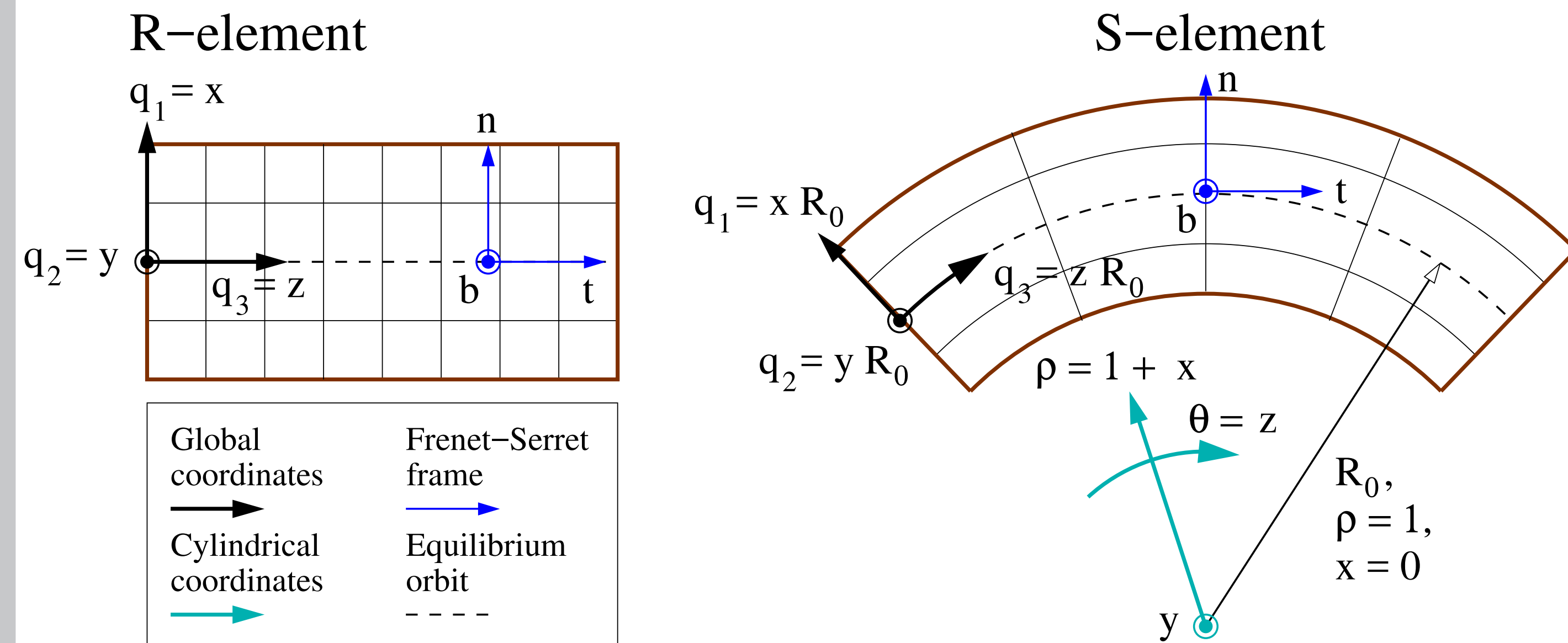


## 1. R- and S- elements



## 2. Multipoles in Cartesian Coordinates — Homogeneous Harmonic Polynomials

$$\Delta_{\perp} \Phi = \frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} = 0$$

$$\Delta_{\perp} \mathbf{A} = \left( \frac{\partial^2 A_z}{\partial x^2} + \frac{\partial^2 A_z}{\partial y^2} \right) \hat{\mathbf{e}}_z = 0$$

Solutions are homogeneous harmonic polynomials of two variables

$$\mathcal{A}_n(x, y) = \Re \mathcal{Z}^n = \frac{1}{2} [(x + i y)^n + (x - i y)^n] = \sum_{k=0}^n \binom{n}{k} x^{n-k} y^k \cos \frac{k \pi}{2}$$

$$\mathcal{B}_n(x, y) = \Im \mathcal{Z}^n = \frac{1}{2i} [(x + i y)^n - (x - i y)^n] = \sum_{k=0}^n \binom{n}{k} x^{n-k} y^k \sin \frac{k \pi}{2}$$

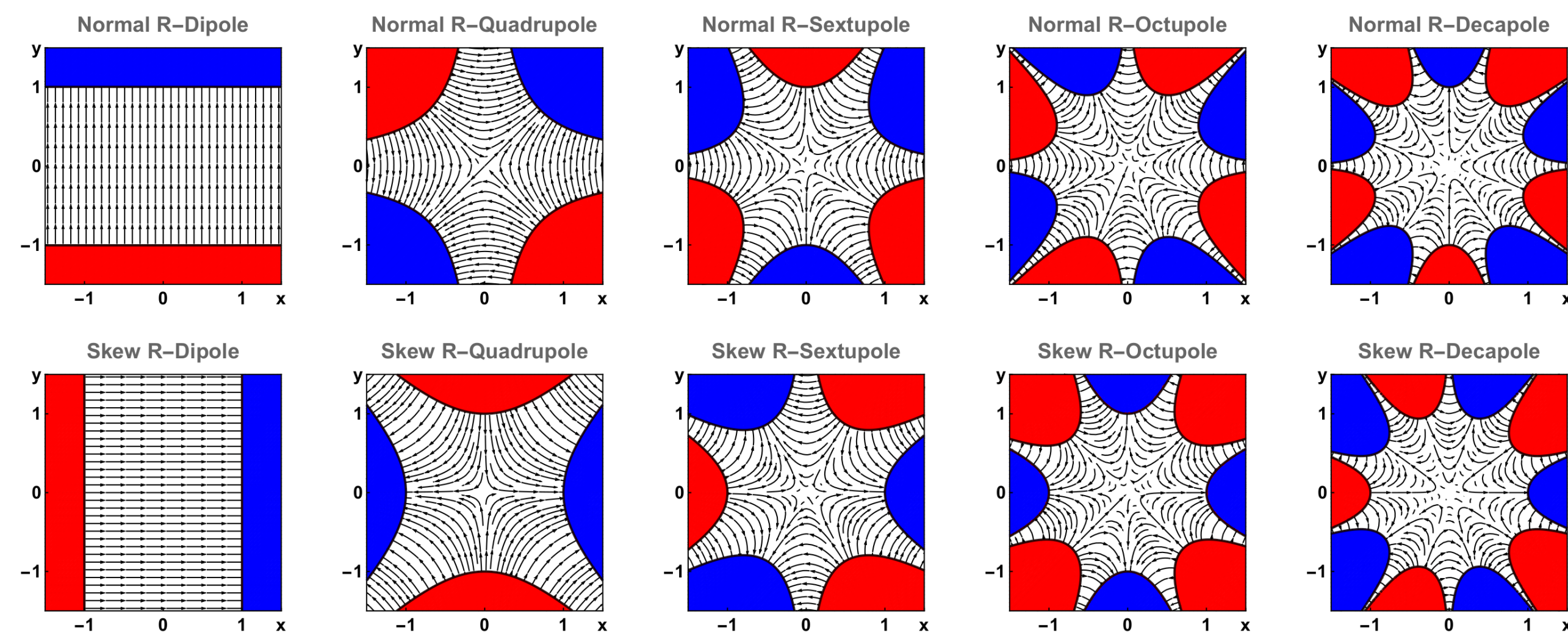
related to each other through the Cauchy-Riemann equation

$$\frac{\partial \mathcal{A}_n}{\partial x} = \frac{\partial \mathcal{B}_n}{\partial y} \quad \frac{\partial \mathcal{A}_n}{\partial y} = -\frac{\partial \mathcal{B}_n}{\partial x}$$

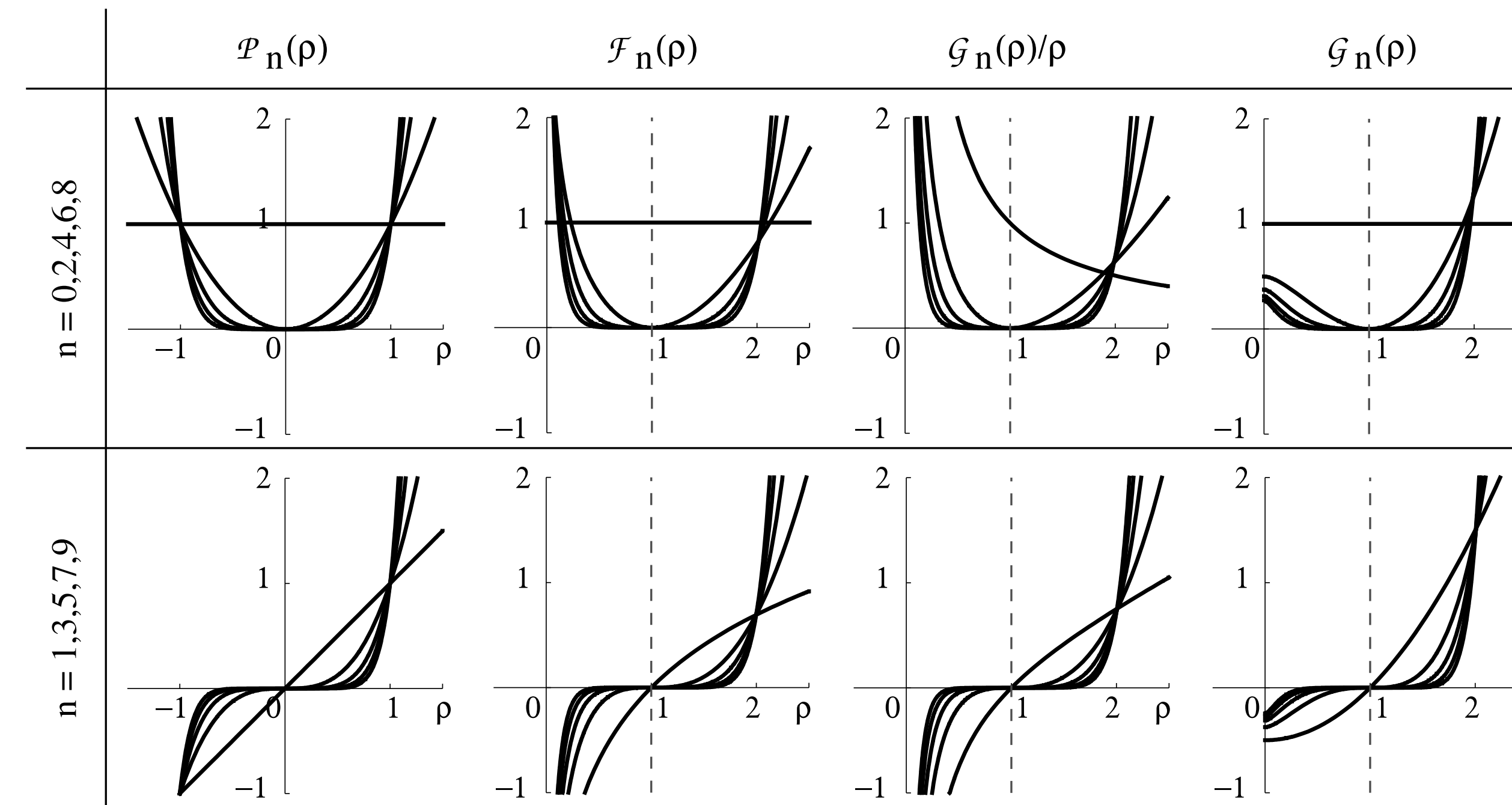
with lowering operator which relates functions of different orders to each other

$$n (\mathcal{A}, \mathcal{B})_{n-1} = \frac{\partial}{\partial x} (\mathcal{A}, \mathcal{B})_n = \pm \frac{\partial}{\partial y} (\mathcal{B}, \mathcal{A})_n \quad \mathcal{A}_0 = 1, \mathcal{B}_0 = 0.$$

Normal	Skew
$\overline{\Phi}^{(n)} = -C_n \frac{\mathcal{B}_n^{(\text{e})}}{n!}$	$\overline{A}_{\theta}^{(n)} = -C_n \frac{\mathcal{A}_n^{(\text{m})}}{n!}$
$\overline{F}_{\rho}^{(n)} = C_n \frac{\mathcal{B}_{n-1}^{(\text{m})}}{(n-1)!}$	$\overline{F}_y^{(n)} = C_n \frac{\mathcal{A}_{n-1}^{(\text{e})}}{(n-1)!}$
$\Phi^{(n)} = -D_n \frac{\mathcal{A}_n^{(\text{e})}}{n!}$	$\underline{A}_{\theta}^{(n)} = D_n \frac{\mathcal{B}_n^{(\text{m})}}{n!}$
$F_{\rho}^{(n)} = D_n \frac{\mathcal{A}_{n-1}^{(\text{m})}}{(n-1)!}$	$\underline{F}_y^{(n)} = -D_n \frac{\mathcal{B}_{n-1}^{(\text{e})}}{(n-1)!}$



## 3. Multipoles in Cylindrical Coordinates and McMillan Harmonics



First step is to restore the symmetry

$$\Delta_{\cap} \Phi = \Delta_{\perp} \Phi + \frac{1}{\rho} \frac{\partial \Phi}{\partial \rho} = \frac{\partial^2 \Phi}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial \Phi}{\partial \rho} + \frac{\partial^2 \Phi}{\partial y^2} = 0$$

$$\Delta_{\cap} \mathbf{A} = \left( \Delta_{\cap} A_{\theta} - \frac{A_{\theta}}{\rho^2} \right) \hat{\mathbf{e}}_{\theta} = \frac{\hat{\mathbf{e}}_{\theta}}{\rho} \left[ \frac{\partial^2}{\partial \rho^2} - \frac{1}{\rho} \frac{\partial}{\partial \rho} + \frac{\partial^2}{\partial y^2} \right] (\rho A_{\theta}) = 0$$

Looking for the solution in a form similar to HHP

$$\Phi = - \sum_{k=0}^n \frac{\mathcal{F}_{n-k}(\rho) y^k}{(n-k)! k!} \left( C_n \sin \frac{k \pi}{2} + D_n \cos \frac{k \pi}{2} \right)$$

$$A_{\theta} = - \sum_{k=0}^n \frac{1}{\rho} \frac{\mathcal{G}_{n-k}(\rho) y^k}{(n-k)! k!} \left( C_n \cos \frac{k \pi}{2} - D_n \sin \frac{k \pi}{2} \right)$$

where functions  $\mathcal{F}_n(\rho)$  and  $\mathcal{G}_n(\rho)$  are determined by two recurrence equations

$$\frac{\partial^2 \mathcal{F}_n(\rho)}{\partial \rho^2} + \frac{1}{\rho} \frac{\partial \mathcal{F}_n(\rho)}{\partial \rho} = n(n-1) \mathcal{F}_{n-2}(\rho)$$

$$\frac{\partial^2 \mathcal{G}_n(\rho)}{\partial \rho^2} - \frac{1}{\rho} \frac{\partial \mathcal{G}_n(\rho)}{\partial \rho} = n(n-1) \mathcal{G}_{n-2}(\rho)$$

one can see that  $\mathcal{F}_n$  and  $\mathcal{G}_n$  are related to each other through

$$\mathcal{G}_{n-1} = \frac{1}{n} \rho \frac{\partial \mathcal{F}_n}{\partial \rho} \quad \text{and} \quad \mathcal{F}_{n-1} = \frac{1}{n \rho} \frac{\partial \mathcal{G}_n}{\partial \rho}$$

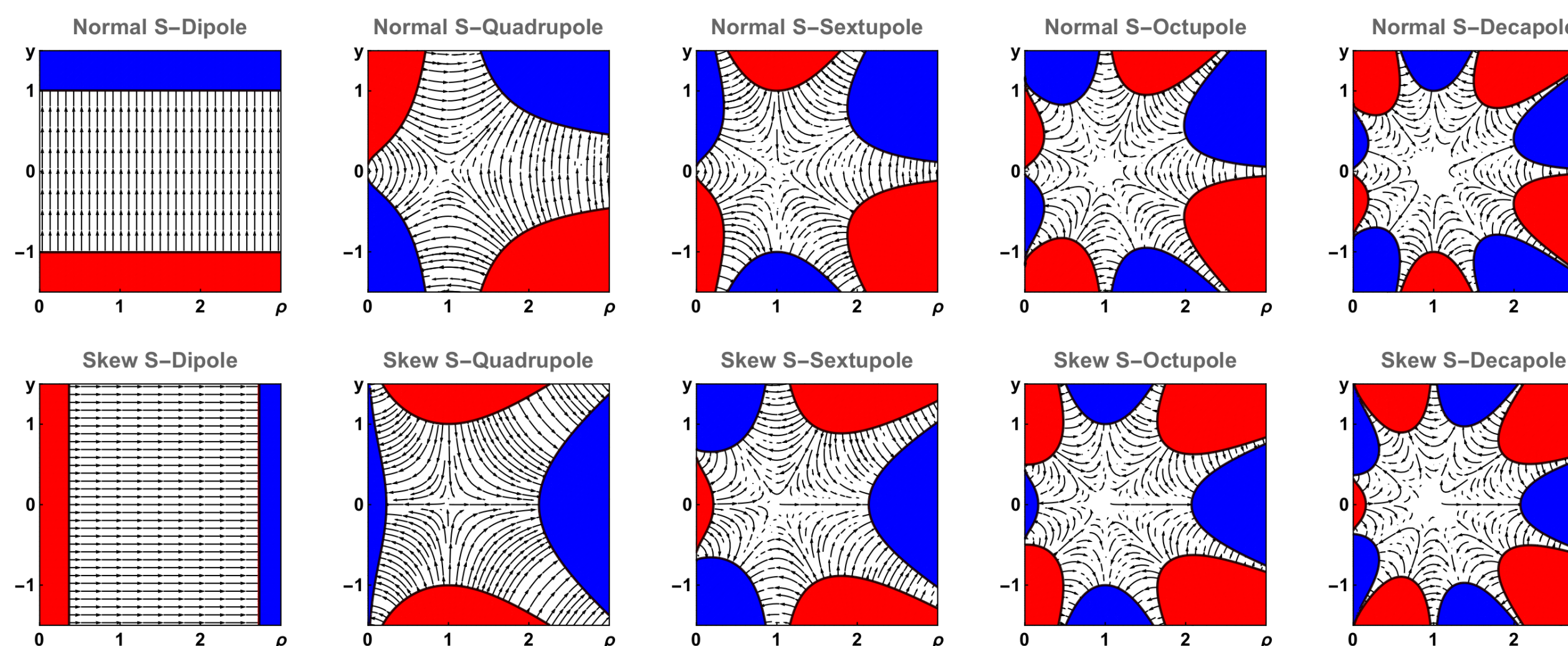
and allow the construction of lowering and corresponding raising operators

$$\mathcal{F}_n = \frac{1}{(n+1)(n+2)} \left[ \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial}{\partial \rho} \right) \right] \mathcal{F}_{n+2} \quad \mathcal{F}_n = n(n-1) \int_1^{\rho} \frac{1}{\rho} \int_1^{\rho} \rho \mathcal{F}_{n-2} d\rho d\rho$$

$$\mathcal{G}_n = \frac{1}{(n+1)(n+2)} \left[ \rho \frac{\partial}{\partial \rho} \left( \frac{1}{\rho} \frac{\partial}{\partial \rho} \right) \right] \mathcal{G}_{n+2} \quad \mathcal{G}_n = n(n-1) \int_1^{\rho} \rho \int_1^{\rho} \frac{1}{\rho} \mathcal{G}_{n-2} d\rho d\rho$$

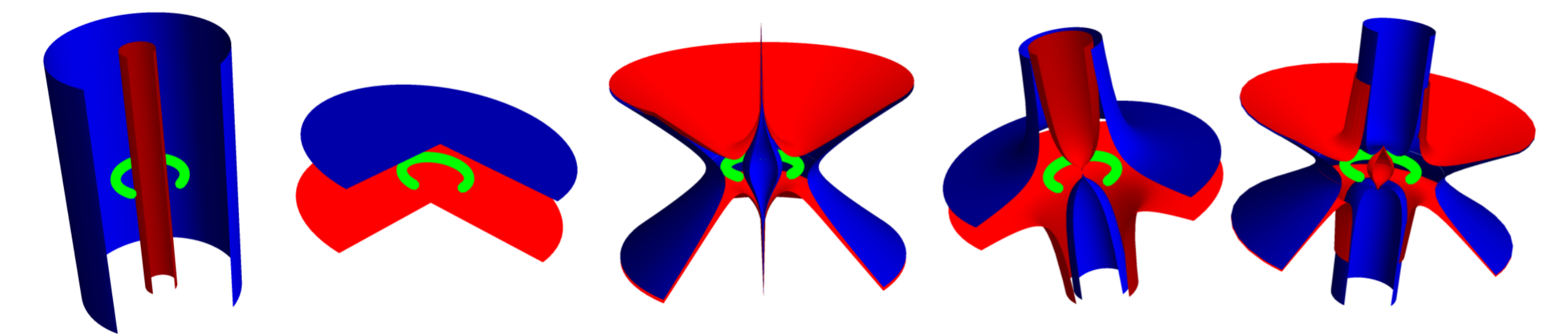
with an additional constraint to terminate recurrences defining lowest orders as

$$\mathcal{F}_0 = 1, \quad \mathcal{F}_1 = \ln \rho, \quad \mathcal{G}_0 = 1, \quad \mathcal{G}_1 = (\rho^2 - 1)/2.$$



## 4. Summary

<https://arxiv.org/pdf/1603.03451v1.pdf>



3D models of 2n-pole sector magnets. From the left to the right: skew S-dipole, normal S-dipole, skew S-quadrupole, normal S-quadrupole and skew S-sextupole.

Table 1: Sector harmonics.

n	
$\mathcal{A}_n^{(\text{e})}$	0 1
1	$\ln \rho$
2	$\left[ \frac{\rho^2 - 1}{2} - y^2 \right] - \ln \rho$
3	$\left[ -3 \frac{\rho^2 - 1}{2} \right] + 3 \left( \frac{\rho^2 + 1}{2} - y^2 \right) \ln \rho$
4	$\left[ \frac{3(\rho^4 + 4\rho^2 - 5)}{8} - 6 \frac{\rho^2 - 1}{2} y^2 + y^4 \right] - 3 \left( \frac{1}{2} + \rho^2 - 2 y^2 \right) \ln \rho$
$\mathcal{A}_n^{(\text{m})}$	0 $\frac{1}{\rho}$
1	$\frac{1}{\rho} \left[ \frac{\rho^2 - 1}{2} \right]$
2	$\frac{1}{\rho} \left\{ \left[ -\frac{\rho^2 - 1}{2} - y^2 \right] + \rho^2 \ln \rho \right\}$
3	$\frac{1}{\rho} \left\{ \left[ \frac{3(\rho^2 + 1)\rho^2 - 1}{4} - 3 \frac{\rho^2 - 1}{2} y^2 \right] - \frac{3}{2} \rho^2 \ln \rho \right\}$
4	$\frac{1}{\rho} \left\{ \left[ -\frac{3(5\rho^4 - 4\rho^2 - 1)}{8} + 6 \frac{\rho^2 - 1}{2} y^2 + y^4 \right] + \frac{3(2 + \rho^2 - 4 y^2)}{2} \rho^2 \ln \rho \right\}$
$\mathcal{B}_n^{(\text{e})}$	0 0
1	y
2	y [2 ln rho]
3	y { [3 \frac{\rho^2 - 1}{2} - y^2] - 3 ln rho }
4	y { [ -12 \frac{\rho^2 - 1}{2} ] + 4 (3 \frac{\rho^2 + 1}{2} - y^2) ln rho }
$\mathcal{B}_n^{(\text{m})}$	0 0
1	$\frac{y}{\rho}$
2	$\frac{y}{\rho} \left[ 2 \frac{\rho^2 - 1}{2} \right]$
3	$\frac{y}{\rho} \left\{ \left[ -3 \frac{\rho^2 - 1}{2} - y^2 \right] + 3 \rho^2 \ln \rho \right\}$
4	$\frac{y}{\rho} \left\{ \left[ 4 \frac{3(\rho^2 + 1)\rho^2 - 1}{4} - 4 \frac{\rho^2 - 1}{2} y^2 \right] - 6 \rho^2 \ln \rho \right\}$

Table 2: Relationship between the coefficients of pure normal and skew sector multipoles, and, power series expansion of field in radial and vertical planes.

$C_n$	$D_n$
$n \quad x = 0 \quad y = 0$	$x = 0 \quad y = 0$
1 $F_y \quad F_y$	$F_x \quad F_x$
2 $\partial_y F_x \quad \partial_x F_y$	$-\partial_y F_y \quad \partial_x F_x + F_x$
3 $-\partial_y^2 F_y \quad \partial_x^2 F_y + \partial_x F_y$	$-\partial_y^2 F_x \quad \partial_x^2 F_x + \partial_x F_x - F_x$
4 $-\partial_y^3 F_x \quad \partial_x^3 F_y + \partial_x^2 F_y - \partial_x F_y$	$\partial_y^3 F_y \quad \partial_x^3 F_x + 2 \partial_x^2 F_x - \partial_x F_x + F_x$