

High Intensity Issues in Storage Rings

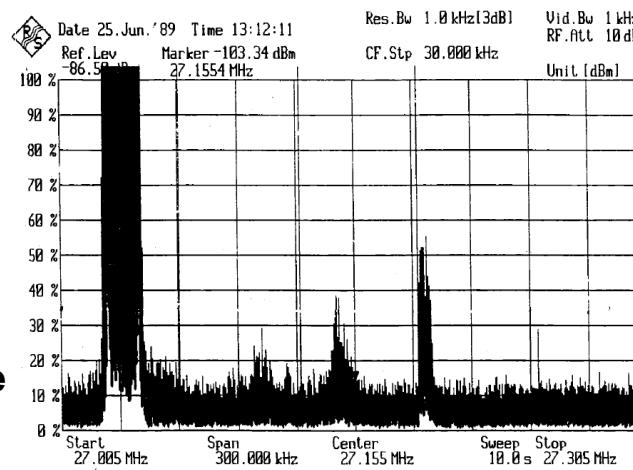
M. Steck
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Storage Ring Department

In storage rings intensity issues are mainly related to high brightness, rather than high intensity

Beam cooling plays an important role

- to prepare high quality beams for experiments**
- including internal targets**
- to support efficient operation**
- to achieve long beam lifetime (minutes to days)**

Test Storage Ring TSR, MPI Heidelberg



rf stacking
deceleration within the
momentum acceptance

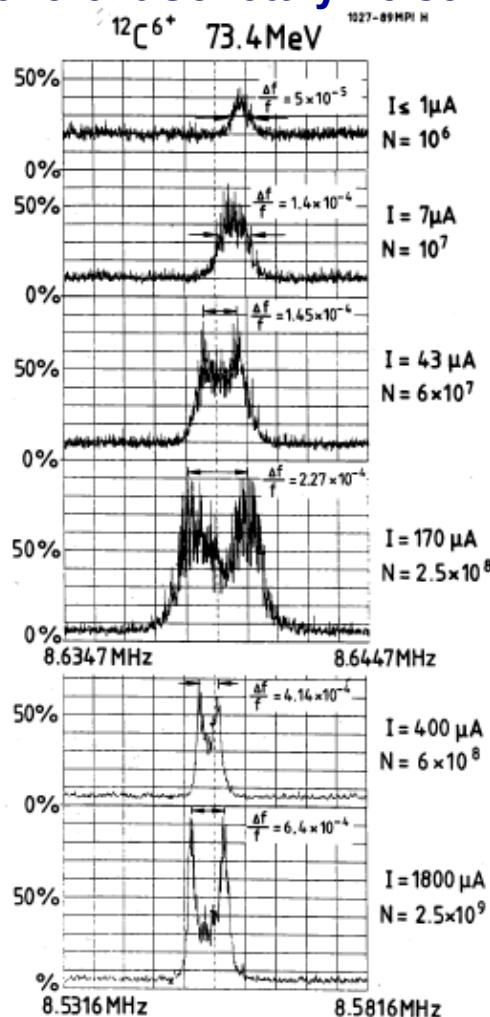
from: https://www.mpi-hd.mpg.de/blaum/storage-rings/tsr/tsr_intensities.de.html

Beam	Energy [MeV]	Intensity [μ A]
p	21	2400
HD ⁺	2	40
⁷ Li ⁺	13.4	23
⁹ Be ⁺	7.3	6
¹² C ⁶⁺	73.3	18000
³² S ¹⁶⁺	195	1600
³⁵ Cl ¹⁷⁺	293	1000
⁵⁶ Fe ¹⁸⁺	240	300
⁶³ Cu ²⁶⁺	510	110
⁸⁰ Se ²⁵⁺	480	110
¹⁹⁷ Au ⁵⁰⁺	695	3

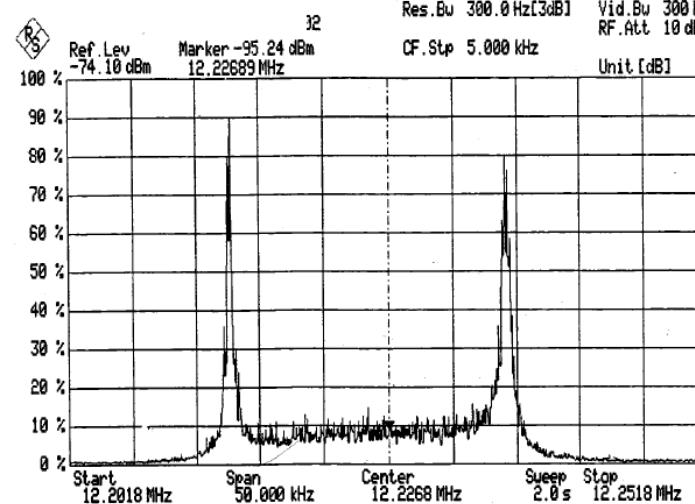
accumulated by multturn injection
and rf stacking
(injected current only 15 μ A)

high intensity cooled beam: achieved by intentional degradation of cooling

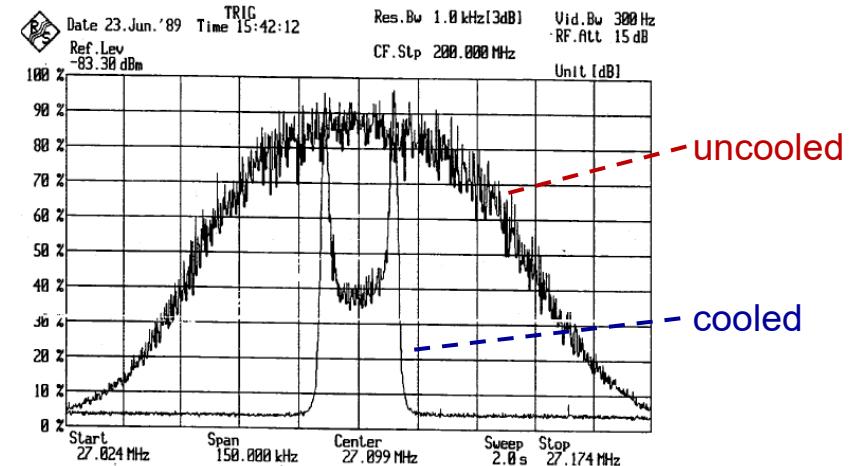
intensity dependence of coherent Schottky noise



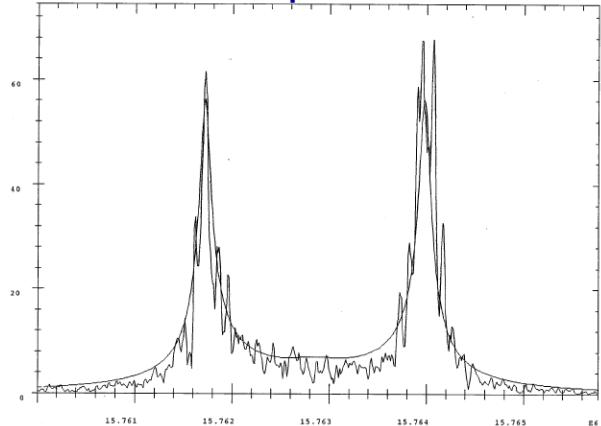
large coherent frequency at maximum intensity



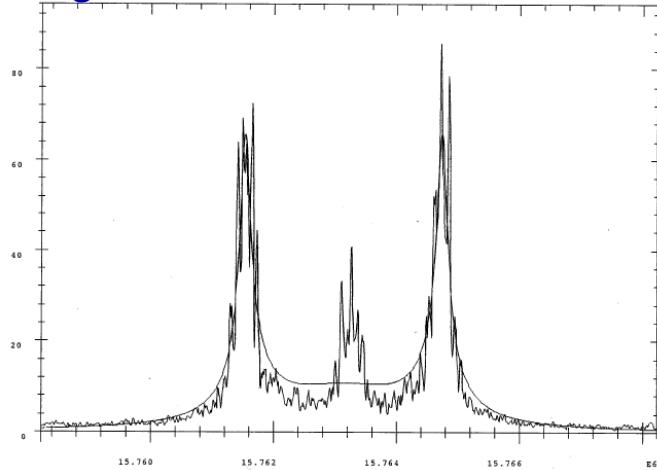
noise suppression of cooled beam signal



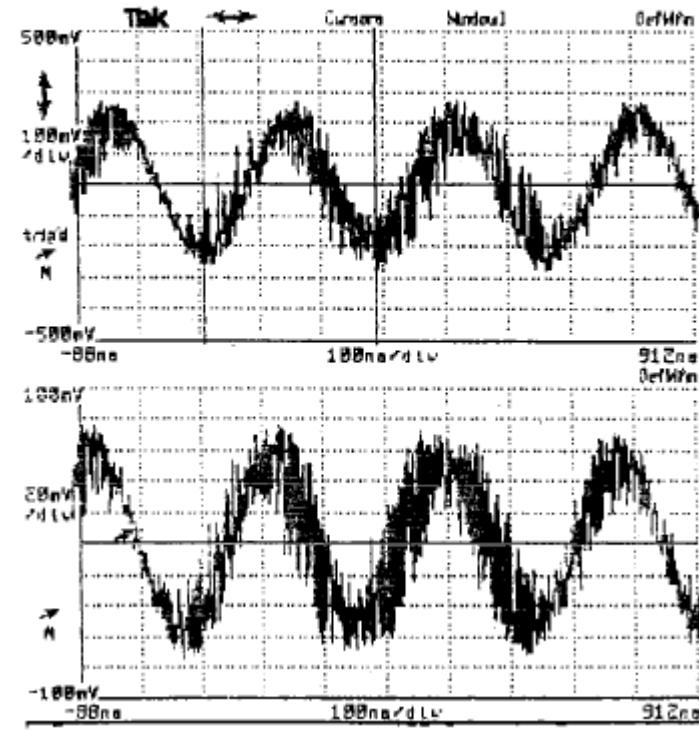
fitting curves of coherent signal
to extract beam parameters



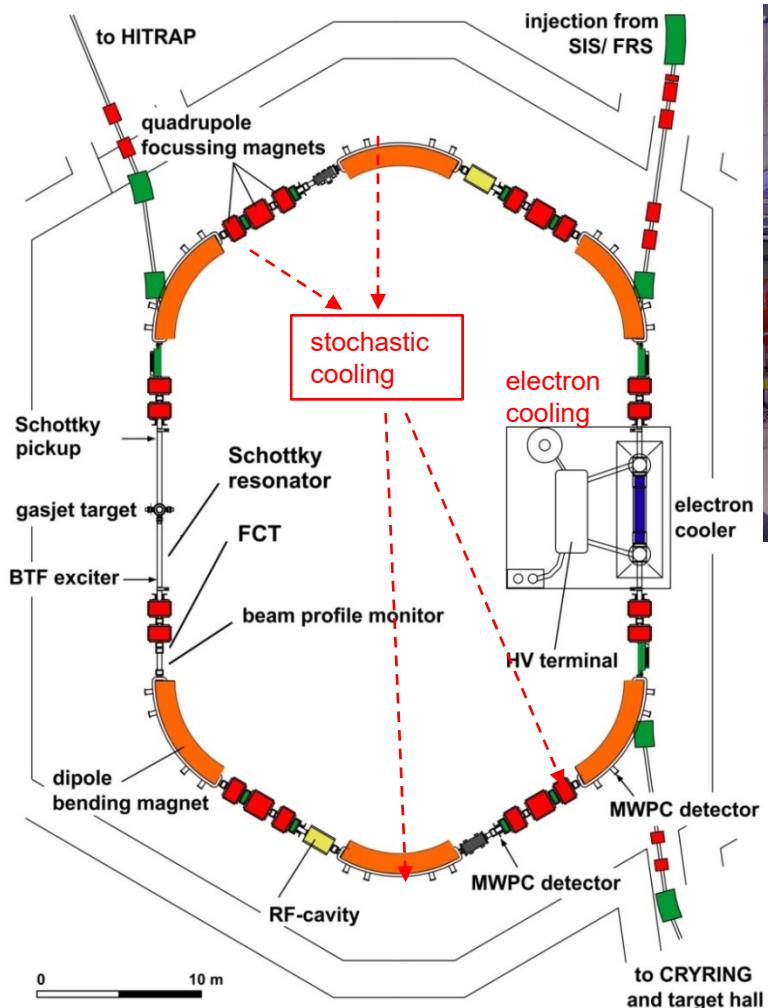
fitting fails for transverse oscillations



pick-up time domain signal revealing
transverse coherent oscillations



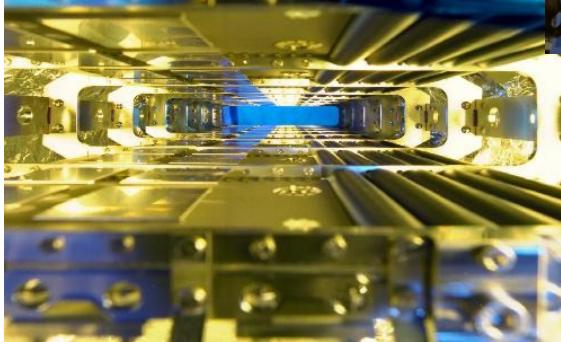
The Heavy Ion Storage Ring ESR



circumference 108.36 m
bending power 10 Tm



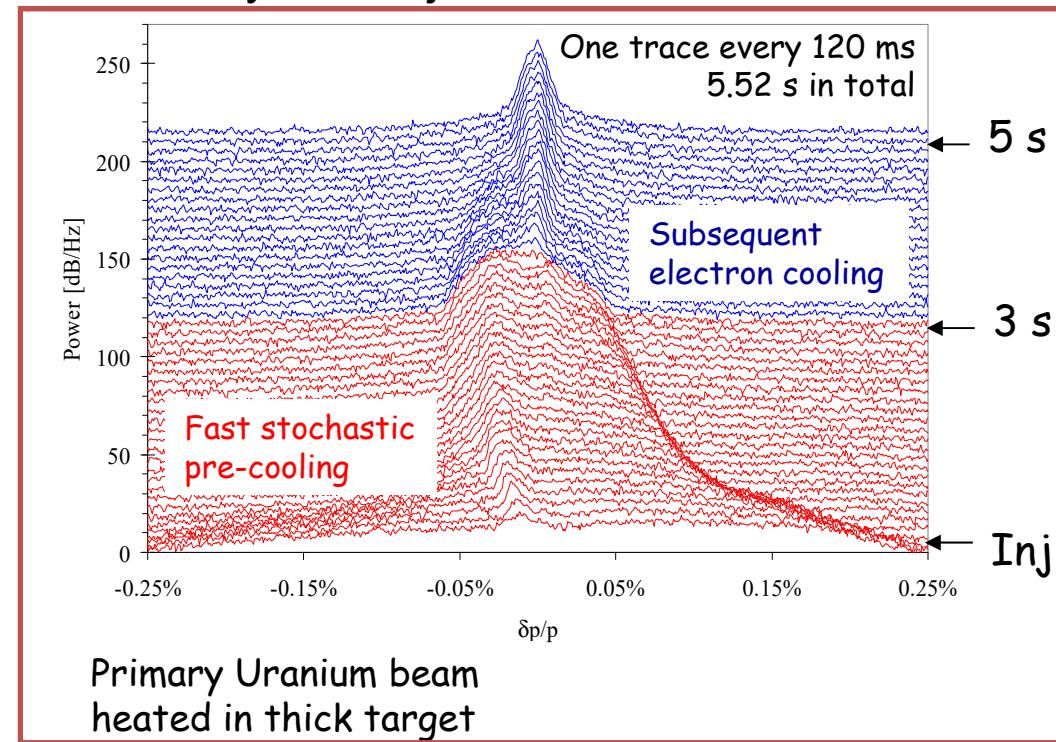
electron cooling
energy 3-430 MeV/u
electron current 0.001-1 A
magnetic field up to 0.2 T



stochastic cooling
energy ≥ 400 MeV/u
cooling acceptance:
 $\delta p/p \leq +/- 0.35\%$
 $\epsilon \leq 10 \pi \text{ mm mrad}$

Combination of Stochastic and Electron Cooling

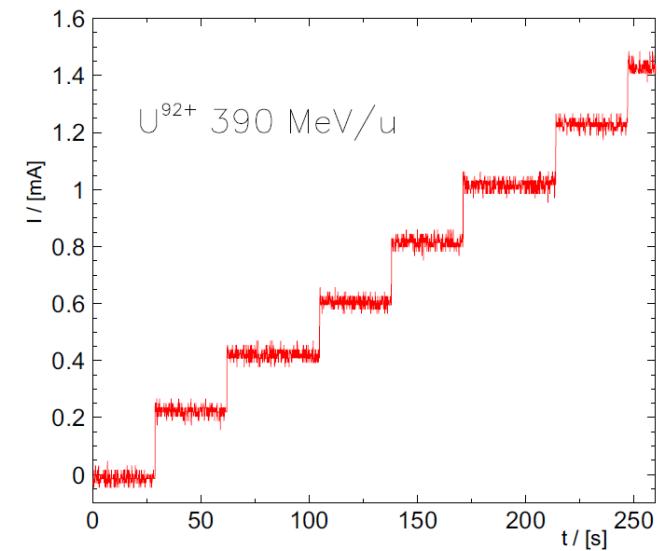
stochastic pre-cooling + final electron cooling
immediately after injection



Stochastic pre-cooling reduces the total cooling time to a few seconds, electron cooling only takes 10 - 60 s

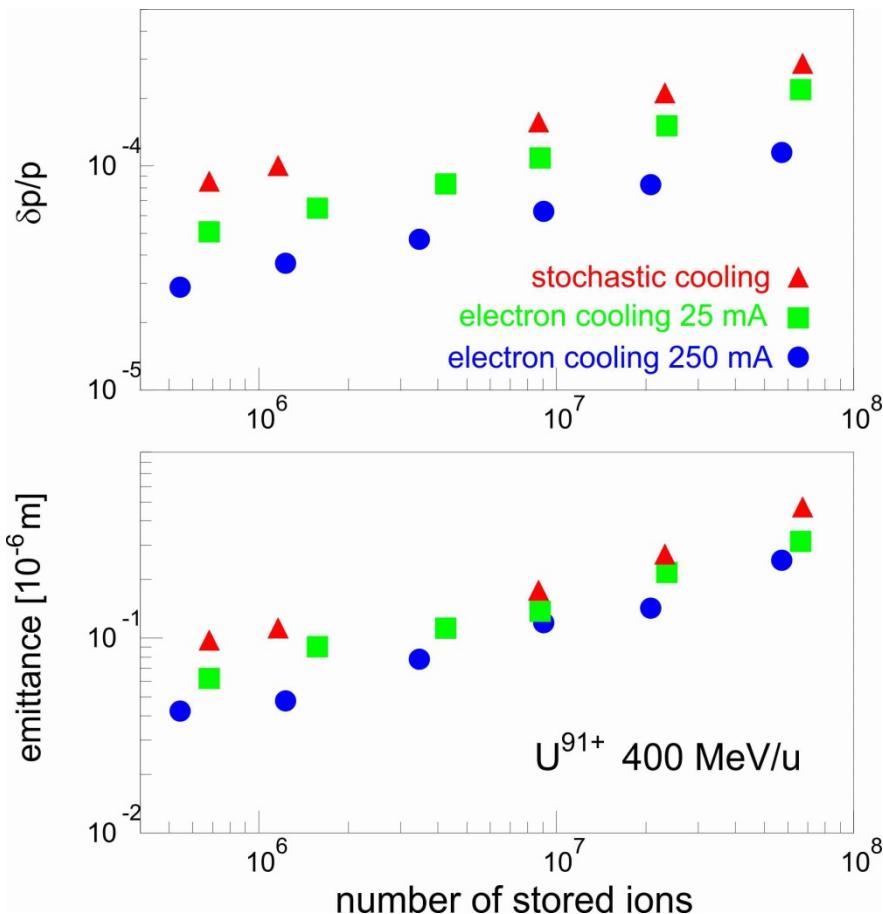
Accumulation of secondary beams

- 1) s.c. on injection orbit
- 2) rf stacking
- 3) electron cooling of stack



Intensity increase for secondary beams

Equilibrium Beam Parameters of Cooled Heavy Ion Beams in the ESR



beam quality is limited by intrabeam scattering

Electron cooling results in smaller momentum spread and smaller emittance compared to stochastic cooling.

The equilibrium is a balance between the cooling rate and the heating rate by intrabeam scattering.

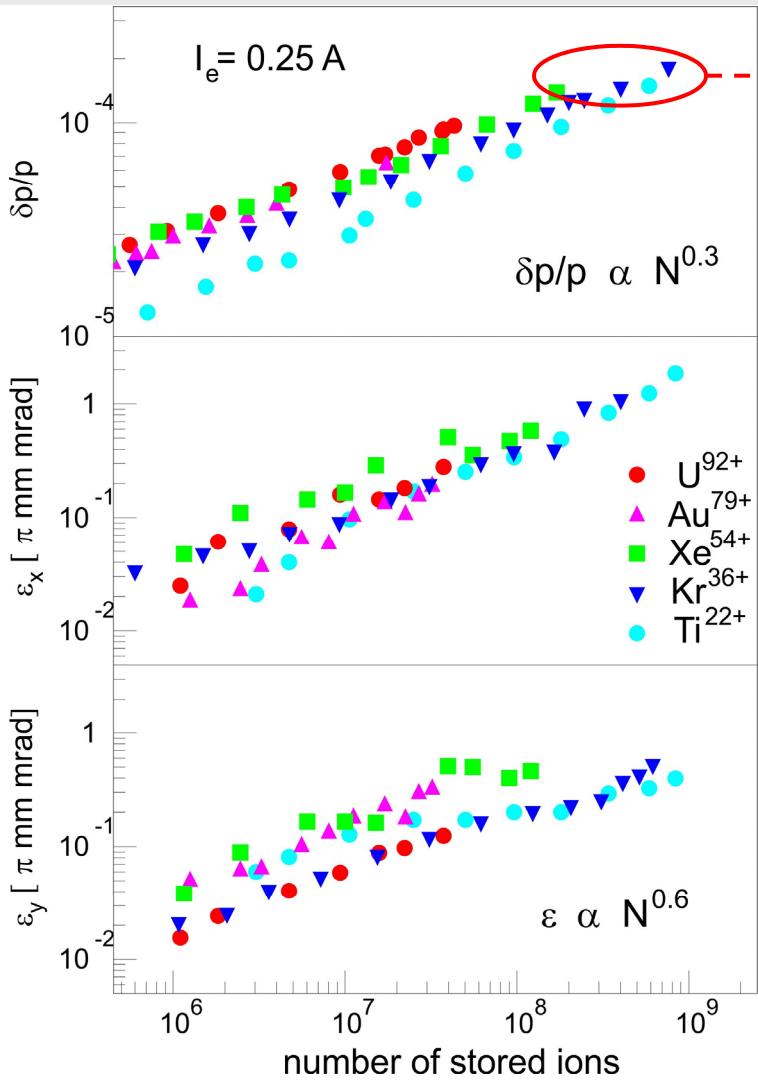
calculated IBS-heating/cooling rate [s⁻¹]

CERN code INTRABT

	longit.	transv.
stoch. cool. (Palmer)	0.9 - 2.2	0.5 - 1.3
elec. cool. [25 mA]	2.0 - 6.0	1.4 - 3.3
elec. cool. [250 mA]	18 - 58	7 - 10

→ Electron cooling is most powerful for cold beams.

stochastic cooling time increases $\propto N_i$



stable current of up to 10 mA could be achieved for lighter ions

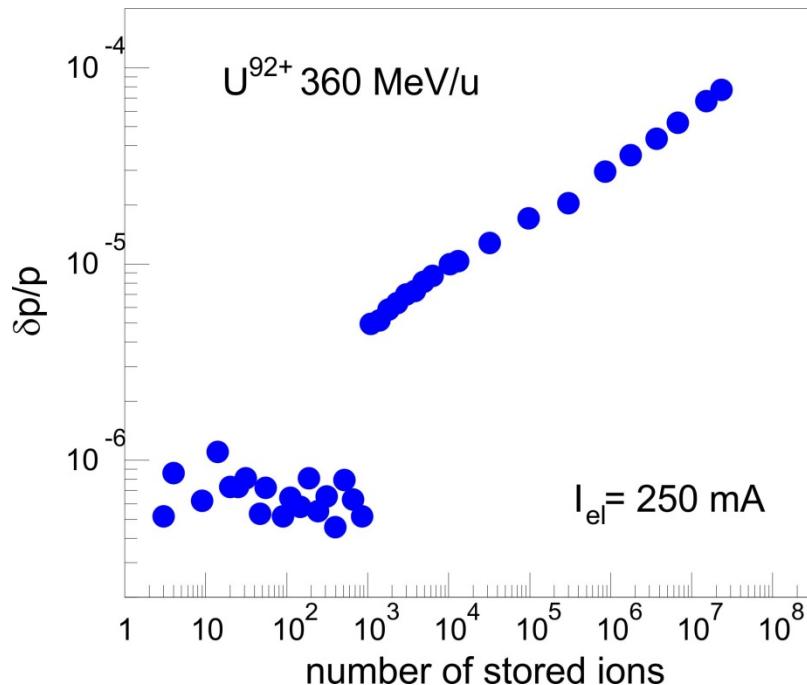
for heavier ions somewhat lower intensities were available

beam parameters governed by intrabeam scattering (IBS)

$$\tau_{\text{IBS}}^{-1} = \frac{q^4 e^4}{(A m_i)^2} \cdot \frac{N}{C \varepsilon_h \varepsilon_v \delta p / p} \cdot \frac{1}{(\gamma^4 \beta^3 c^3)} \cdot 4\pi L_{\text{c}}^{\text{IBS}}$$

in equilibrium with electron cooling

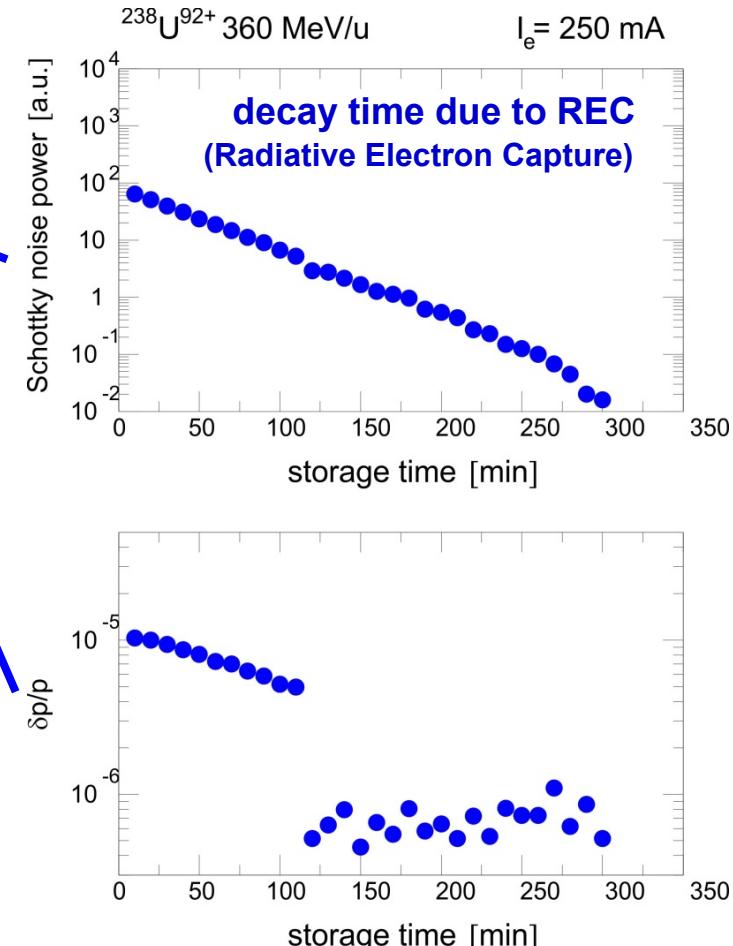
$$\tau_{\text{ec}}^{-1} \propto \frac{q^2}{A} \frac{n_e \eta}{\gamma^2}$$

reduction of momentum spread

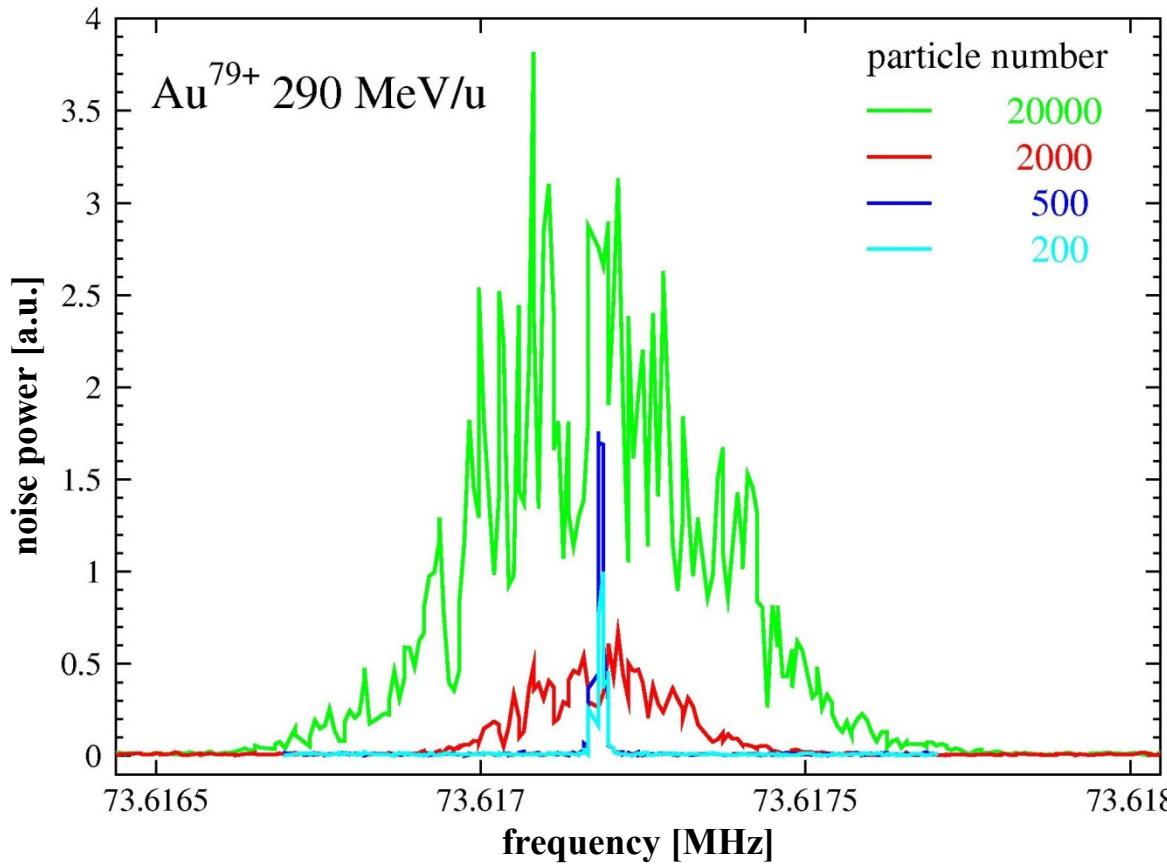
sudden reduction of the momentum spread
for less than about one thousand stored ions

⇒ linear ordering in ion string

temporal evolution of Schottky noise allows independent determination of particle number

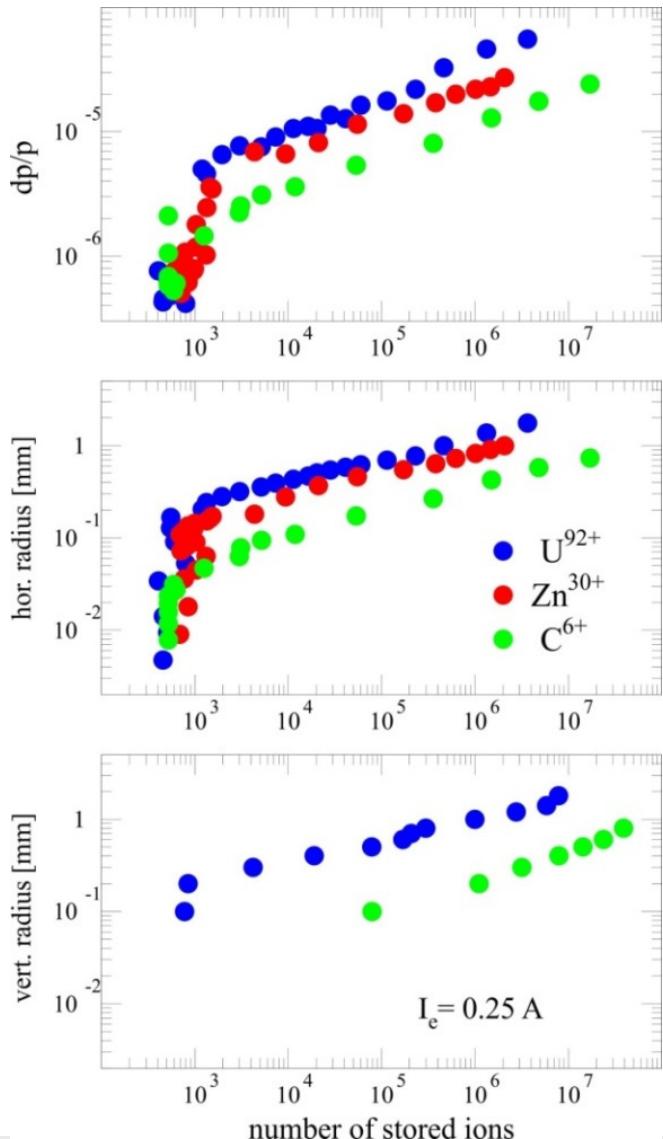


Schottky Noise Signal at Transition to Ordered Beam



The frequency (momentum) spread can be derived from the distribution.
The integrated noise power is proportional to the particle number.

Plasma Parameter



Plasma parameter $\Gamma = \frac{U_{\text{Coul}}}{k_B T} = \frac{q^2 e^2}{4\pi \epsilon_0 a k_B T}$

$$a_{\perp} = x_{\text{rms}}$$

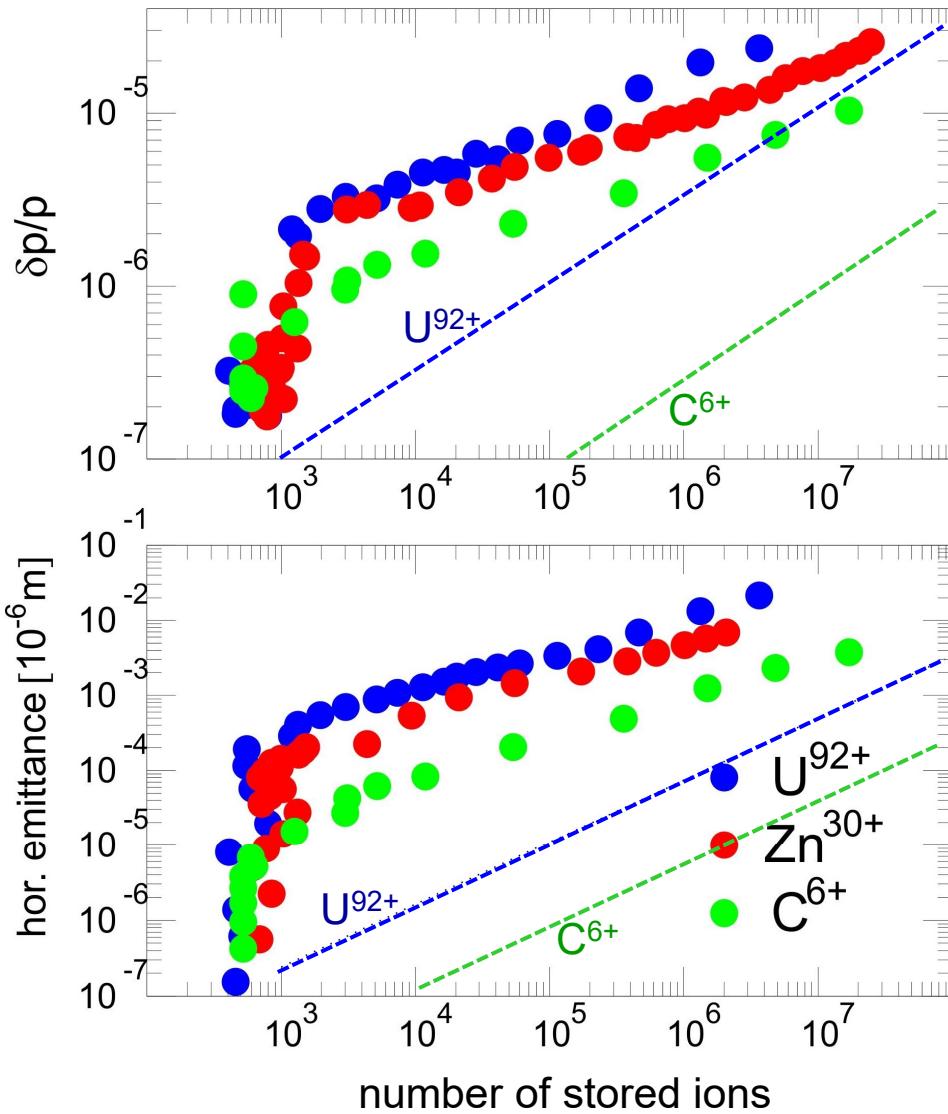
$$a_{\parallel} = \gamma C/N$$

typical values (meV) of beam parameters below transition:

	C ⁶⁺	Zn ³⁰⁺	U ⁹²⁺
$k_B T_x$	0.14	1.0	0.93
$k_B T_{\parallel}$	0.26	1.5	5.0
U_{\perp}	194	3960	84260
U_{\parallel}	0.002	0.1	0.63

$$\frac{1}{2} k_B T_{\perp} = \frac{1}{2} m v_{\perp}^2 = \frac{1}{2} m c^2 \beta^2 \gamma^2 \theta_{\perp}^2$$

$$\frac{1}{2} k_B T_{\parallel} = \frac{1}{2} m v_{\parallel}^2 = \frac{1}{2} m c^2 \beta^2 \left(\frac{\delta p_{\parallel}}{p} \right)^2$$



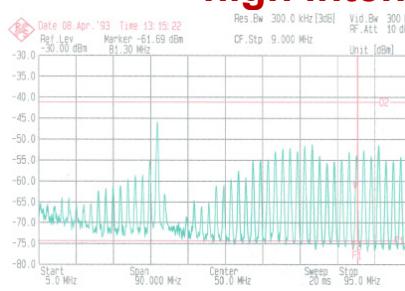
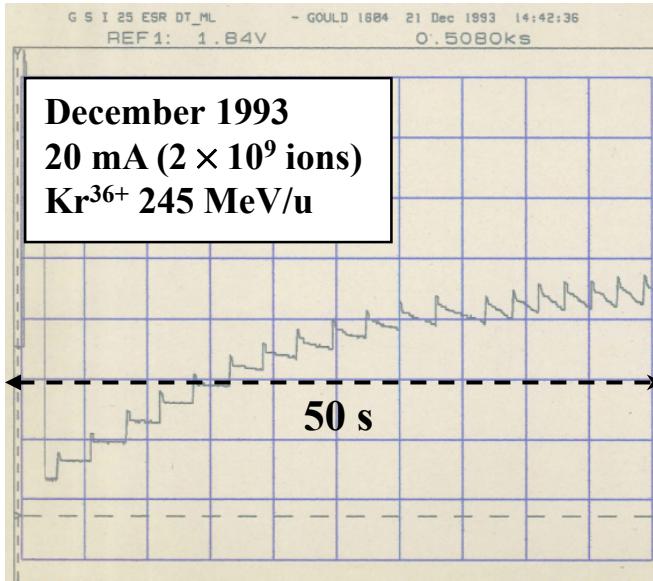
Keil-Schnell criterion

$$\left(\frac{\delta p}{p}\right)^2 \geq \frac{N q^2 e^2}{|\eta| \beta \gamma A m_0 c c} \frac{\pi}{\ln 2} \left| \frac{Z_{||}}{m} \right|$$

space charge limit ($\Delta Q = 0.01$)

$$\Delta Q = - \frac{r_p N i q^2}{2\pi A \beta^2 \gamma^3 \epsilon_x}$$

High Intensity Cooled ESR Beams



transverse coherent Schottky signal 0-90 MHz

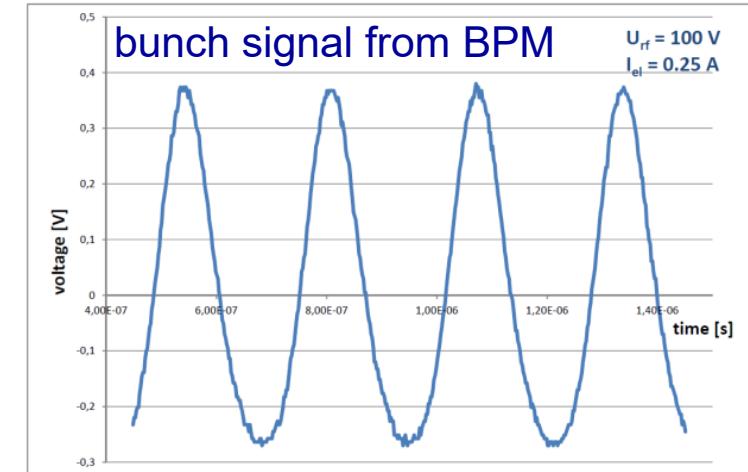


longitudinal Schottky signal with central coherent peak

observations and attempts to apply feedback were not conclusive

December 2005

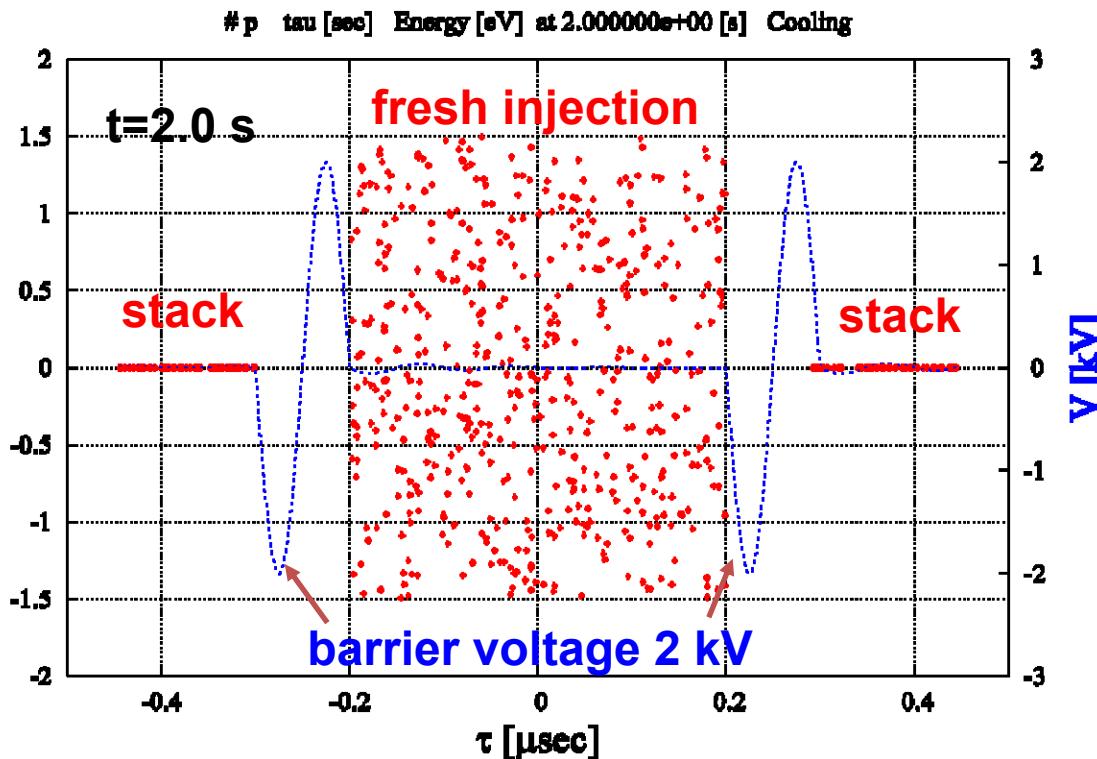
18 mA Xe⁵⁴⁺ 350 MeV/u bunched ($h=2$, $B=0.40$) corresponding to over 40 mA coasting beam



what matters?
cooling \Rightarrow emittance, momentum spread orbit, tune, chromaticity, impedances

basic idea: confine stored beam to a fraction of the circumference, inject into gap
apply strong electron cooling to merge the two beam components

⇒ fast increase of intensity (for low intensity RIBs)



proposed for the NESR
Storage Ring of FAIR

$^{132}\text{Sn}^{50+}$

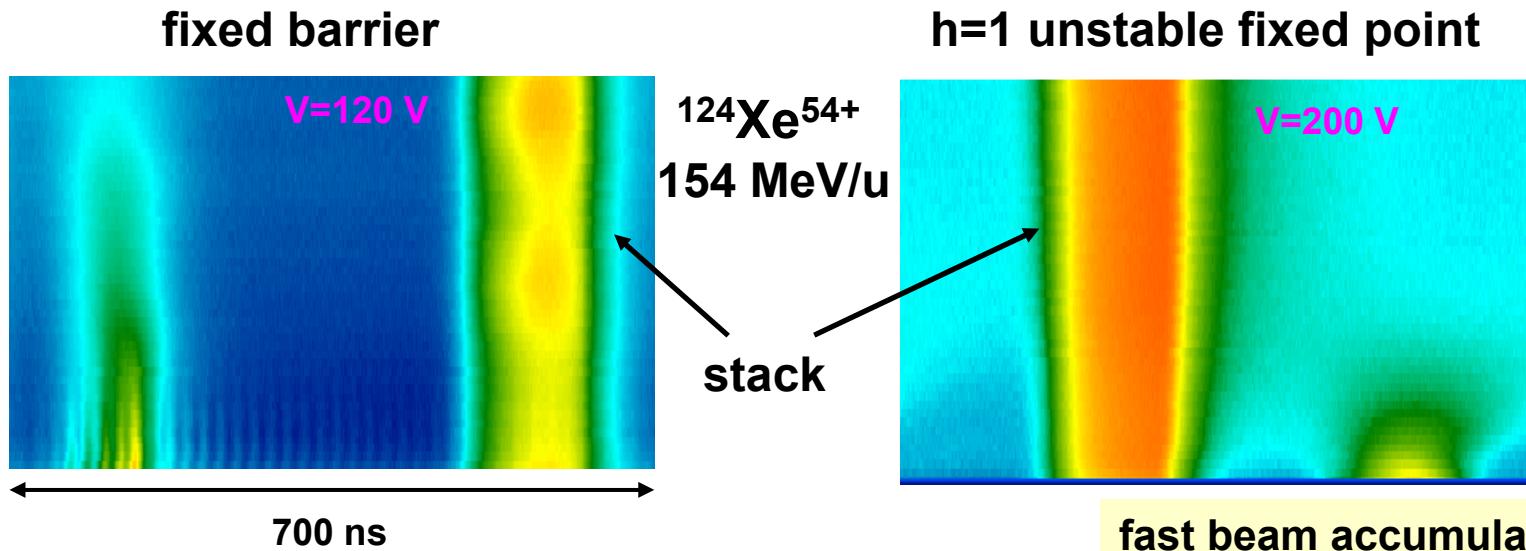
$E_k = 740 \text{ MeV/u}$

Longitudinal stacking
with Barrier Buckets

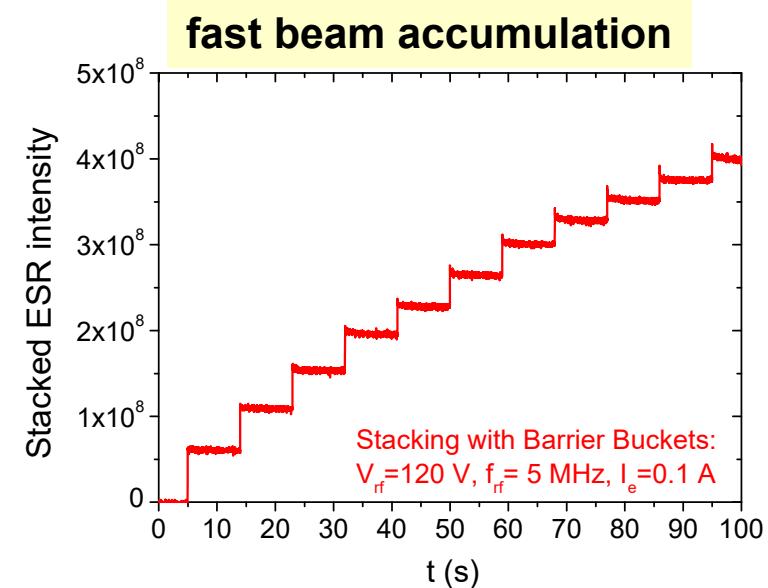
(simulation by T. Katayama)

revolution time 0.9 μs ⇒ rather short bunches of stored cooled beam

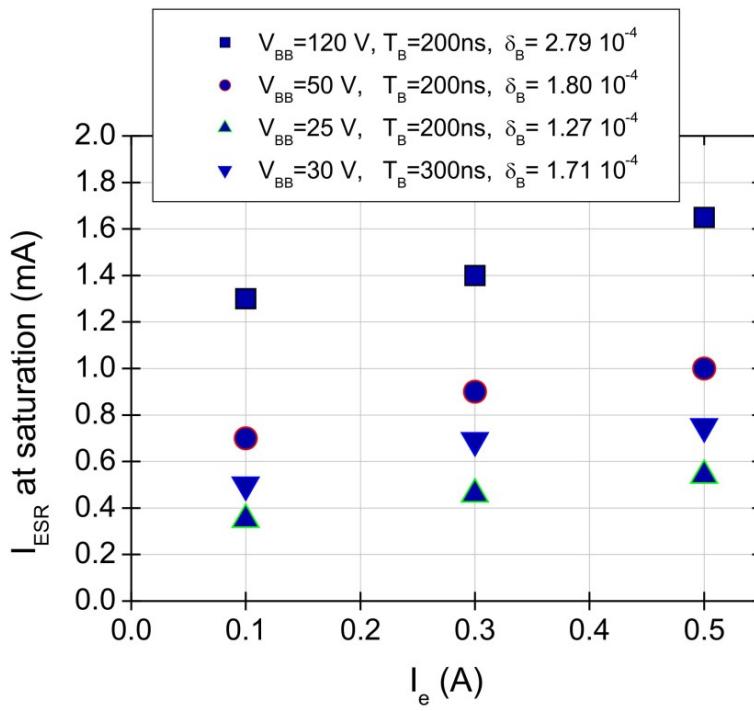
200 turns/frame



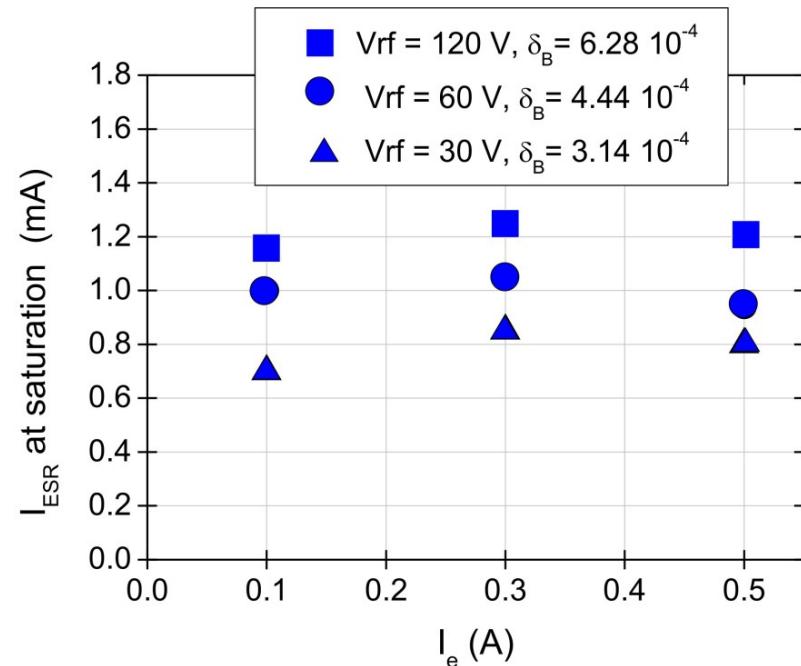
accumulation schemes work well:
cooling times close to expectations
efficient accumulation
high quality timing and kicker pulses required
Intensity limits: rf voltage and instabilities



Stacking with Barrier Buckets

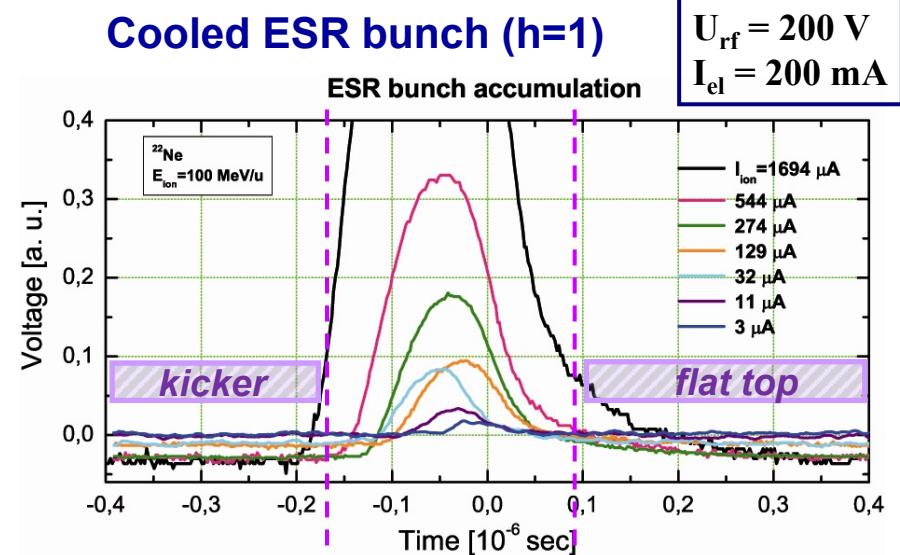
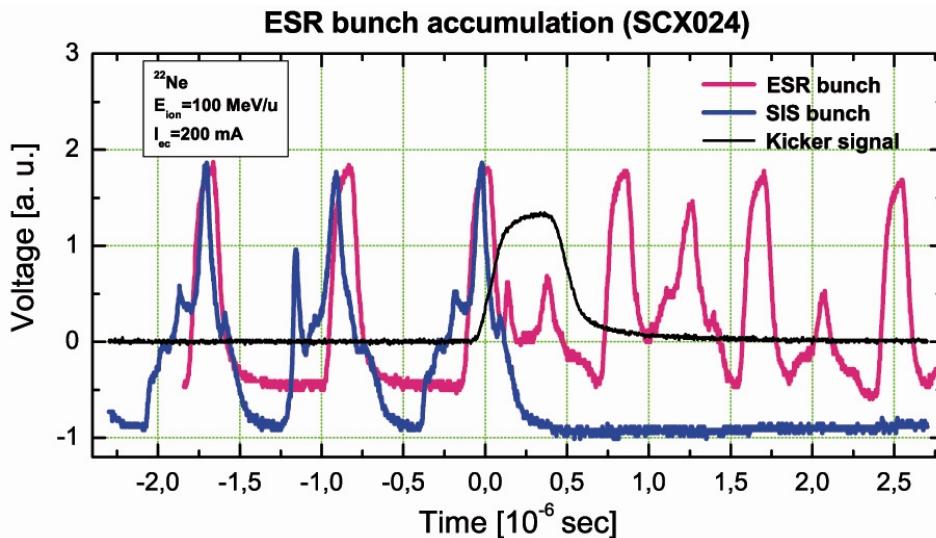


Stacking at Unstable Fixed Point $h=1$

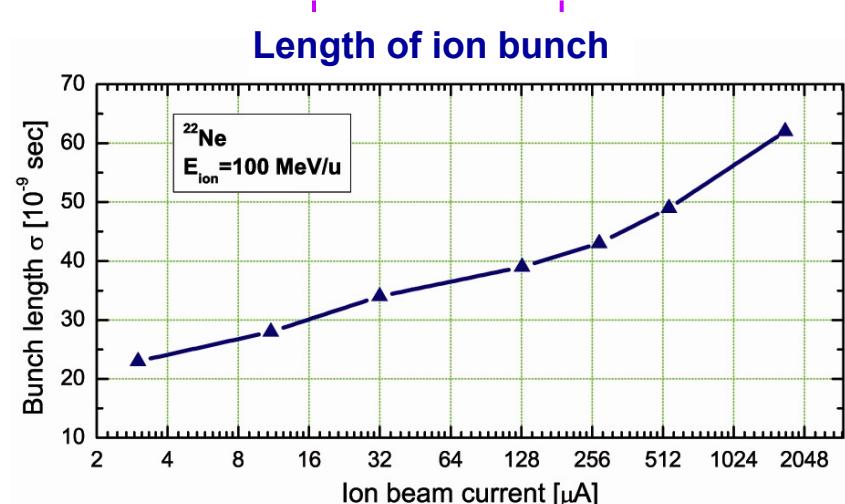
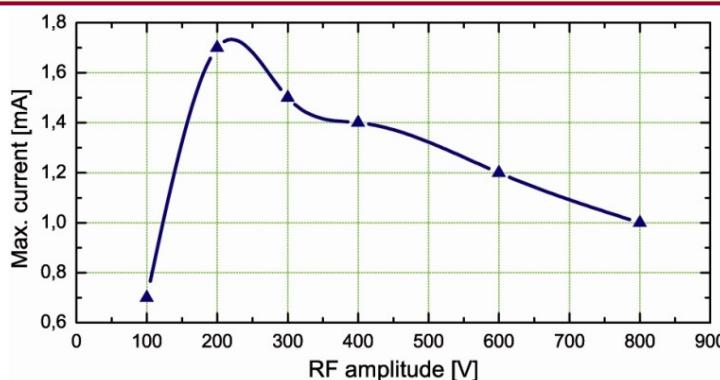


barrier height and cooling rate determine the intensity limit

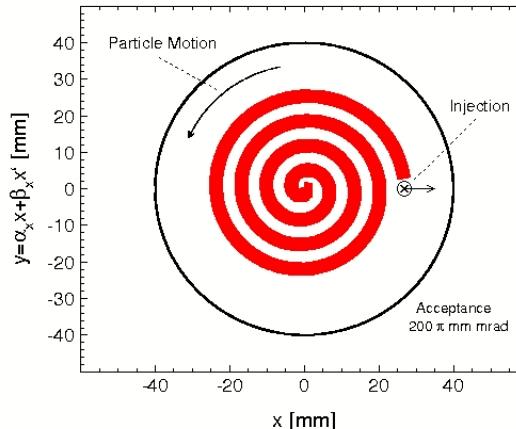
Accumulation of $^{22}\text{Ne}^{10+}$ at 100 MeV/u



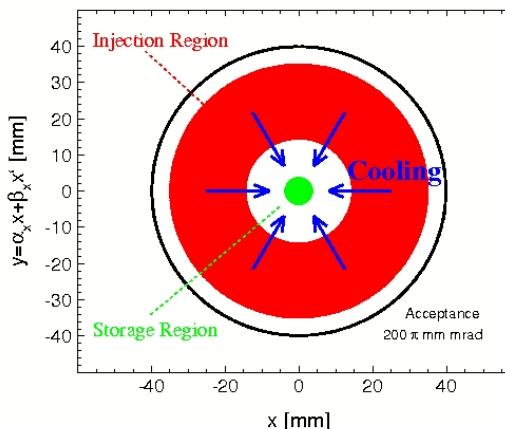
Intensity limit is a trade off between bunch length ($\leq 200 \text{ ns}$) and local ion current ($\leq 6-8 \text{ mA}$)



standard multiturn injection

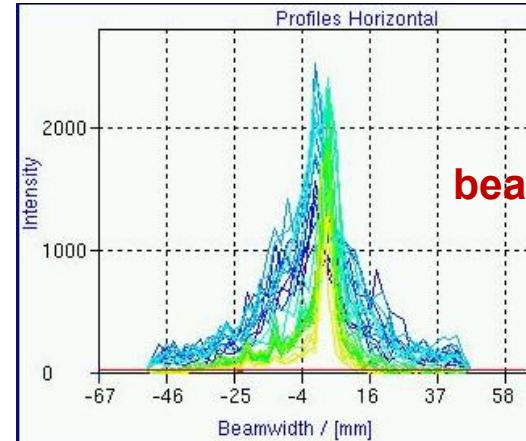


fast accumulation by multiple multturn injection with electron cooling

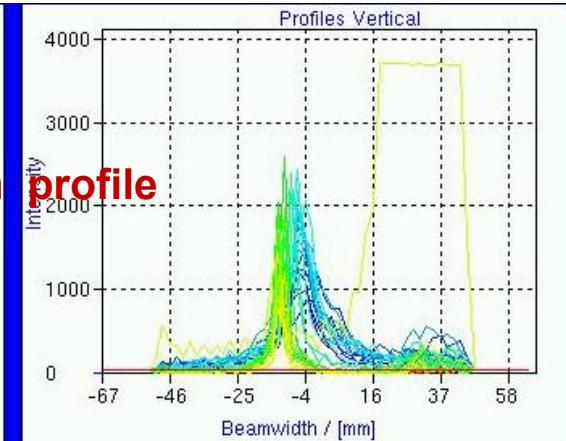


fast transverse cooling

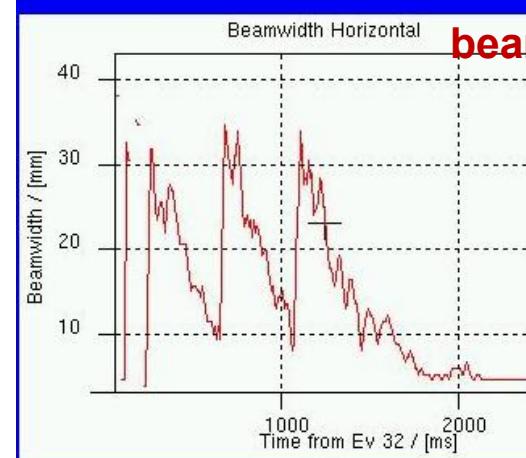
horizontal



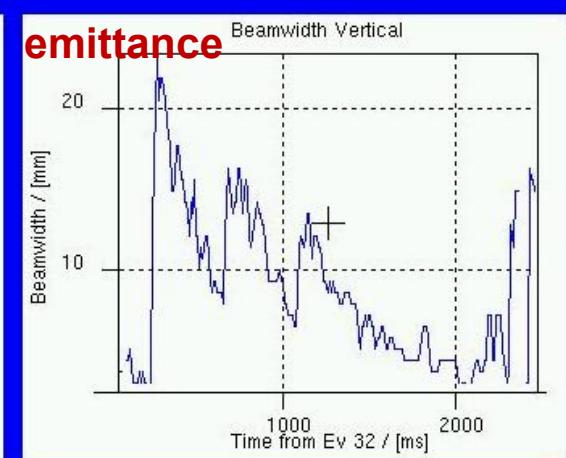
vertical



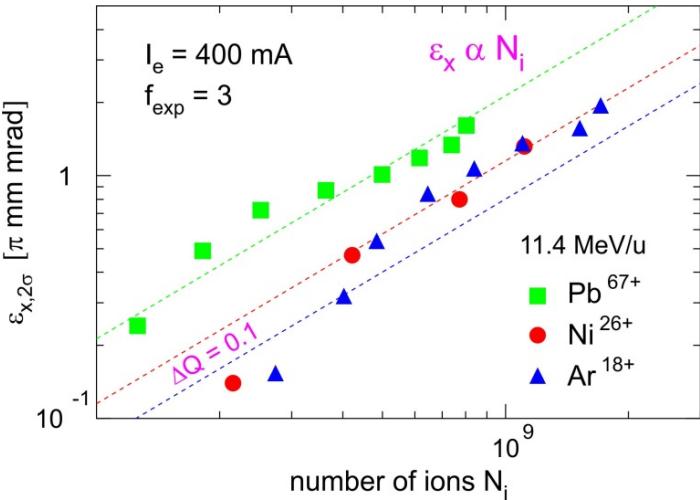
beam emittance



beam emittance

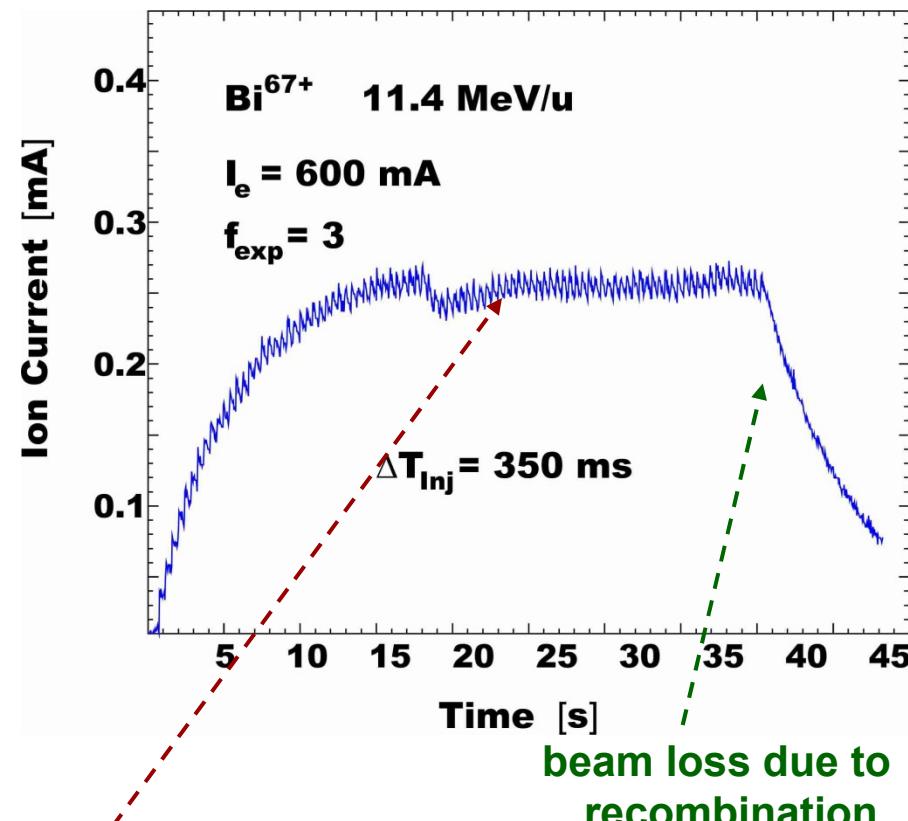
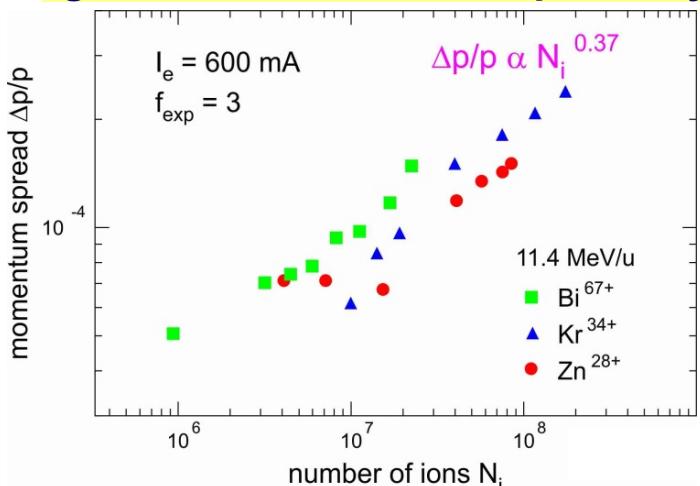


emittance growth by space charge



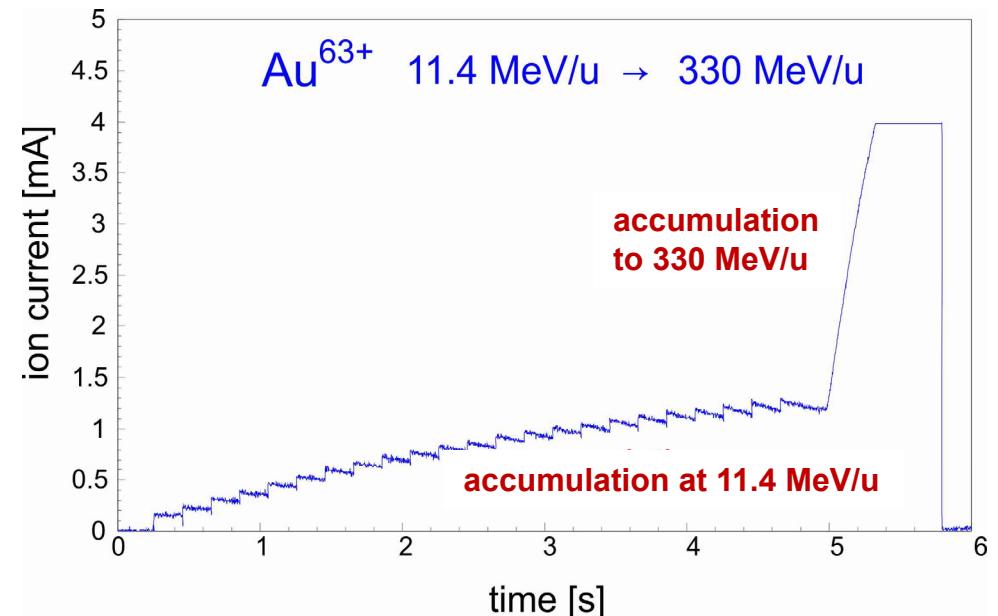
beam lifetime (recombination, vacuum)
must be long compared to cooling time

growth of momentum spread by IBS

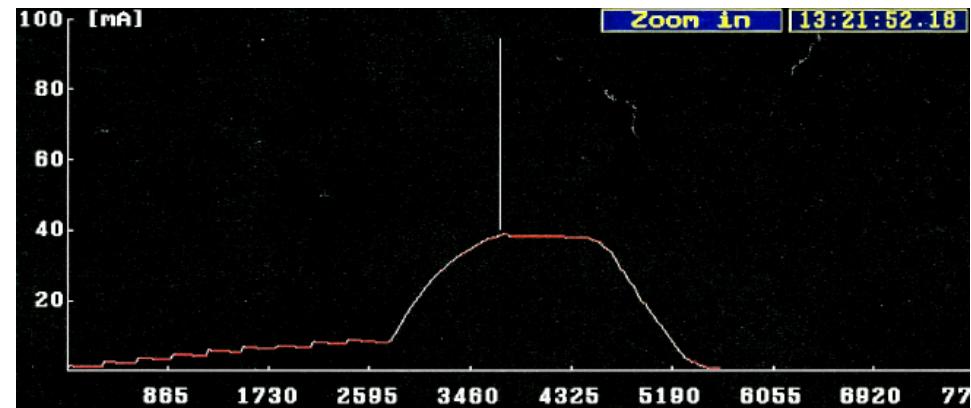


saturation due to equilibrium
injection rate – recombination rate

high intensity beams of
a few milliAmperes have been
accumulated routinely in SIS18



during operation with
high intensity uranium beams
a limit of around 10 mA
(3×10^9 ions of $^{238}\text{U}^{73+}$ at 11.4 MeV/u)
was encountered,
accompanied by
collective transverse beam oscillations



CELSIUS Uppsala (Beam Cooling Montreux 1993)

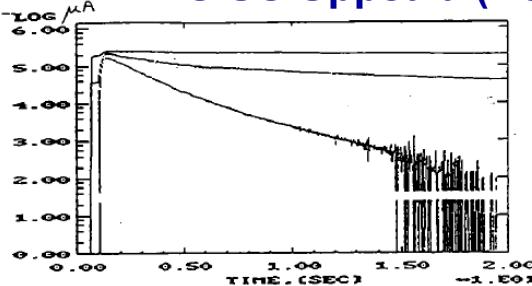


Fig. 3. Measurement of intensity (shown as $10 \log$ of μA) vs. time (from zero to 20 s) for a bunched beam of 48 MeV protons (rf. voltage about 200 V), without electrons (best lifetime), with 400 mA electron beam with 4 cm diameter, and with 100 mA of electron beam with 2 cm diameter (worst lifetime).

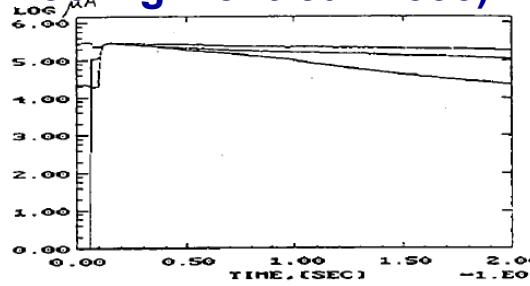
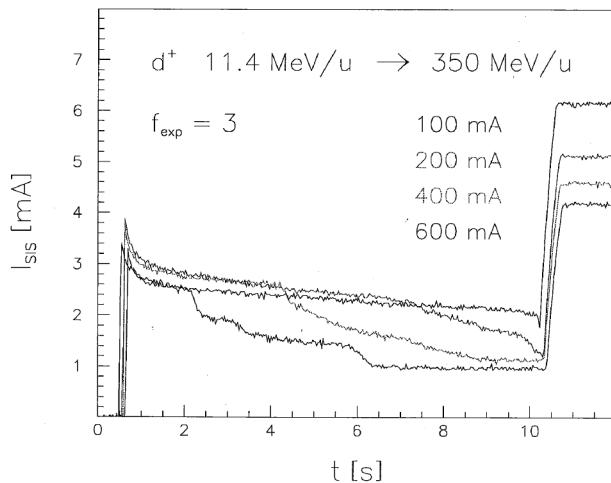


Fig. 4. As fig. 3 but for coasting beam of 48 MeV protons. The best lifetime is without electrons. The next curve is with 400 mA electron beam with 4 cm diameter. The worst lifetime is with 100 mA electron beam with 2 cm diameter.

electron cooling of deuterons in SIS18



Losses during cooling of the initially hot injected ion beam

possible reasons:
transverse instabilities
non-linear resonances
impedance of electron beam
heating due to non-linear fields of the electron beam

COSY Jülich (annual report 2003)

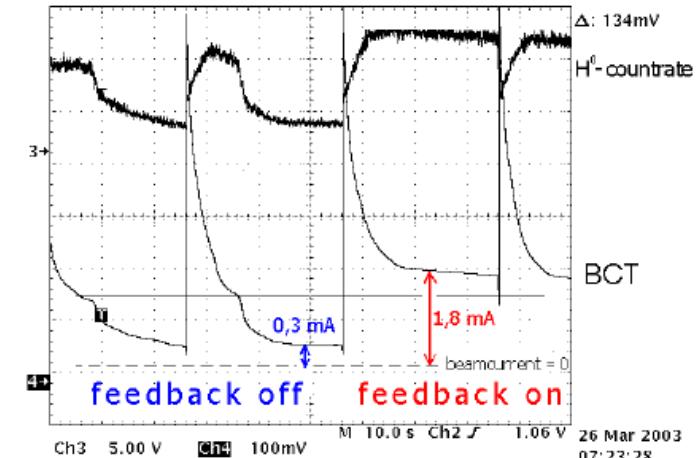
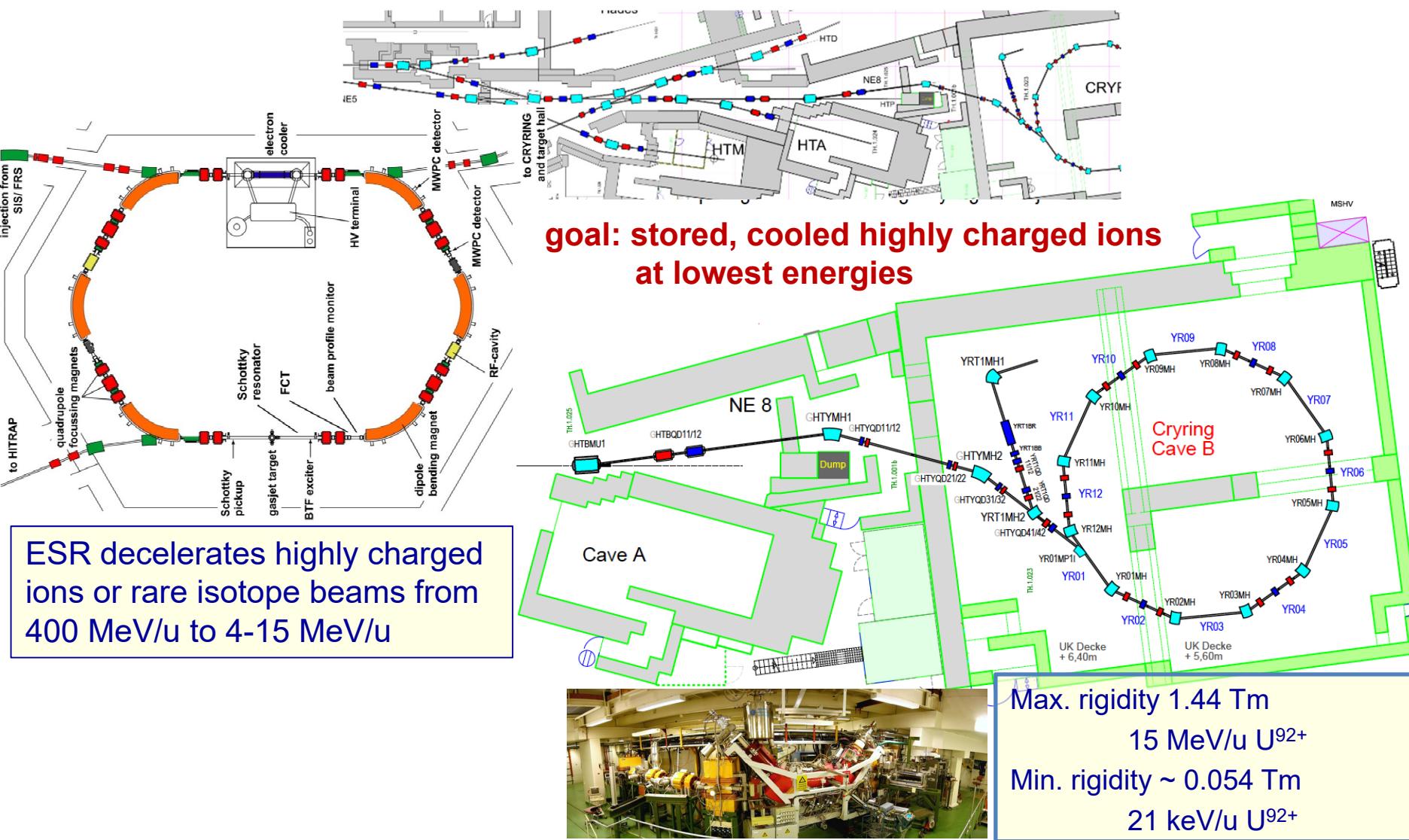
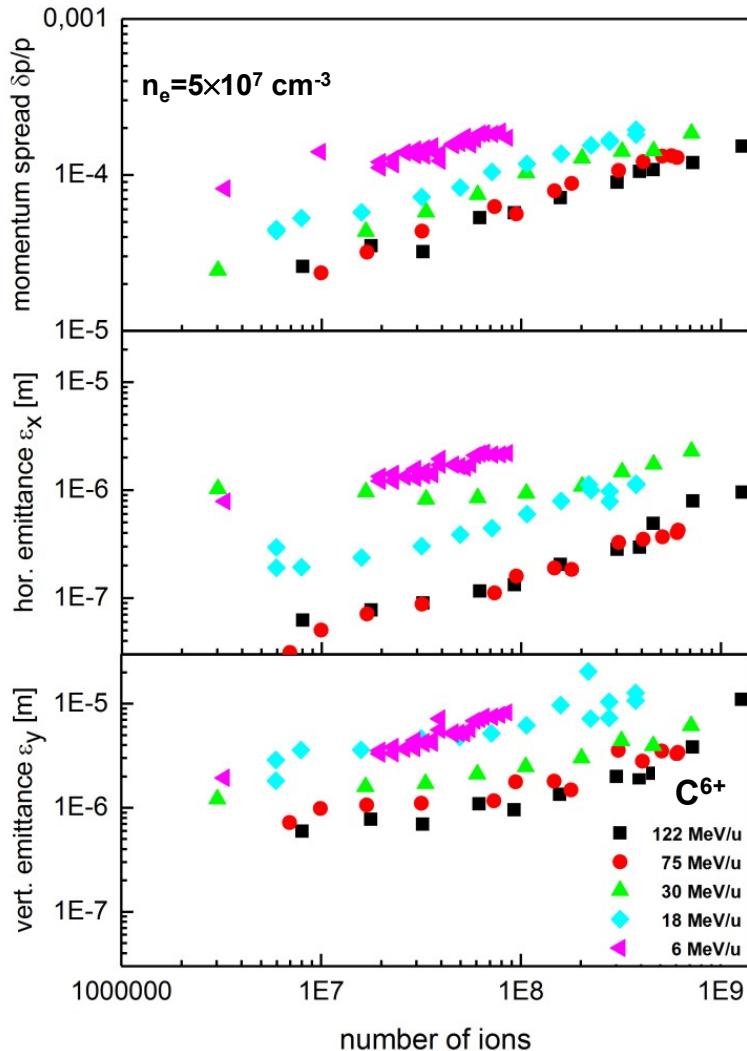


Fig. 3: BCT and H^0 -countrate signals without FB and with FB switched on. Time scale 10 s/div.





Intrabeam Scattering rate

$$\tau_{IBS}^{-1} = \frac{Q^4 e^4}{(Am_i)^2} \cdot \frac{N}{C\epsilon_h \epsilon_v \delta p / p} \cdot \frac{1}{(\gamma^4 \beta^3 c^3)} \cdot 4\pi L_C^{IBS}$$

in case of equipartitioning:
emittance increases \approx proportional $\beta\gamma$

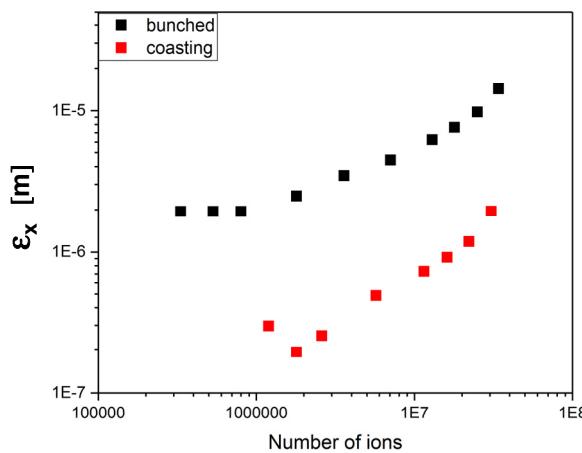
tune shift

$$\Delta Q_x = \frac{r_p Z^2 N g}{\pi A \beta^2 \gamma^3 B (\epsilon_x + \sqrt{\epsilon_x \epsilon_y Q_x / Q_y})}$$

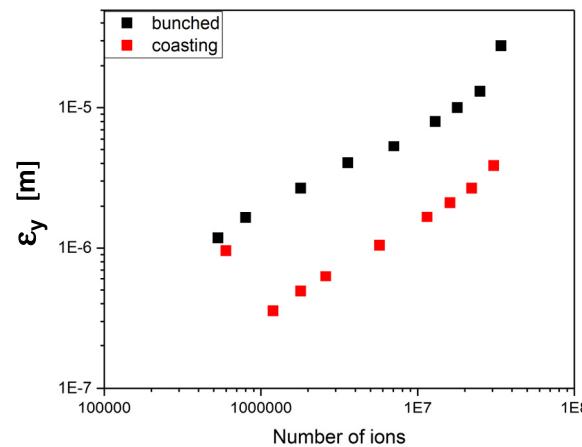
increases proportional $1/(\beta^2 \gamma^3)$

space charge effects are increased
for low energy cooled beams

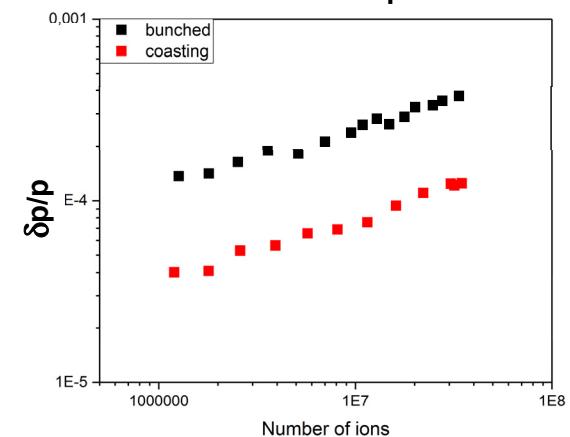
horizontal emittance



vertical emittance



momentum spread

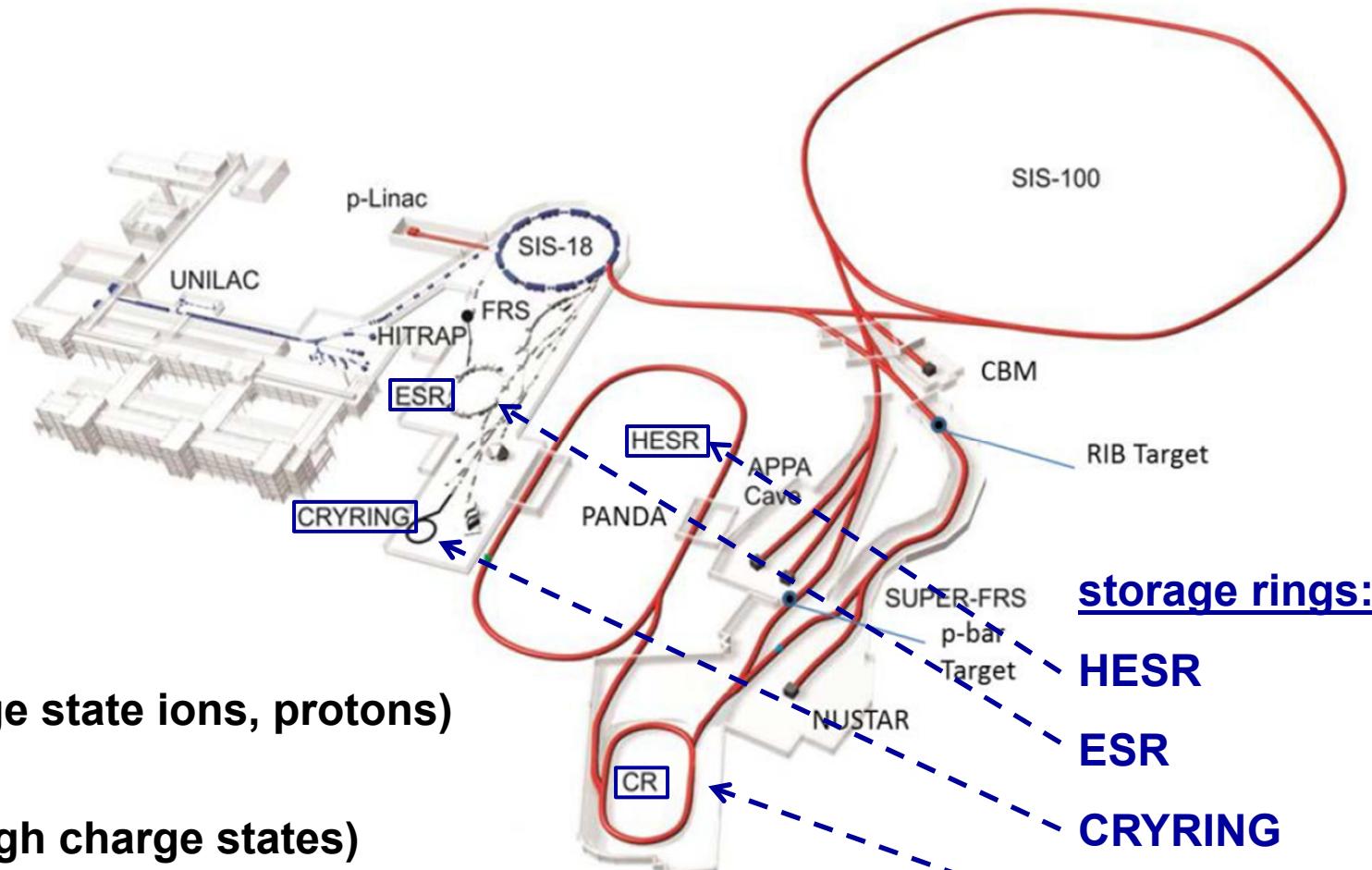


$^{238}\text{U}^{89+}$ 75 MeV/u cooled by $I_{\text{el}}=0.2$ A

increase of emittance (factor of 5-8) is not proportional to the bunching factor (15-50)

space charge effects are increased for bunched beams

and finally: bunching the beam for deceleration or transfer
is the most critical situation with respect to space charge issues

**goals:**

higher intensity
(heavy low charge state ions, protons)

higher energy
(heavy ions in high charge states)

production of antiprotons and RIBs

high quality secondary beams (cooling)

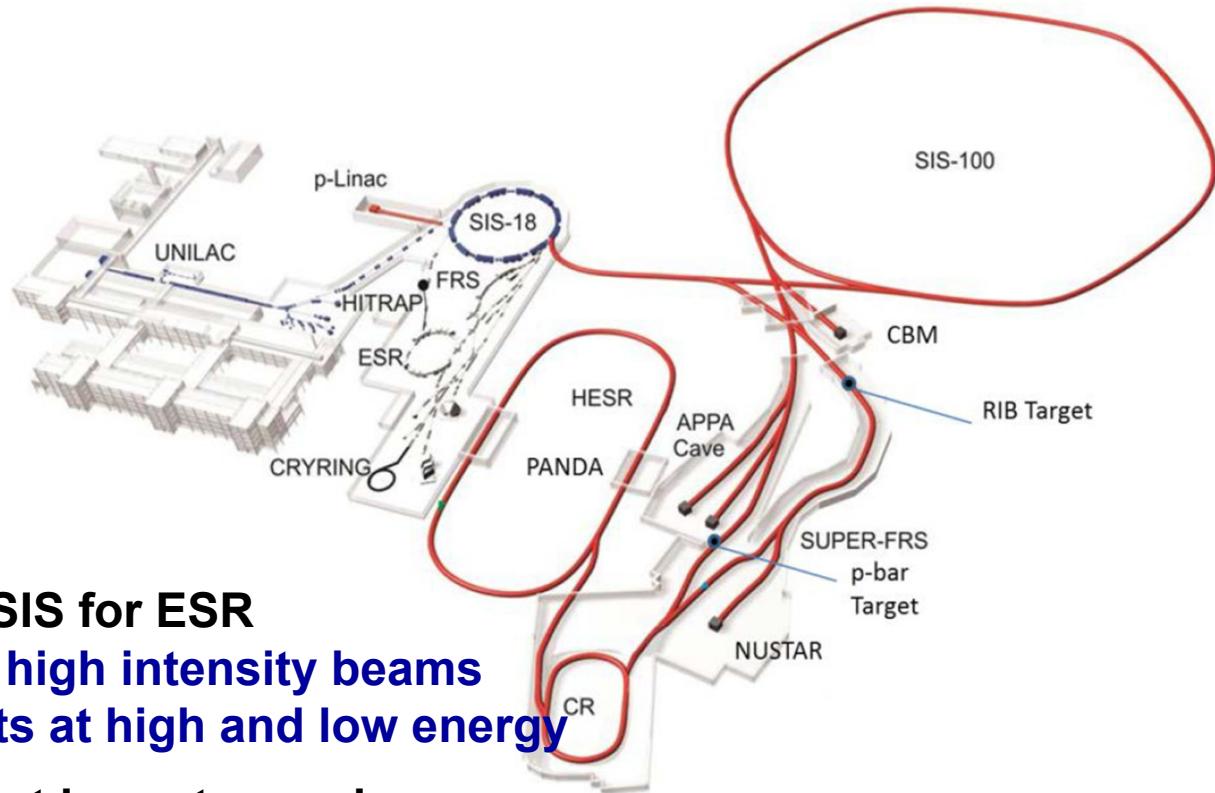
storage rings:

HESR

ESR

CRYRING

Collector Ring



- high intensity beams from SIS for ESR
⇒ studies of the cooling of high intensity beams
nuclear physics experiments at high and low energy
- highly charged ions stored at lowest energies
⇒ space charge issues in ESR and CRYRING
- availability of dedicated barrier bucket system for accumulation
of high intensity beams in ESR and HESR
- continue experiments with extremely cold ion beams in ESR and CRYRING