

Estimation of longitudinal profiles of ion bunches in the LHC using Schottky-based diagnostics



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Paper ID: MOCO03



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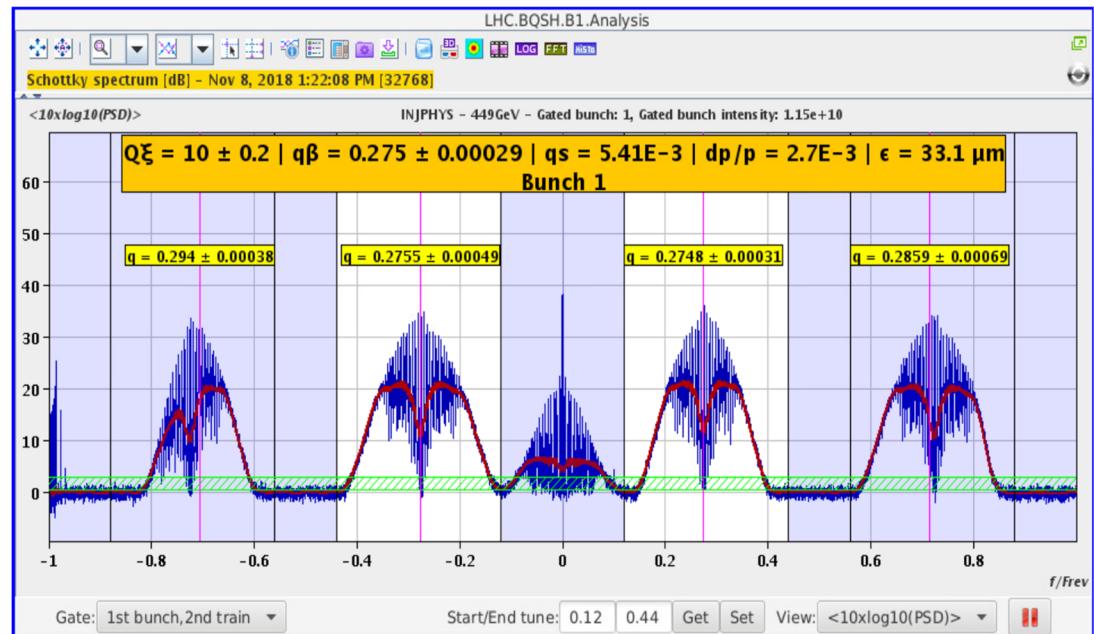
Schottky-based diagnostics

Traditionally used for:

- Average momentum
- Momentum spread
- Tune
- Chromaticity
- Intensity
- Emittance

Another application:

- Longitudinal bunch profiles





Motivation



What we want to achieve:

- Better understanding of the LHC Schottky spectra
 - Be able to simulate Schottky spectra
- Explore other capabilities of LHC Schottky monitors
 - Estimation of longitudinal bunch profiles



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Overview



- Longitudinal bunch profile
- LHC Schottky system
- Longitudinal Schottky spectrum in the LHC
- Results
- Conclusions and outlook



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Overview



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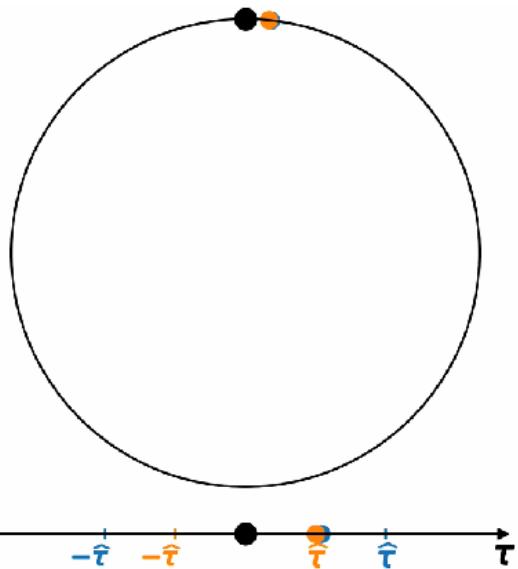


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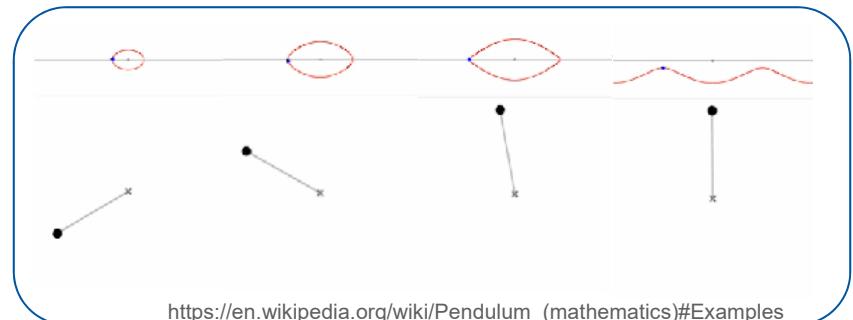
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Synchrotron motion – brief review

$$\tau_i = \hat{\tau}_i \sin(\omega_{s_i} t + \phi_{s_i}) \quad [\text{Boussard1986}]$$



$$\omega_{s_i} \simeq \left(1 - \frac{(h\omega_0 \hat{\tau}_i)^2}{16} \right) \omega_{s_0}$$



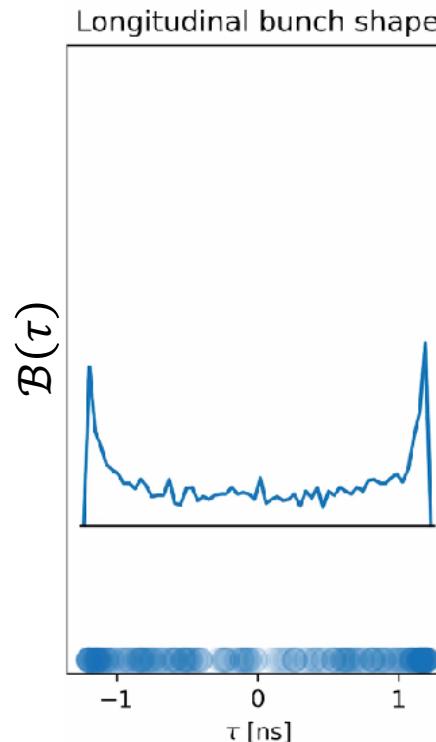
[https://en.wikipedia.org/wiki/Pendulum_\(mathematics\)#Examples](https://en.wikipedia.org/wiki/Pendulum_(mathematics)#Examples)

- Each particle performs simple harmonic motion
- Synchrotron frequency, ω_{s_i} , depends on synchrotron amplitude, $\hat{\tau}_i$

Synchrotron motion – single particle distribution

$$g_{\tau}(\tau|\hat{\tau}) = \frac{1}{\pi\sqrt{\hat{\tau}^2 - \tau^2}}$$

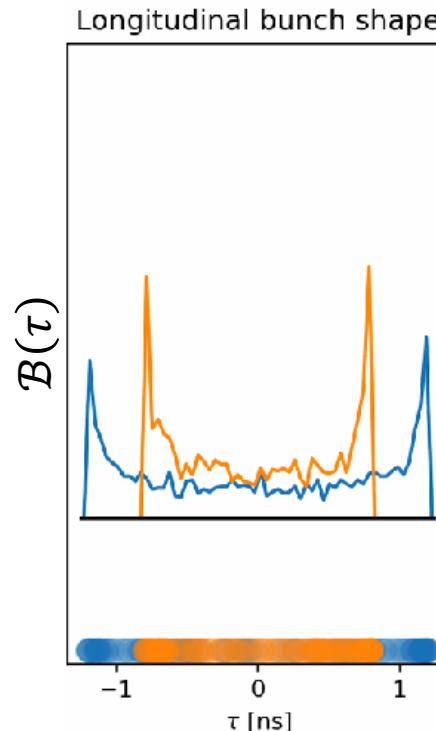
Single particle distribution



Synchrotron motion – single particle distribution

$$g_{\tau}(\tau|\hat{\tau}) = \frac{1}{\pi\sqrt{\hat{\tau}^2 - \tau^2}}$$

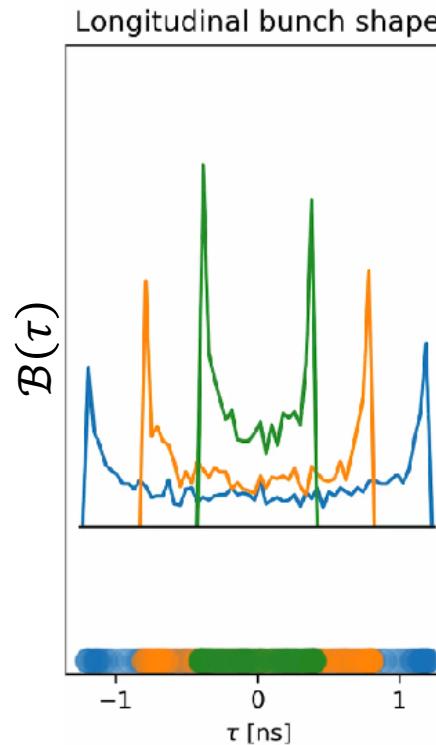
Single particle distribution



Synchrotron motion – single particle distribution

$$g_{\tau}(\tau|\hat{\tau}) = \frac{1}{\pi\sqrt{\hat{\tau}^2 - \tau^2}}$$

Single particle distribution



Longitudinal bunch profile

Without intra-bunch coherent motion:

$$\mathcal{B}(\tau) = \int_{|\tau|}^{\infty} \frac{1}{\pi\sqrt{\hat{\tau}^2 - \tau^2}} g_{\hat{\tau}}(\hat{\tau}) d\hat{\tau}$$

Longitudinal bunch profile

$g_{\hat{\tau}}(\hat{\tau})$ - Amplitude distribution

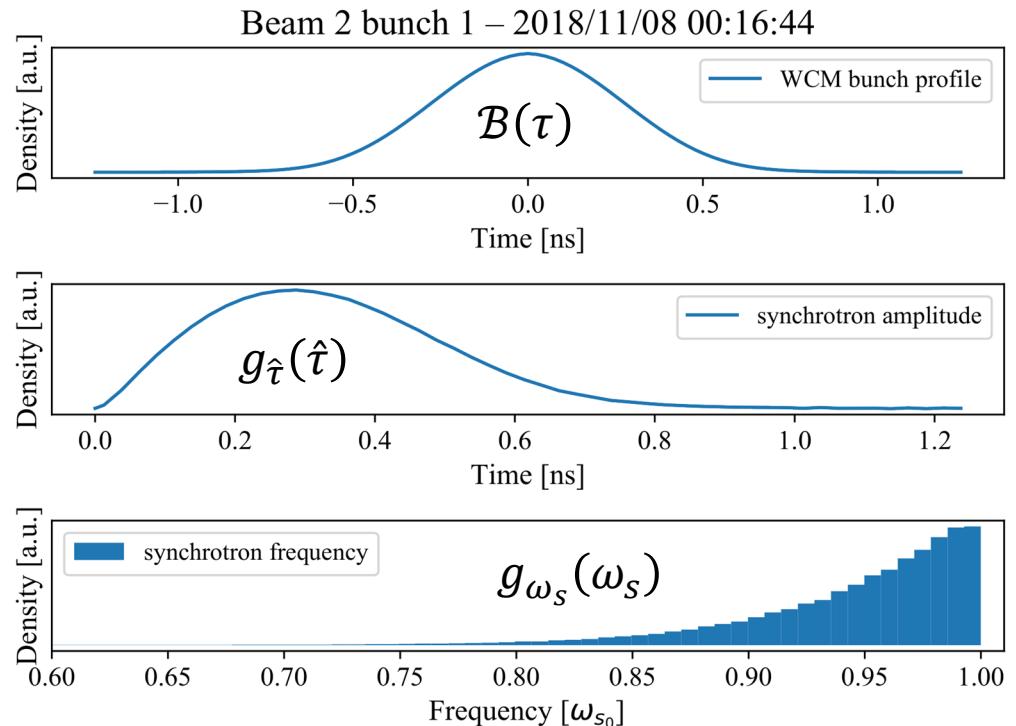
$g_{\tau}(\tau|\hat{\tau}) = \frac{1}{\pi\sqrt{\hat{\tau}^2 - \tau^2}}$

Single particle distribution

Equivalence of distributions

$$\mathcal{B}(\tau) = \int_{|\tau|}^{\infty} \frac{g_{\hat{\tau}}(\hat{\tau})}{\pi \sqrt{\hat{\tau}^2 - \tau^2}} d\hat{\tau}$$

$$\omega_s \approx \left(1 - \frac{(h\omega_0 \hat{\tau})^2}{16}\right) \omega_{s_0}$$





Overview



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- **LHC Schottky system**
- Longitudinal Schottky spectrum in the LHC
- Results
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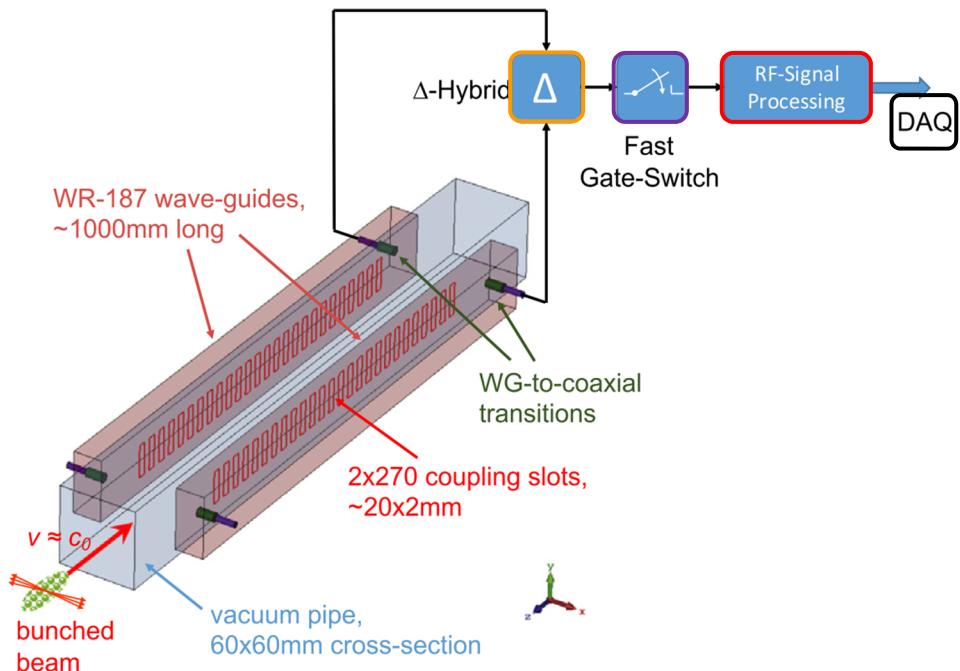


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LHC Schottky systems



- 4 transverse (Δ) systems (H/V, B1/B2)
- Longitudinal (Σ) signal not fully suppressed

- Flexible bunch gating

- Down-mix 4.8GHz → 11kHz
- Final BW ≈ 15kHz

- ADC + FFT

$F_s = 44\text{kHz}$
 FFT len ~ 1.5s
 Averaging ~ 1.5min
 Update rate = 1Hz

Figure courtesy of M. Betz et al, NIM 2017



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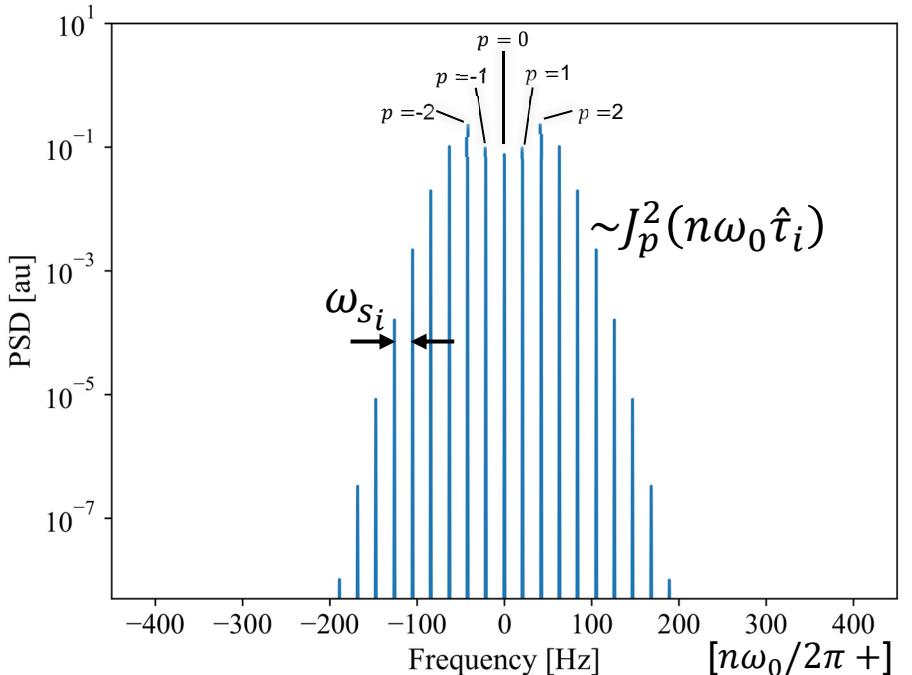
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Single particle longitudinal Schottky spectrum

Single particle contribution
to sum (Σ) signal

$$s_i(t) = \sum_{p=-\infty}^{\infty} j^p J_p(n\omega_0 \hat{t}_i) e^{j[(n\omega_0 + p\omega_{s_i})t + p\phi_{s_i}]}$$

[Boussard 1986]



Longitudinal Schottky spectrum

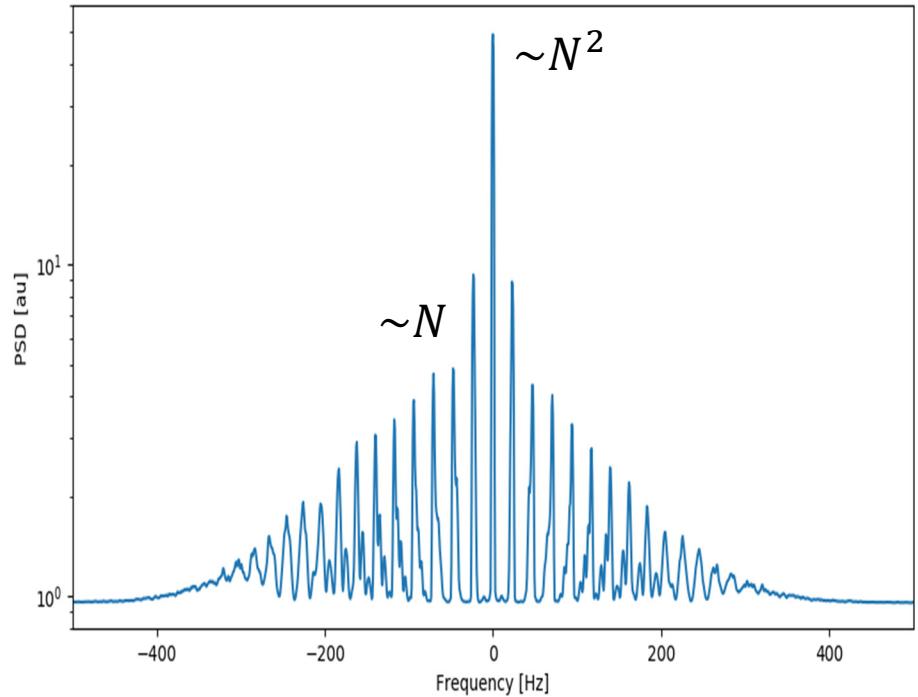
Total Σ signal: $s(t) = \sum_{i=1}^N s_i(t)$

Instantaneous Schottky spectrum is not deterministic!
 ➤ depends on the instantaneous phases → random
 ... so we average ...

No intra-bunch coherent motion

$$\langle P(\omega) \rangle \sim N \langle P_i(\omega, \hat{t}_i) \rangle \sim N \int_0^{\infty} g_{\hat{t}}(\hat{t}) P(\omega, \hat{t}) d\hat{t}$$

Synch amp pdf
Single particle spectrum





Matrix formalism

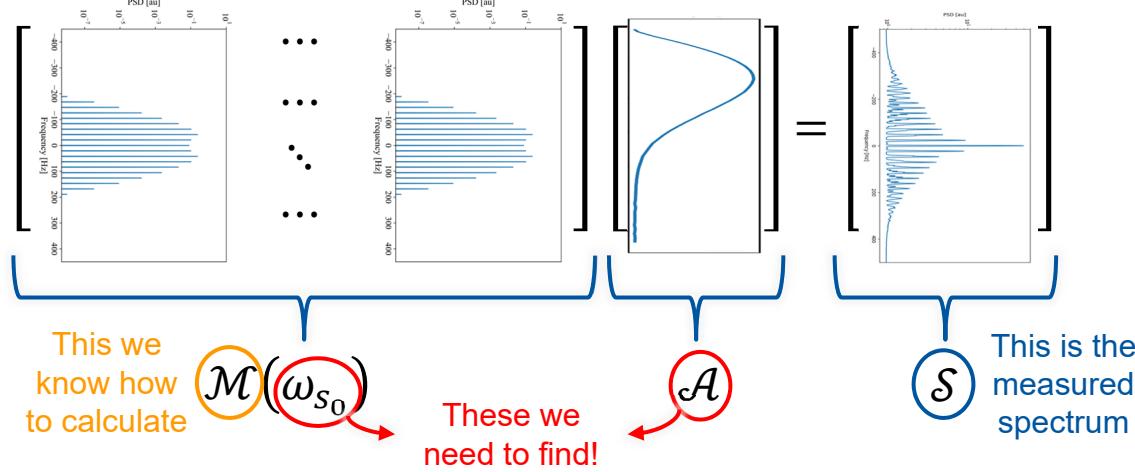
$$\left[\begin{array}{ccc} P(\omega_1, \hat{\tau}_1) & \cdots & P(\omega_1, \hat{\tau}_n) \\ P(\omega_2, \hat{\tau}_1) & \cdots & P(\omega_2, \hat{\tau}_n) \\ \vdots & \ddots & \vdots \\ P(\omega_m, \hat{\tau}_1) & \cdots & P(\omega_m, \hat{\tau}_n) \end{array} \right] \left[\begin{array}{c} g_{\hat{\tau}}(\hat{\tau}_1) \\ g_{\hat{\tau}}(\hat{\tau}_2) \\ \vdots \\ g_{\hat{\tau}}(\hat{\tau}_n) \end{array} \right] = \left[\begin{array}{c} \langle P(\omega_1) \rangle \\ \langle P(\omega_2) \rangle \\ \vdots \\ \langle P(\omega_m) \rangle \end{array} \right]$$

$\mathcal{M}(\omega_{s_0})$ \mathcal{A} \mathcal{S}

- \mathcal{M} – cols are the single particle spectra
- \mathcal{A} – synchrotron amplitude distribution
- \mathcal{S} – averaged power spectral density



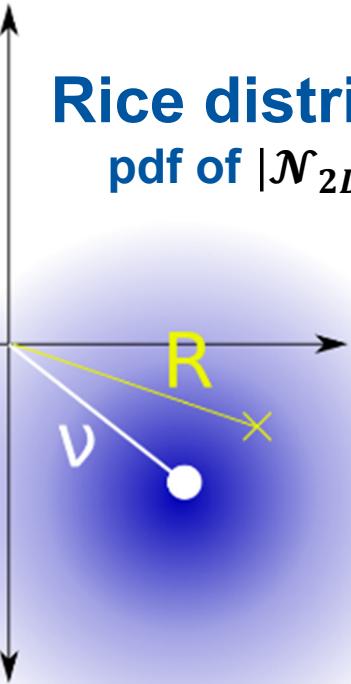
Matrix formalism



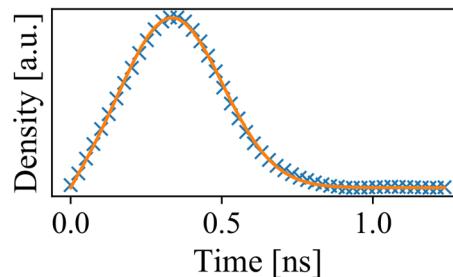
- \mathcal{M} – cols are the single particle spectra
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- \mathcal{S} – averaged power spectral density

Amplitude distribution in the LHC

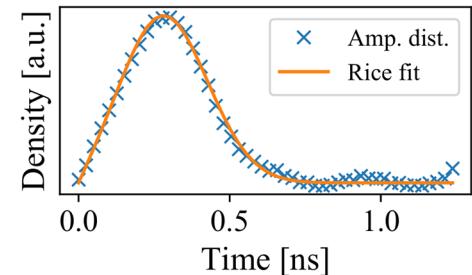
Rice distribution
 pdf of $|\mathcal{N}_{2D}(\nu, \sigma)|$



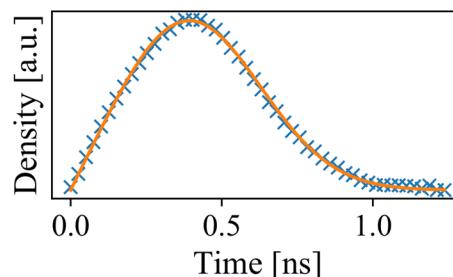
Beam 2 bunch 31171 RAMP
 2018/11/07 17:49:41



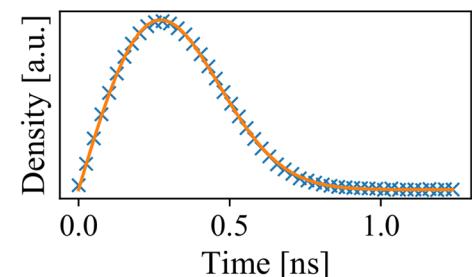
Beam 1 bunch 1 STABLE
 2018/10/26 09:44:36



Beam 1 bunch 1 INJPHYS
 2018/11/27 23:25:52



Beam 2 bunch 1 ADJUST
 2018/11/08 00:16:44





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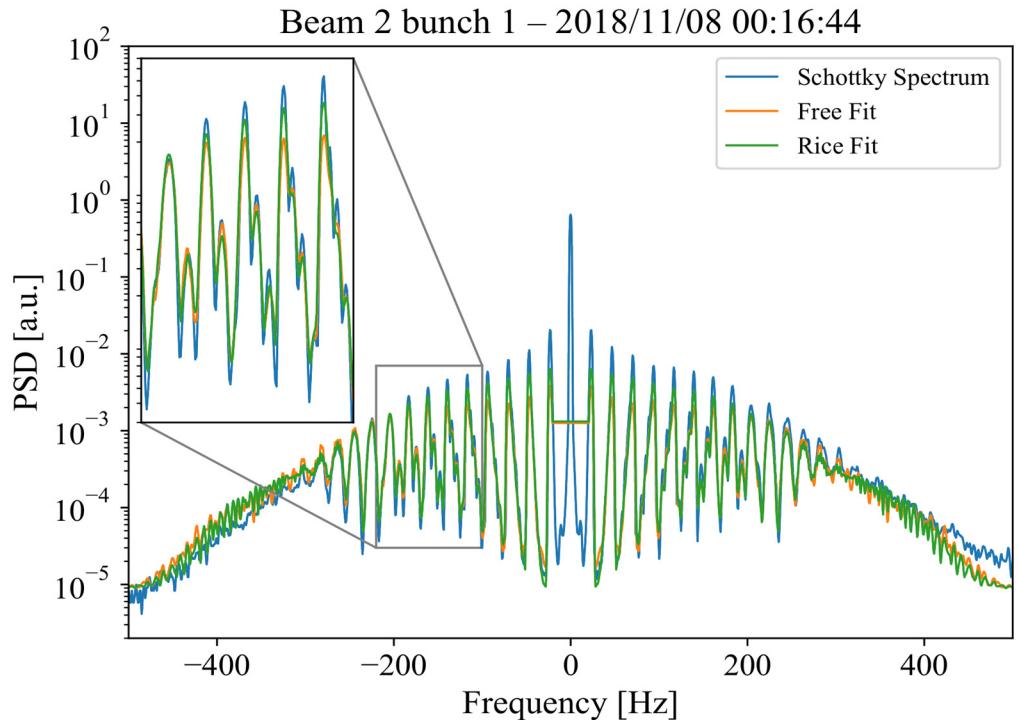
Fitting the spectrum

Minimise cost function:

$$C(\omega_{s_0}, \mathcal{A}) = |\log[\mathcal{M}(\omega_{s_0})\mathcal{A}] - \log[\mathcal{S}]|^2$$

Beam parameters:

- Pb⁸²⁺
- E = 6.37 TeV
- I_{beam} = 9.4×10^{10} charges
- I_{gate} = 5×10^9 charges



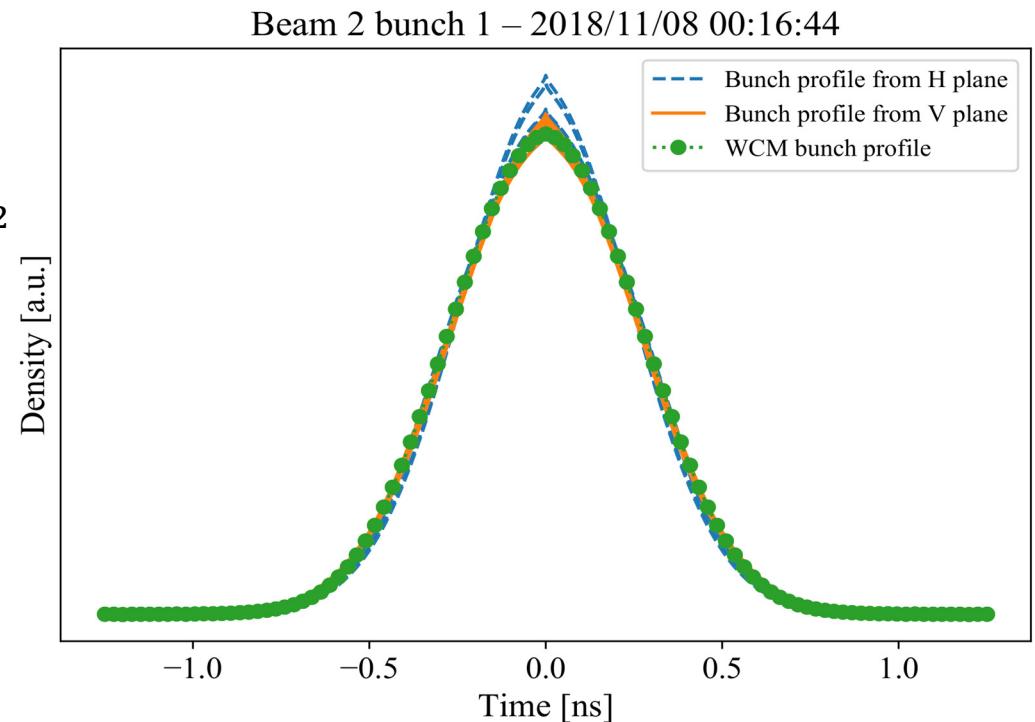
Estimated bunch profiles

Minimise cost function:

$$C(\omega_{s_0}, \mathcal{A}) = \left| \log[\mathcal{M}(\omega_{s_0})\mathcal{A}] - \log[\mathcal{S}] \right|^2$$

Beam parameters:

- Pb⁸²⁺
- E = 6.37 TeV
- I_{beam} = 9.4 × 10¹⁰ charges
- I_{gate} = 5 × 10⁹ charges





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Conclusions

- Fitted spectrum similar to experimental one
- Profiles from H&V channels are consistent
- Profiles match WCM measurements
- Reconstruction of longitudinal profile from narrowband measurements
- Developed mathematical framework for simulating Schottky spectra





Outlook



- Important step towards a better understanding of the LHC Schottky spectra
 - Mathematical framework can be adapted to transverse signals (e.g. estimate chromaticity and tune)
 - Obtained bunch profiles can be compared with WCMs and serve as a validation of the estimation of other parameters
- Extend mathematical framework to include the treatment of coherent motion
 - Can be applied e.g. during injection oscillations



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ACKNOWLEDGEMENTS



- T. Argyropoulos
- O. R. Jones
- T. Lefèvre
- T. Levens
- O. Marqversen
- M. Wendt

Thank you for your attention!



SPARE SLIDES



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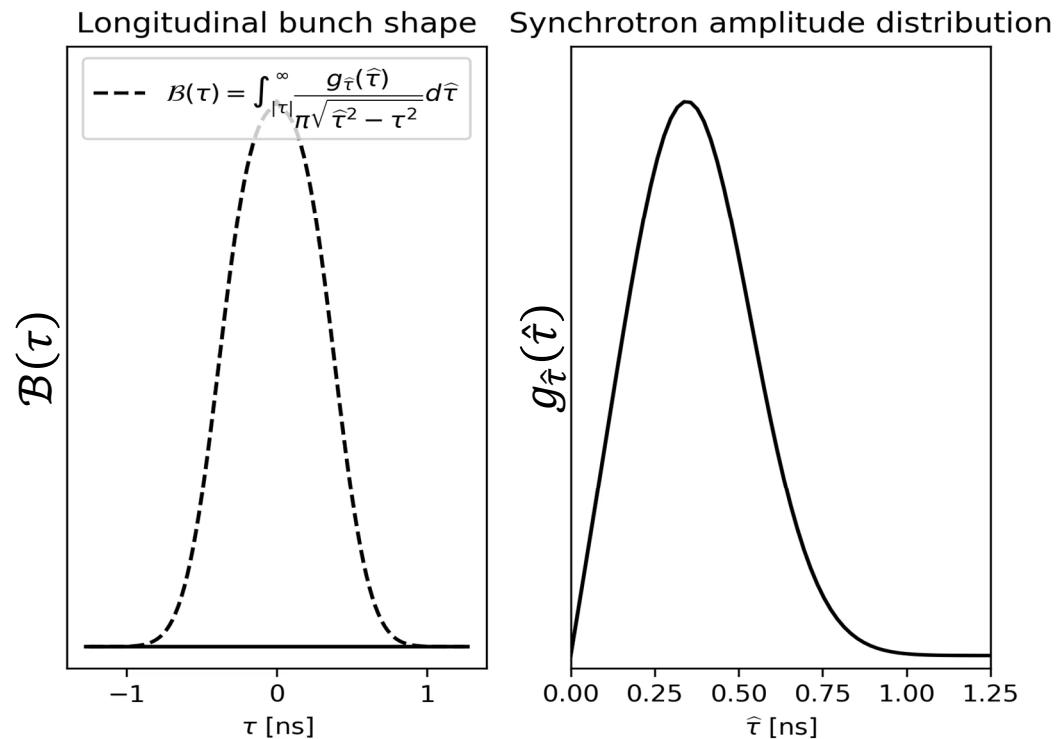
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Typical LHC bunch profile



$$\mathcal{B}(\tau) = \int_{|\tau|}^{\infty} \frac{g_{\hat{\tau}}(\hat{\tau})}{\pi\sqrt{\hat{\tau}^2 - \tau^2}} d\hat{\tau}$$



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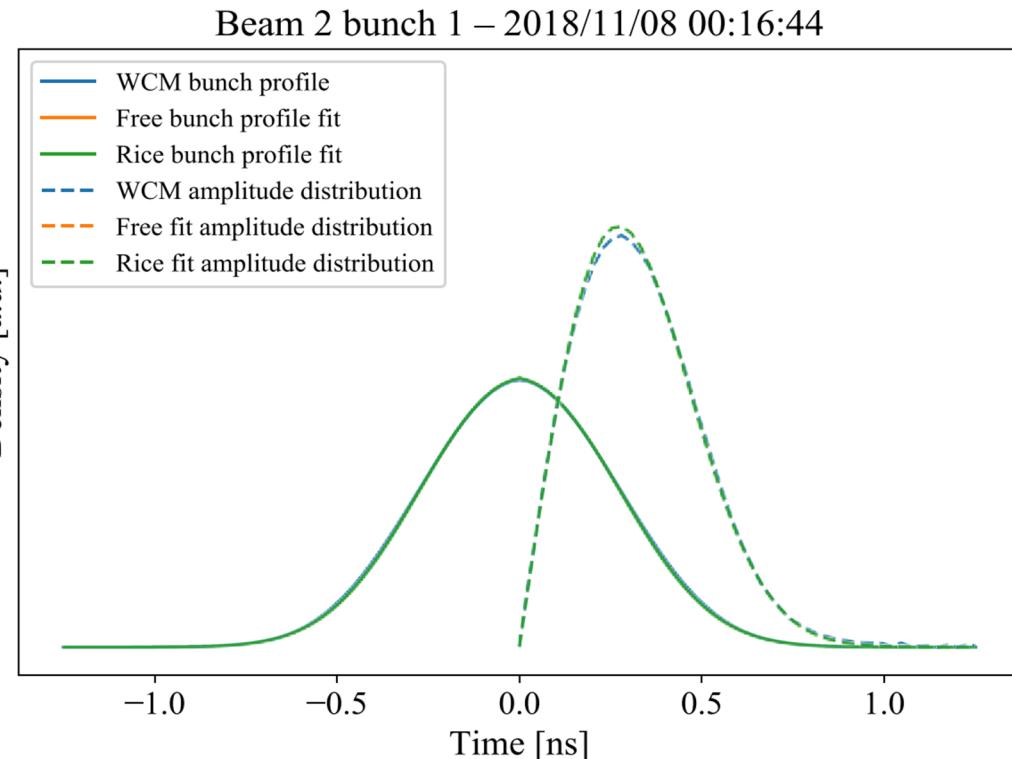
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Optimisation process

Minimise cost function:

$$C(\omega_{s_0}, \mathcal{A}) = |\log[\mathcal{M}(\omega_{s_0})\mathcal{A}] - \log[\mathcal{S}]|^2$$

Optimization process uses a *differential evolution* algorithm for global minimisation

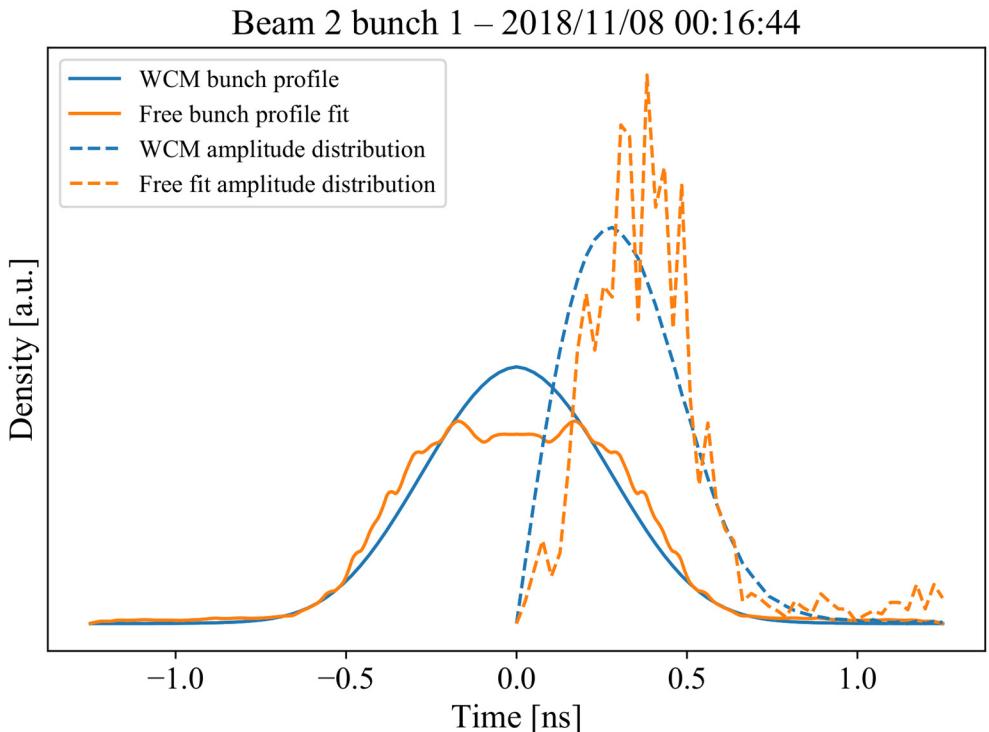


Free fit optimisation process

Minimise cost function:

$$C(\omega_{s_0}, \mathcal{A}) = |\log[\mathcal{M}(\omega_{s_0})\mathcal{A}] - \log[\mathcal{S}]|^2$$

Optimization process uses a *differential evolution* algorithm for global minimisation

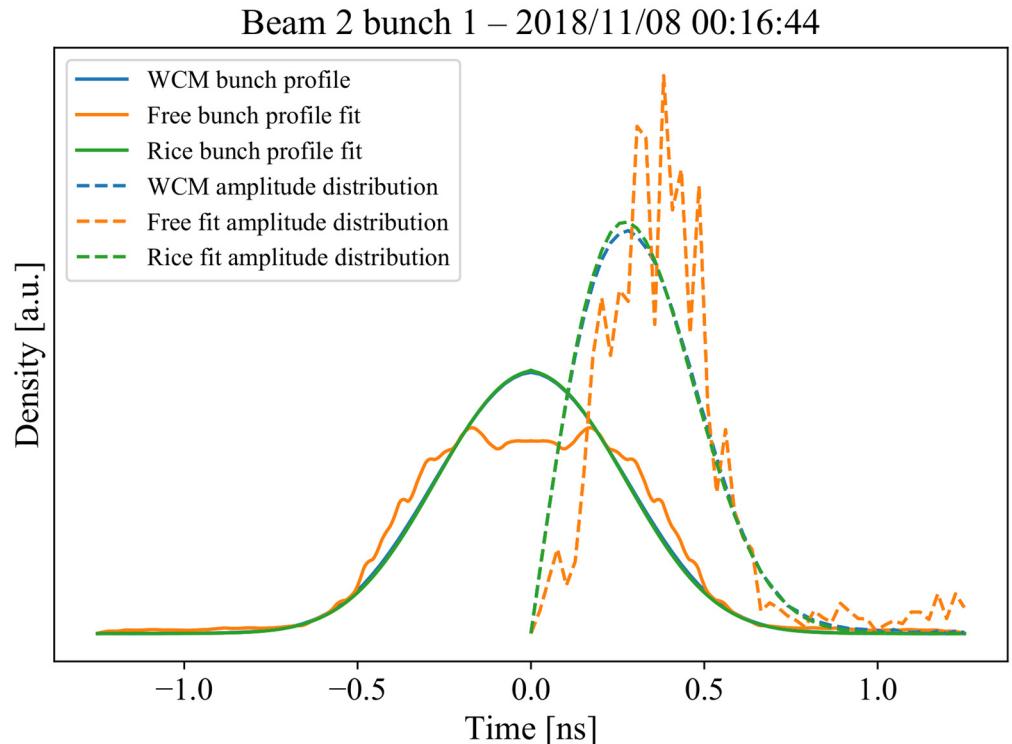


Optimisation process

Minimise cost function:

$$C(\omega_{s_0}, \mathcal{A}) = |\log[\mathcal{M}(\omega_{s_0})\mathcal{A}] - \log[\mathcal{S}]|^2$$

Optimization process uses a *differential evolution* algorithm for global minimisation





Matrix formalism

$$\underbrace{\begin{bmatrix} P(\omega_1, \hat{\tau}_1) & \dots & P(\omega_1, \hat{\tau}_n) \\ P(\omega_2, \hat{\tau}_1) & \dots & P(\omega_2, \hat{\tau}_n) \\ \vdots & \ddots & \vdots \\ P(\omega_m, \hat{\tau}_1) & \dots & P(\omega_n, \hat{\tau}_n) \end{bmatrix}}_{\mathcal{M}} \cdot \underbrace{\begin{bmatrix} g(\hat{\tau}_1) \\ g(\hat{\tau}_2) \\ \vdots \\ g(\hat{\tau}_n) \end{bmatrix}}_{\mathcal{A}} = \underbrace{\begin{bmatrix} \langle P(\omega_1) \rangle \\ \langle P(\omega_2) \rangle \\ \vdots \\ \langle P(\omega_m) \rangle \end{bmatrix}}_{\mathcal{S}}$$

$\mathcal{M}(:, i)$ - spectrum of single particle with amplitude τ_i

\mathcal{A} - synchrotron amplitude distribution

\mathcal{S} - cumulative spectrum





Matrix formalism



$$\underbrace{\begin{bmatrix} P(\omega_1, \hat{\tau}_1) & \cdots & P(\omega_1, \hat{\tau}_n) \\ P(\omega_2, \hat{\tau}_1) & \cdots & P(\omega_2, \hat{\tau}_n) \\ \vdots & \ddots & \vdots \\ P(\omega_m, \hat{\tau}_1) & \cdots & P(\omega_m, \hat{\tau}_n) \end{bmatrix}}_{\mathcal{M}} \cdot \underbrace{\begin{bmatrix} g(\hat{\tau}_1) \\ g(\hat{\tau}_2) \\ \vdots \\ g(\hat{\tau}_n) \end{bmatrix}}_{\mathcal{A}} = \underbrace{\begin{bmatrix} \langle P(\omega_1) \rangle \\ \langle P(\omega_2) \rangle \\ \vdots \\ \langle P(\omega_m) \rangle \end{bmatrix}}_{\mathcal{S}}$$

Precisely:

$$\mathcal{M} = \mathcal{M}(\omega_{s_0})$$





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- **Longitudinal bunch profile**
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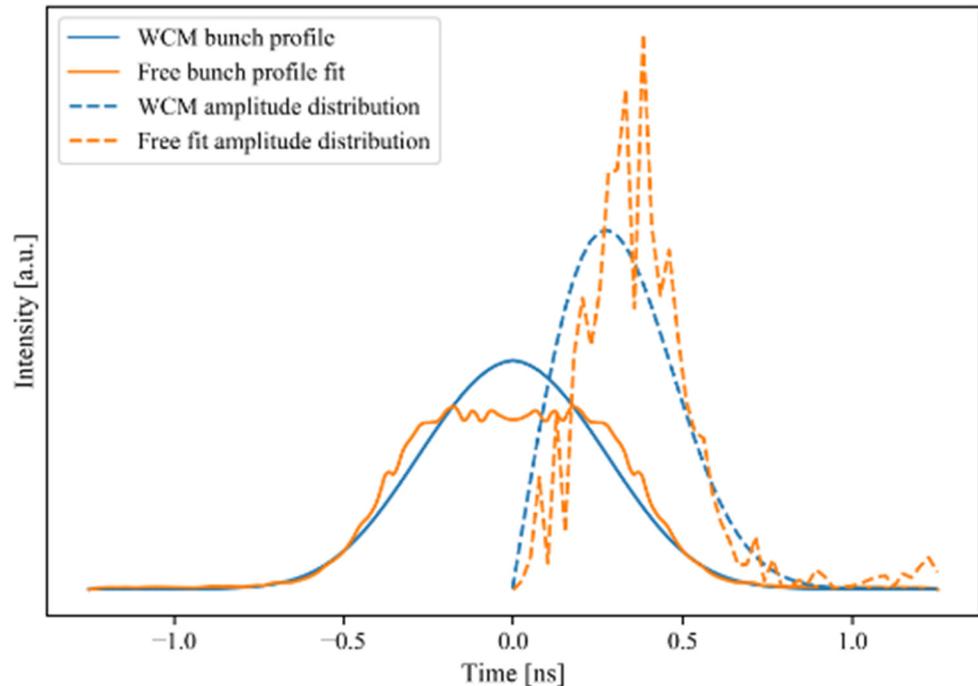


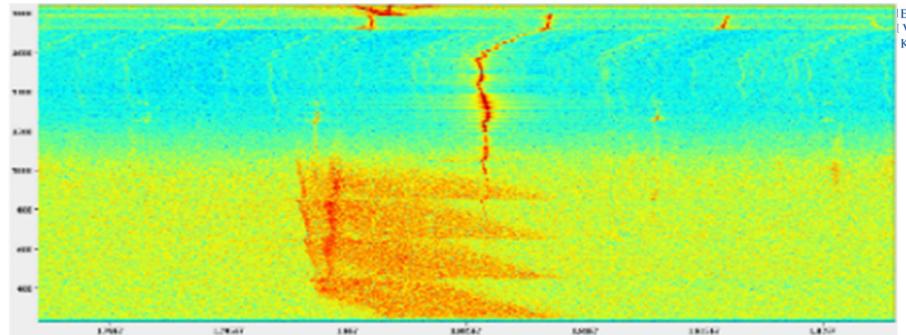
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Beam 2 bunch 1 – 2018/11/08 00:16:44





$$\begin{bmatrix} P(\omega_1, \hat{\tau}_1) & \dots & P(\omega_1, \hat{\tau}_n) \\ P(\omega_2, \hat{\tau}_1) & \dots & P(\omega_2, \hat{\tau}_n) \\ \vdots & \ddots & \vdots \\ P(\omega_m, \hat{\tau}_1) & \dots & P(\omega_m, \hat{\tau}_n) \end{bmatrix} \begin{bmatrix} g_{\hat{\tau}}(\hat{\tau}_1) \\ g_{\hat{\tau}}(\hat{\tau}_2) \\ \vdots \\ g_{\hat{\tau}}(\hat{\tau}_n) \end{bmatrix} = \begin{bmatrix} \langle P(\omega_1) \rangle \\ \langle P(\omega_2) \rangle \\ \vdots \\ \langle P(\omega_m) \rangle \end{bmatrix}$$



Notation



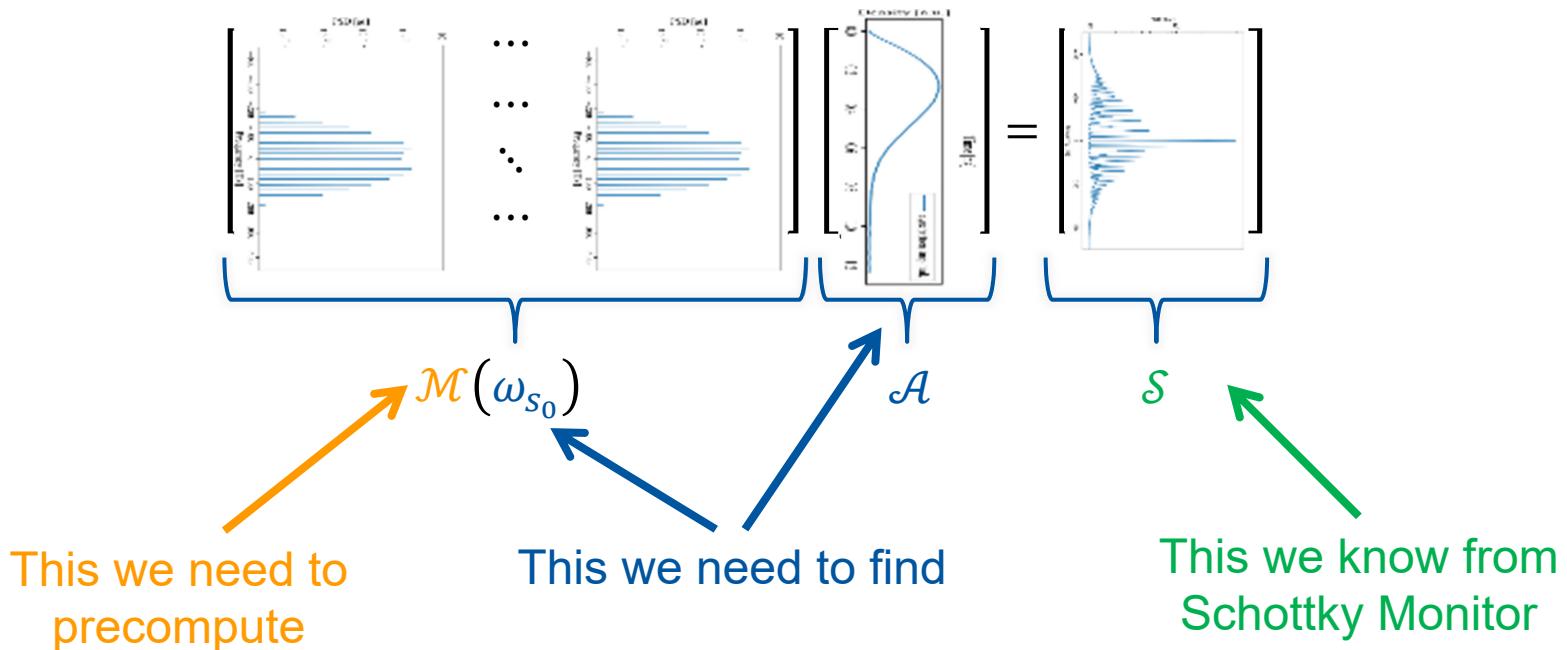
- τ = longitudinal time difference w.r.t. synchronous particle,
- $\hat{\tau}$ = time amplitude of synchrotron oscillations,
- ϕ_s = initial synchrotron phase,
- ω_{s_0} = nominal synchrotron frequency,
- ω_s = synchrotron frequency.

Assumptions

- Simple harmonic motion for synchrotron oscillations (% of RF bucket~2.5ns)
 - $\tau = \tau(\hat{\tau}, \phi_s) = \hat{\tau} \cos(\omega_s t + \phi_s)$ [Boussard 1986]
- No coherent intra-bunch motion



Matrix formalism



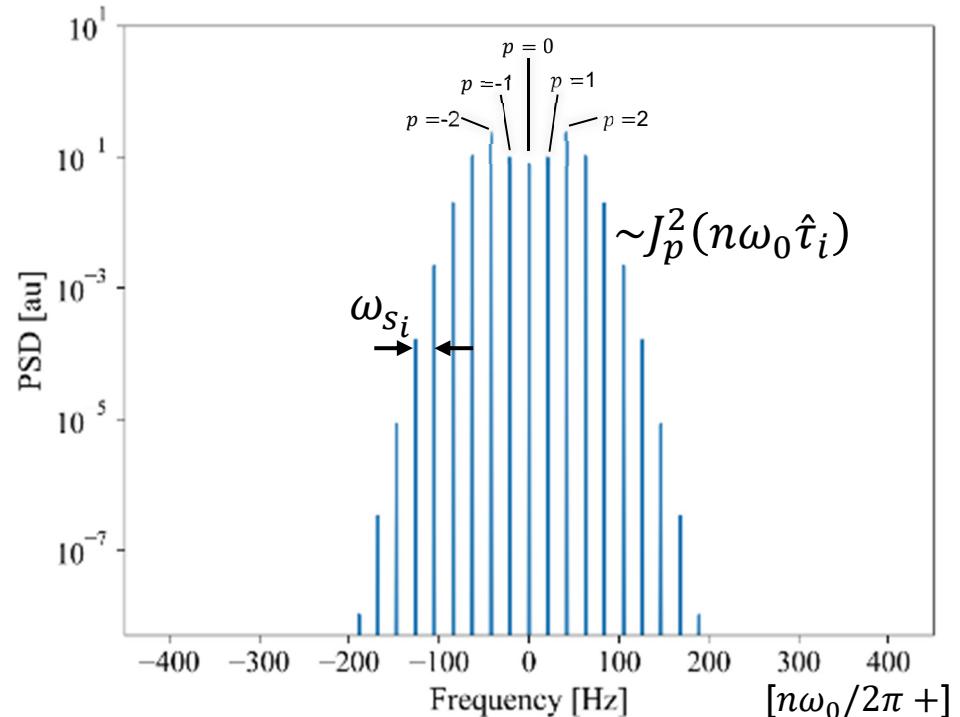
Single particle Schottky spectrum

Single particle contribution
to sum (Σ) signal
[Boussard1986]

$$s_i(t) = \sum_{p=-\infty}^{\infty} j^p J_p(n\omega_0 \hat{\tau}_i) e^{j[(n\omega_0 + p\omega_{s_i})t + p\phi_{s_i}]}$$

As ω_{s_i} and $\hat{\tau}_i$ are related, we have that:

$$P_i(\omega) = P_i(\omega, \hat{\tau}_i)$$





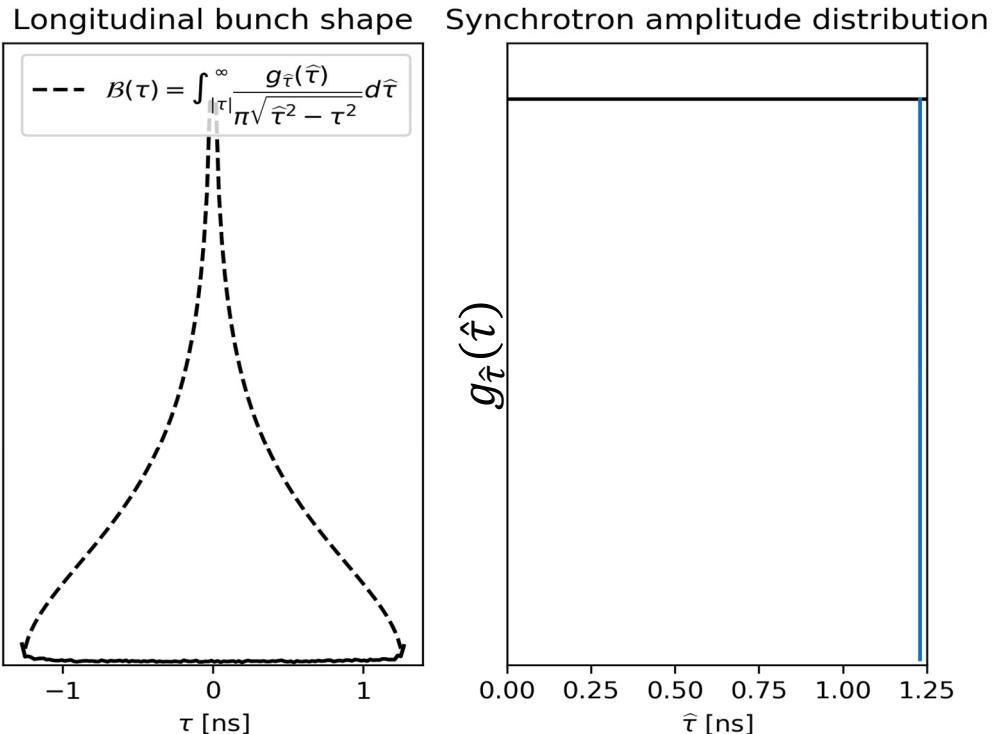
Bunch profile numerical example

$$g_{\tau}(\tau|\hat{\tau}) = \frac{1}{\pi\sqrt{\hat{\tau}^2 - \tau^2}}$$

Single particle distribution

$$\mathcal{B}(\tau) = \int_{|\tau|}^{\infty} \frac{1}{\pi\sqrt{\hat{\tau}^2 - \tau^2}} g_{\hat{\tau}}(\hat{\tau}) d\hat{\tau}$$

Longitudinal bunch profile

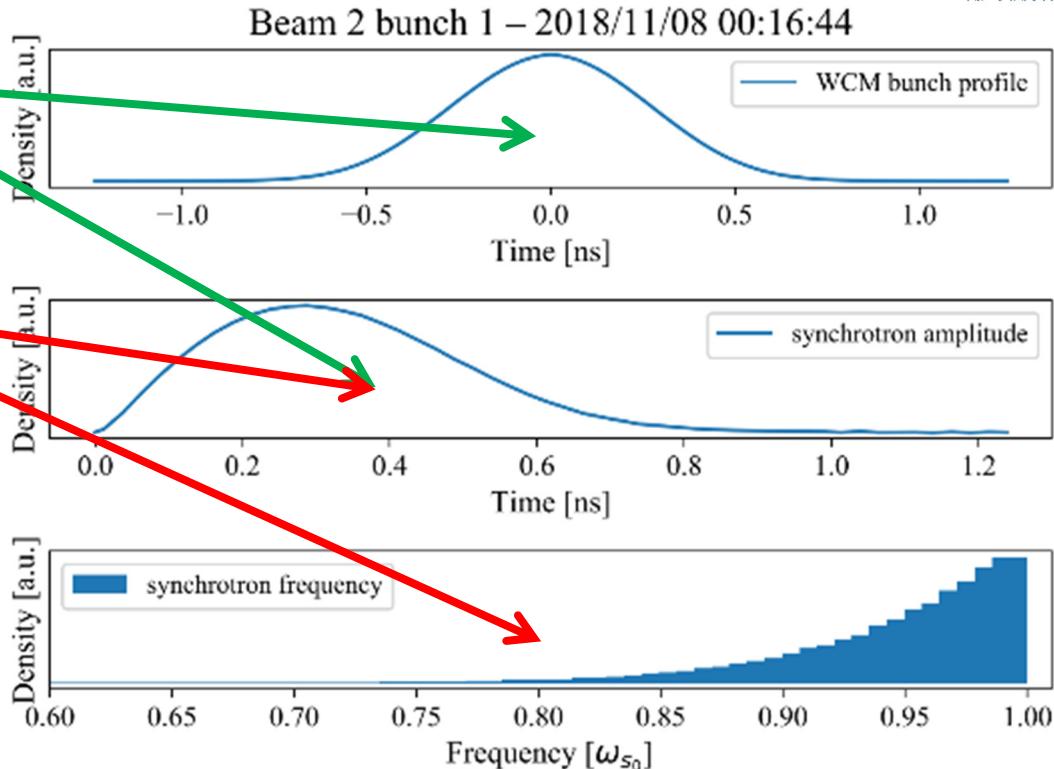
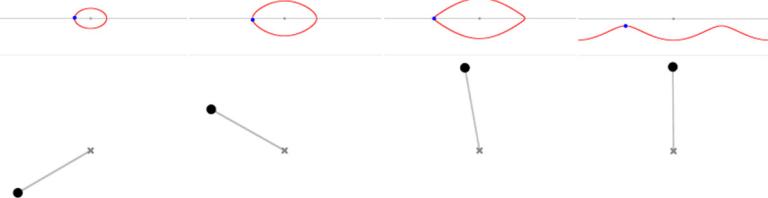


Equivalence of distributions

$$\mathcal{B}(\tau) = \int_{|\tau|}^{\infty} \frac{g_{\hat{\tau}}(\hat{\tau})}{\pi \sqrt{\hat{\tau}^2 - \tau^2}} d\hat{\tau}$$

$$\omega_s = \frac{\pi}{2K} \left[\sin\left(\frac{h\omega_0 \hat{\tau}}{2}\right) \right] \omega_{s_0}$$

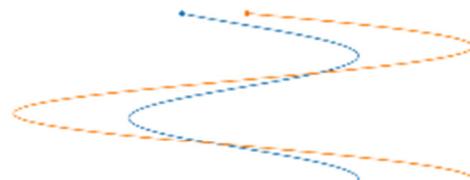
From the theory of pendulum:



Synchrotron motion

Harmonic oscillator with amplitude dependent frequency

$$\tau_i = \hat{\tau}_i \cos(\omega_{s_i} t + \phi_{s_i}) \quad [\text{Boussard 1986}]$$



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Longitudinal profile



How to relate the distribution of synchrotron amplitudes
with the longitudinal bunch profile ...

$$g_{\hat{\tau}}(\hat{\tau}) \leftrightarrow \mathcal{B}(\tau)$$



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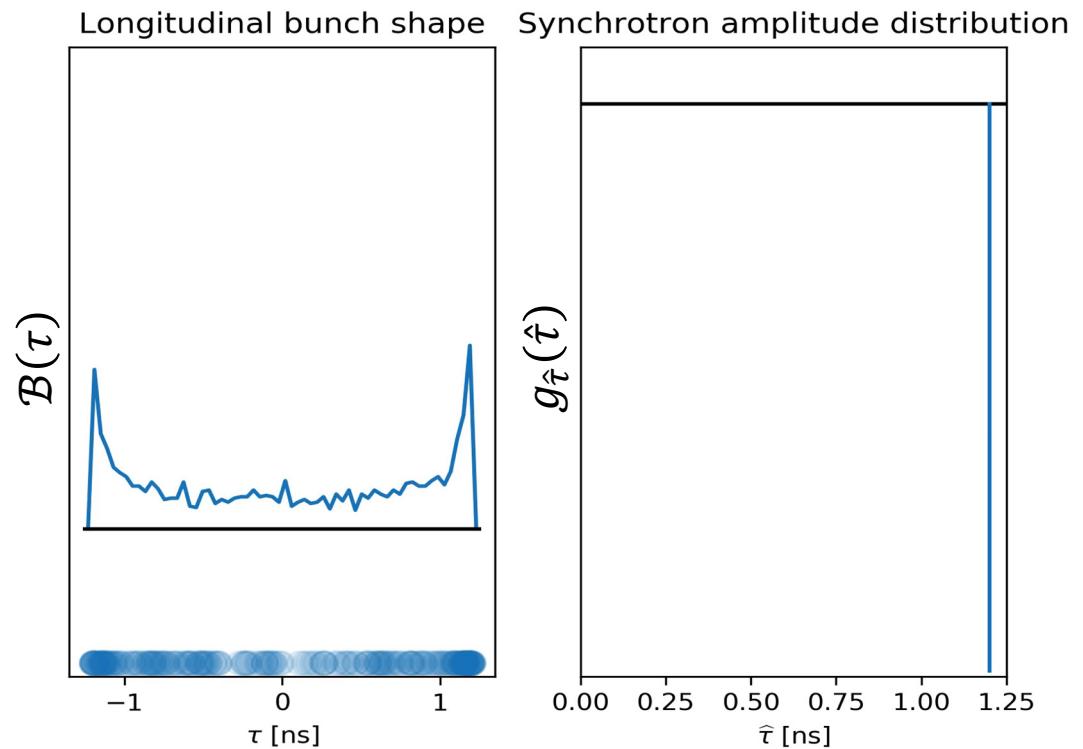
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Synchrotron motion – single particle distribution



$$g_{\tau}(\tau|\hat{\tau}) = \frac{1}{\pi\sqrt{\hat{\tau}^2 - \tau^2}}$$

Single particle distribution



Longitudinal Schottky spectrum

Total pick-up signal: $s(t) = \sum_{i=1}^N s_i(t)$

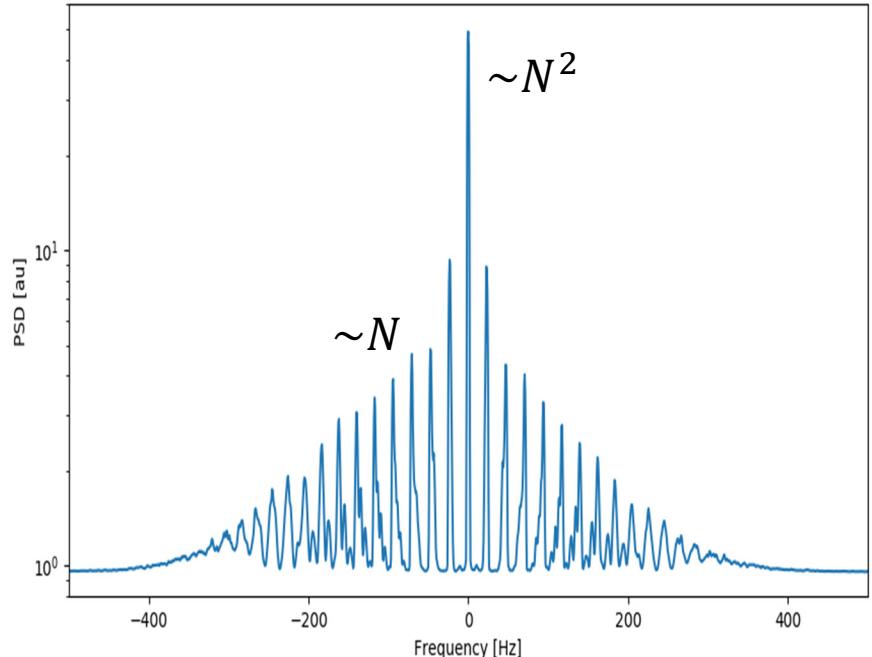
Instantaneous Schottky spectrum is not deterministic!

- depends on ϕ_s
- changes with time
- ... so we average ...

No intra-bunch coherent motion

$$\langle P(\omega) \rangle \propto N \langle P_i(\omega, \hat{\tau}_i) \rangle \sim N \sum_{\hat{\tau}} g_{\hat{\tau}}(\hat{\tau}) P(\omega, \hat{\tau})$$

Synch amp pdf Single particle spectrum





$$\omega_{s_i} = \frac{\pi}{2K \left[\sin \left(\frac{h\omega_0 \hat{t}_i}{2} \right) \right]} \omega_{s_0}$$

$$g_{\phi_s}(\phi_s) = \mathcal{U}[-\pi, \pi]$$
$$g_{\phi_s, \hat{\tau}}(\phi_s, \hat{\tau}) = g_{\phi_s}(\phi_s)g_{\hat{\tau}}(\hat{\tau})$$



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REFERENCES

[Boussard1986]



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Outlook



- Important step towards understanding the Schottky spectra
 - Mathematical framework can be adapted to transverse signals (e.g. estimate chromaticity and tune)
 - Obtained bunch profiles can be compared with WCMs and serve as a validation of the estimation of other parameters
- Mathematical treatment of coherent motion
 - Evolution of bunch profiles @ injection
- Apply method to proton bunches (higher intensities and shorter bunch lengths)



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