



Beam Loss Monitoring for Demanding Environments

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Melbourne

Special thanks for their input to Sanja Damjanovic and Eduardo Nebot

Contributions from: Bernhard Auchmann, Marcin Bartosik, Mark Boland, Bernd Dehning, Maria Kastriotou, Anton Lechner, Sophie Mallows, Giuseppe Venturini, Arden Warner, Kay Wittenburg, Manuel Zingl

Overview

- Considerations for the Design of a BLM System
 - Machine Protection
 - Short vs. Distributed Detectors
 - Machine Size and Radiation
- Background Sources
 - Synchrotron Radiation
 - Distant Losses
 - Accelerating Structures
- Summary

plus examples

Design Considerations

Beam Loss Monitors Roles

Beam diagnostics —
operation and commissioning

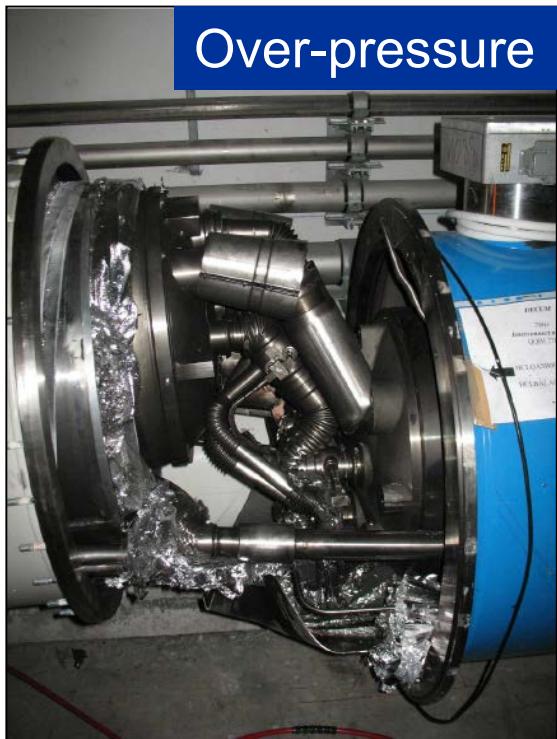
Quench / damage protection

Activation / aging / human exposure

Machine Protection

Energy stored in the Magnets — Release of 600 MJ

- LHC 2008 incident **without beam**
 - Electrical arc provoked a He pressure wave **damaging ≈ 600 m of LHC**
- LHC magnets at 7 TeV: **10 GJ**



Arcing @ interconnection



Magnet displacement

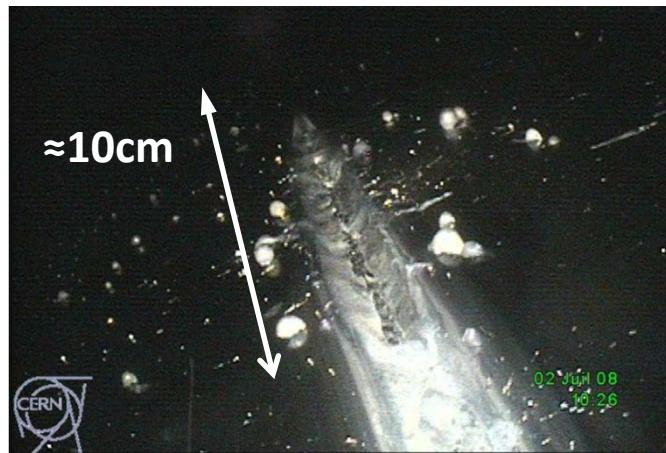


Energy stored in the Beams — Uncontrolled Losses

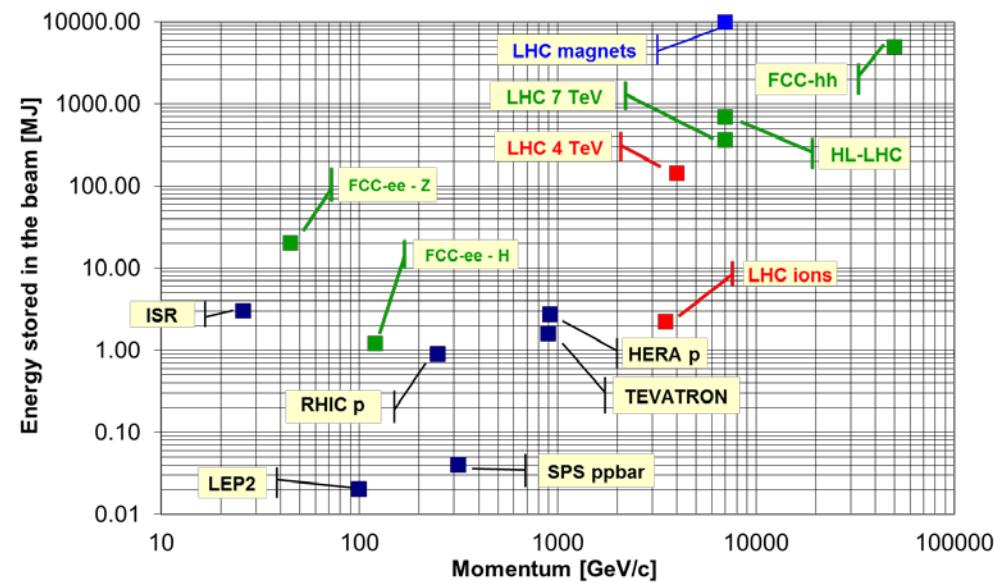
- LHC at 7 TeV 360 MJ:
 - Pilot bunch of 5×10^9 close to damage level
 - Loss of 3×10^{-7} of nominal beam over 10 ms can create a quench

1MJ can heat and melt 1.5 kg of copper

- SPS incident in June 2008,
400 GeV beam with **2 MJ**
(*J. Wenninger, CERN-BE-2009-003-OP*)



World record: LHC
140 MJ @ 4 TeV

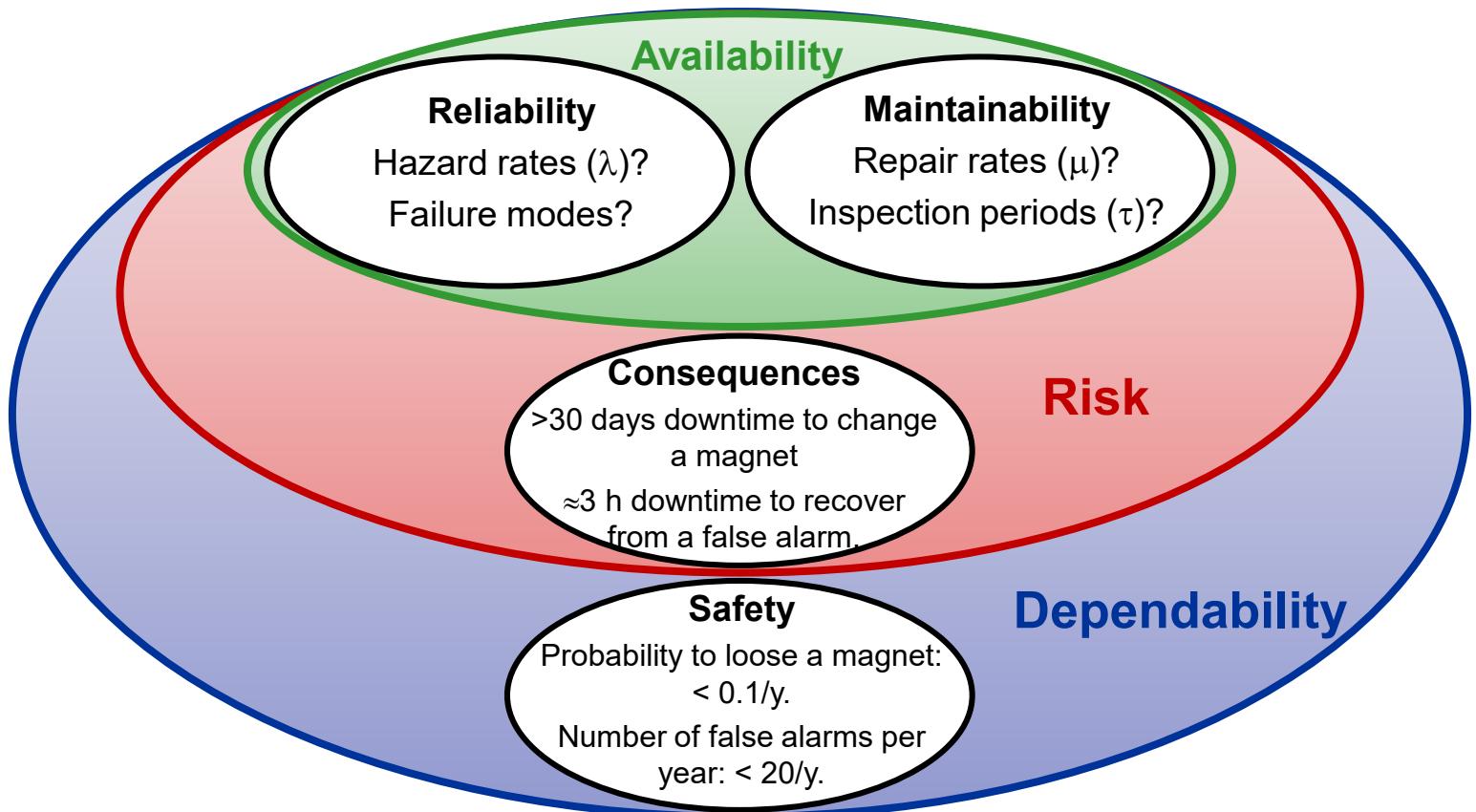


Machine Protection by BLM

- LHC
 - Safely extract beam when loss exceeds threshold (on any of \approx 3600 detectors)
 - \approx 1.5 million thresholds depend on
 - Detector location
 - Beam energy
 - Integration time (40 μ s–84s)
- CLIC
 - **Prevent subsequent injection** when potentially dangerous beam losses are detected (“**next cycle permit**”)
 - Damage to beam-line components determined by power density (not by beam power) of the beams

Dependability (colloquially: reliability) analysis

- Machine protection system must be integrated in the machine design
- Dependability (reliability, availability, maintainability and safety) analysis → allowances for
 - Probability of component damage due to malfunctioning
 - Downtime due to false alarms
 - Downtime due to maintenance



Short vs. Distributed Detectors

Short localised detectors vs. Long distributed detectors

Time resolution
Coverage
Position resolution
Position \leftrightarrow Magnitude
Loss magnitude

Machine protection

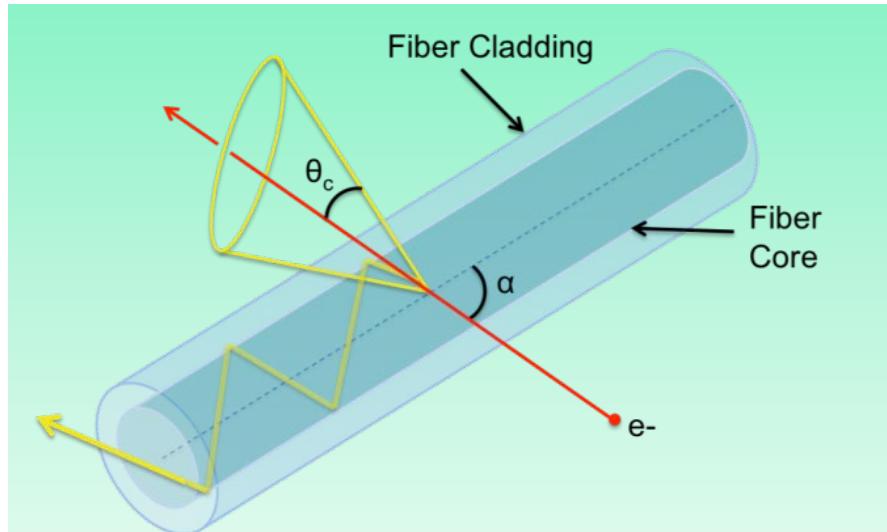
Coverage
Cost

Short beams
Position resolution
Time resolution
Loss magnitude (calibration including attenuation)

Long trains
Position \leftrightarrow Time \leftrightarrow Magnitude
R&D needed for machine protection

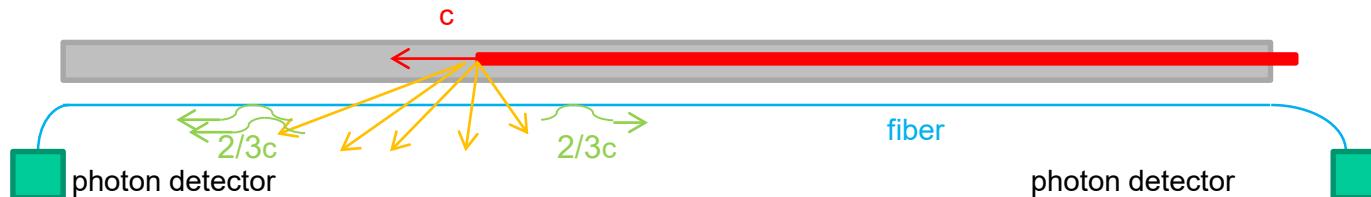
Cherenkov Fibers BLMs — Detection Principle

- When a charged particle with $v > c$ enters the fiber, photons are produced along Cherenkov cone



$$\cos \theta_{\text{Cherenkov}} = \frac{1}{n_{\text{core}}} \beta$$

- Optical fibers for loss measurement are increasingly popular
- Overview: *T. Obina , Y. Yano, IBIC 2013;*
L. Fröhlich et al., DIPAC 2011; F. Wulf, M. Körfer, DIPAC 2009; ...



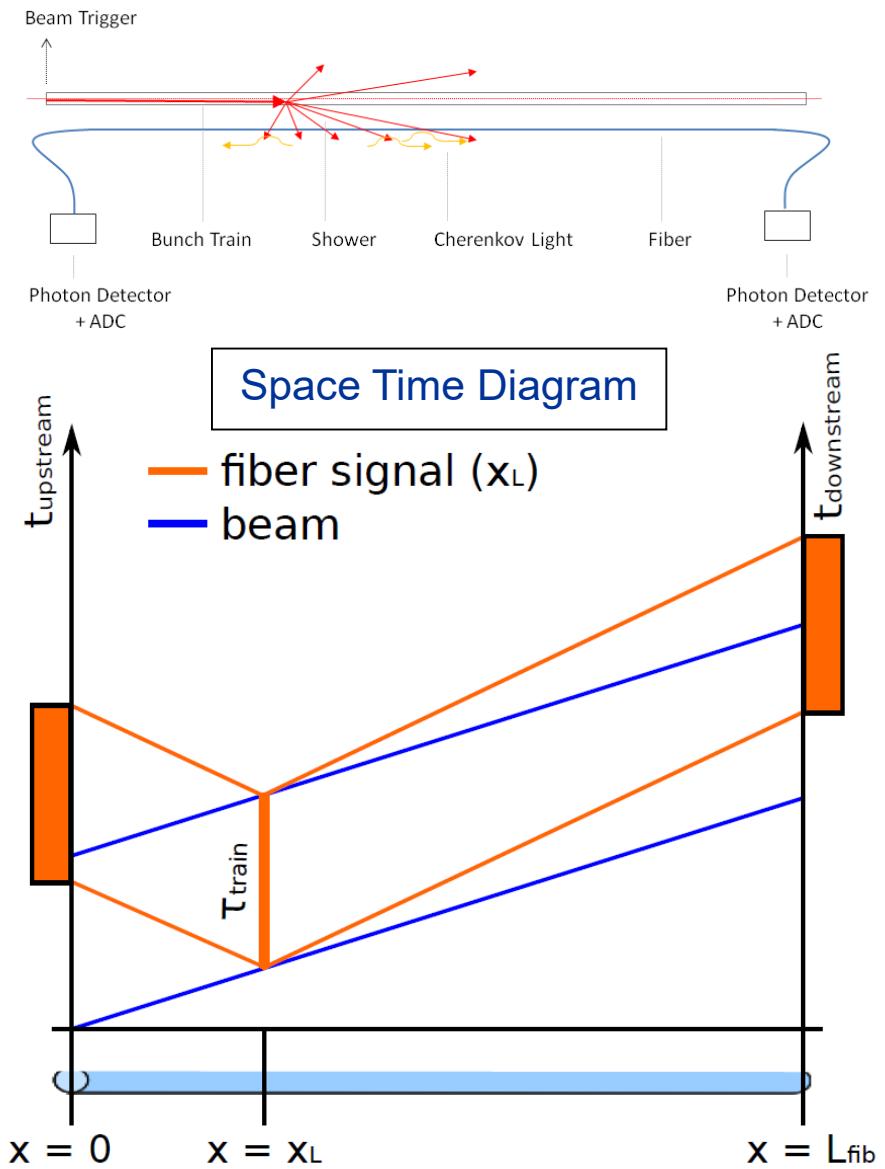
Cherenkov Fiber BLMs — Pros and Cons

- Insensitive to photons
- Insensitive to E and B fields
- Small and fast
- Relatively radiation hard (depending on type)
- Combination fiber / readout can adapt to a wide dose range
- Time resolution 1 ns
- Position resolution 0.5 – 1 m
- (Radiation induced) attenuation and dispersion (multi mode fibers)
- Position \leftrightarrow time \leftrightarrow magnitude
 - Ideas:
 - Couple with localised high time resolution measurement
 - Match pattern to a catalogue of known loss scenarios (experience, simulation)

Position Resolution for Long Bunches

- Case: single loss location — constant for all bunches
- Simulations:
 - Starting point of the losses can be determined from the signal's rising edge, with:
 - < 1m longitudinal resolution
 - ≈ 1ns time resolution
- First results from installation at the Australian Synchrotron and at CTF3 (CLIC Test Facility):

E. Nebot: "Position Resolution of Optical Fibre-Based Beam Loss Monitors Using Long Electron Pulses"



UFO Losses

Position ↔ magnitude

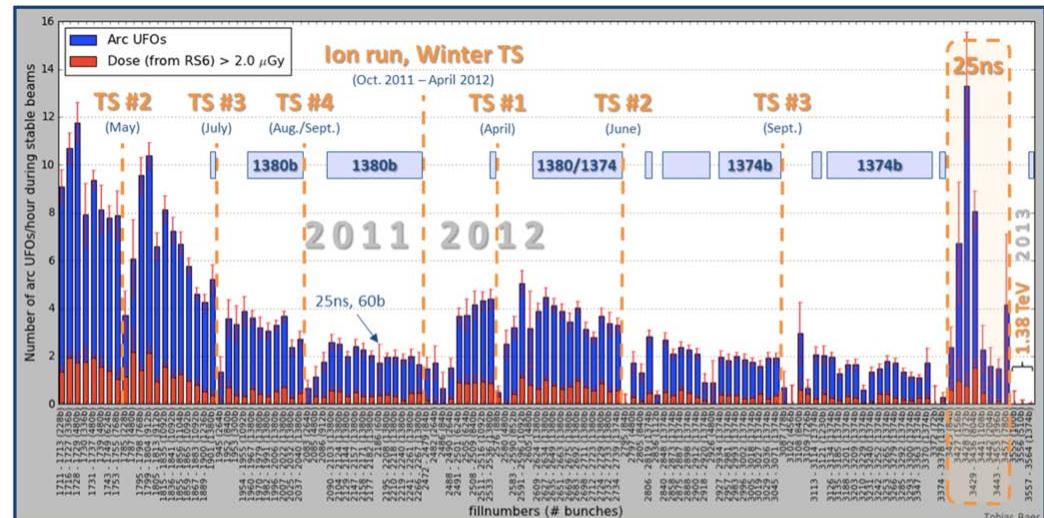
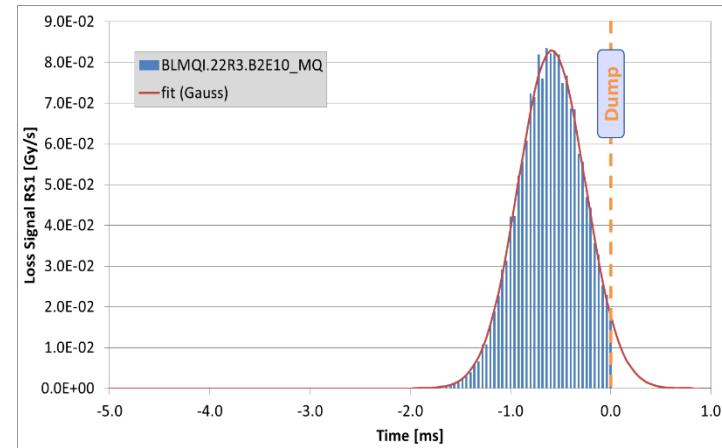
Determine position and magnitude
with simulations

Improve longitudinal coverage

UFOs — Causing Quenches at 6.5 TeV



- Fast and localised losses all around the ring believed to be caused by macro particles interacting with the beam
- “UFO”: Unidentified Falling Objects
- No quenches at 4 TeV
 - Less heat deposited
 - Lower magnetic field
 - Conservative BLM thresholds
- 6.5 TeV: thresholds set to quench limit
→ quenches occurred



Tobias Baer

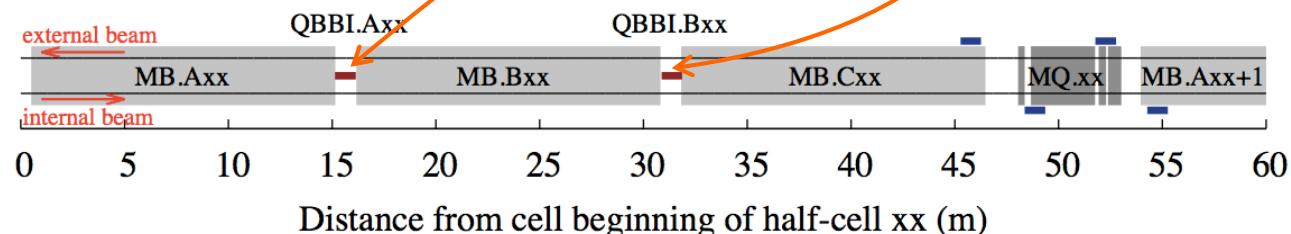
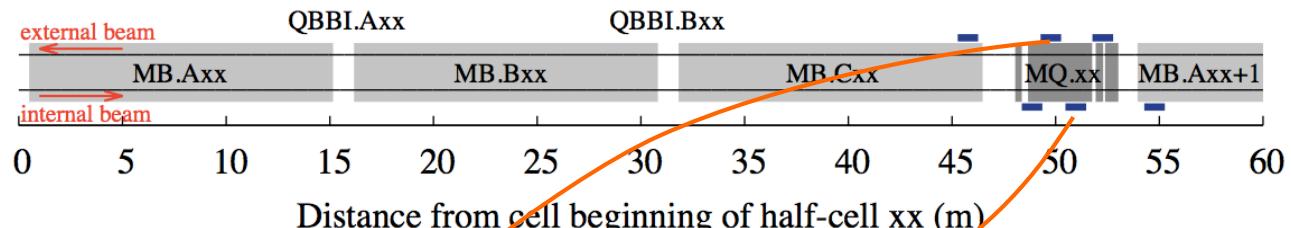
Relocation of 1/3 of Arc Detectors (Long Shutdown 1)



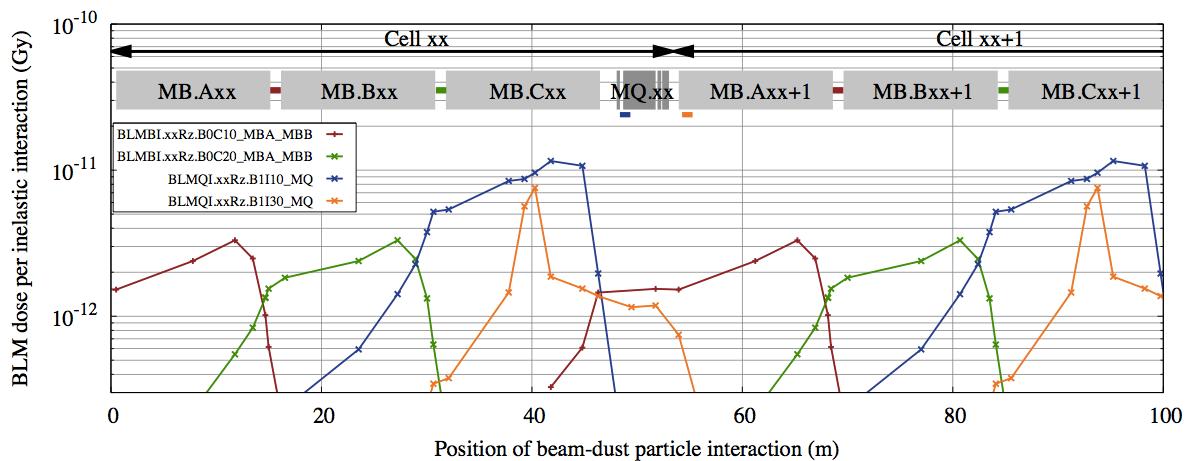
Pre-LS1



Post-LS1



- Coverage post-LS1: increases sensitivity by a factor 30 → 100% coverage can be achieved

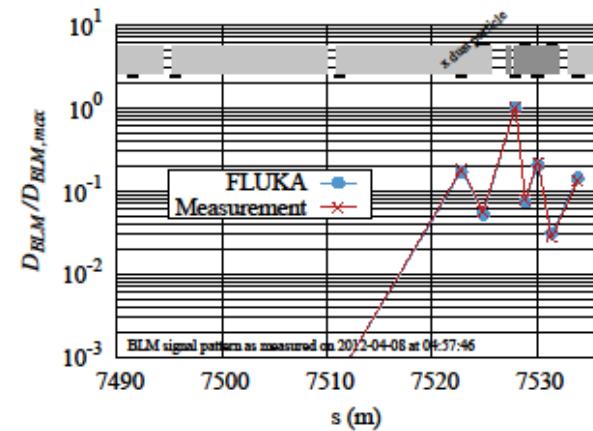
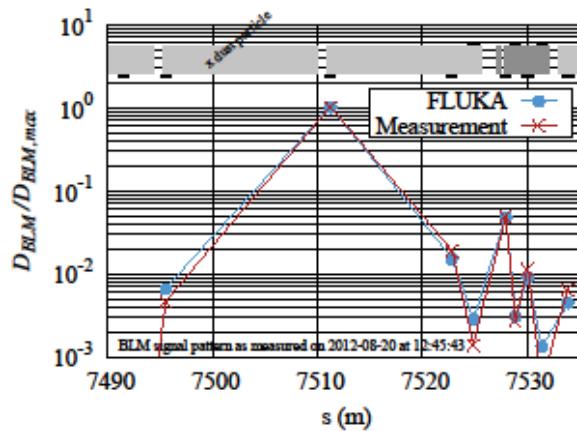
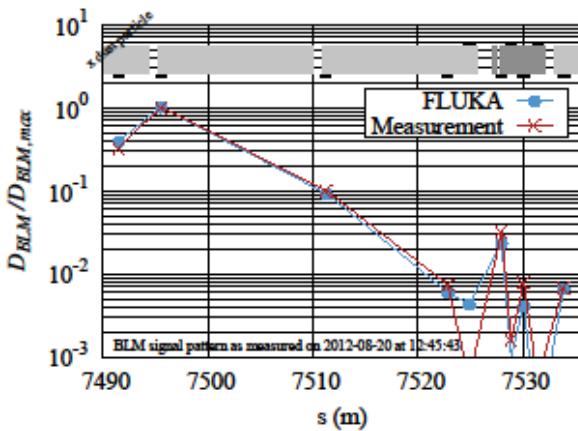


A. Lechner, Workshop on Beam-Induced Quenches, CERN, 2014

UFO Losses: Comparison Simulation — Measurement

Anton Lechner

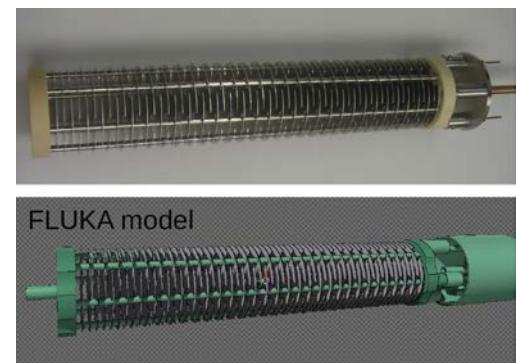
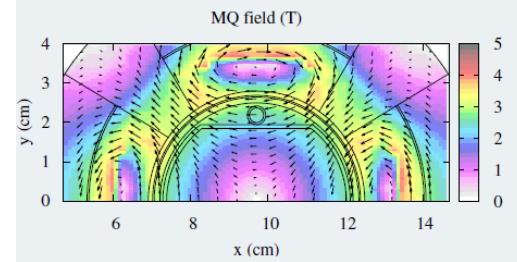
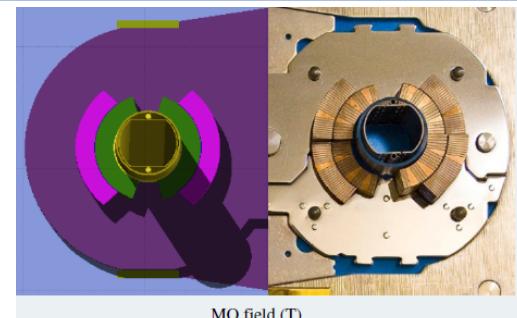
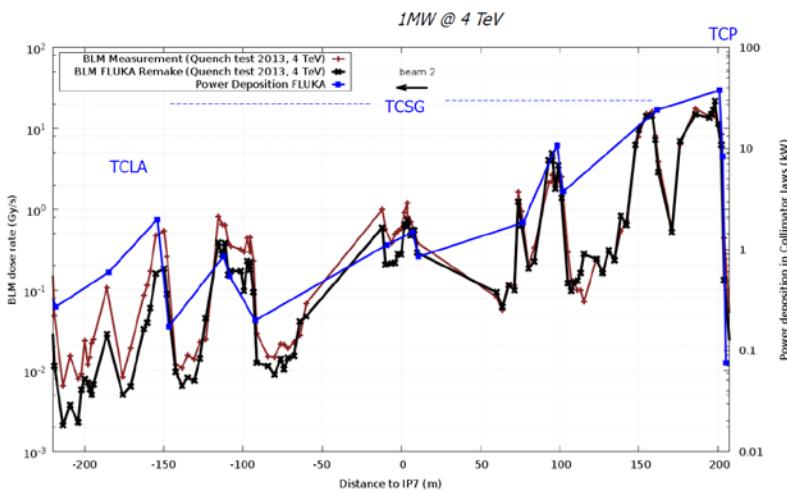
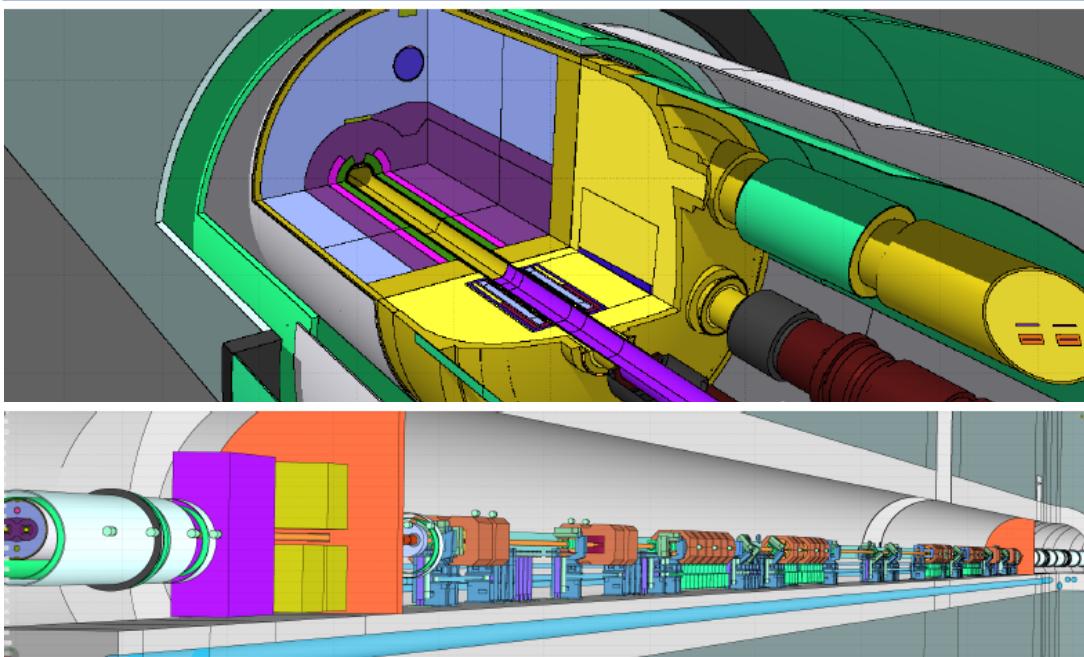
- If several detectors record the loss: Determine the loss position and magnitude with the help of simulations
 - Loss position: ± 1 m
 - Number of inelastic proton-dust particle interactions: factor 2



- $1-4 \times 10^6$ inelastic proton-dust particle interactions in this cell
- Other cells 10–100 times higher

A. Lechner, Workshop on Beam-Induced Quenches, CERN, 2014

Require Very Accurate Models



**Francesco Cerutti, Anton Lechner,
Eleftherios Skordis (FLUKA team)**

High Time Resolution

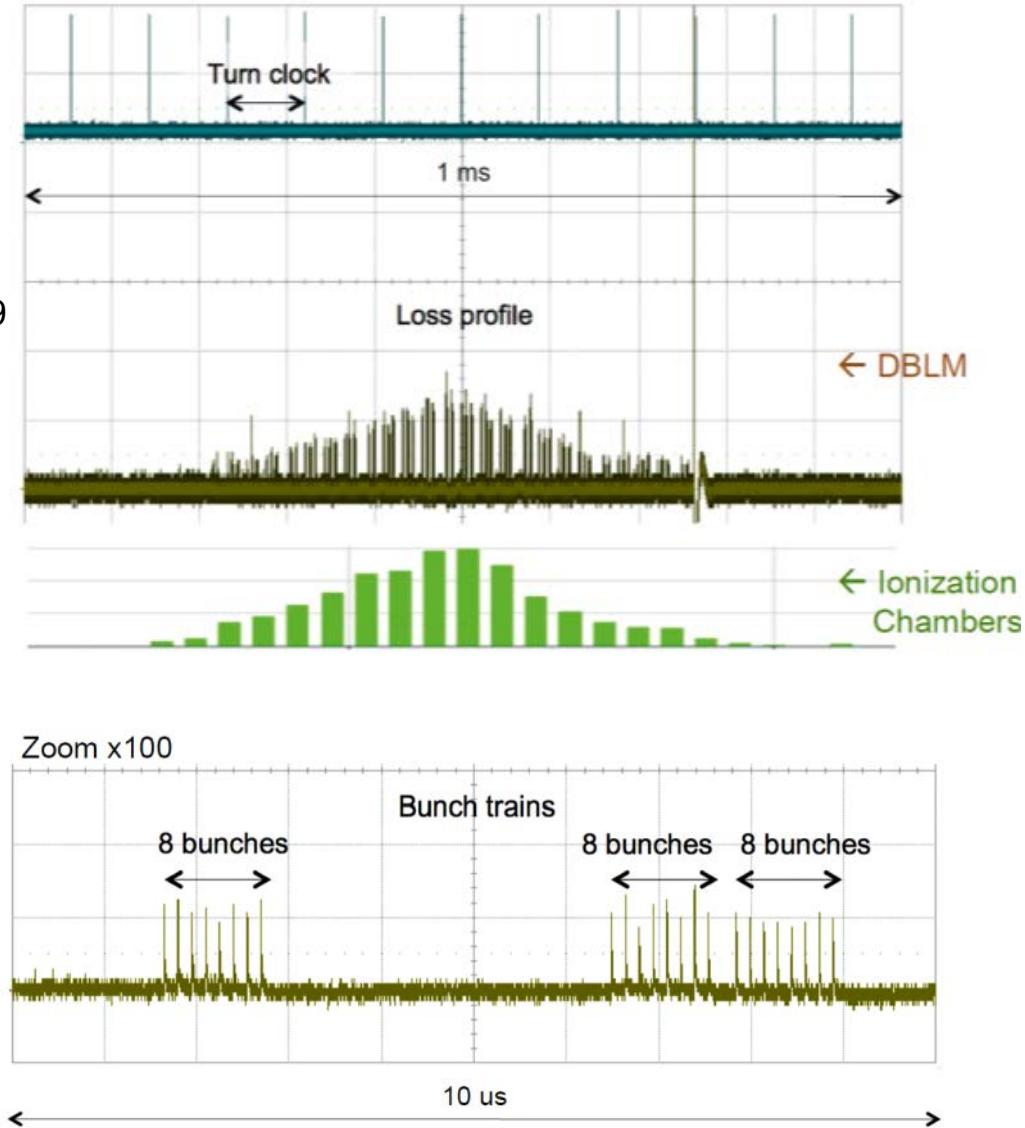
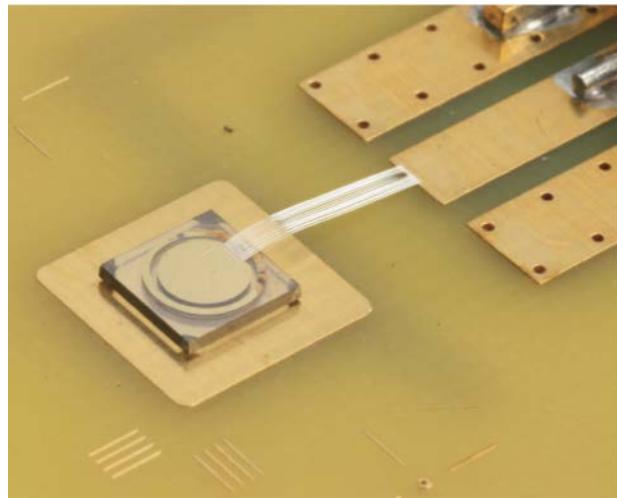
Require very small detector

→ Diamond detector LHC
(few ns)

→ Small Cherenkov + ultrafast readout:
Could they reach $\approx 100\text{ps}$?

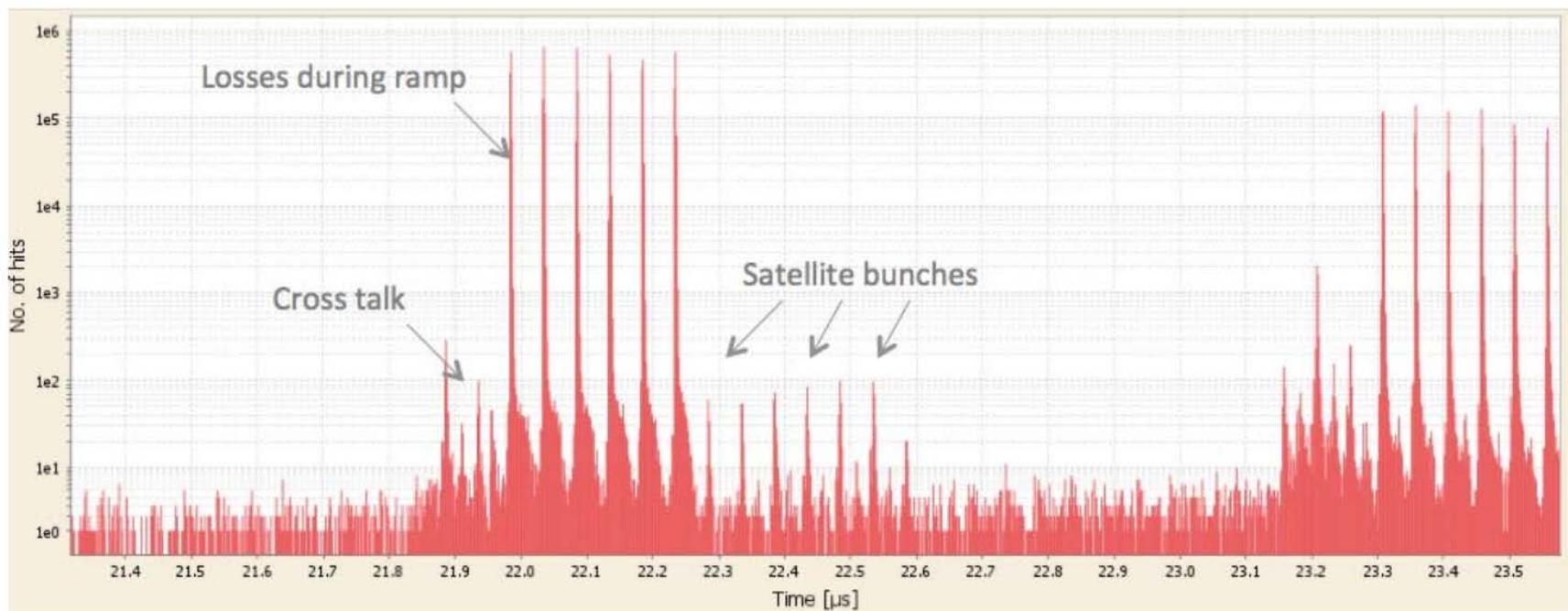
LHC Diamond Detectors

- Fast and sensitive
- Small and radiation hard
- Used in LHC to distinguish bunch-by-bunch losses
- Dynamic range of monitor: 10^9
- Temporal resolution: few ns



Diamond: Arrival Time Histogram During Ramp

- 50 ns bunch spacing
- Loss signal at 25 ns is from opposite beam (“cross talk”)
→ sub 25 ns resolution required to resolve



B. Dehning

Machine Size and Radiation

Machine Size Implications

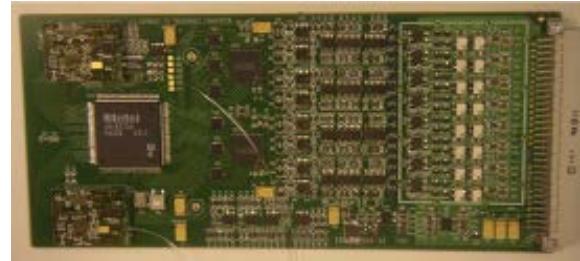
Short detectors: cost, dependability per channel, data treatment

Front-end electronics in the tunnel
→ radiation tolerant

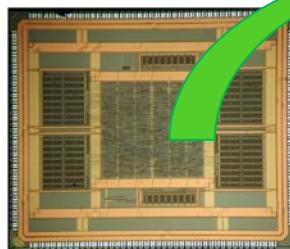
Low noise, low loss signal transmission: optical techniques

Radiation Tolerant Readout

- LHC BLM front-end: charge-to-frequency converter
 - 500 Gy certified — ok for arcs
 - Insertion regions:
up to 300–800 m long cables



- New development: radiation hard Application Specific Integrated Circuit (ASIC)
 - Dynamic range 10^6
 - Bipolar input current
 - Certified up to 100 kGy



G. Venturini

Background

Background Implications

Limit sensitivity to primary beam losses

Can compromise machine protection

Synchrotron Radiation

Synchrotron Radiation Background

- FLUKA simulation study for CLIC damping ring (DR) arc
- Aim: worst case estimate using a simplified geometry
- Impacting on vacuum chamber: wall thickness 1.5 mm
- Detectors: 10 cm and 40 cm from beam pipe
- Many CLIC DR parameters similar to third generation light sources
(but: very small longitudinal normalised emittance of 6 keV m)

| | Elettra | ALBA | DLS | ESRF | APS | Spring-8 | ASLS | CLIC DR |
|--------------|---------|------|-----|------|------|----------|------|------------------------|
| Energy [GeV] | 2 | 3 | 3 | 6 | 7 | 8 | 3 | 2.86 |
| Circumf. [m] | 259 | 269 | 562 | 845 | 1104 | 1436 | 216 | 427.5 |
| Lattice type | DBA | DBA | DBA | DBA | DBA | DBA | DBA | TME (arc) / FODO (LSS) |
| Current [mA] | 300 | 400 | 300 | 200 | 100 | 100 | 200 | 200 |

(table based on: A. Wolski, *Synchrotron Light Machines*, CAS 2012)



CLIC DR and Australian Synchrotron Comparison

| | E [GeV] | Intensity [electrons] | bunches | Pulse length [ns] | Circumf. [m] | F_{rev} [MHz] | Bunch spacing [ns] | $\gamma \epsilon_x$ [nm rad] | $\gamma \epsilon_y$ [nm rad] |
|----------------|--------------|--------------------------|---------|-------------------------|-----------------|---------------------------|--------------------------|------------------------------------|------------------------------------|
| CLIC DR | 2.86 | 1.28×10^{12} | 312 | 156 | 427.5 | 0.73 | 0.5 | 472 | 4.8 |
| ASLS | 3.0 | 0.9×10^{12} | 300 | 600 | 216 | 1.38 | 2 | 58708 | < 5 |

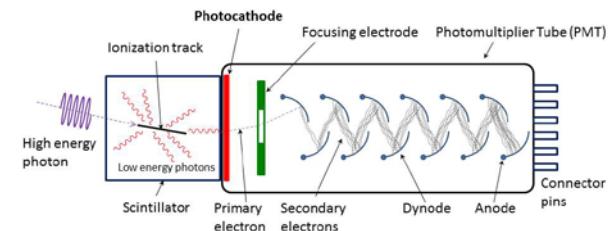
| | Dipole | | | | | Wiggler | |
|----------------|---------|-------------------------|------------|---------|---------------------------|---------|-------------------------|
| | B [T] | E_{crit} [keV] | ρ [m] | L [m] | $\Delta \theta$ [mrad] | B [T] | E_{crit} [keV] |
| CLIC DR | 1 | 5.4 | 9.5 | 0.6 | 60 | 2.5 | 14 |
| ASLS | 1.3 | 7.8 | 7.7 | 1.7 | 220 | 1.9 | 12 |

- LHC type ionization chambers (IC)

- N_2 at 1.1 bar
- 1.5 liter, 50 cm long

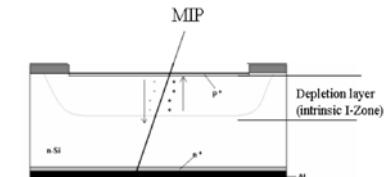
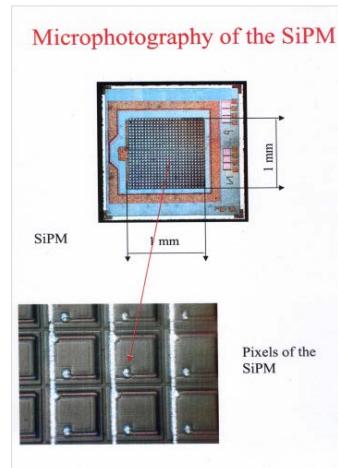


- Plastic NE102 scintillators, 0.8 liter ($25 \times 16 \times 2 \text{ cm}^3$) + photomultiplier (gain 10^4)



[File](#) by Qwerty123uiop / [CC BY-SA 3.0](#)

- Small quartz Cherenkov crystals, 1 cm^3 + photomultiplier (gain 10^4) OR + SiPM (gain 10^5)



(K. Wittenburg, CAS)

- Silicon PIN diodes in current mode, 1 cm^2 , depletion layer $100 \mu\text{m}$

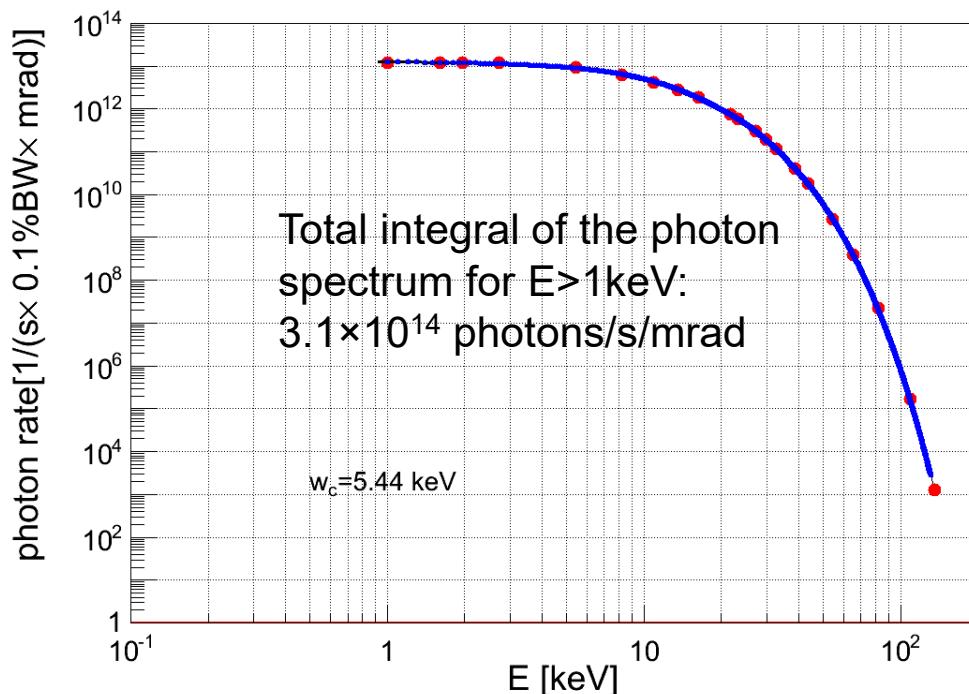
Synchrotron Radiation Spectrum

Sanja Damjanovic

- Spectrum is not hard enough to produce electrons above the Cherenkov threshold in quartz (≈ 190 keV)

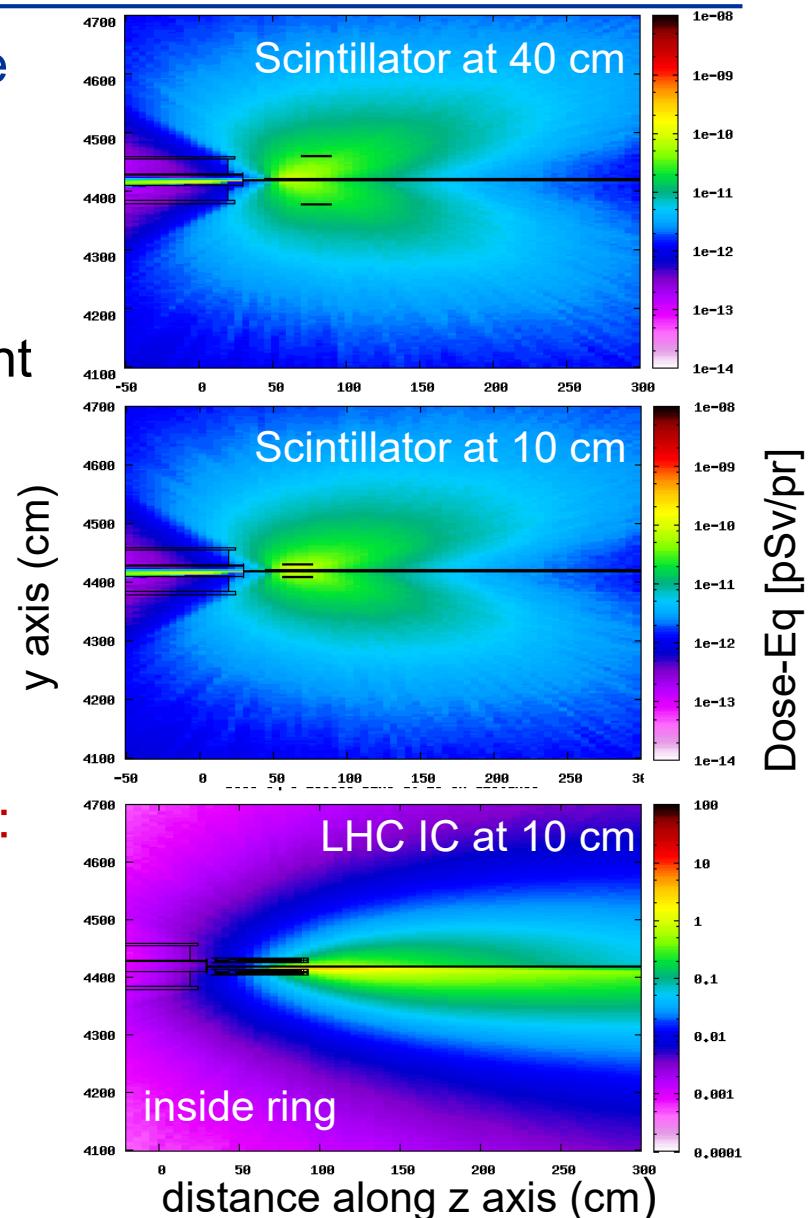


Cherenkov detectors are insensitive to the synchrotron radiation induced charged particle showers



$$N(\lambda) = 2.457 \times 10^{13} E[\text{GeV}] I[\text{A}] G_1 \left[\frac{\text{photons}}{\text{s} \cdot 0.1\% \text{BW} \cdot \text{mrad} \theta} \right]$$

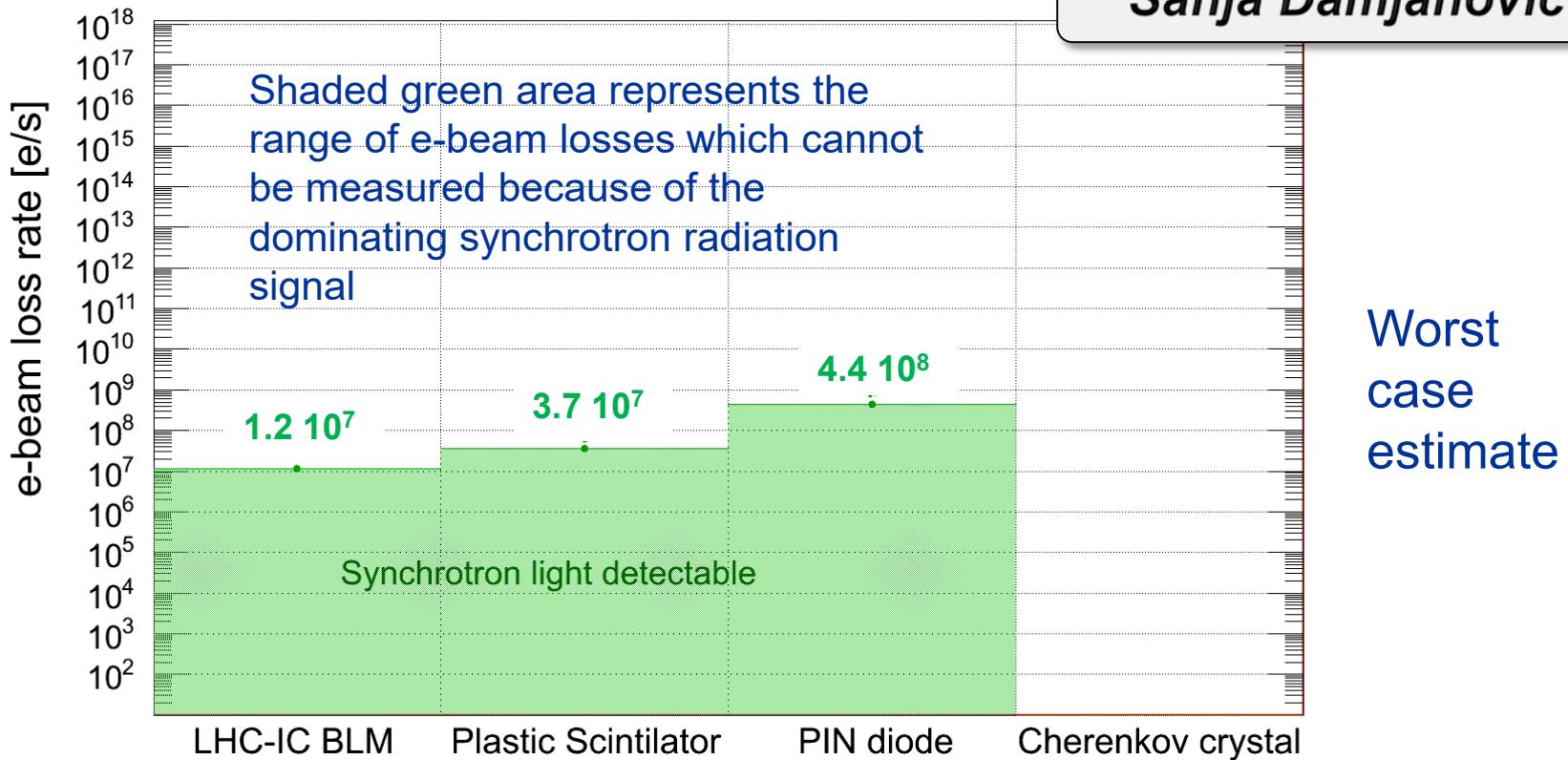
- Synchrotron radiation at 10 cm, outside of the ring, maximizing the signal:
 - IC 80 pA
 - Scintillator 64 A
 - PIN 300 pA < typ. dark current
 - Quartz insensitive



- Electron loss at same detector location:
 - IC 1.5×10^{-6} GeV/e
 - Scintillator 1.0×10^{-13} Gy/e
 - PIN 1.4×10^{-13} Gy/e
 - Quartz 9.3×10^{-14} Gy/e

Sensitivity Limits to Beam Loss due to Synchrotron Rad.

Sanja Damjanovic

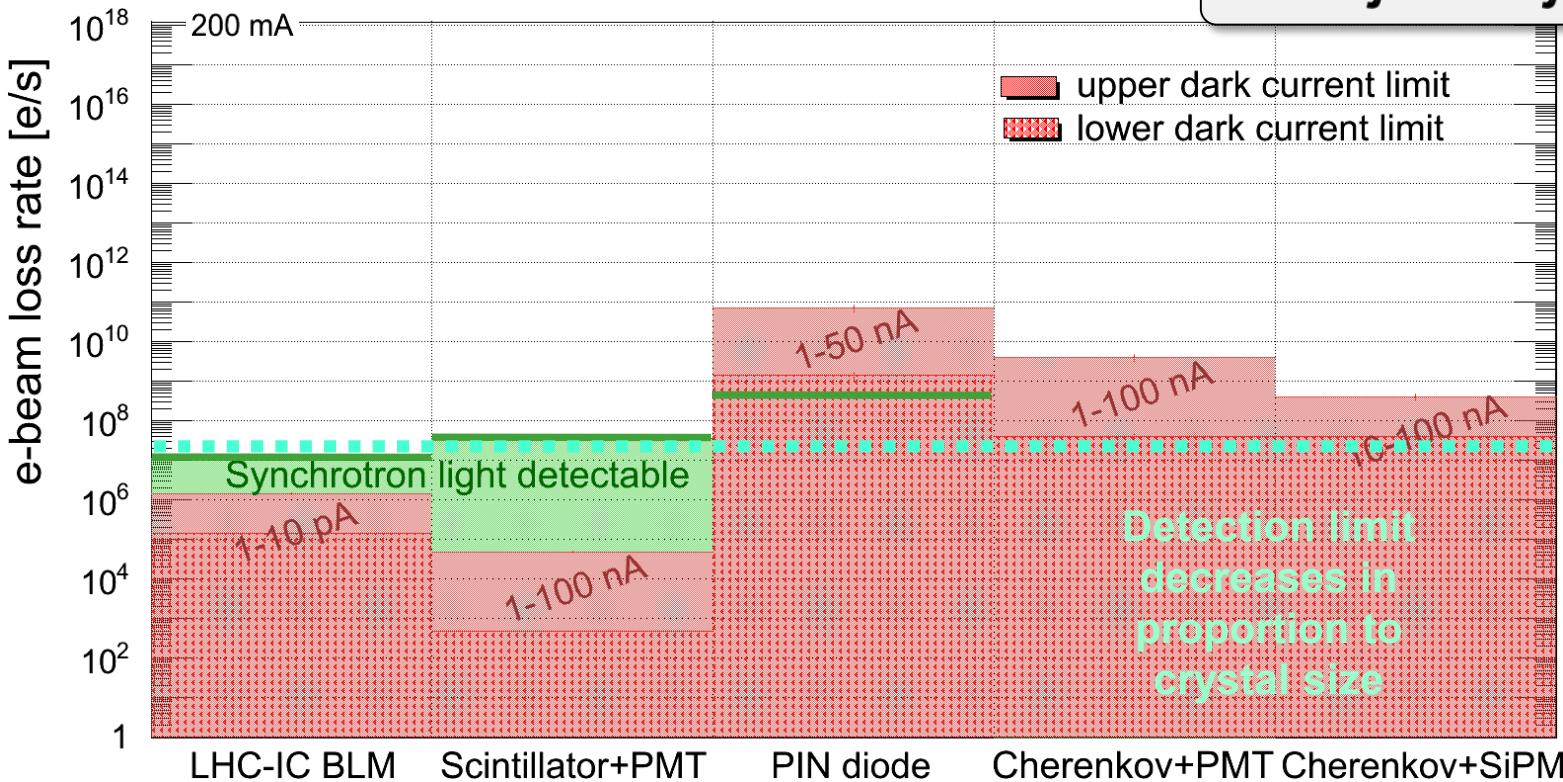


- e-beam loss rates generating detector currents equal to the currents produced by synchrotron light:

| | |
|----------------------|-----------------------|
| LHC-IC | 1.2×10^7 e/s |
| Plastic scintillator | 3.7×10^7 e/s |
| PIN diode | 4.4×10^8 e/s |

Sensitivity Limits due to Detector Dark Current

Sanja Damjanovic

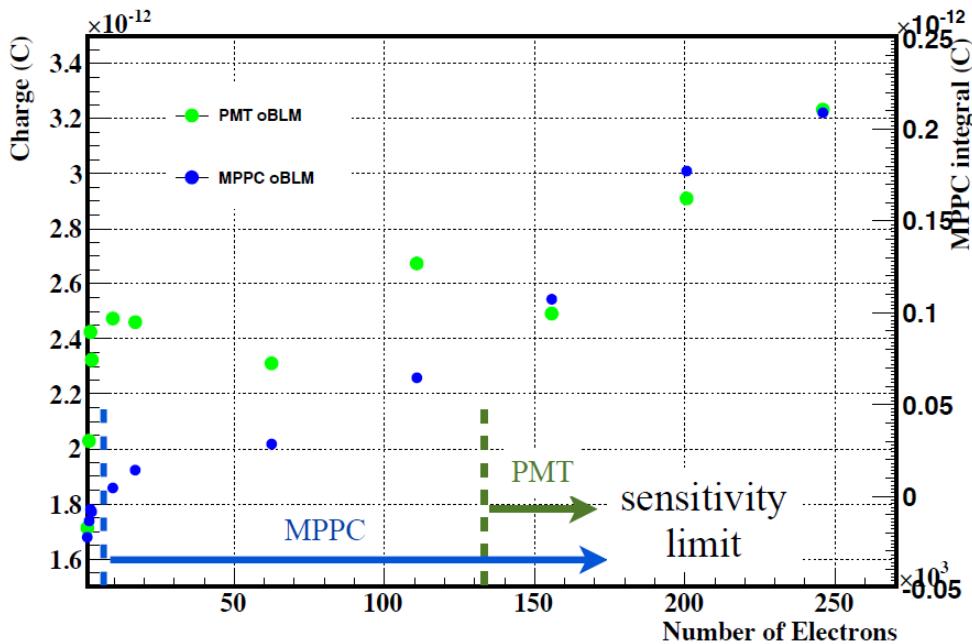


- Primary e-beam losses would be detectable for beam loss rates larger than:

| | |
|----------------------|------------------------------------|
| LHC-IC | $\geq 1.5 \times 10^5 \text{ e/s}$ |
| Plastic scintillator | $\geq 5 \times 10^2 \text{ e/s}$ |
| PIN diode | $\geq 1.5 \times 10^9 \text{ e/s}$ |
| Cherenkov crystal | $\geq 4 \times 10^7 \text{ e/s}$ |

Optical Fiber Loss Measurements, Australian Synchrotron

Eduardo Nebot et al.

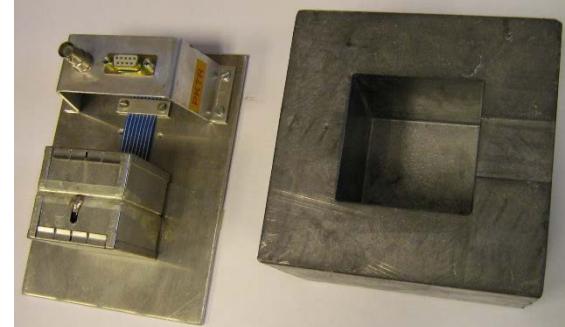
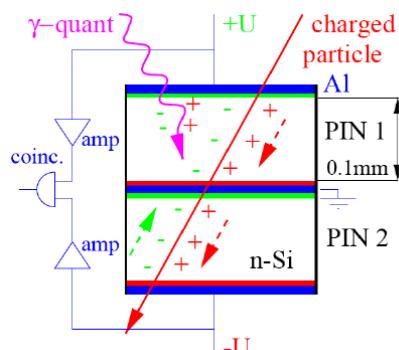


CLIC DR CRD:
 $2 \times 10^7 \text{ e/s/m}$

- Sensitive to $\approx 1 \times 10^4$ electrons lost in single location with MPPC (Multi-Pixel Photon Counter)
- Linear response up to 1×10^9 electrons with PMT
- Combination of photon-sensors \rightarrow dynamic range $\approx 10^5$

E. Nebot et al., "Measurement of Beam Losses at the Australian Synchrotron",
IBIC 2014

- Quench protection
- Superconducting 6.3 km p-ring (920 GeV/c)
- e-ring (30 GeV/c)
- Synchrotron radiation:
 - $E_{\text{crit}} = 88 \text{ keV}$
 - Dose rate $\approx 10^4 \text{ Gy/year}$
- Two silicon PIN diodes in coincidence counting mode + shielding box
 - Efficiency charge particle $> 30\%$
 - Efficiency photon 3.5×10^{-5}
- 10.4 MHz count rate (96 ns bunch spacing)
- Integration time: 5.2 ms — speed limited by low count rate
- Dynamic range of up to 10^9
- Good calibration: measured lifetime by current decay and losses agree within factor 2



- Photon coincidence counting:
 - Statistically
 - Electron created (photoelectric absorption, Compton effect) reaches second diode energy
 - → add thin metal layer between the diodes
- Electrons which lose energy are lost at dispersive aperture restrictions (close to horizontal quadrupoles)

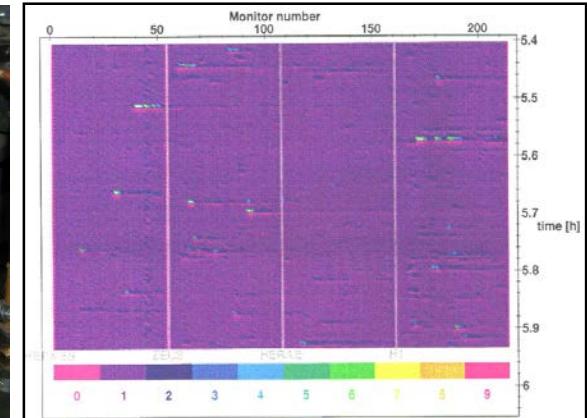
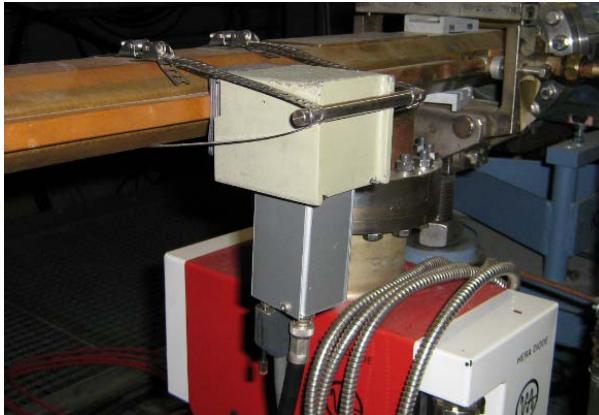
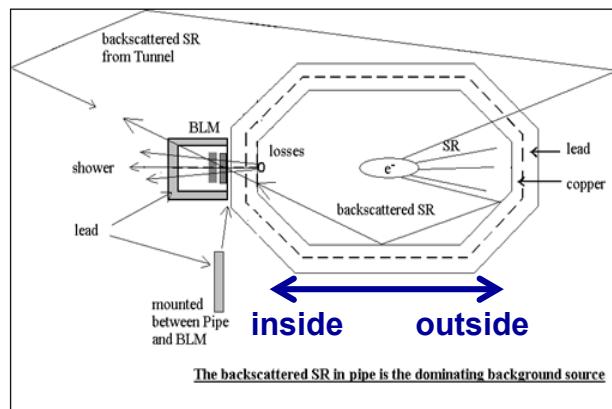
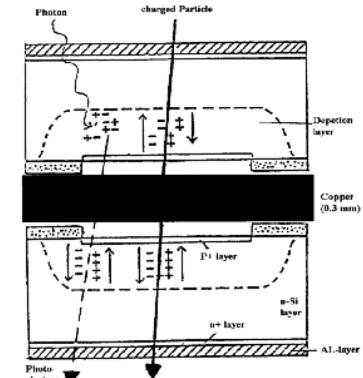


Fig. 2: The directions of synchrotron radiation and losses

Distant Losses

Distant Losses

Collimation regions

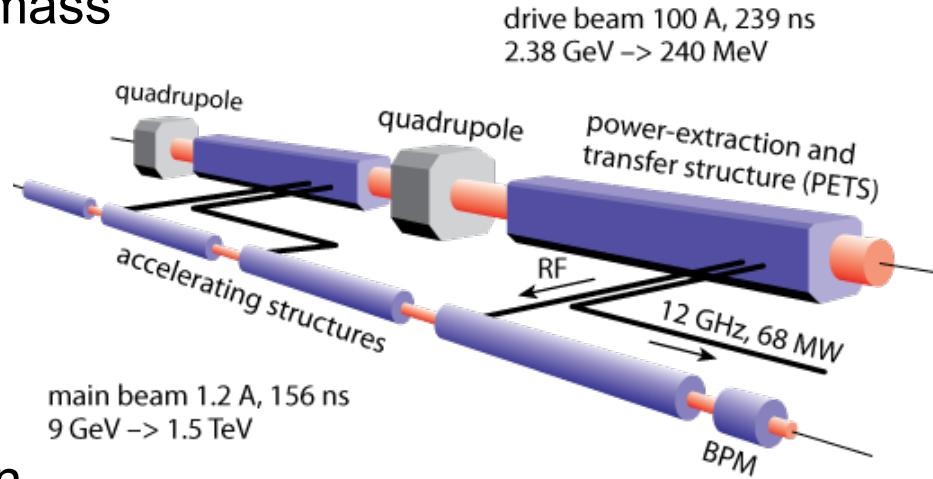
Injection and extraction

Around interaction points

The CLIC (Compact Linear Collider) Two-Beam-Module



- Future e^+e^- collider, centre of mass energy of 3 TeV
- High accelerating gradients (100 MV/m)
→ novel two beam acceleration method
- High intensity Drive Beam decelerated in Power-Extraction and Transfer Structures (PETS)
- RF power at 12 GHz is transferred to the Main Beam



| | Energy range [GeV] | Pulse length | Electrons per train |
|-------------------|------------------------|---------------------------------------|-----------------------|
| Drive Beam | $2.4 \rightarrow 0.24$ | $244 \text{ ns} \approx 80 \text{ m}$ | 1.53×10^{14} |
| Main Beam | $9 \rightarrow 1500$ | $156 \text{ ns} \approx 50 \text{ m}$ | 1.16×10^{12} |

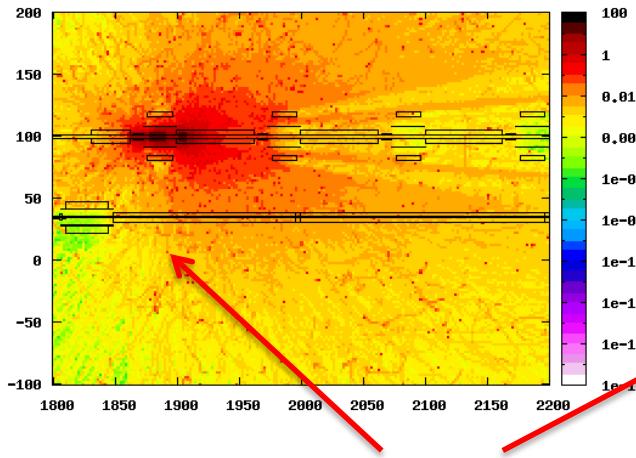
Distinguish Losses CLIC Two-Beam Module – “crosstalk”

- Dose (Gy) from FLUKA simulations:

Sophie Mallows

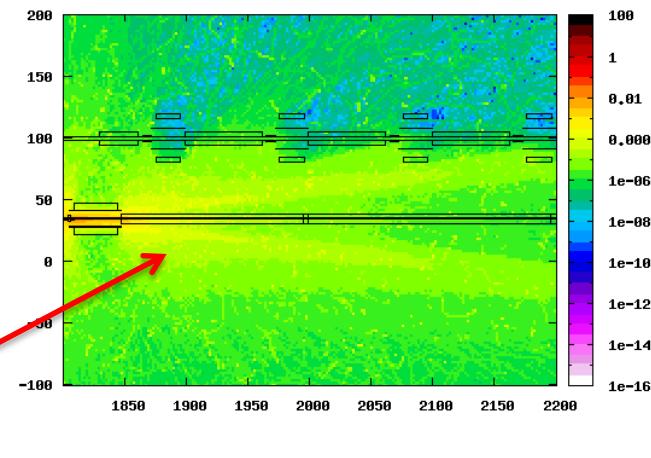
Destructive Drive Beam loss

1.0% of bunch train hits single aperture restriction



Destructive Main Beam loss

0.01% of bunch train hits single aperture restriction



- At the very beginning of the Main Beam: Destructive Drive Beam loss provokes similar signal as destructive Main Beam loss in the region close to Main Beam quadrupole
- Not a machine protection issue — dangerous loss would never go unnoticed**

| | Energy |
|------------|---------|
| Drive Beam | 2.4 GeV |
| Main Beam | 9 GeV |

First Measurements CTF3 Two-Beam Module – “crosstalk”

Maria Kastriotou



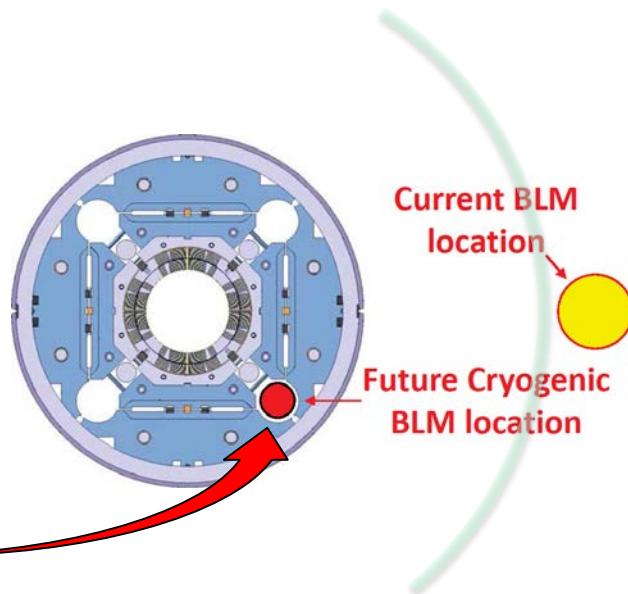
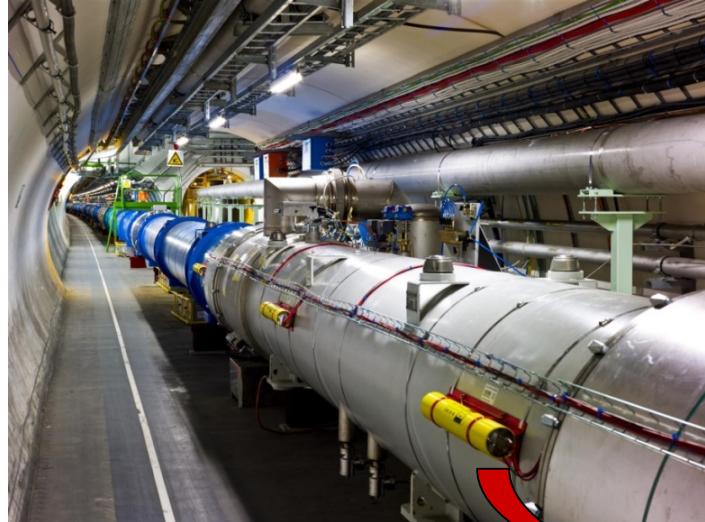
- Little ionization chambers
- Optical fibers
- Drive / Main Beam: 120 / 200 MeV
- First results show 1–5% “crosstalk” on the other beam

M. Kastriotou et al., “BLM Crosstalk Studies at the CLIC Two-Beam Module”, MOPB045

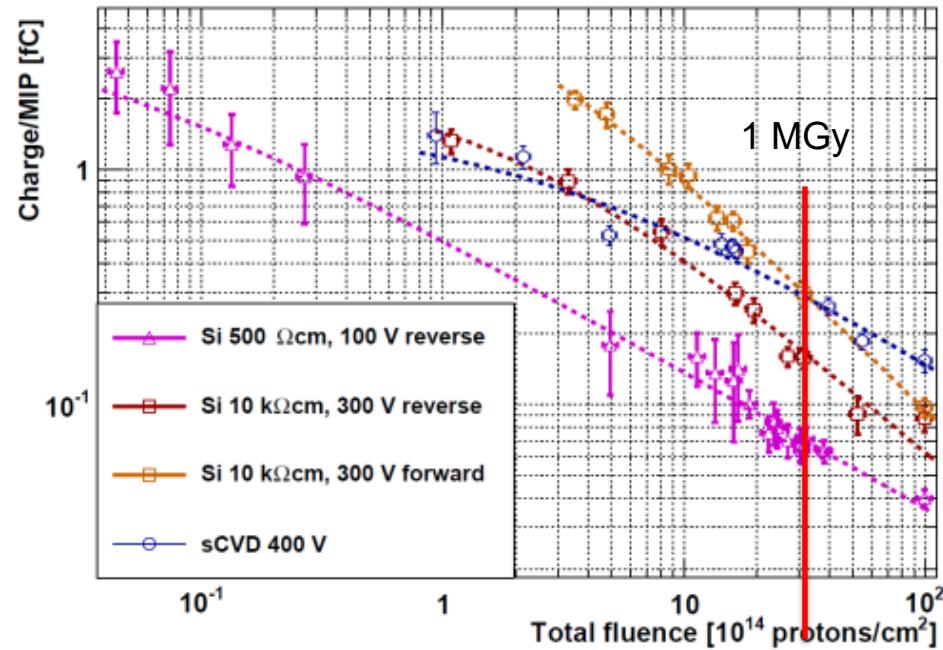
- 10^{-3} losses of either beam → unacceptable luminosity losses due to beam loading variations
- FLUKA simulations show: $BLM_{DriveBeam}$ up to $100 \times BLM_{MainBeam}$
- How to measure Main Beam losses in the presence of Drive Beam losses?



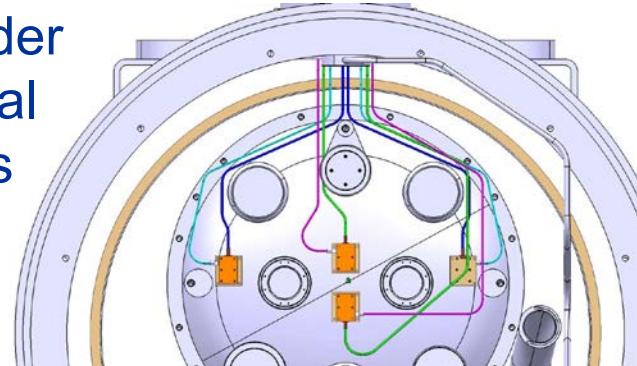
- Loss monitor closer to loss location → avoid that the signal is dominated by other radiation sources (e.g. physics debris at the insertion region triplet magnets)
- HL-LHC triplets: move detector inside the cryostat
 - Operate in liquid helium at 1.9 K and at 2 T
 - Dose of 2 MGy in 20 years, without access
 - Fast pressure rise (1.1 to 20 bar) in case of quench



- Since 2011 extensive classification program of the detector parameters in cryogenic temperature under radiation
- Single crystal chemical vapour deposition (scCVD) diamond
- $p^+ \text{-} n \text{-} n^+$ silicon
- Both materials can operate up to 2 MGy:
 - Diamond/silicon sensitivity reduction 14% / 25%
 - Leakage current of Si irradiated at cold much less than at room temperature



Now testing both detectors in LHC cryostat at $\approx 20 \text{ K}$ under operational conditions



M. Bartosik, "Cryogenic Beam Loss Monitors for the Superconducting Magnets of the LHC", MOPB042

Accelerating Structures

Dark current (field emission) and
voltage breakdown (electric arcs)
→ electrons and X-rays

Limits the sensitivity to primary
beam losses

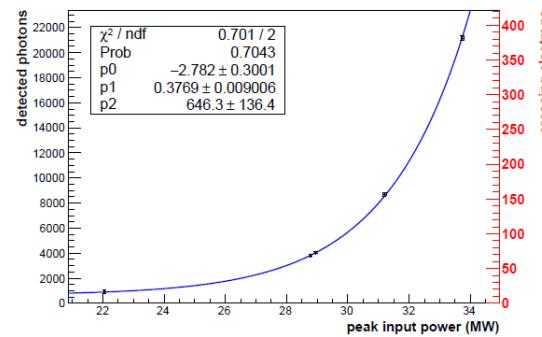
Can be monitored with BLM

- Test at CTF3 of unloaded cavity
- Optical fiber 900 μm at 2.5 cm distance (30 cm exposure)
- 200ns RF pulses: Measured dark current and breakdown as function of cavity input power
- Extrapolated to 40 MW (unloaded) and 60 MW (loaded) structure
- Detected photons: 2×10^5 (dark current), 6×10^6 (breakdown) at 40 MW
- → very high electron background

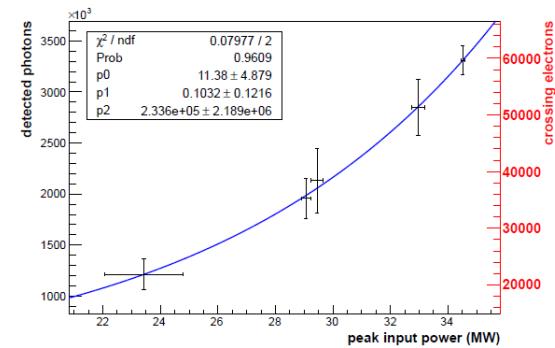
M. Kastriotou et al., "RF cavity induced sensitivity limitations on Beam Loss Monitors", LA3NET 2015, Physics Procedia, to be published



(a) The BLM (red fibre) above the CLIC structure (in the middle)



(a) Electron field emission



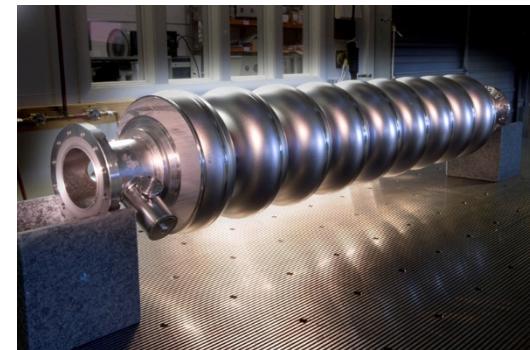
(b) RF breakdown

Fig. 5. Detected photons and estimation of the number of electrons with $\beta = 1$ that will give the respective signal when crossing the BLM

Cryogenic Loss Monitors (CLM) at Fermilab

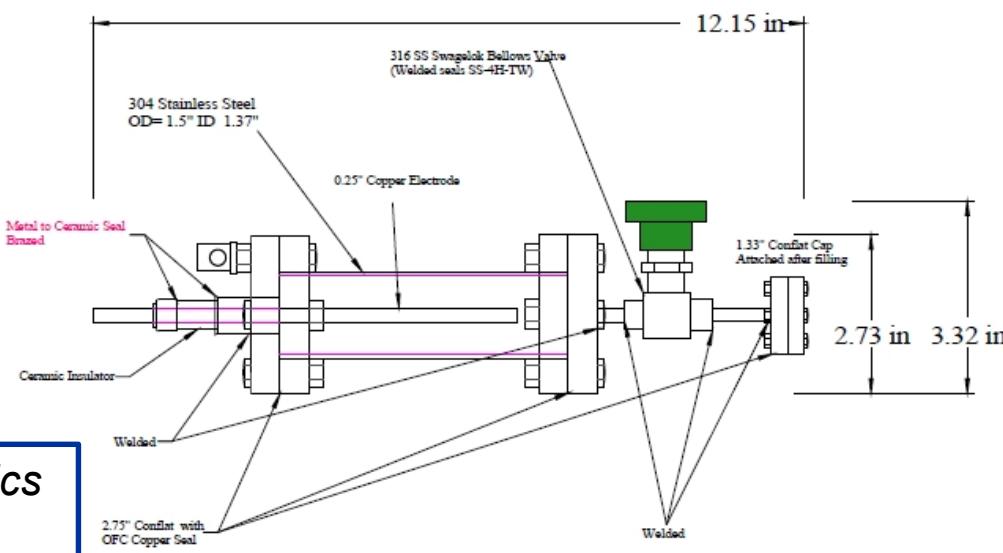
Arden Warner

- Superconducting RF (SRF) cavities measure primary beam losses and dark current losses at Fermilab's Advanced Superconducting Test Accelerator (ASTA)
- 5 K (inside SRF cryo-module) to 350 K
- Stainless steel vessel, coaxial design, 120 cm³, He-gas 1.0–1.5 bar
- Sensitivity: 1.9 pA/(rad/h), max. dose rate: 30 krad/h
- Readout via current-to-frequency converter and FPGA-TDC
- Housing at -95 V



Fill port

A. Warner and J. Wu, Physics
Procedia 37 (2012) 2031.



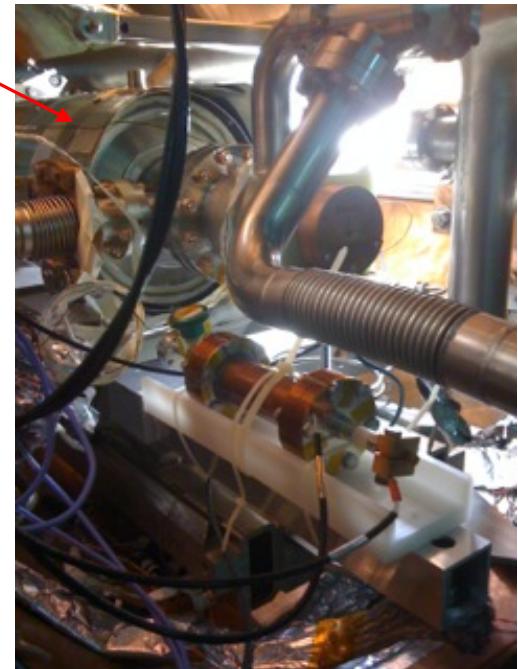
- Dark current measurements at the A0-photo-injector test accelerator and the Horizontal Test Stand (HTS)

Loss due to Dark current background at A0-photo-injector;
measured to be ≈ 400 nA downstream of bend magnet



40 μ s RF gate (dark current only no photo-electrons injected)

Test cavity



HTS installation

A. Warner and J. Wu, "Cryogenic loss monitors with FPGA TDC signal processing", Physics Procedia 37 (2012) 2031.

Summary

- Pro — contra of localized vs. distributed systems
 - Position resolution
 - Time resolution
 - Absolute measurement
- Challenges for machine protection, large machines and radiation environment
- Background sources limiting sensitivity to primary beam losses
 - Synchrotron radiation
 - Distant beam losses and physics collision debris
 - RF accelerating structures

**Thank you
for your
Attention**

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