

SIMULATION STUDY ON TRANSVERSE LASER COOLING AND CRYSTALLIZATION OF HEAVY-ION BEAMS AT THE COOLER STORAGE RING S-LSR

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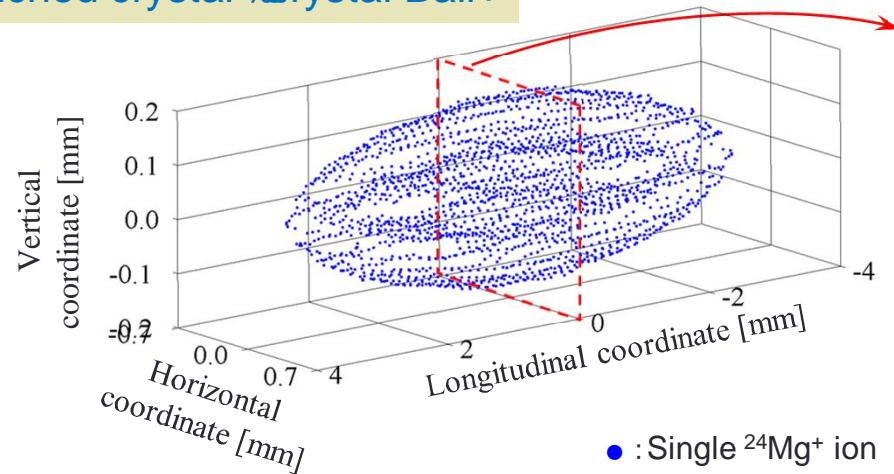
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Outline

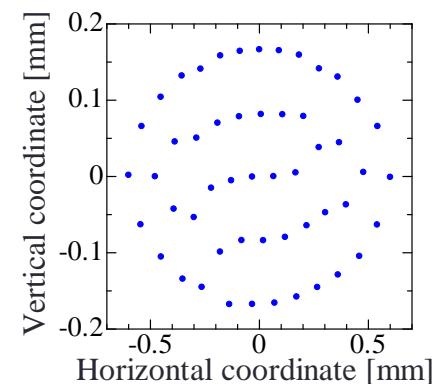
- „ Intro: Purpose of the present study
- „ Molecular dynamics (MD) simulation
- „ Simulation conditions
- „ MD results
 - „ Three-dimensional laser cooling
 - „ Crystallization
- „ Summary

Beam Crystallization

3D bunched crystal %Crystal Ball+



Cross-sectional profile
in the bunch center



- „ Coulomb crystalline state of an ion beam strongly cooled in a storage ring
- „ Characteristics:
 - „ Ultralow emittance
 - „ Coulomb coupling constant > 170
 - „ Periodic oscillation with the external focusing force
 - „ Stable after removing the cooling force

Purpose of the Present Study

- “ Feasibility of beam crystallization was already predicted if the ring and laser conditions were sufficient. (PRL2004, PRSTAB2005)
- “ However, laser-cooling conditions have been limited in the recent experiments at S-LSR.
 - “ Single laser beam, low power, and fixed detuning.
- “ To show numerically how to attain a low-emittance beam using **Resonant Coupling** and **Laser Cooling** by assuming actual parameters at S-LSR.
 - “ Optimization of a cooling laser for high cooling efficiency
(To be presented at NA-PAC13)
 - “ **Fast 3D cooling of low-current beams**
 - “ **Feasibility of beam crystallization**
- “ Numerical study using a Molecular Dynamics (MD) simulation technique.

Molecular Dynamics (MD) Simulation

“ The most reliable simulation technique for the study of beam cooling and crystallization.

“ Hamiltonian

$$H = \frac{p_x^2 + p_y^2 + p_z^2}{2} - \frac{\gamma}{\rho} x p_z + \frac{x^2}{2\rho^2} - \frac{K(s)}{2} (x^2 - y^2) + \frac{r_p}{\beta^2 \gamma^2} \phi.$$

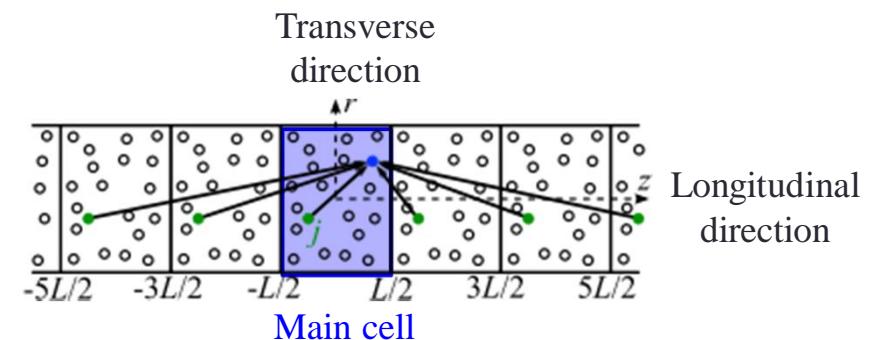
“ Motion of real particles is integrated in a symplectic manner.
 “ Coulomb potential --- Periodic boundary condition imposed

$$\phi = \sum_j (\phi_{short}^{(j)} + \phi_{long}^{(j)}).$$

$$\phi_{short}^{(j)} = \frac{1}{\sqrt{(x - x_j)^2 + (y - y_j)^2 + (z - z_j)^2}}$$

$$\phi_{long}^{(j)} = \frac{2}{L} \int_0^\infty \frac{\cosh(kz^{(j)}/L) J_0(kr^{(j)}/L) - 1}{e^k - 1} dk$$

where $z^{(j)} = |z - z_j|$ and $r^{(j)} = \sqrt{(x - x_j)^2 + (y - y_j)^2}$.



For a bunched beam, L can be set as a bucket length (C/h).

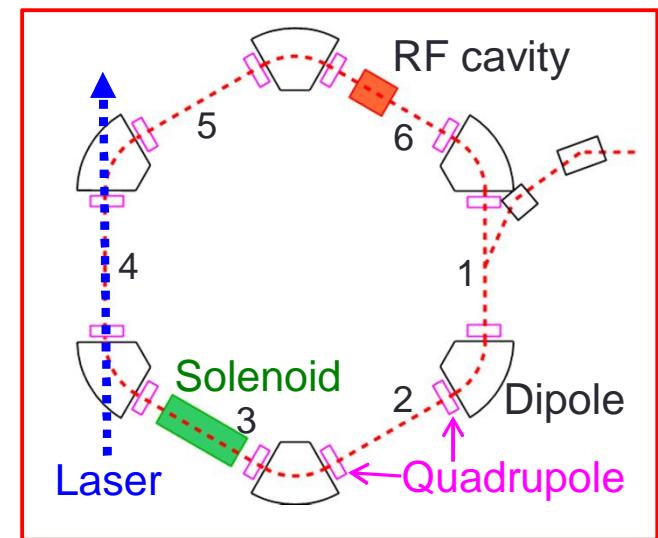
Resonant Coupling for 3D Cooling

- “ A possible scheme for efficient transverse cooling
 - “ H. Okamoto, D. Mohl, and A. M. Sessler, (PRL1993, PRE1994)
- “ First, introduce a coupling source in the ring.
 - “ RF cavity placed where the dispersion is finite for X-Z coupling
 - “ Solenoid magnet for X-Y coupling
- “ Then, operate the ring at a difference resonant condition;
 - $\nu_x - \nu_z \approx \text{integer}$ for X-Z coupling
 - $\nu_x - \nu_y \approx \text{integer}$ for X-Y coupling

MD Simulation Conditions (1)

„ Machine (S-LSR at Kyoto Univ.)	
„ Circumference	22.56 m
„ Superperiodicity	6
„ Lattice	
„ Tunes	Case-I (v_x, v_y, v_z)~(2.07, 1.12, 0.07) Case-II (v_x, v_y, v_z)~(2.07, 1.07, 0.07)
„ RF bunching voltage	~40 V
„ Harmonic number	100
„ Adiabatic capture	5,000 turns (0.2sec)
„ Beam	
„ Ion species	40-keV $^{24}\text{Mg}^+$
„ Lorentz factors	$\beta=1.89\times10^{-3}$, $\gamma=1.00000179$
„ Revolution frequency (period)	25 kHz (40 μsec)
„ Initial RMS emittance ($\varepsilon_x=\varepsilon_y$)	$1\times10^{-9} \pi \text{ m.rad}$ (Normalized)
„ Initial dp/p (rms)	$5\times10^{-7} \pi \text{ m.rad}$ (Un-normalized) 3×10^{-4}

Schematic view of S-LSR



From the
measurement
result

MD Simulation Conditions (2)

“Laser (1 co-propagating laser)

“ Power	8mW
“ Spot radius w (2sigma)	0.66 mm (Peak Saturation Power~4.6)
“ Detuning Δ (fixed)	-200 MHz
“ Cooling time	3 sec

} From the experiment

These parameters are rather limited as compared to past experiments in TSR & ASTRID.

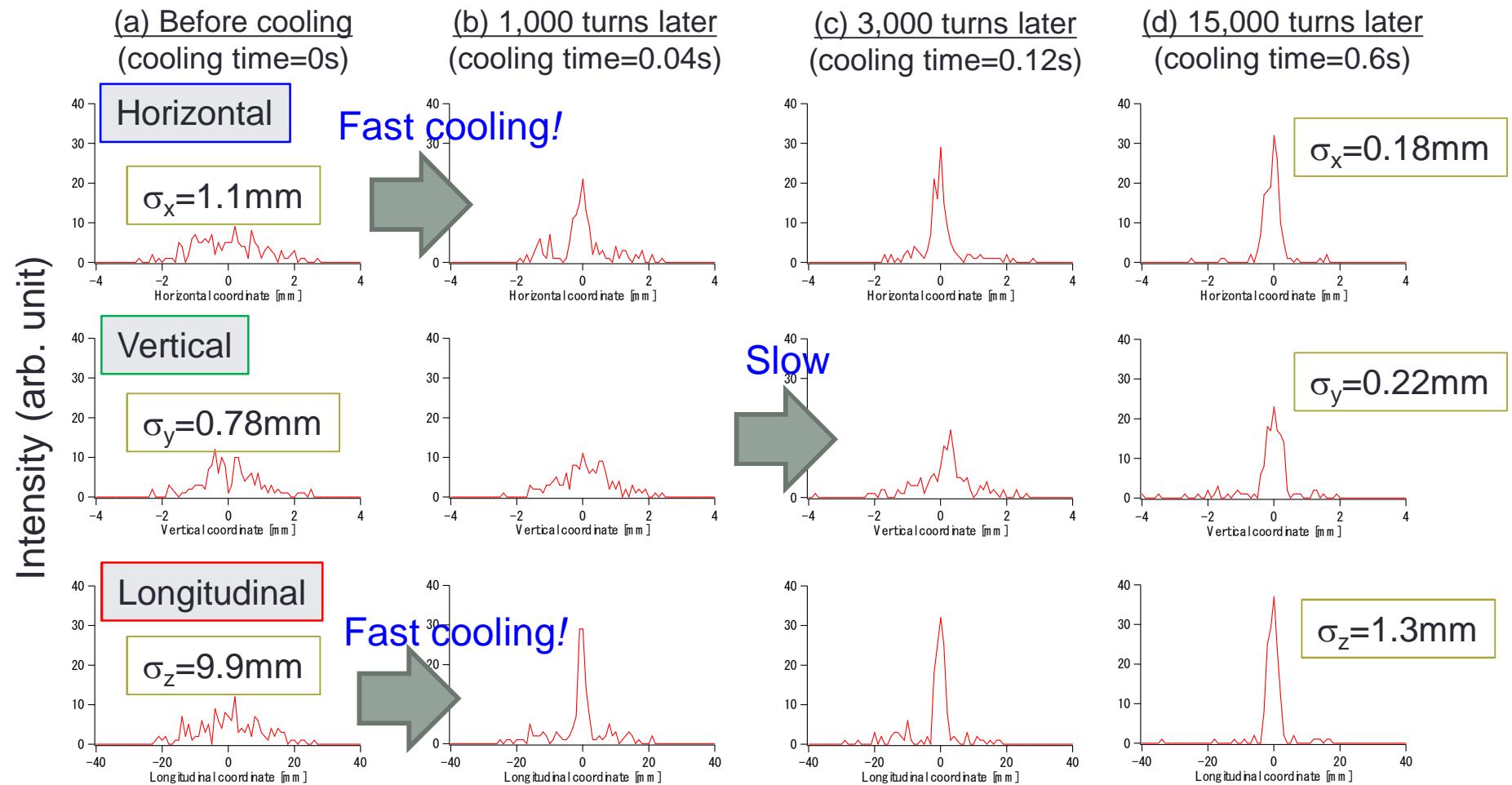
$$\text{Cooling force: } F = \frac{1}{2} \hbar k_L \Gamma \frac{S_L}{1 + S_L + (2\Delta/\Gamma)^2}$$

$$\text{Saturation parameter : } S_L = S_0 \exp\left[-\frac{2(x^2 + y^2)}{w^2}\right]$$

$$\text{Laser detuning : } \Delta \approx \omega\gamma \left[1 - \beta \left(1 + \frac{\delta p}{p} \right) \right] - \omega_0$$

MD Results (1: Time evolution)

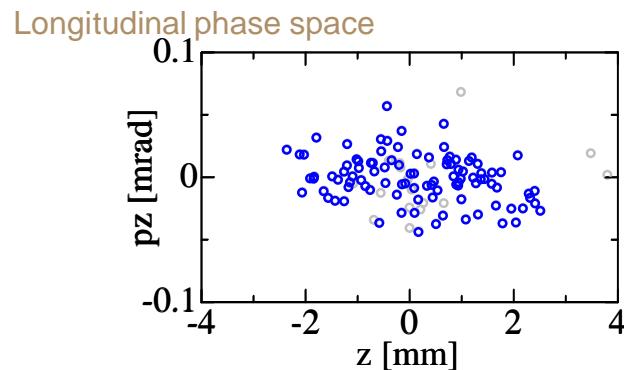
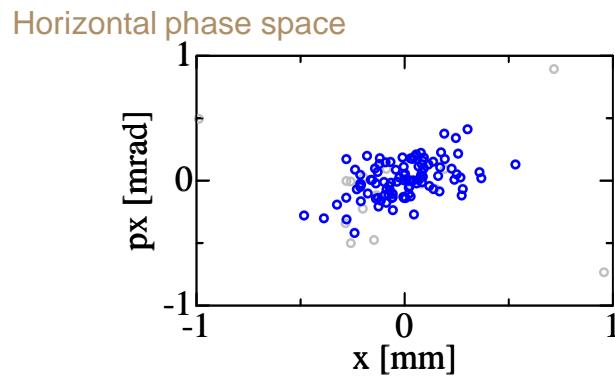
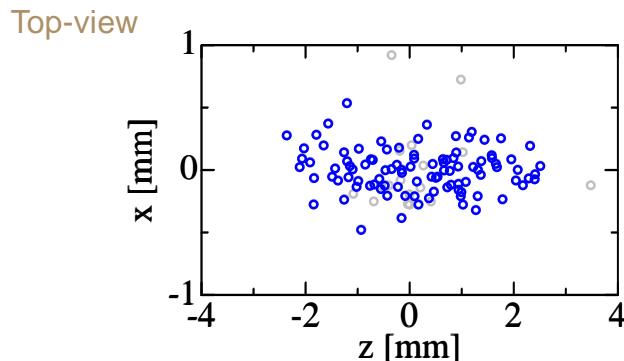
Case-I
 (v_x, v_y, v_z)
 $\sim(2.07, 1.12, 0.07)$



The vertical direction is cooled through the Coulomb interaction between ions, although no artificial cooling force is introduced.

MD Results (2: Equilibrium state)

Case-I
 (v_x, v_y, v_z)
 $\sim(2.07, 1.12, 0.07)$



- “ The ion number of the cooled part (**blue ions** in the picture) is about 100. Namely, the cooling efficiency is about 70%.

- “ Horizontal

- “ Norm. rms $\varepsilon=4.6 \times 10^{-11} [\pi m \cdot rad]$
- “ $T_x=18[K]$
- “ **Radius $\sigma=0.18 mm$**

- “ Vertical

- “ Norm. rms $\varepsilon=2.6 \times 10^{-11} [\pi m \cdot rad]$
- “ $T_y=3.8[K]$
- “ **Radius $\sigma=0.22 mm$**

- “ Longitudinal

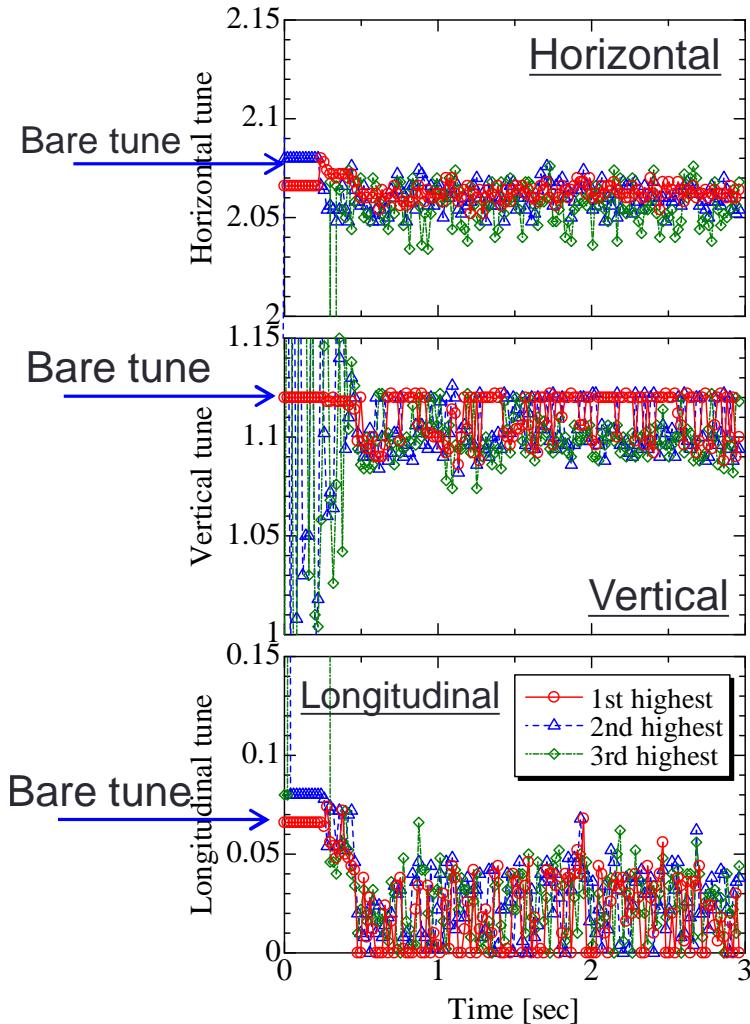
- “ Rms $dp/p = 2.2 \times 10^{-5}$
- “ $T_z=0.45[K]$
- “ **Radius $\sigma=1.3 mm$**

These values agree well with the observation result in S-LSR!!

The beam is three-dimensionally cooled, but the ordered configuration cannot be seen.

Case-I
 (v_x, v_y, v_z)
 $\sim(2.07, 1.12, 0.07)$

MD Results (3: Tune shift)

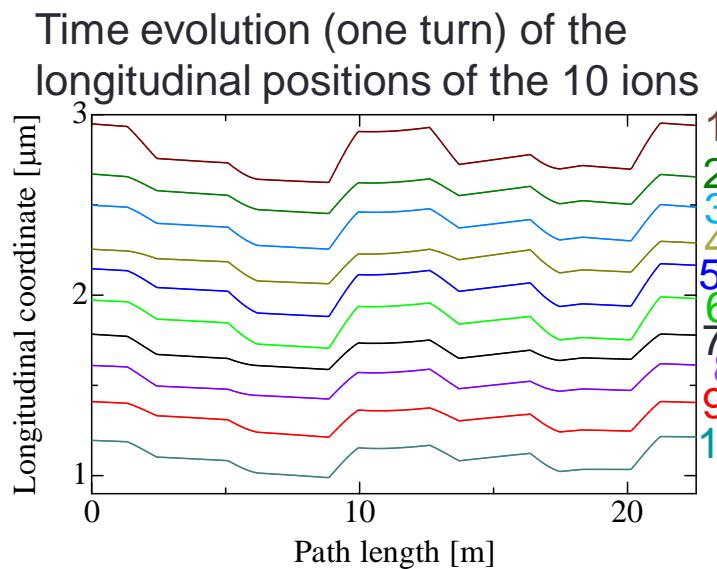
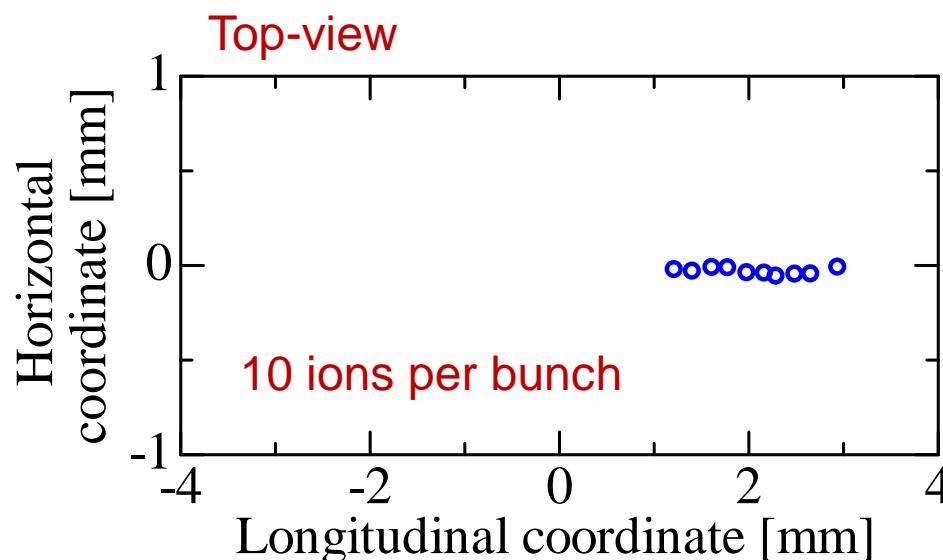


- “ The orbits of several ions are Fourier-transformed to see the time evolution of tunes in all three directions.
- “ The three highest peaks in the power spectrum (right pictures) are plotted.
- “ Result: tune shift
 - “ Horizontal 2.07 --> 2.05~2.06
 - “ Vertical 1.12 --> 1.09~1.10
 - “ Longitudinal 0.07 --> 0.00~0.04
- “ The synchrotron tune is almost damped by laser cooling.
- “ The beam is still oscillating in the transverse direction.

The laser-cooled beam is three-dimensionally space-charge-dominated.

MD Results (4: Crystallization)

Case-I
 (v_x, v_y, v_z)
 $\sim(2.07, 1.12, 0.07)$



- Even with the limited laser-cooling condition, **1D string crystal** can be formed when the beam current is sufficiently low and detuning is small.

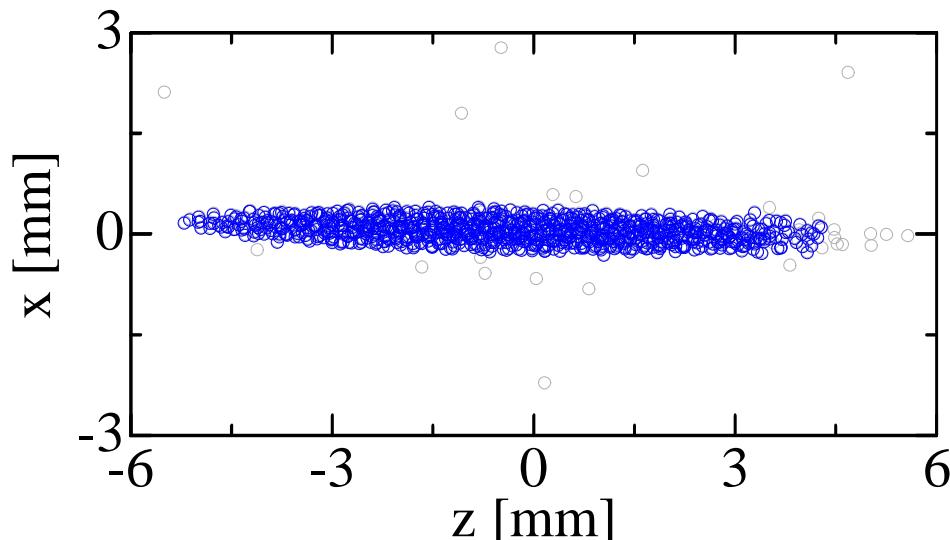
- Each ion does not pass by neighboring ions.
- The synchrotron oscillation is fully depressed.

Note that the bunch is positioned forward because the beam is pushed by the co-propagating laser.

Beam crystallization is feasible at S-LSR!!

MD Results (5: Ideal case)

Case-II
 (v_x, v_y, v_z)
 $\sim(2.07, 1.07, 0.07)$



- „ 1000 ions per bunch
- „ 3D-resonant tunes:
 - „ $(v_x, v_y, v_z)=(2.07, 1.07, 0.07)$
 - „ Weak solenoid $B=80G$
- „ Laser conditions:
 - „ 2 lasers (co- and counter-propagating)
 - „ High power (100mW)
 - „ Frequency scanned (-4GHz to -40MHz for 1sec)

- „ More than 90% ions are laser-cooled.
- „ Transverse norm. rms emittance $\sim 1 \times 10^{-11} \pi m.rad$ ($T_{x,y} \sim 10K$)
- „ Longitudinal momentum $dp/p \sim 1 \times 10^{-5}$ ($T_z \sim 0.1K$)

The highest-quality heavy-ion beam can be formed just by improving the laser system in S-LSR!!

Summary

- „ 3D laser cooling of the heavy-ion beam in S-LSR was studied using the MD simulation technique.
- „ The three-dimensionally low-temperature bunched ion beam was generated through resonant coupling.
- „ The MD result agreed well with the observation result in the recent experiment in S-LSR.
- „ Beam crystallization (1D string at low line density) is possible even in a limited cooling situation.
- „ An ultra-low-emittance bunched beam can be formed at a high intensity by a combination of powerful laser cooling and resonant coupling.