantum excitation, and beam-induced voltages in both cavities. The simulations run until the cavities and beam reach equilibrium, after which we extract steady-state bunch lengths and centroids for comparison to the results obtained from **ELEGANT**.

Figure 2 presents a comparative analysis of steady-state bunch density distributions and rms bunch lengths as calculated by **SPACE** and **ELEGANT**. The density plots clearly illustrate uniform bunch lengthening across the entire bunch train. For the optimal 35 kHz detuning frequency, there is a small but noticeable difference of approximately 3 ps between the two codes. This discrepancy arises primarily from the distinct numerical methods: particle tracking with transient modulation in **ELEGANT** and self-consistent solution of the Vlasov–Fokker–Planck equation in **SPACE**. Despite this minor difference, both simulation approaches validate the effectiveness of passive 3HC operation for achieving significant bunch lengthening in NSLSII-U.

In normal NSLS-II operations, the storage ring is typically filled with 1200 bunches, leaving a gap of 120 RF buckets. To examine the impact of this non-uniform filling pattern, simulations were performed using **ELEGANT** for two different detuning frequencies of the harmonic cavity ($\Delta f_H = 32$ kHz and $\Delta f_H = 35$ kHz). As shown in Fig. 3, the bunch length (σ_t) exhibits a clear transient behavior along the bunch train. Bunches near the head of the train have shorter lengths, which gradually increase toward the tail, reaching a maximum before slightly decreasing near the end. A corresponding shift in the bunch centroid positions (Δc_t) is also observed, starting with positive values at the head and transitioning to negative values toward the tail. Such transient effects in bunch lengthening and centroid shifts due to filling gaps have been previously reported and are consistent with other studies [6]. These preliminary simulations highlight the importance of careful consideration of filling patterns in the operation of passive harmonic cavities. Further comparisons, including **SPACE** simulations, are underway to fully resolve observed discrepancies and ensure robustness of these observations.

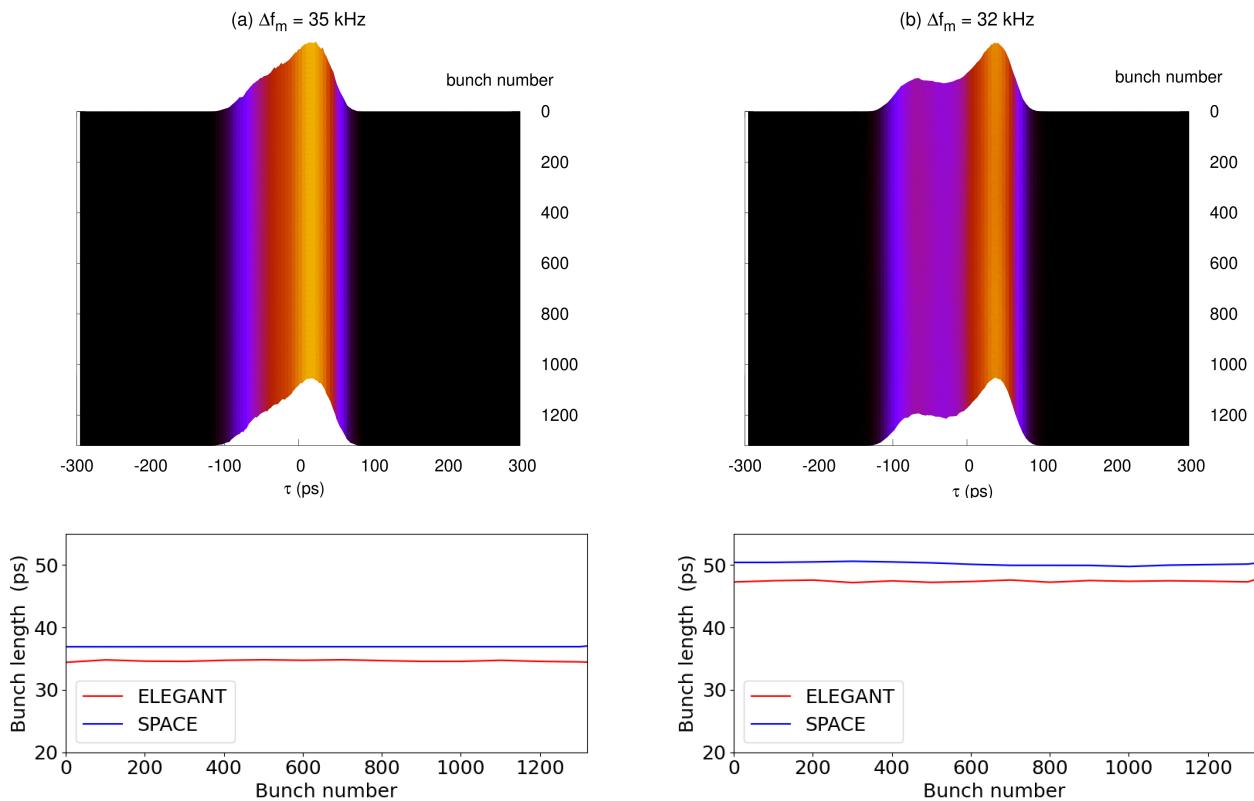


Figure 2: Comparison of steady-state longitudinal density distributions (a, b) computed with SPACE and bunch lengths (c, d) obtained from SPACE and ELEGANT simulations. Both simulations yield consistent results, with a small bunch length difference of approximately 3 ps.

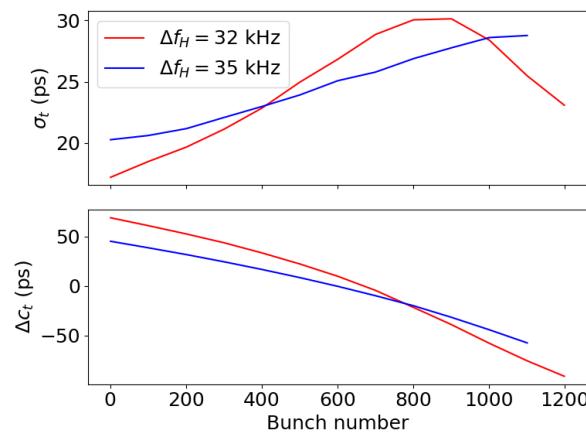


Figure 3: Simulated bunch length σ_t (top) and centroid shift Δc_t (bottom) along a 1200-bunch train with a 120-bucket gap, calculated using ELEGANT for two harmonic cavity detuning frequencies, $\Delta f_H = 32$ kHz (red) and $\Delta f_H = 35$ kHz (blue).

CONCLUSION

Passive third-harmonic cavities lengthen NSLSII-U bunches by a factor of 4–6 under uniform filling, but superconducting R/Q limits prevent the perfectly flat profile predicted for an ideal quartic potential. Introducing the 120-bucket ion gap used in routine operations reduces the attainable

lengthening from $u \approx 4.4$ to ≈ 3.5 at 35 kHz detuning. This sensitivity to filling pattern underscores the need for multi-harmonic RF schemes [7], which are now being evaluated and will be benchmarked with both ELEGANT and SPACE in future work.

ACKNOWLEDGMENT

This work has been supported by DOE under Contract No. DE-SC0012704.

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