

# ***High Gain FEL based on a Transverse Gradient Undulator***

***Panagiotis Baxevanis, Zhirong Huang, Yuantao Ding and  
Ronald Ruth (SLAC)***

***August 28, 2013***



# Overview

- The **transverse gradient undulator (TGU)** is a concept for reducing the sensitivity of the FEL gain with respect to the energy spread of the electron beam.

- Previous studies include:

low gain regime: T. Smith et. al., J. Appl. Phys. 50, 4580 (1979).

N. Kroll et. al., IEEE Journal of Quan. Electro. QE-17, 1496 (1981).

high gain regime (1-D theory and 3-D simulations): Z. Huang, Y. Ding and C. Schroeder, PRL 109, 204801 (2012)

- In this talk, we present a **3-D theory** of a **high-gain** TGU-based FEL and discuss some numerical examples.

# Introduction

- The FEL performance depends critically on the energy spread of the driving electron beam.
- ✓ Efficient lasing requires

$$\frac{\sigma_\delta}{\rho} < 1$$

Big energy spread results in a large spread of the resonant wavelength

$$\rho = \left( \frac{K_0^2 J J^2}{16 \gamma_0^3 k_u^2 \sigma_x \sigma_y I_A} \right)^{1/3}$$

is the FEL parameter.

- ✓ For beams from plasma wakefield accelerators, it is hard to satisfy this condition (low emittance and high peak current but also large energy spread).

# Transverse Gradient Undulator (TGU)

- ✓ Use a dispersive element to spread out the beam:

$$\delta = \frac{\gamma - \gamma_0}{\gamma_0} = \frac{x}{\eta} \rightarrow \gamma = \gamma_0(1 + x/\eta)$$

- ✓ Introduce a linear field gradient by canting the poles:

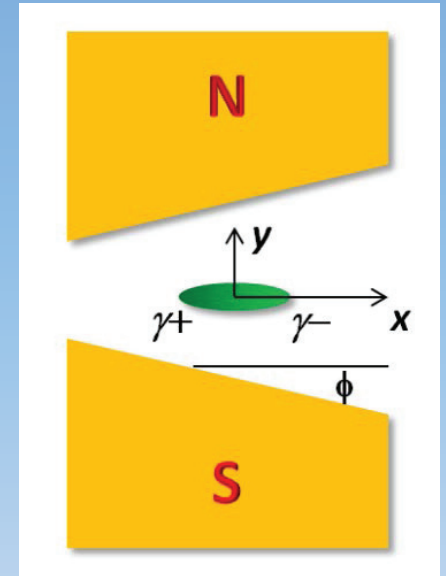
$$K = K_0(1 + ax)$$

- If we select

$$a = \frac{2 + K_0^2}{\eta K_0^2}$$

all particles in the beam satisfy the resonance condition

$$\lambda_r = \lambda_u \frac{1 + K^2/2}{2\gamma^2}$$



# Self-consistent 3-D theory

- FEL phase equation in the **parallel beam** limit

$$\theta' = \frac{d\theta}{dz} = 2k_u \left( \delta - \frac{x}{\eta} \right) \quad (\text{extra term due to the field gradient})$$

- Linearized, frequency-domain, *Vlasov-Maxwell equations*

$$\frac{\partial f_\nu}{\partial z} + i\nu\theta' f_\nu = -\kappa_1 \frac{\partial f_0}{\partial \delta} E_\nu e^{-i\Delta\nu k_u z}$$

$$\Delta\nu = \nu - 1$$

$$\nu = \omega/\omega_r$$

$$\left( \frac{\partial}{\partial z} + \frac{\nabla_\perp^2}{2ik_r} \right) E_\nu = -\kappa_2 e^{i\Delta\nu k_u z} \int_{-\infty}^{\infty} d\delta f_\nu$$

- The background distribution for the dispersed beam:

$$f_0 = \frac{N_b/l_b}{(2\pi)^{3/2} \sigma_x \sigma_y \sigma_\delta} \exp \left( -\frac{(x - \eta\delta)^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2} \right) \exp \left( -\frac{\delta^2}{2\sigma_\delta^2} \right)$$

✓ We seek the guided FEL eigenmodes:

$$E_\nu(\mathbf{x}, z) \propto A(\mathbf{x}) e^{i\mu z}$$

✓ This leads to an **eigenmode equation**

$$\left( \mu - \frac{\nabla_\perp^2}{2k_r} \right) A(\mathbf{x}) = -8\rho_T^3 k_u^3 A(\mathbf{x}) \exp \left( -\frac{x^2}{2\sigma_T^2} - \frac{y^2}{2\sigma_y^2} \right) \\ \times \int_{-\infty}^0 d\xi \xi e^{i(\mu - \Delta\nu k_u)\xi} e^{-2\sigma_{ef}^2 k_u^2 \xi^2} \exp \left( -2ik_u \xi \frac{\sigma_x^2}{\sigma_T^2} \frac{x}{\eta} \right)$$

$$\sigma_T = (\sigma_x^2 + \eta^2 \sigma_\delta^2)^{1/2}$$

total horizontal beam size

$$\rho_T = \rho \left( 1 + \frac{\eta^2 \sigma_\delta^2}{\sigma_x^2} \right)^{-1/6}$$

attenuated FEL parameter and

$$\sigma_{ef} = \sigma_\delta \left( 1 + \frac{\eta^2 \sigma_\delta^2}{\sigma_x^2} \right)^{-1/2}$$

**effective energy spread.** An efficient TGU requires  $\eta\sigma_\delta/\sigma_x \gg 1$ , so  $\sigma_{ef} \approx \sigma_x/\eta$

- ✓ Assuming a trial solution of the form

$$A(\mathbf{x}) = \exp(-a_x x^2 + bx) \exp(-a_y y^2)$$

we use a **variational technique** to obtain an approximation to the growth rate of the fundamental FEL eigenmode.

- ✓ The final dispersion relations can be expressed in a **fully analytical form (allows for fast calculations)**.
- ✓ The eigenmode analysis can be extended in order to take into account **emittance in the y-direction** as well as the undulator **natural focusing**.
- ✓ However, we disregard **natural focusing in the x-direction** ( $\sim (\gamma\eta)^{-1}$ , typically weak) and assume that a small net bending ( $\sim B_0/\gamma$ ) has been corrected.

# Soft X-ray FEL with 1 GeV LPA beam

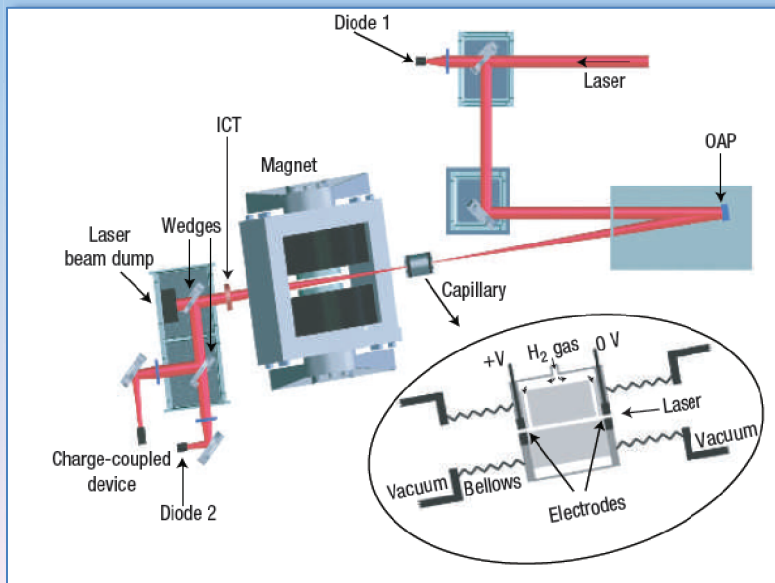
## LETTERS

### GeV electron beams from a centimetre-scale accelerator

W. P. LEEMANS<sup>1\*</sup>, B. NAGLER<sup>1</sup>, A. J. GONSALVES<sup>2</sup>, Cs. TÓTH<sup>1</sup>, K. NAKAMURA<sup>1,3</sup>, C. G. R. GEDDES<sup>1</sup>, E. ESAREY<sup>1\*</sup>, C. B. SCHROEDER<sup>1</sup> AND S. M. HOOKER<sup>2</sup>

Nat. Phys. 2, 696 (2006)

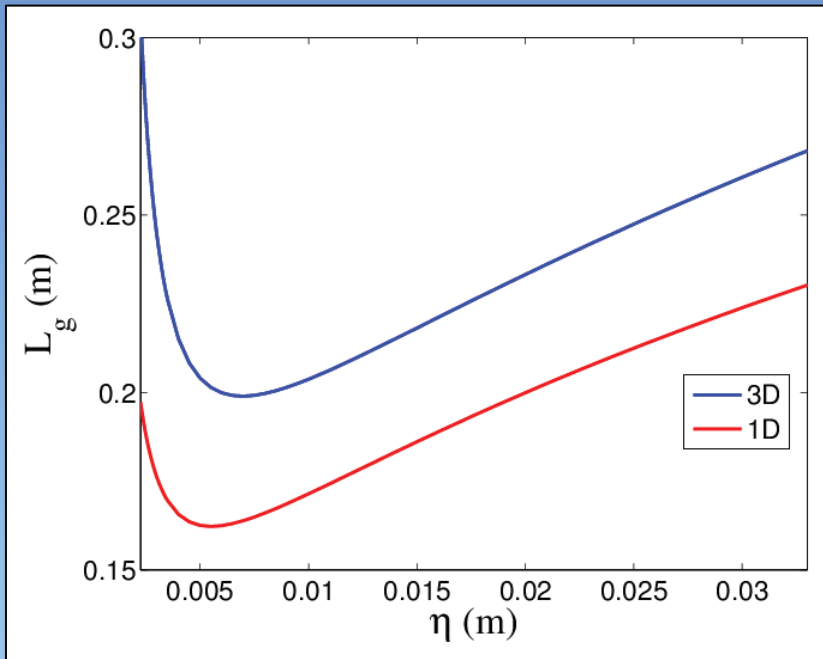
Laser-plasma accelerators (LPAs) have demonstrated the capability to produce e-beams in the GeV range.



- 1 GeV,  $\lambda_u = 1$  cm,  $K_0 = 2$   
radiation wavelength  $\lambda_r = 3.9$  nm
- 10 kA peak current,
- 0.1 mm-mrad norm. emittance
- 1% energy spread

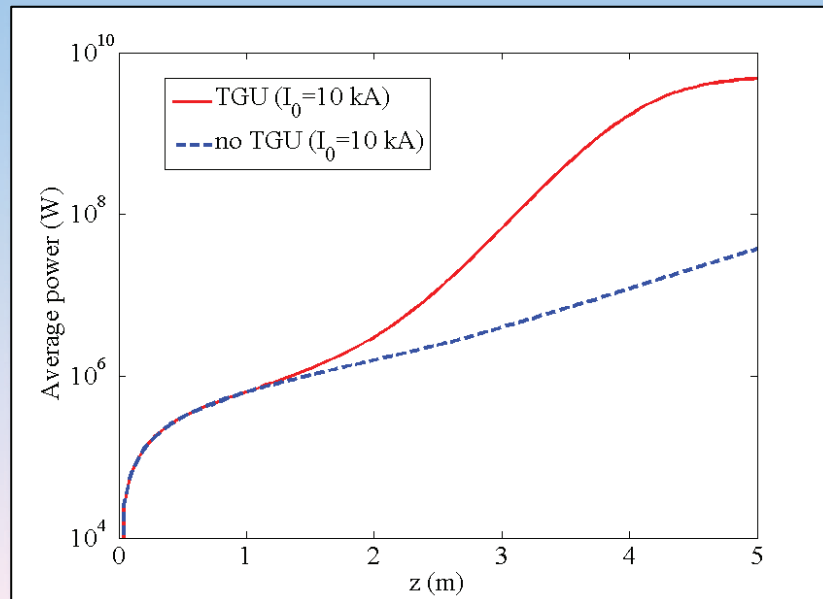
$$\sigma_x = \sigma_y \sim 10 \mu\text{m}$$





- ✓ We calculate the **frequency-optimized gain length** as a function of dispersion using the parallel beam theory.
- ✓ Results are compared with the **1D formula**

$$L_g \approx \frac{\lambda_u}{4\pi\sqrt{3}\rho_T} \left[ 1 + \frac{\sigma_{ef}^2}{\rho_T^2} \right]$$

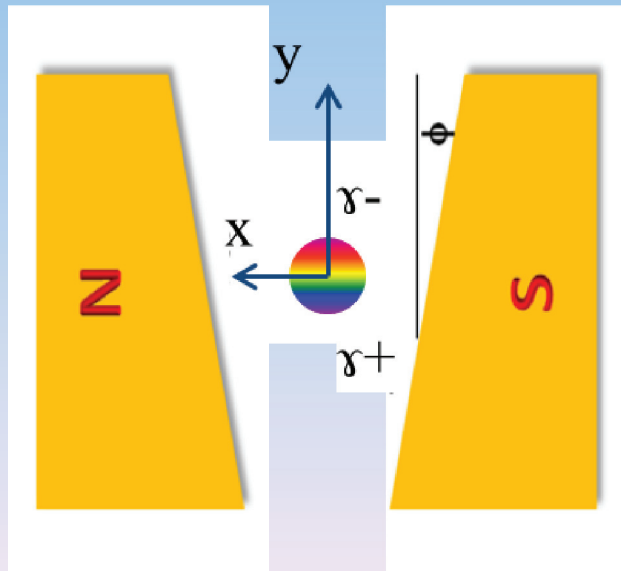


For a **1 cm** dispersion

- ✓ e-beam size: 100  $\mu\text{m}$  x 10  $\mu\text{m}$
- ✓ optimum gain length  $\sim 21$  cm
- ✓ Saturation within 5 m

# Soft X-ray FEL in an Ultimate Storage Ring

- A single pass, high-gain FEL based on an ultimate storage ring (USR) may be possible using a TGU.
- ✓ The undulator is placed in a bypass next to the ring.
- ✓ Rotate by  $90^\circ$  to take advantage of the very low vertical emittance.

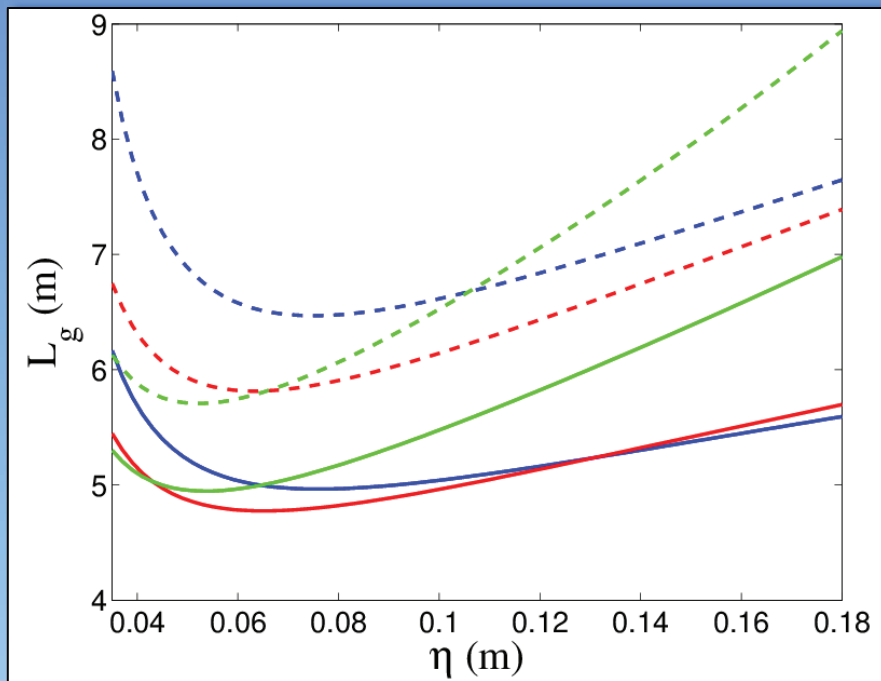


## PEP-X as an example

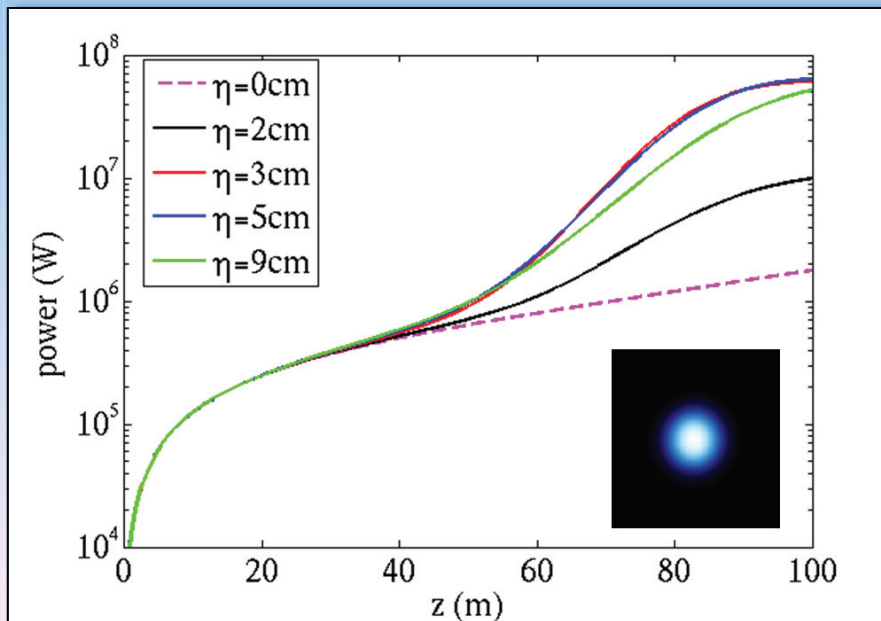
(Y. Ding et al. , IPAC-13)

- 4.5 GeV,  $\lambda_u = 2$  cm,  $K_0 = 3.68$   
radiation wavelength  $\lambda_r = 1$  nm
- 200 A peak current  
0.0123/1.23 mm-mrad n.emittance  
0.15% energy spread

$$\sigma_x \sim 10 \mu\text{m} , \sigma_y \sim 40 \mu\text{m}$$



Gain length vs. dispersion for several values of the detuning ( $\Delta\nu/(2\rho) = 0.0/-0.2/-0.4$ ) from the parallel beam model (solid lines) and including emittance (dashed).



For a dispersion  $\sim 3-5$  cm

- ✓ Saturation within 100 m.
- ✓ Roughly a round electron/radiation beam (better transverse coherence).

# Summary

- A self-consistent 3D theory has been developed for a TGU-based, high-gain FEL.
- We use a variational method to determine the properties of the fundamental FEL mode.
- In the **parallel beam** limit, we derived fully analytical relations.
- This model has been extended to include **vertical emittance** and **natural focusing** (relevant for USR-based TGU FELs).
- The theory agrees with simulation as well as other analytical results and can be used for optimization studies.