



DEVELOPMENT OF BEAM POSITION MONITOR USING CHERENKOV DIFFRACTION RADIATION

Ken-ichi Nanbu

Research Center for Electron Photon Science, Tohoku University

International Beam Instrumentation Conference 2019
Malmö, Sweden
8-12, September 2019



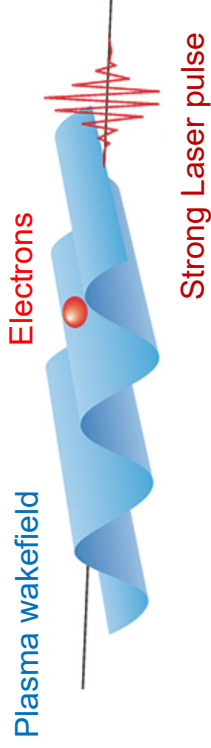
OUTLINE

- Background & Motivation
- Cherenkov diffraction radiation
- Azimuthal distribution of Cherenkov diffraction radiation
- Considerations in BPM development
 - Position derivation
 - Effect of beam size
 - Intensity
- Summary

Background

Laser-wakefield accelerators (LWFAs):

Laser-plasma acceleration

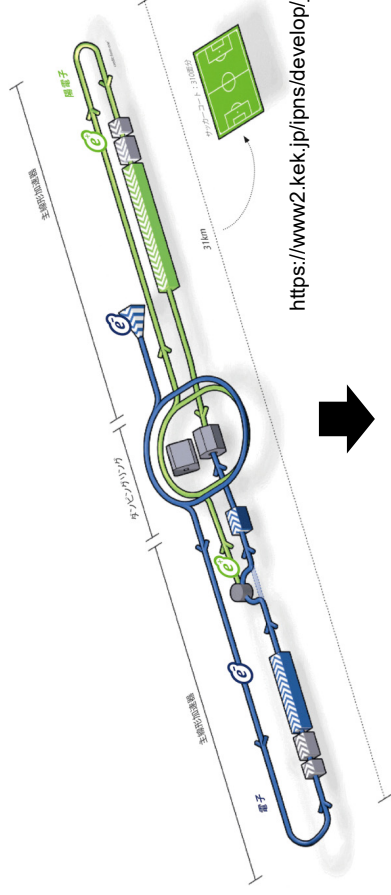


Applications

- XFEL
- Collider

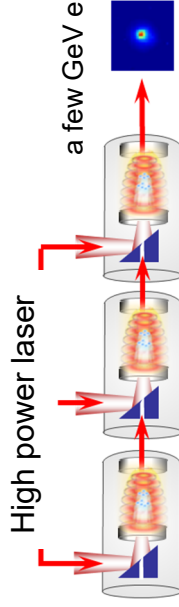
Predicted maximum field strength:

$$1 \text{ GeV} / 1 \text{ cm} = 100 \text{ GeV} / \text{m}$$



500 GeV beam would be achieved within 500 cm

Background



Beam generation acceleration

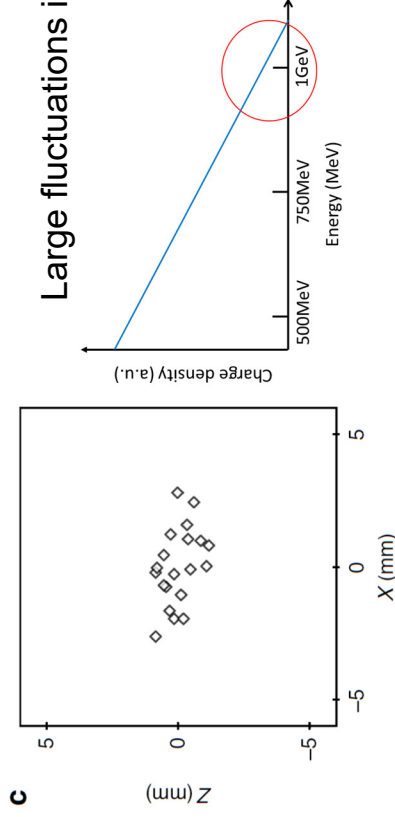
ImPACT project (Sano PM)

<http://www.jst.go.jp/impact/sano/>

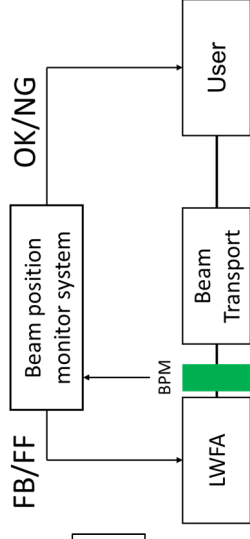
- Recently, dramatically improved their beam quality
- Small source size, Low emittance
- Large energy spread
- Still poor pointing stability
- Low beam repetition rate



Large fluctuations in beam position would be unacceptable for many beam users



Position monitoring

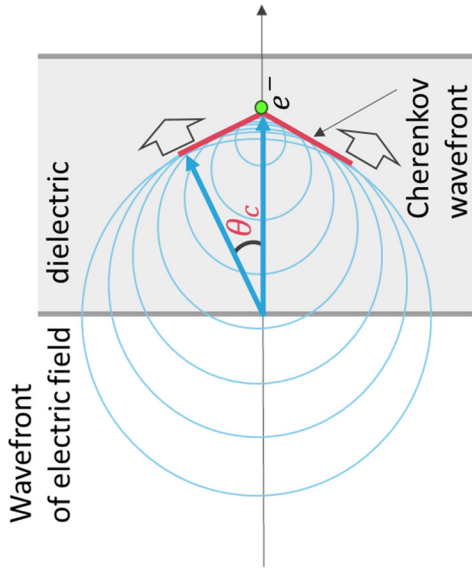


Shot-to-shot measured pointing stability

T. André et al., NATURE COMMUNICATIONS (2018) 9:1334

Non-destructive beam monitor is required in practical use of the LWFA

Cherenkov radiation



Cherenkov angle: θ_c

$$\cos \theta_c = \frac{1}{\beta n}$$

$$\beta = c/v$$

c : Speed of light

v : Electron velocity

n : Refractive index

$$(\beta \simeq 1, \beta n > 1)$$

Number of generated photons: N_0

(Frank-Tamm formula)

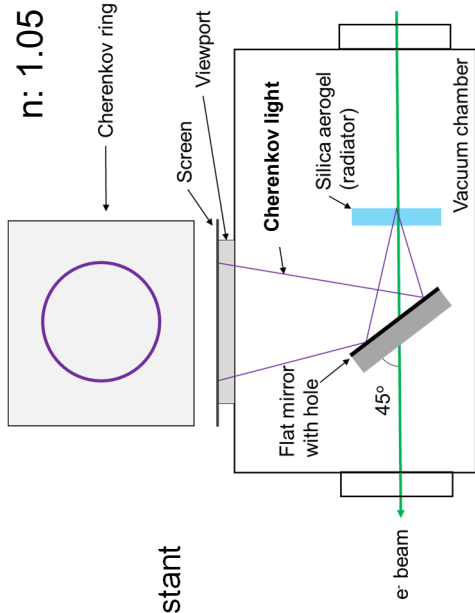
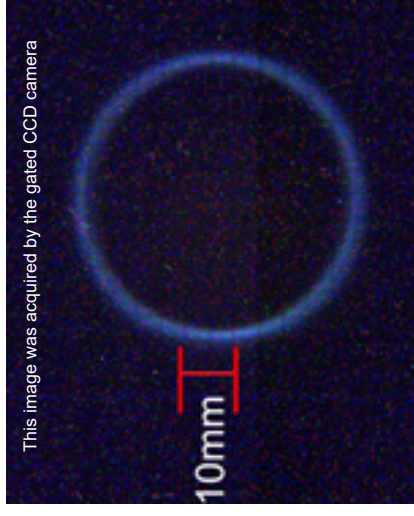
$$N_0 = 2\pi\alpha L \left(1 - \frac{1}{\beta^2 n^2(\omega)} \right) \left(\frac{1}{\lambda_1} - \frac{1}{\lambda_2} \right)$$

α : Fine structure constant

λ : Wavelength

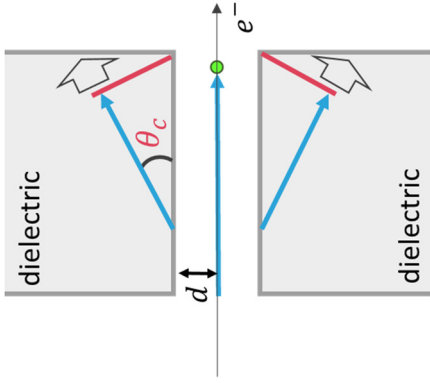
L : Radiator length

Jelley, J.V. " Čerenkov Radiation And Its Applications" 1958, pp.4-5



Cherenkov diffraction radiation (ChDR)

Cherenkov radiation from a nearby dielectric



The ChDR intensity can be expressed by multiplying the intensity of ordinary Cherenkov radiation by a coupling factor, K .

- Coupling factor

$$K = \exp\left(-4\pi \frac{d}{\beta\gamma\lambda}\right)$$

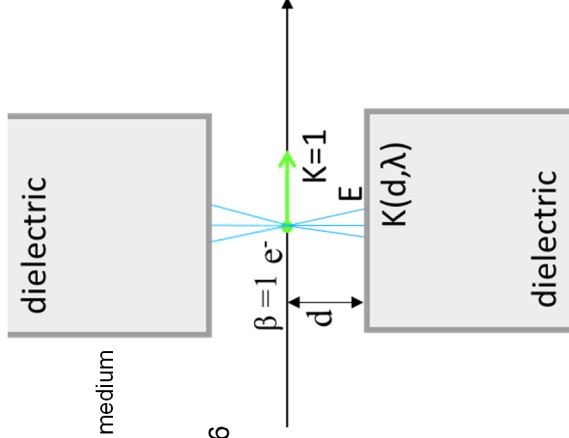
d : distance between the electron and the surface of the medium
 γ : Lorentz factor $\gamma=(1-\beta^2)^{-1/2}$

R. Ulrich, Z. Phys. 194, 180 (1966)

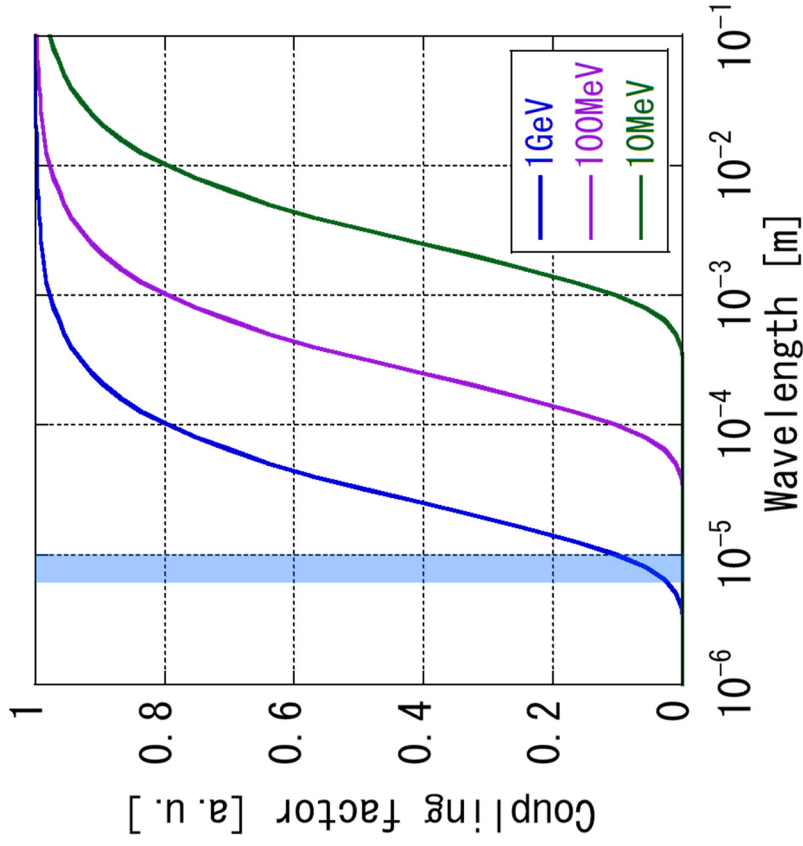
T. Takahashi et al., Phys. Rev. E62(2000) 8606

- Number of generated photons (ChDR): N

$$N = N_0 K$$



Characteristic of coupling factor



Large photon yield ➡ Higher K

1 GeV: Infrared region

100 MeV: Terahertz region

10 MeV: Millimeter-wave region

d and λ have to be chosen with respect to beam energy

Variation of the coupling factor for different beam energies ($d = 3.5$ mm)

Azimuthal distribution of intensity

Distance d between the charged particle and surface of medium is calculated by

$$d = \sqrt{R^2 + r_0^2 - 2r_0R \cos(\varphi - \varphi_0)}$$

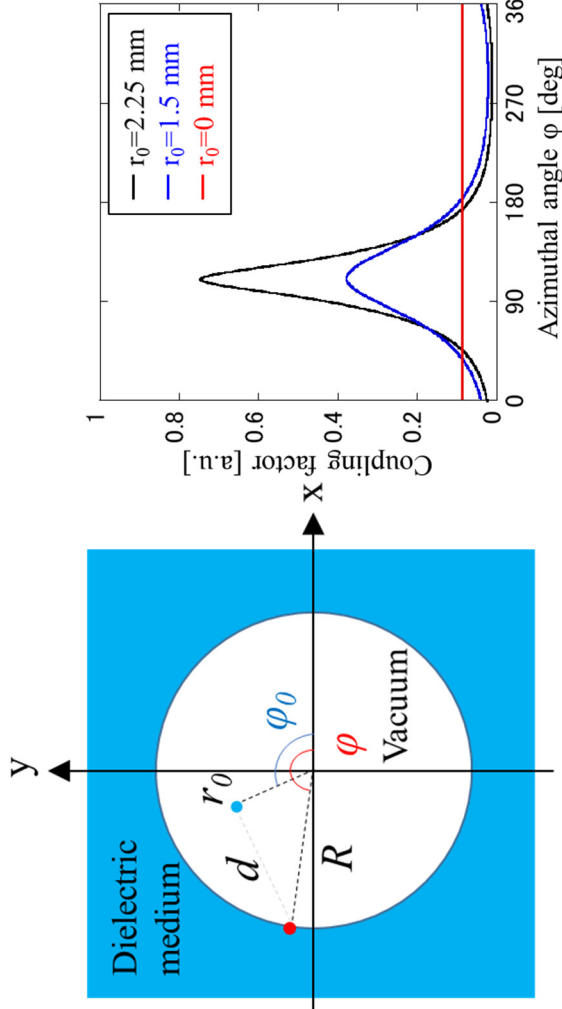
(r_0, φ_0) are the polar coordinates of the beam position

(R, φ) are the polar coordinates of the position, which defines the intersection of the observation direction with surface of medium.

$$K(r_0, \varphi_0, \varphi) = \exp\left(-4\pi \frac{\sqrt{R^2 + r_0^2 - 2r_0R \cos(\varphi - \varphi_0)}}{\beta\gamma\lambda}\right)$$

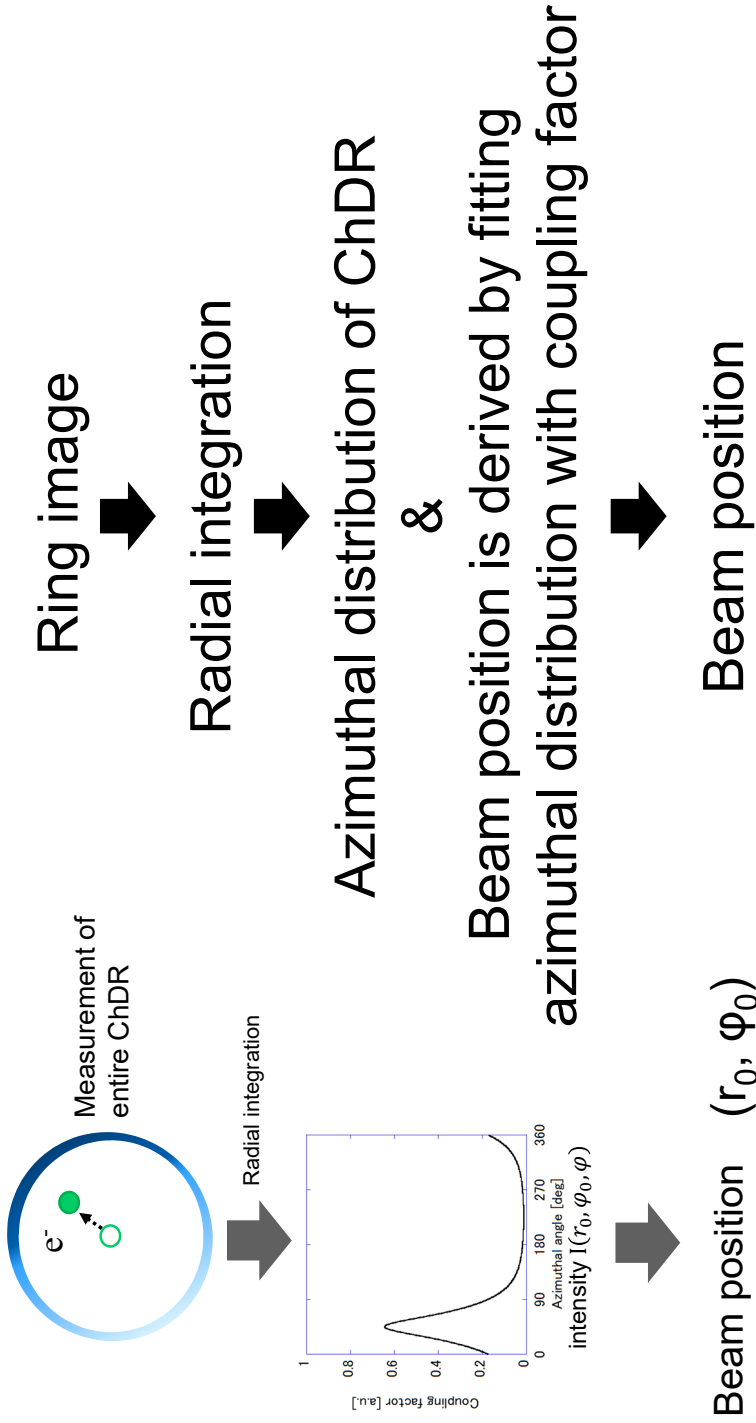
Azimuthal distribution of intensity

$$I(r_0, \varphi_0, \varphi) = N_0 \exp\left(-4\pi \frac{\sqrt{R^2 + r_0^2 - 2r_0R \cos(\varphi - \varphi_0)}}{\beta\gamma\lambda}\right)$$

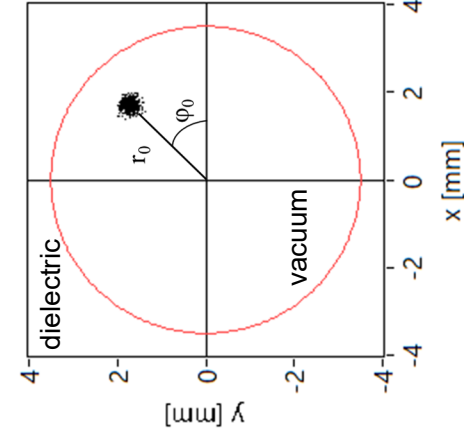


Position derivation

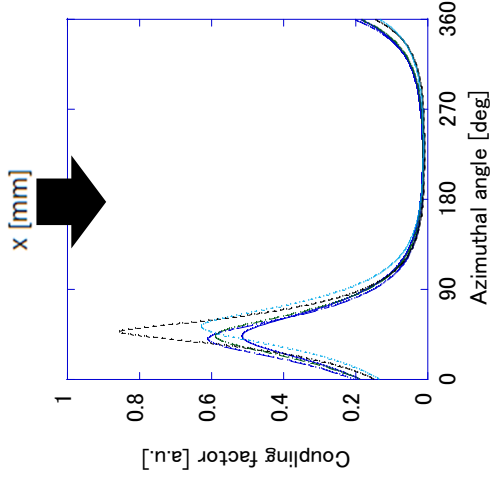
■ Concept



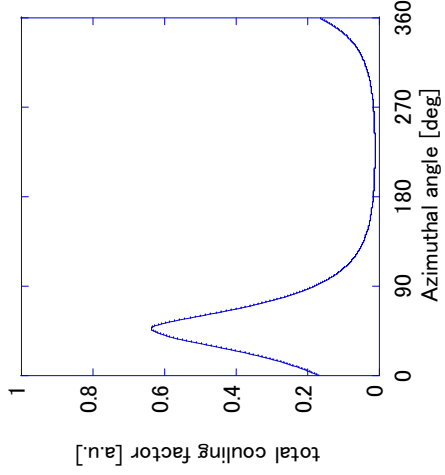
Position derivation (finite beam size)



- Beam size, σ_b :
 - Gaussian distribution
 - σ_b : 0, 25, 50, 100, 200, 300 μm
 - N: 10000
- Beam position, r_0 : 0 ~ 2 mm
- Each coupling factor is calculated from the given position of macro particles
- A total coupling factor, which is defined as the average of all the coupling factors is calculated
- Beam position is derived by fitting the total coupling factor with $K(\varphi)$



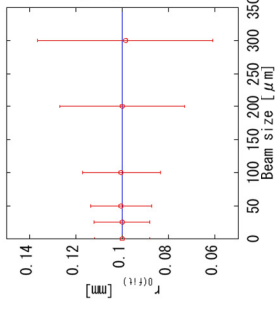
ave. 



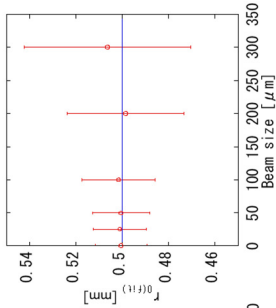
Beam position (r_0, φ_0)

Effect of beam size (Numerical simulation)

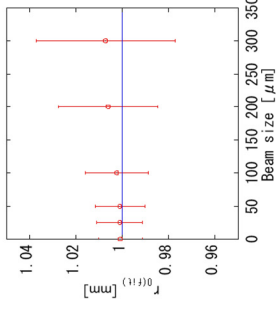
$r_0 = 0.1 \text{ mm}$



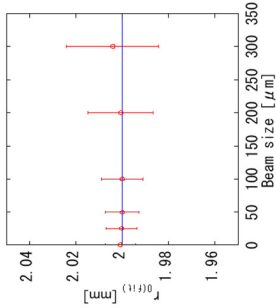
$r_0 = 0.5 \text{ mm}$



$r_0 = 1.0 \text{ mm}$



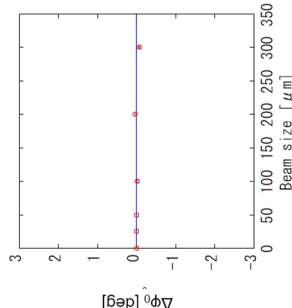
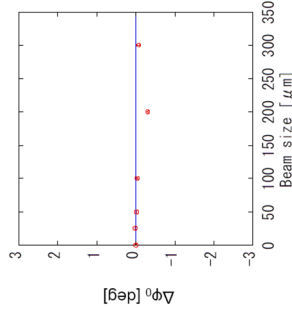
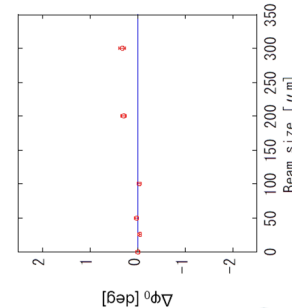
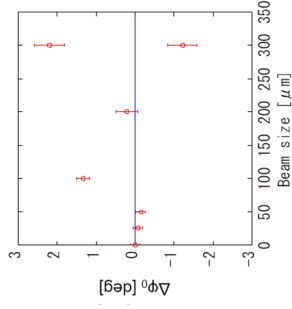
$r_0 = 2.0 \text{ mm}$



r_0 :

Large beam size → Large position error

Beam size should be small



φ_0 :

Small r_0 → Large error

※ The error bars represent the fitting error

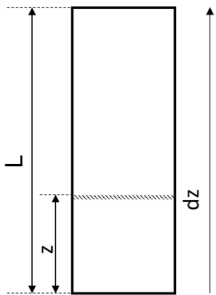
■ 1 GeV beam (R: 3.5 mm, $\lambda = 10 \text{ } \mu\text{m}$):

$$\sigma_b \leq 100 \text{ } \mu\text{m}$$

$$\Rightarrow \Delta r_0 \cong 17 \text{ } \mu\text{m}$$

Requirement for hollow dielectric as radiator

ChDR $I(\omega)$ Intensity



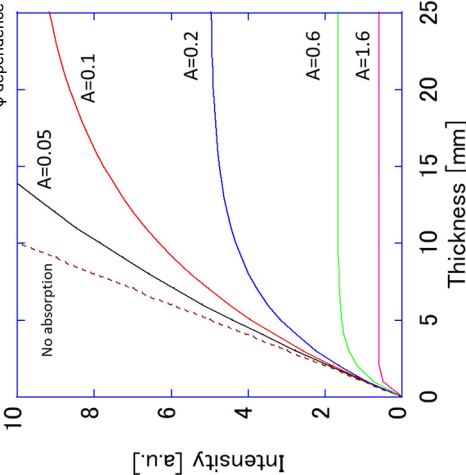
$$N_0 = 2\pi\alpha L \left(1 - \frac{1}{\beta^2 n^2}\right) \equiv \eta L$$

⇒ Photon yield : $\int_0^L \eta e^{(-A(L-z))} dz$

A is attenuation coefficient
z: Emission point

$$I(r_0, \varphi_0, \varphi) = \frac{\eta}{A} (1 - e^{-AL}) \cdot \exp\left(-4\pi \frac{\sqrt{R^2 + r_0^2 - 2r_0 R \cos(\varphi - \varphi_0)}}{\beta \gamma \lambda}\right)$$

A, and L have no azimuthal dependence
ϕ dependence of intensity comes only from coupling factor, K



Photon yields for various attenuation coefficients

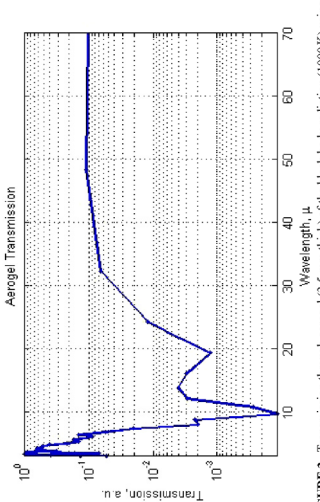


FIGURE 2. Transmission through aerogel (3.5 mm thick) of the black body radiation (1000 K) using BLIS interferometer.

n = 1.008

R. Tikhoplav et al., AIP Conference Proceedings 1086, 610 (2009)

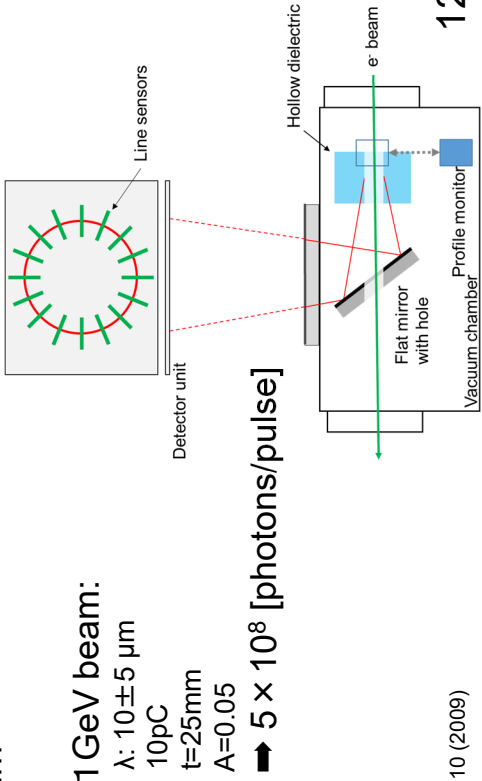
- Hollow dielectric medium

- Large photon yield ➡ Larger n, Small absorption coefficient
- Cherenkov angle has to be small ➡ Smaller n



Smaller absorption coefficient, appropriate refractive index

Entire ChDR ring can be observed employing appropriate dielectric medium



Summary

- Laser-wakefield accelerators have attracted much attention and have been developed actively towards practical applications worldwide.
- Single-shot and non-destructive beam monitor will provide the opportunity for effective use in practical application of LWFA.
- By measuring ChDR, single-shot and non-destructive beam diagnostics can be realized.
- The observed wavelength and the size of the hollow dielectric have restrictions related to the beam energy.
- The beam size affects the position accuracy.
- Since the observation wavelength changes depending on the beam energy, it is necessary to optimize the refractive index, thickness, and transmittance of the radiator. As a future prospect, the development of more effective radiators is considered to be an important theme.

Acknowledgement

This work was supported by JSPS KAKENHI Grant Numbers 18K11915.



Thank you very much for your attention