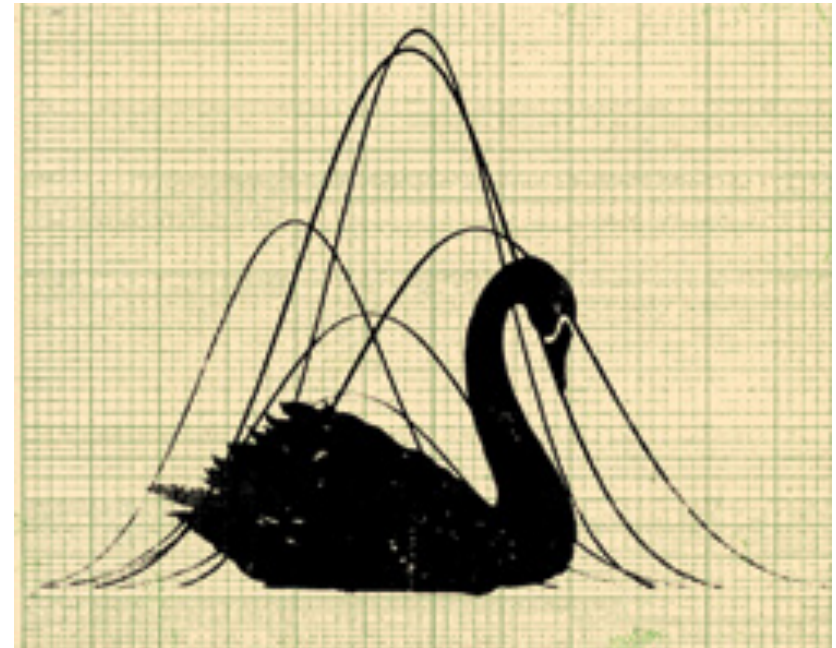


Experimental characterization of the laser heater effects on a seeded FEL

E. Ferrari^{1,2}, G. Penco¹, E. Allaria¹, S. Spampinati¹,
L. Giannessi^{1,3}, W. Fawley¹, Z. Huang⁴

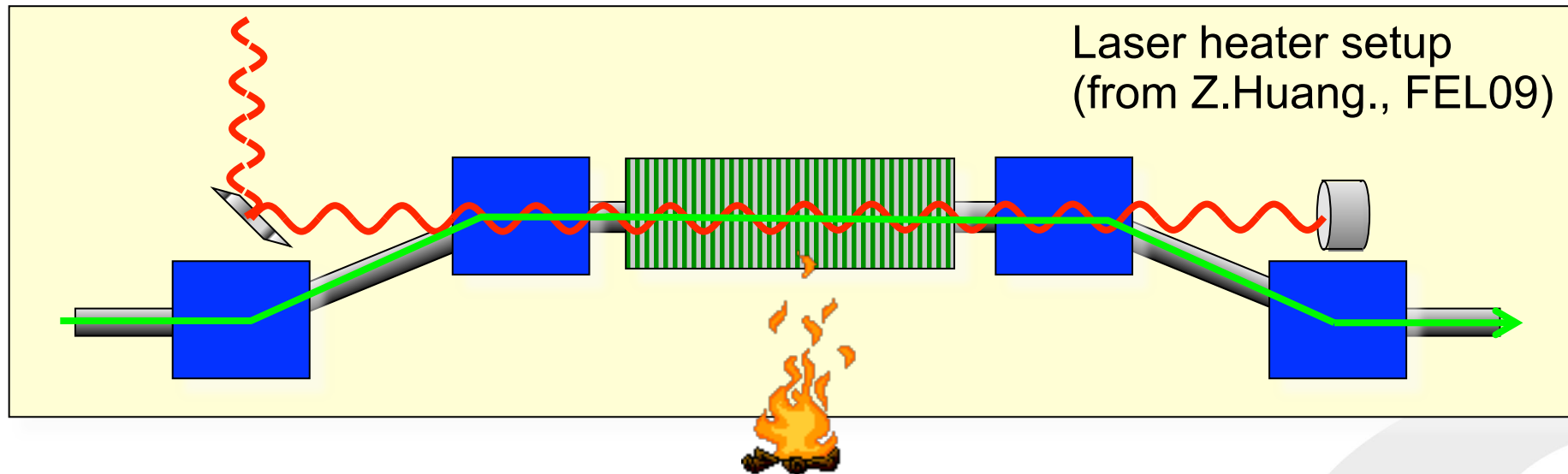
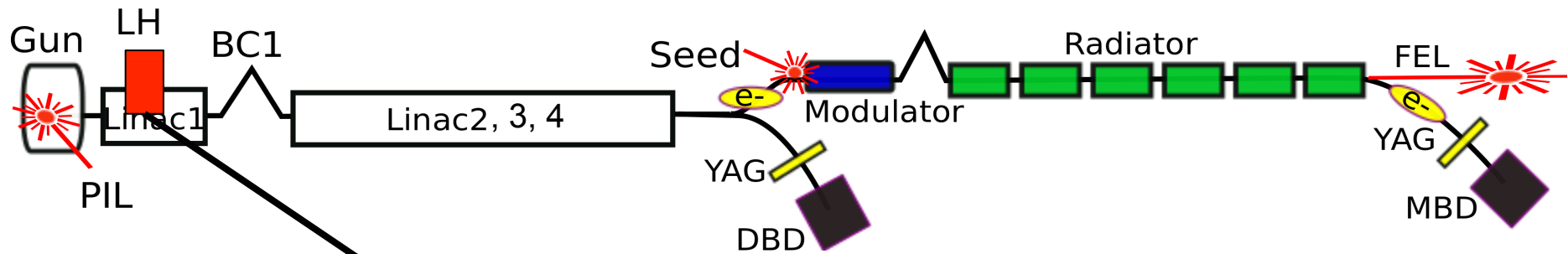
1. Elettra - Sincrotrone Trieste;
2. Universita' degli Studi di Trieste;
3. ENEA C.R. Frascati
4. SLAC National Accelerator Laboratory

- ★ Fermi experimental setup
- ★ Impact of the laser heater on seeded FEL
- ★ Non-Gaussian effects



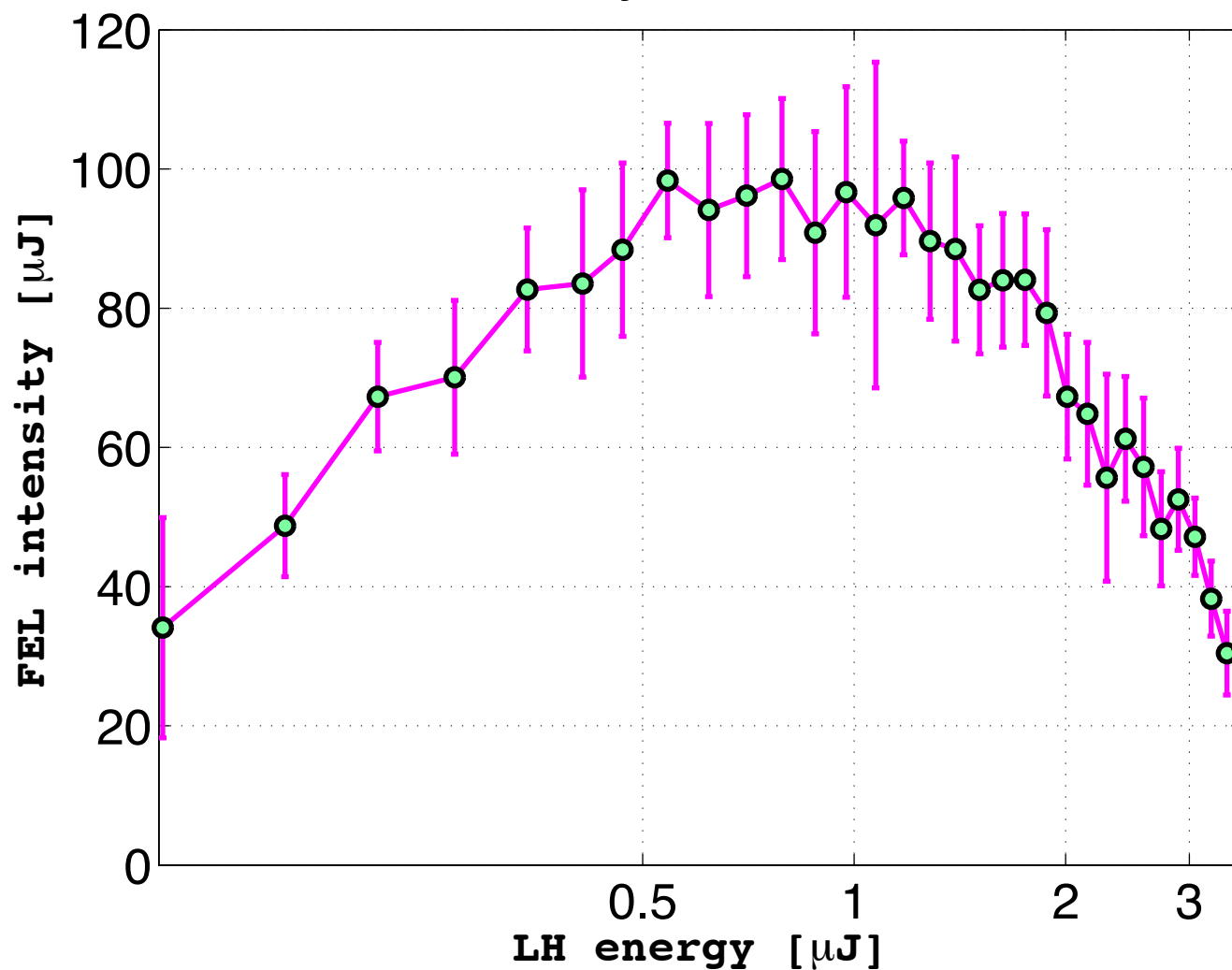
Experimental setup

Fermi scheme



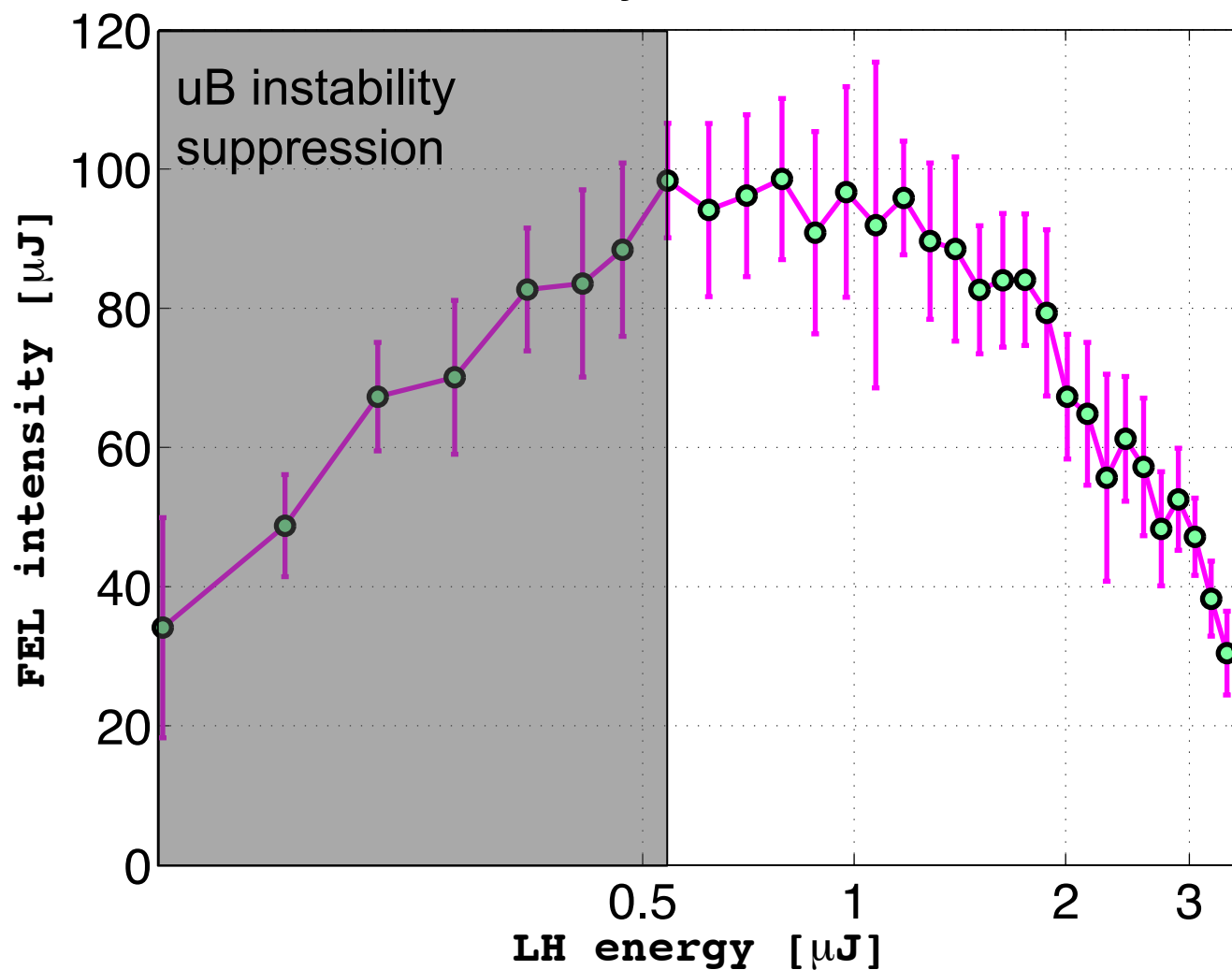
Laser heater – FEL intensity

FEL intensity vs. Laser Heater



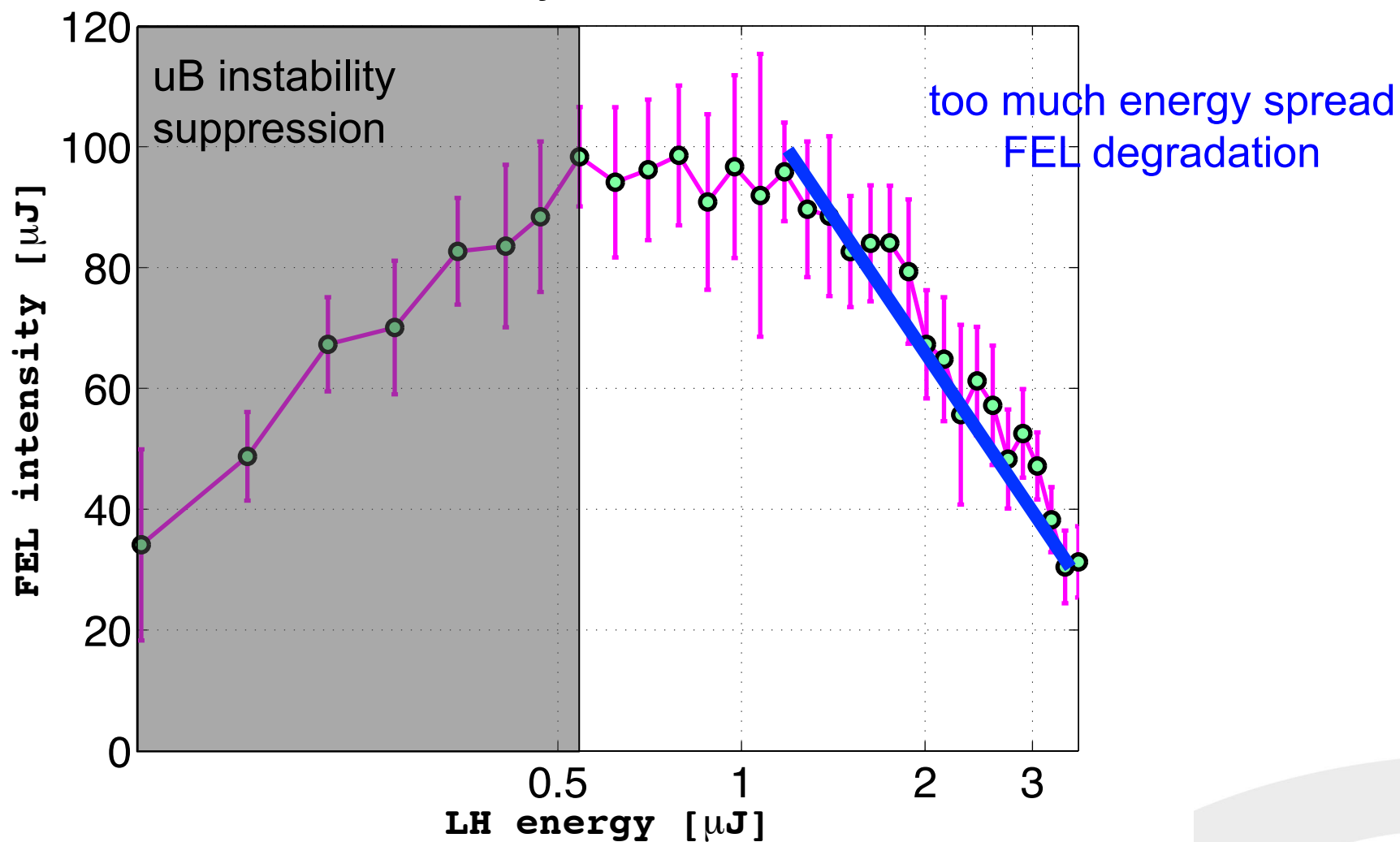
Laser heater – FEL intensity

FEL intensity vs. Laser Heater



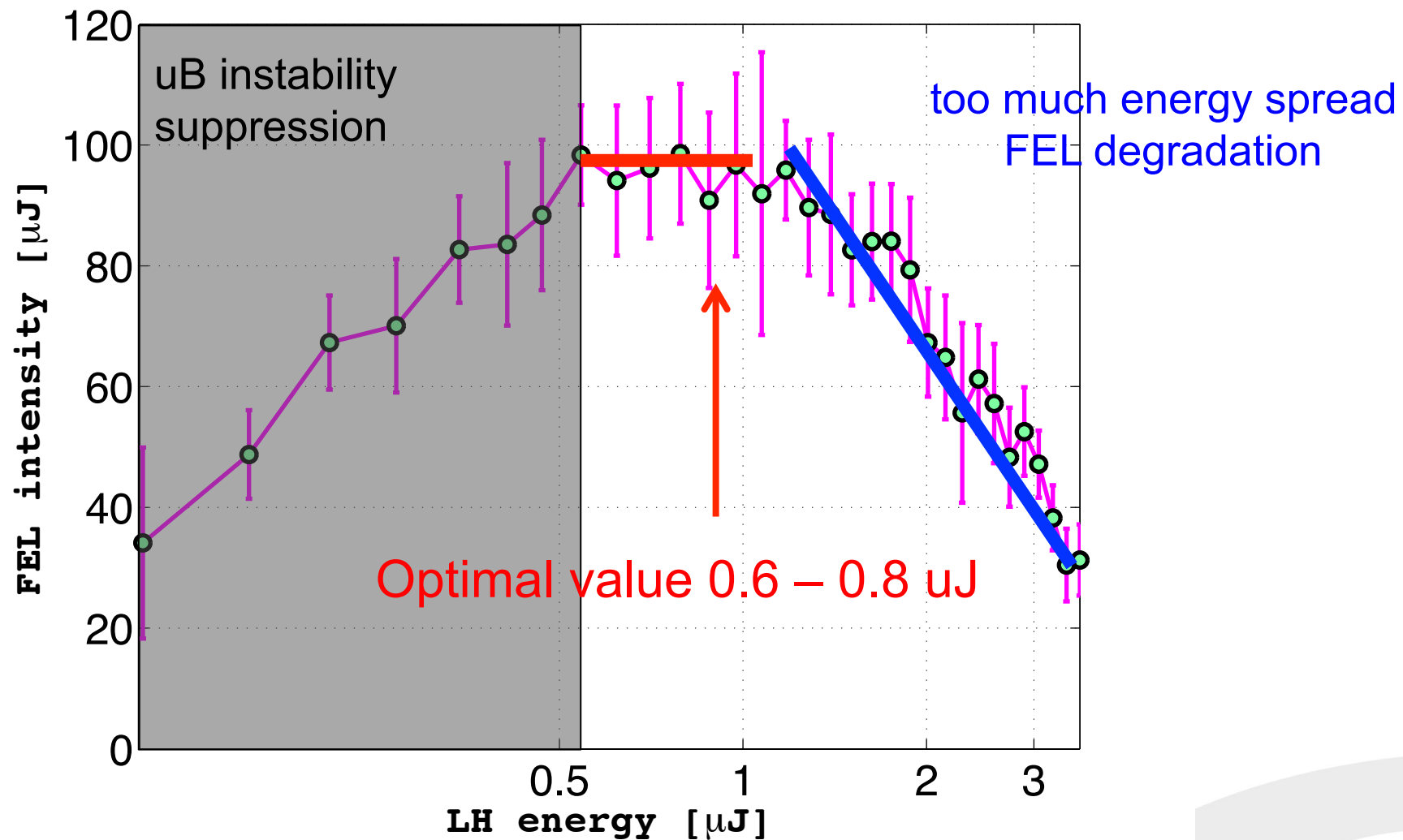
Laser heater – FEL intensity

FEL intensity vs. Laser Heater



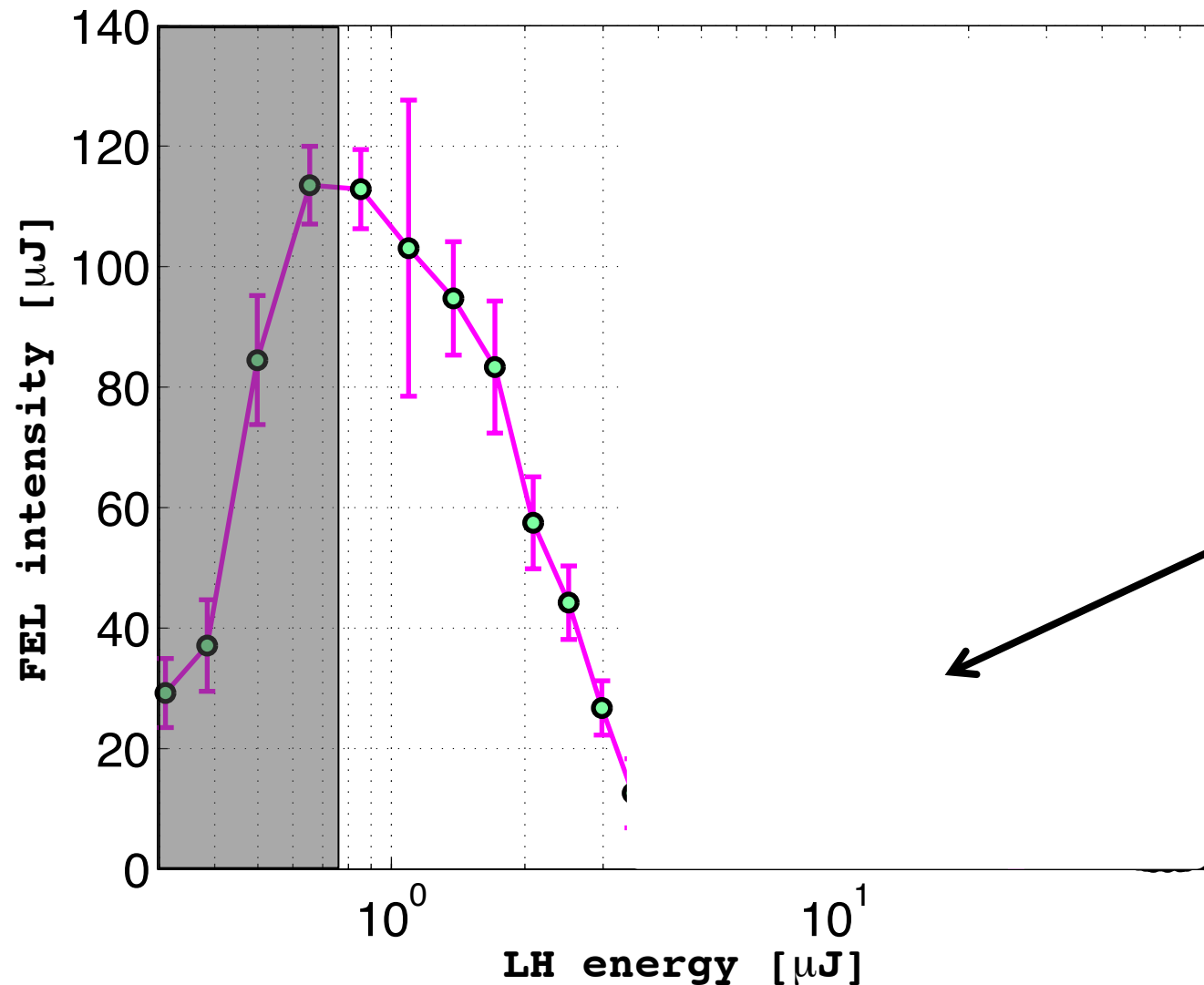
Laser heater – FEL intensity

FEL intensity vs. Laser Heater



But if we enlarge the scan range...

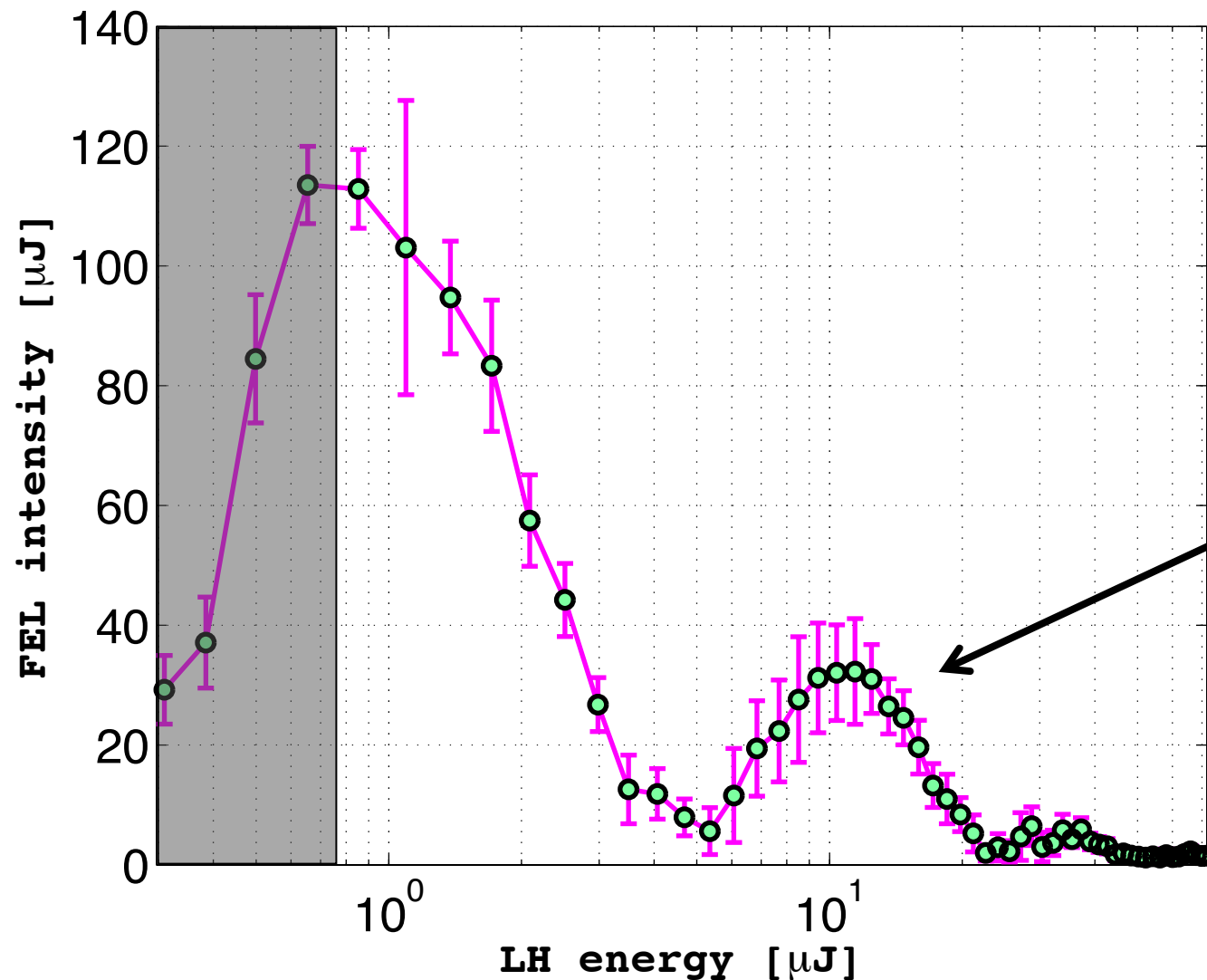
FEL intensity vs. heating at 32 nm



Appearance of local maxima at large heating (x10 with respect to optimal)

But if we enlarge the scan range...

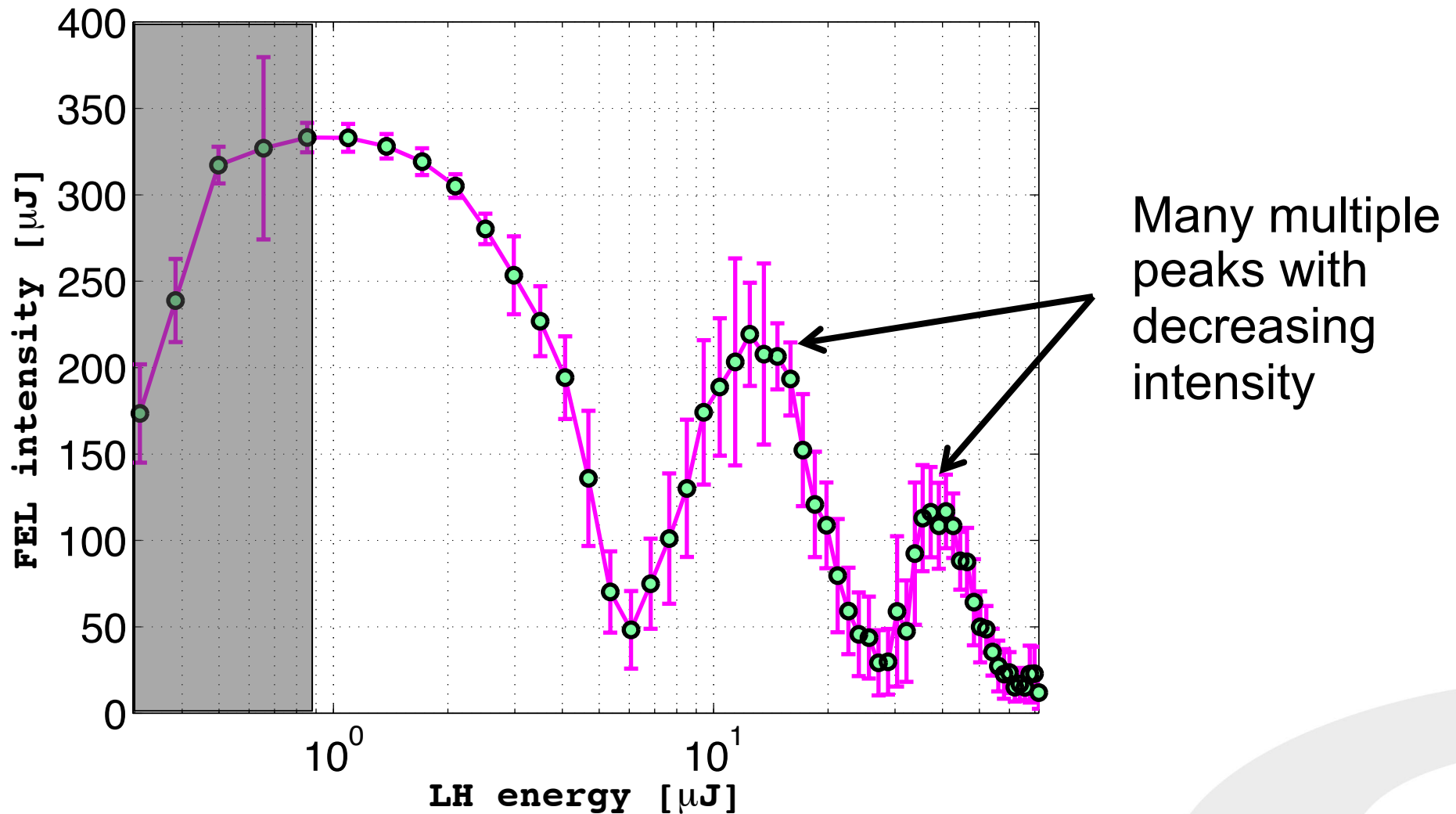
FEL intensity vs. heating at 32 nm



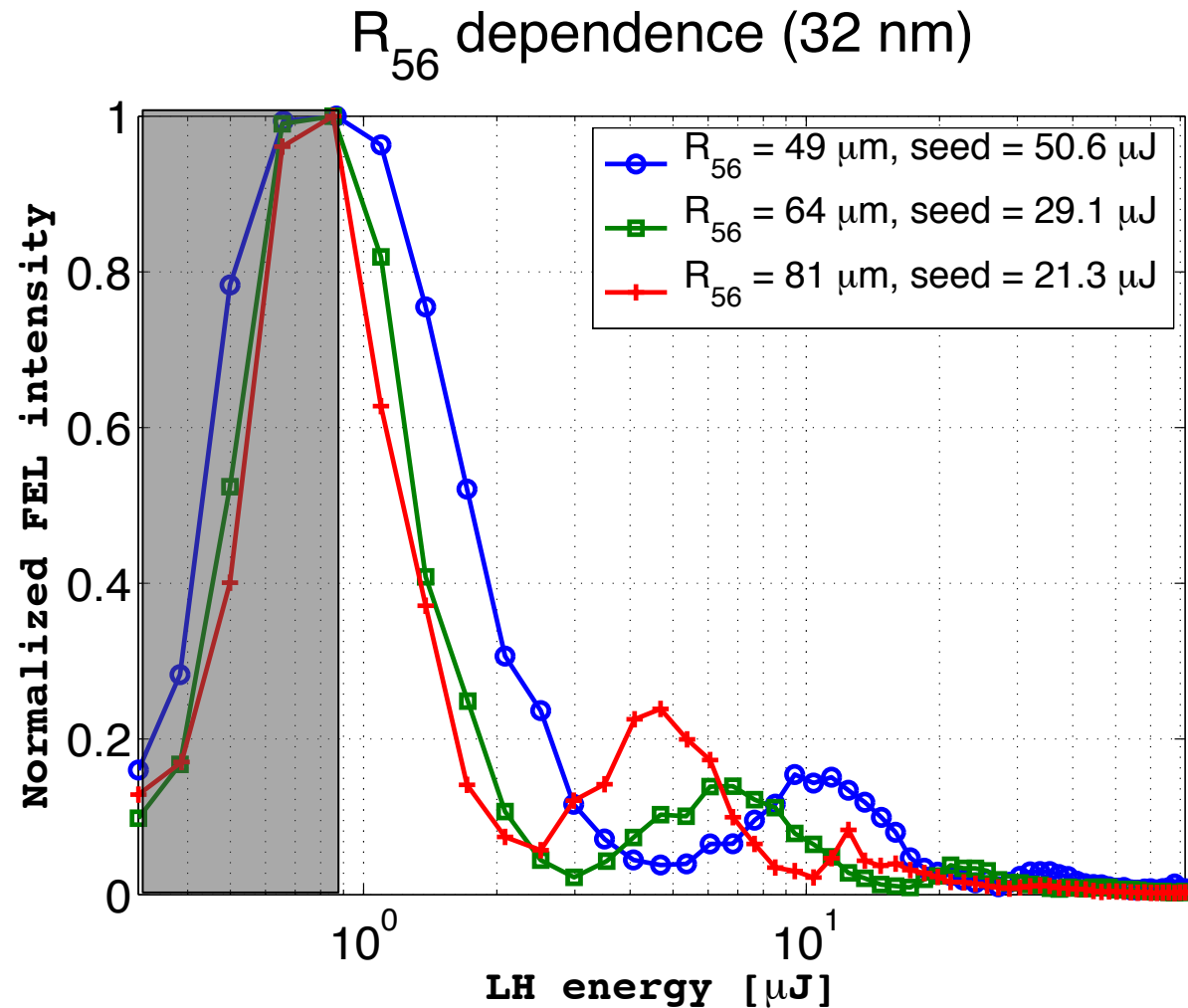
Appearance of local maxima at large heating (x10 with respect to optimal)

The effect can be dramatic!

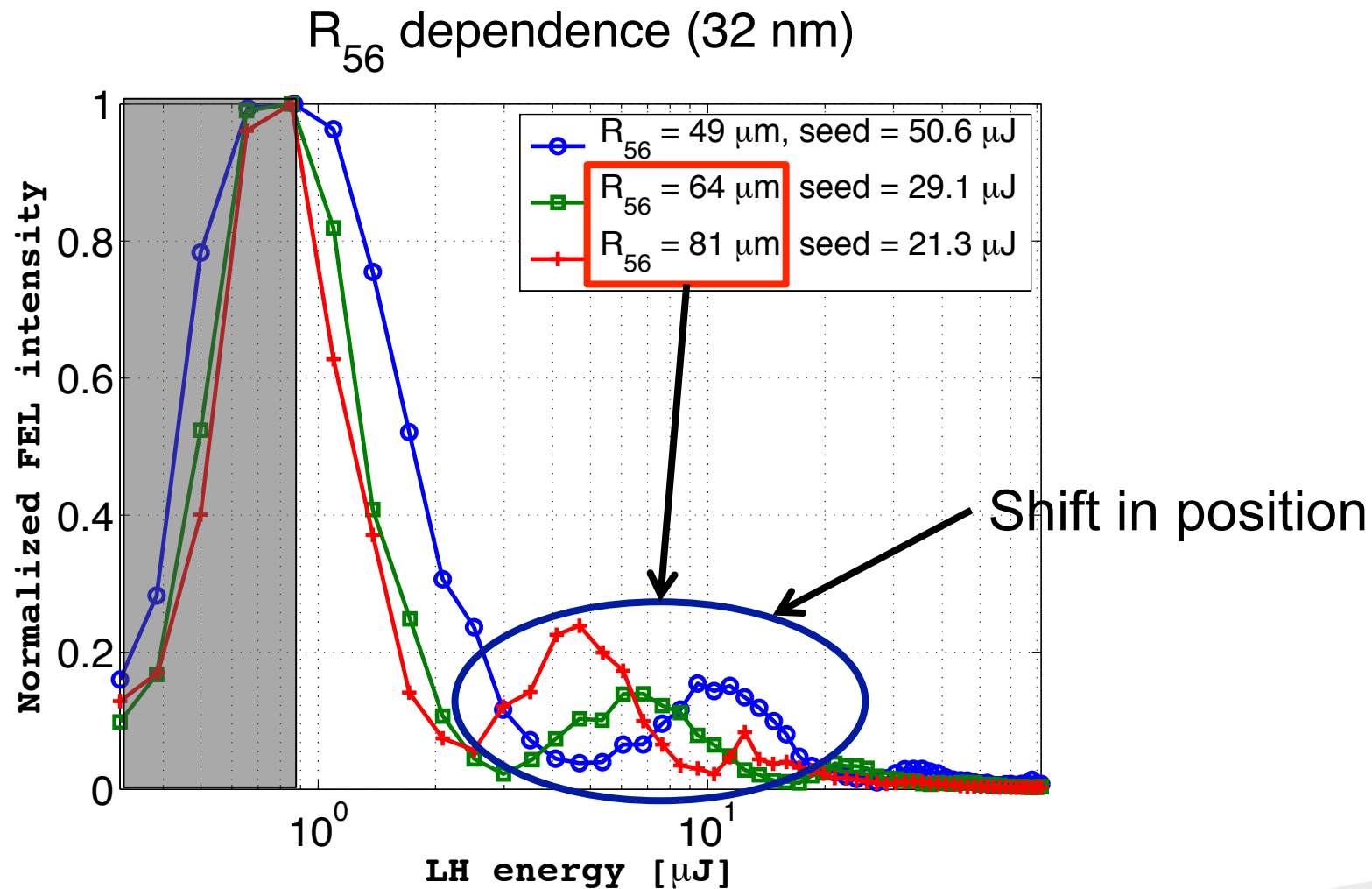
FEL intensity vs. heating at 52 nm



Dependence on R56 and seed

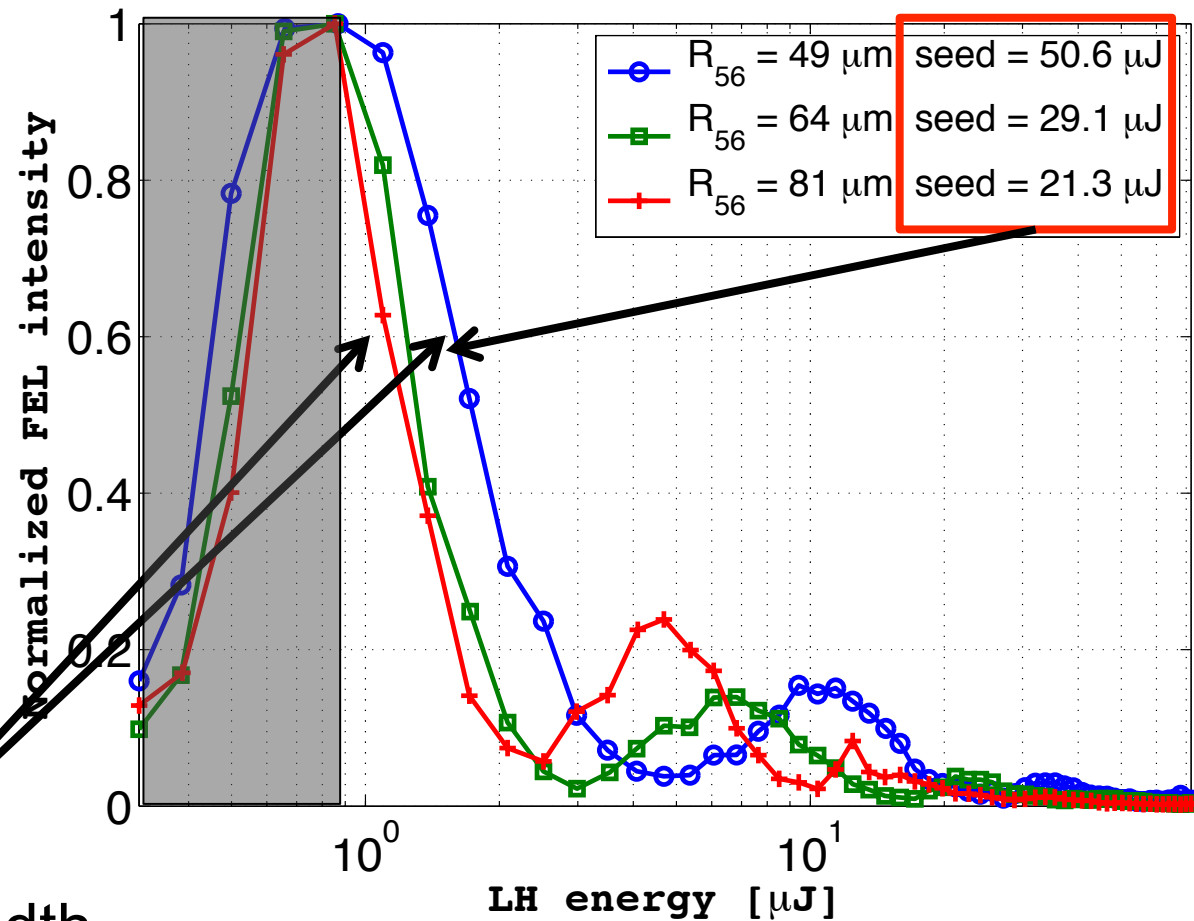


Dependence on R56 and seed

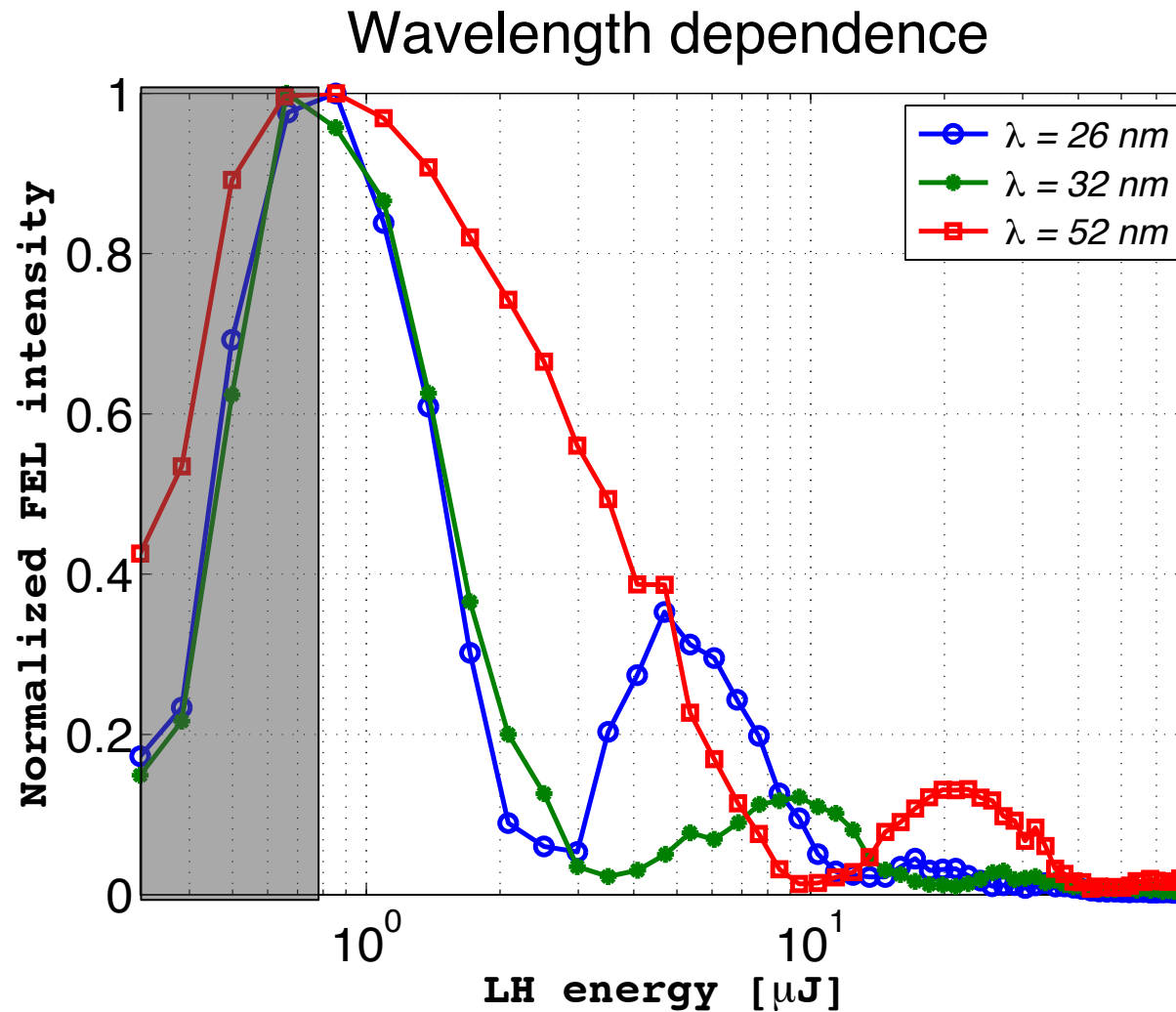


Dependence on R56 and seed

R_{56} dependence (32 nm)



Dependence on FEL wavelength



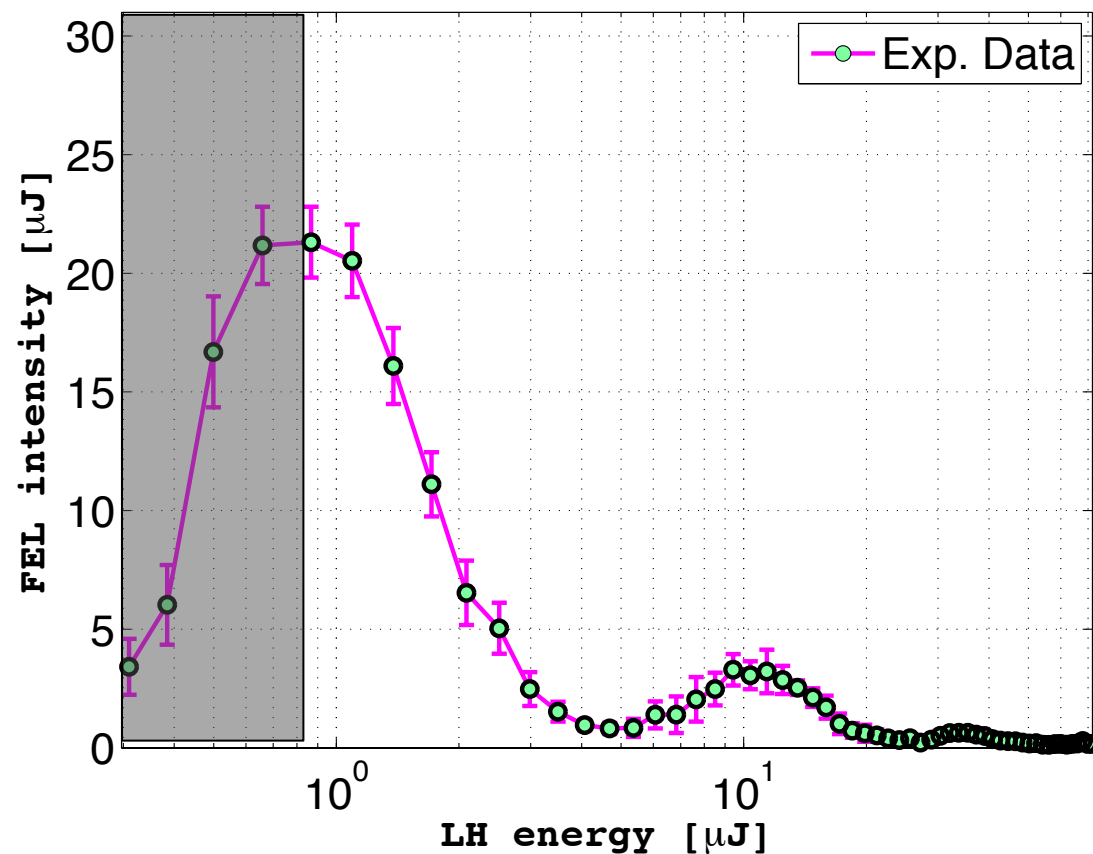
The number, position and relative intensity of secondary peaks can be tuned

How to explain this behaviour?

FEL vs. LH - Without gain

Without gain, the FEL intensity is almost **proportional** to the square of the bunching b_m

FEL intensity vs. heating (no gain, 32 nm)

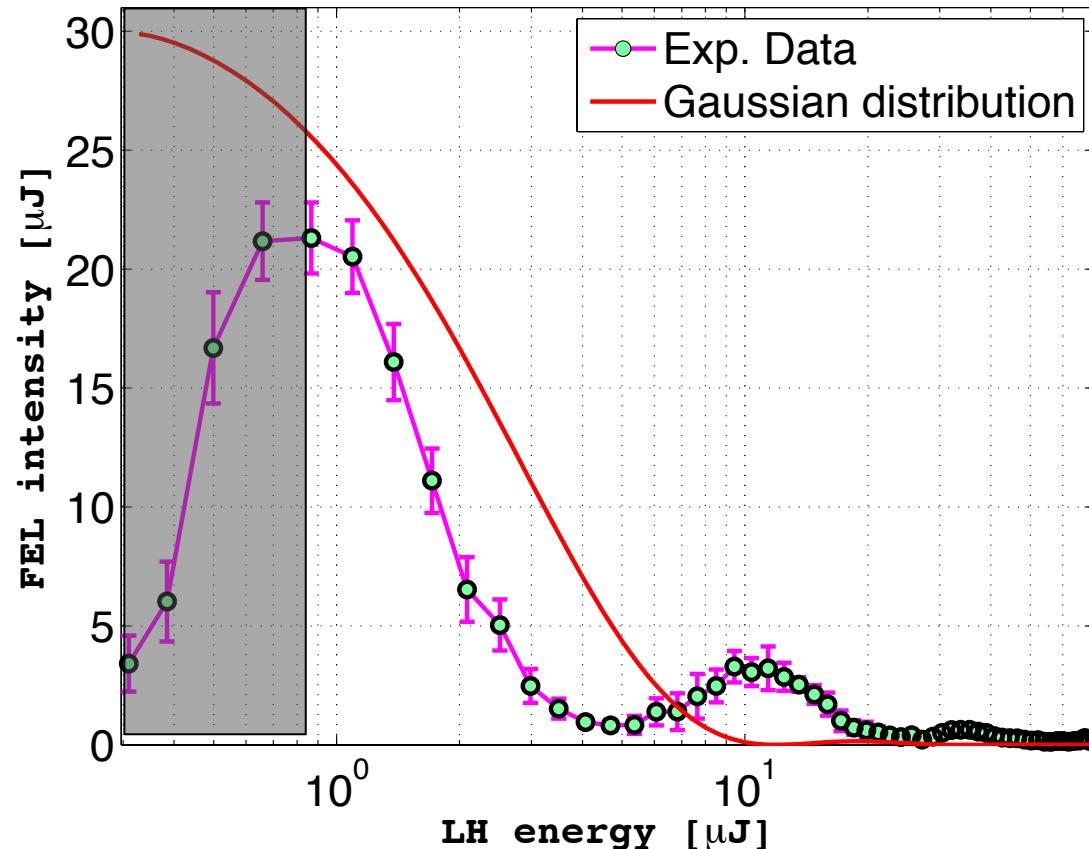


Coherent emission from
three radiators only

Bunching

$$(1) \quad b_m = \exp\left(-\frac{1}{2}m^2D^2\sigma_\gamma^2\right)J_m(mD\Delta\gamma)$$

FEL intensity vs. heating (no gain, 32 nm)



$$D = \frac{2\pi R_{56}}{\gamma_0 \lambda} \quad \text{Dispersion}$$

R_{56} Momentum compaction

γ_0 e^- energy

λ FEL wavelength

m Harmonic number

σ_γ Energy spread (rms)

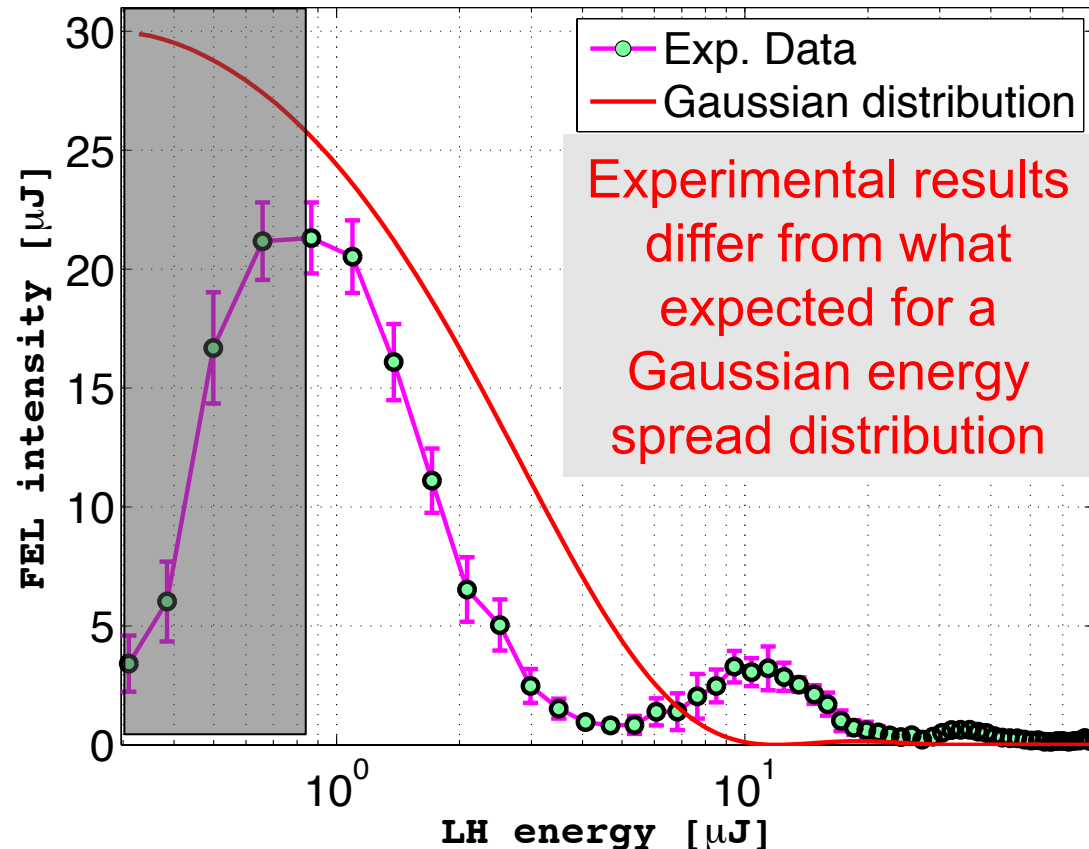
J_m m-th order Bessel

$\Delta\gamma$ FEL energy modulation

Bunching

$$(1) \quad b_m = \exp\left(-\frac{1}{2}m^2D^2\sigma_\gamma^2\right)J_m(mD\Delta\gamma)$$

FEL intensity vs. heating (no gain, 32 nm)



$$D = \frac{2\pi R_{56}}{\gamma_0 \lambda} \quad \text{Dispersion}$$

R_{56} Momentum compaction

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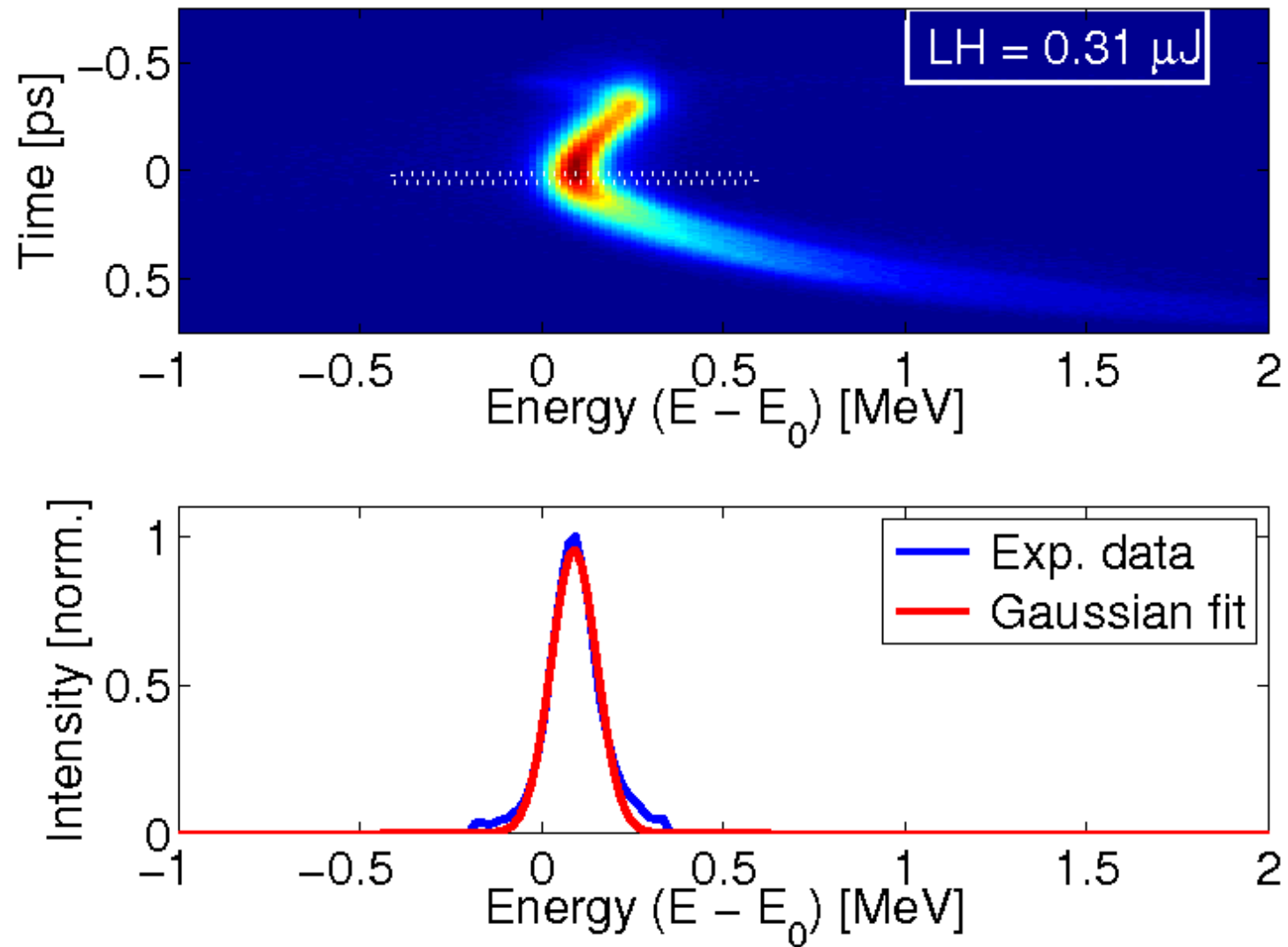
m Harmonic number

σ_γ Energy spread (rms)

J_m m-th order Bessel

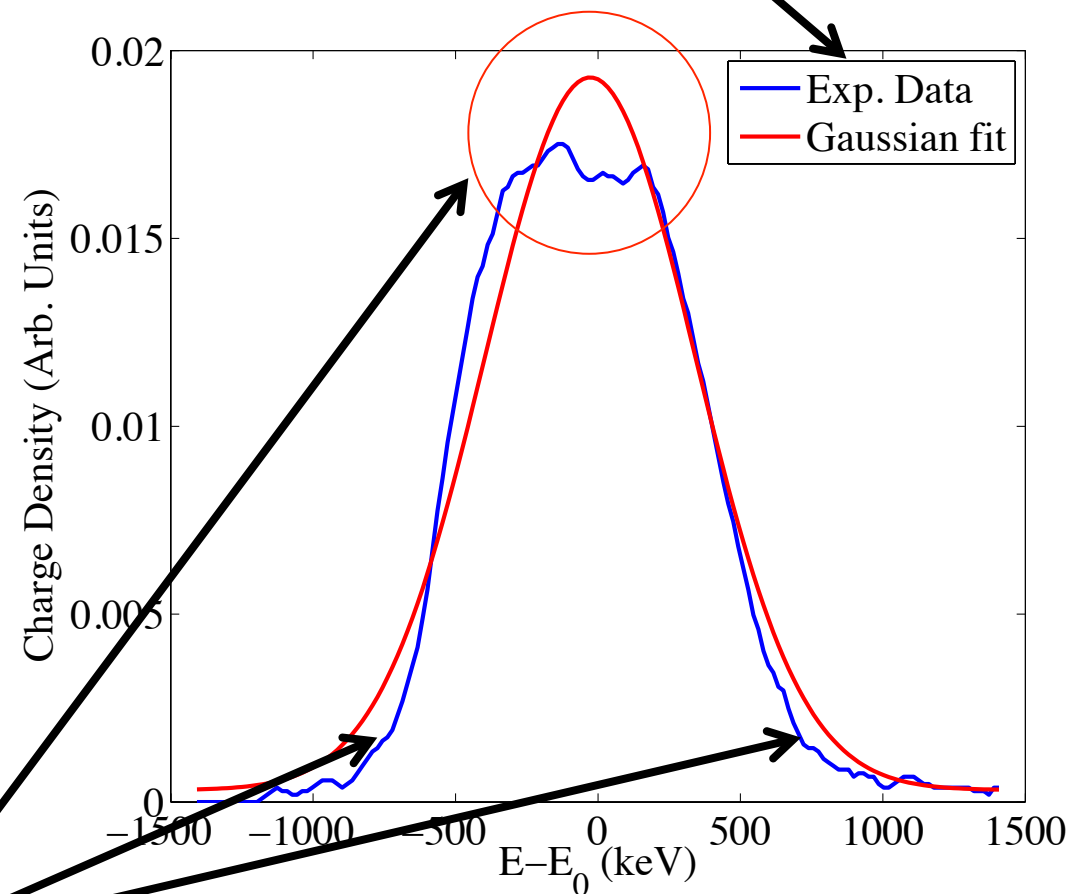
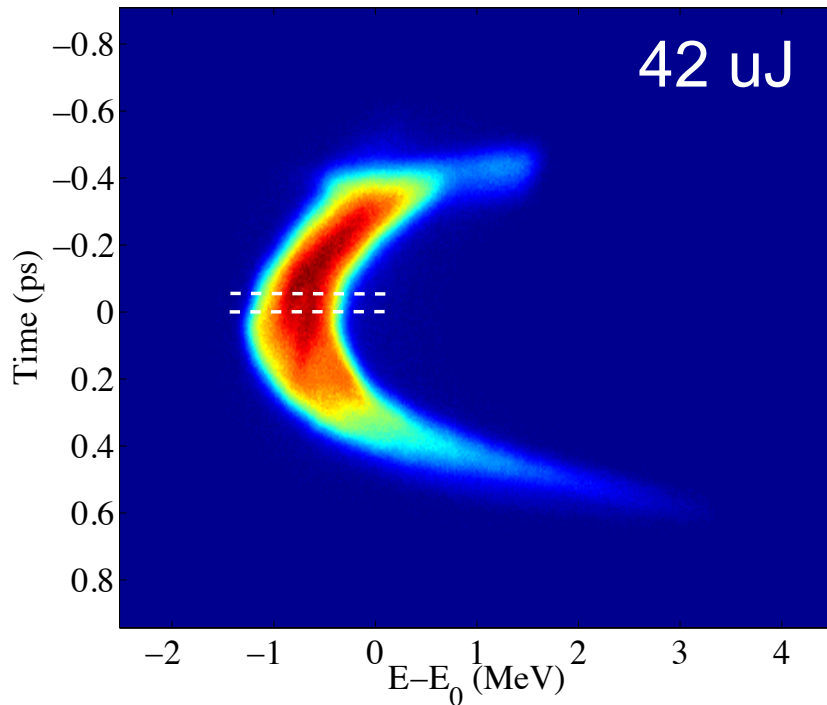
$\Delta\gamma$ FEL energy modulation

Longitudinal Phase Space and heating



Non-Gaussian energy spread

Same area and same second moment



Significant differences

Bunching with non-Gaussian energy spread

$$(2) \quad b_m = \exp\left(-\frac{1}{2}m^2 D^2 \sigma_\gamma^2\right) J_m(mD\Delta\gamma) S_L\left(mD\Delta\gamma, \frac{\sigma_r}{\sigma_x}\right) \quad \text{Z. Huang, PRSTAB 7, 074401 (2004)}$$

$$S_L(A, B) = \int R dR \exp\left(-\frac{R^2}{2}\right) J_0\left[A \exp\left(-\frac{R^2}{4B^2}\right)\right]$$

bunching suppression factor

$$D = \frac{2\pi R_{56}}{\gamma_0 \lambda}$$

m : Harmonic number

σ_γ Energy spread (rms)

J_m m -th order Bessel

$\Delta\gamma$ FEL energy modulation

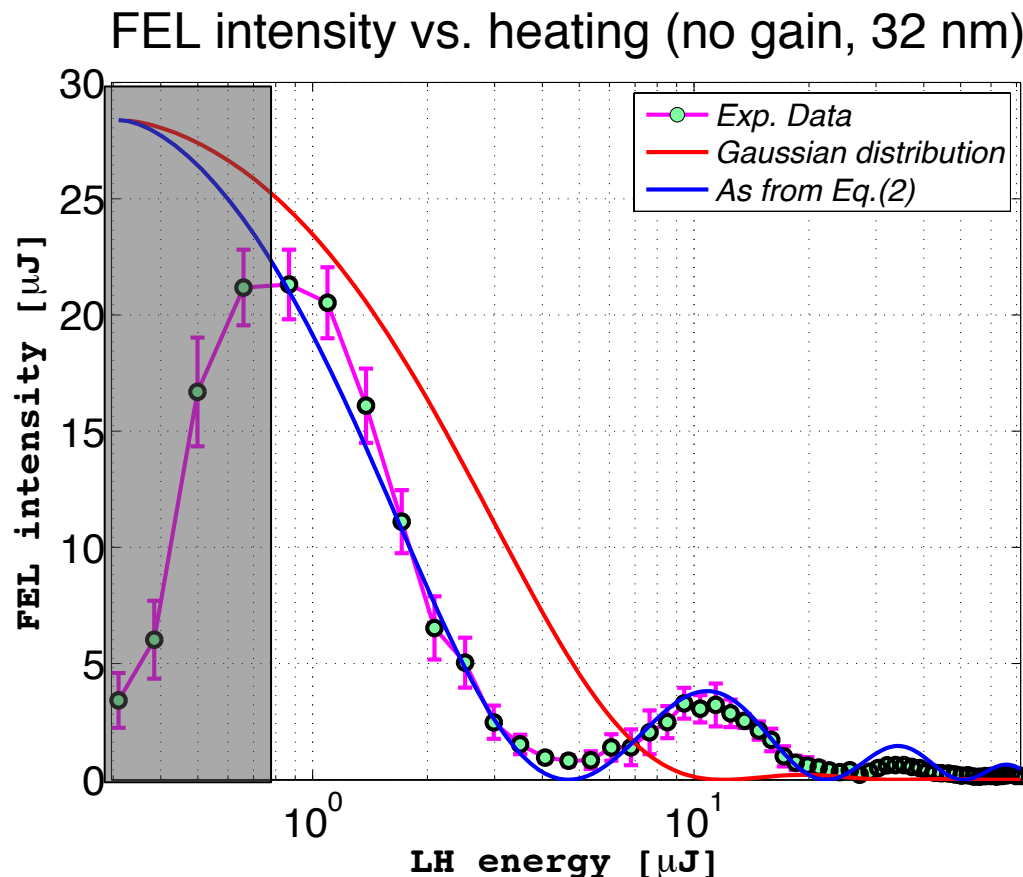
$$S_L(A, B) = \begin{cases} J_0(A), & \text{if } B \gg 1 \\ \frac{2J_1(A)}{A}, & \text{if } B = 1 \end{cases}$$

σ_r laser spot size (in LH)

σ_x e^- spot size (in LH)

Bunching with non-Gaussian energy spread

$$(2) \quad b_m = \exp\left(-\frac{1}{2}m^2D^2\sigma_\gamma^2\right)J_m(mD\Delta\gamma)S_L\left(mD\Delta\gamma, \frac{\sigma_r}{\sigma_x}\right) \quad \text{Z. Huang, PRSTAB 7, 074401 (2004)}$$



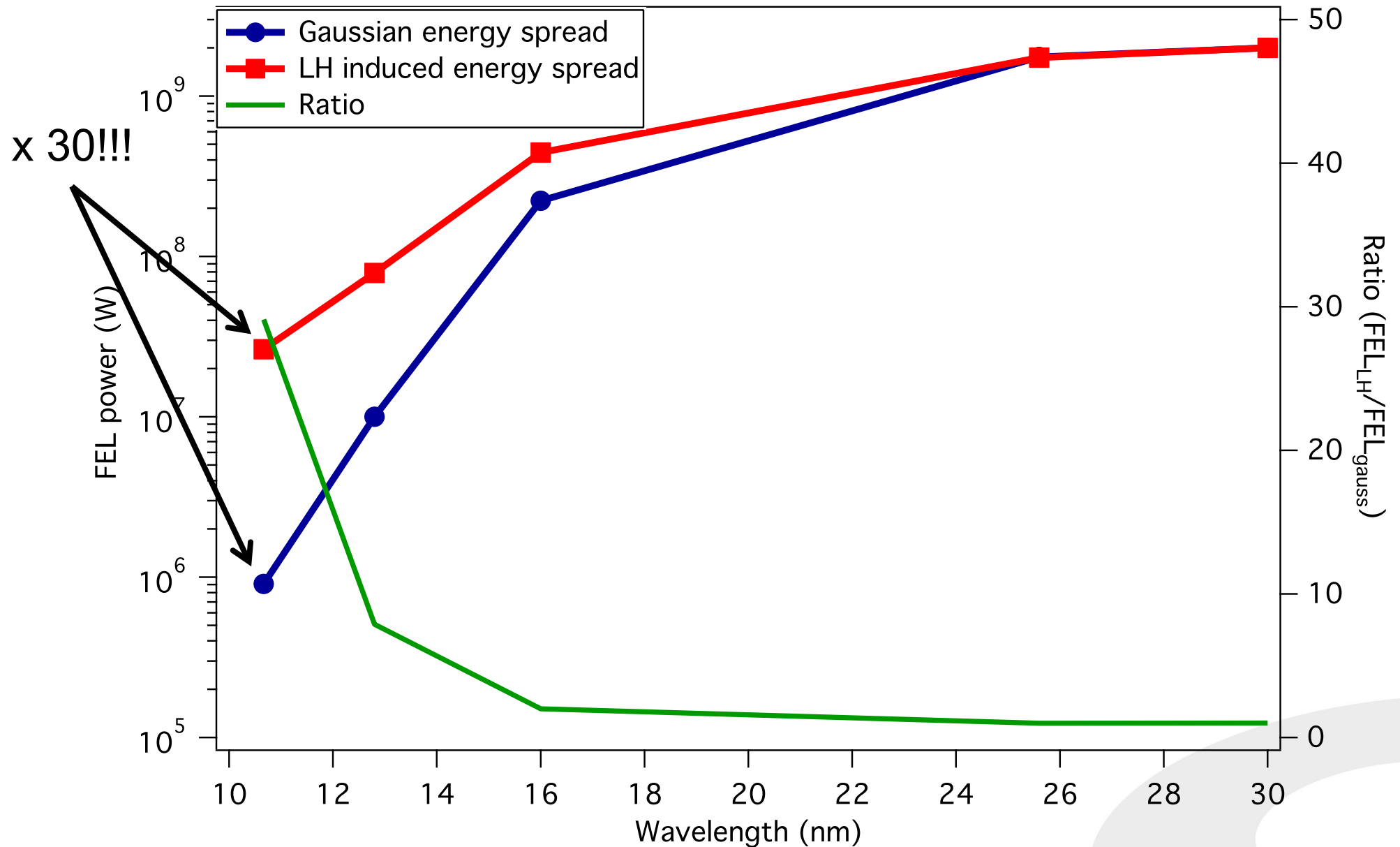
bunching suppression factor

$$S_L(A, B) = \begin{cases} J_0(A), & \text{if } B \gg 1 \\ \frac{2J_1(A)}{A}, & \text{if } B = 1 \end{cases}$$

σ_r laser spot size (in LH)

σ_x e^- spot size (in LH)

Simulated impact on high-harmonic emission



Summary

- ★ A Laser Heater is routinely used in FEL operations at FERMI.
- ★ The **non-Gaussian** distribution of the energy spread induced by the Laser Heater has been shown to be preserved up to the linac end and the undulators.
- ★ The shape of the slice energy spread distribution has a **significant** impact on FEL intensity, as it ultimately determines the bunching.
- ★ In particular, several FEL **local maxima** as a function of LH intensity have been observed, and can be controlled by tuning the machine parameters.
- ★ The unexpected behavior is well reproduced by previously developed LH theory.

- ★ Preliminary numerical simulations show that the non-Gaussian energy spread can **increase** the FEL power at high harmonic (i.e. shorter wavelength) in a HGHG FEL.
- ★ The significant increase in emission power could potentially **extend** the operation range of the **single cascade** HGHG scheme.

We acknowledge the support of the
FERMI COMMISSIONING TEAM



Thanks for your attention!

