

Iván Mauricio Burbano Aldana

Perimeter Scholars International

Quantum Field Theory II

Homework 2

1. a) We have that the numerator on the r.h.s. of (2) has the perturbative expansion

$$\begin{aligned}
 & \int \mathcal{D}\phi \exp\left(-\frac{1}{\hbar} S_0(\phi)\right) \exp\left(-\frac{g}{3!\hbar} \int d^d x \phi(x)^3\right) \phi(x_1) \dots \phi(x_N) \\
 &= \sum_{n=0}^{\infty} \frac{1}{n!} \left(-\frac{g}{3!\hbar}\right)^n \int \prod_{i=1}^n d^d z_i \int \mathcal{D}\phi \exp\left(-\frac{1}{\hbar} S_0(\phi)\right) \prod_{i=1}^n \phi(z_i)^3 \prod_{i=1}^N \phi(x_i) \\
 &= \sum_{n=0}^{\infty} \left(-\frac{g}{\hbar}\right)^n \frac{1}{3!^n n!} \int \prod_{i=1}^n d^d z_i \langle \prod_{i=1}^n \phi(z_i)^3 \prod_{i=1}^N \phi(x_i) \rangle_0 Z_0,
 \end{aligned}$$

where $S_0(\phi) = \int d^d x \left(\frac{1}{2} (\nabla \phi(x))^2 + \frac{1}{2} m^2 \phi(x)^2 \right)$,

$$Z_0 = \int D\phi \exp\left(-\frac{1}{\hbar} S_0(\phi)\right), \quad \text{and}$$

$$\langle A \rangle_0 = \frac{1}{Z_0} \int D\phi \exp\left(-\frac{1}{\hbar} S_0(\phi)\right) A(\phi).$$

Then, we see there is a diagrammatic expansion

$$\langle \phi(x_1) \dots \phi(x_N) \rangle = \sum_{\substack{\text{Diagrams } G \text{ with} \\ N \text{ external vertices} \\ \text{without vacuum subdiagrams}}} c(G)$$

where $c(G)$ is calculated as follows

$$1. \quad \begin{array}{c} \bullet \\ \text{---} \\ \bullet \\ x \quad y \end{array} = \hbar G_0(x-y)$$

$$2. \quad \begin{array}{c} \diagup \\ \bullet \\ \diagdown \\ | \\ \bullet \\ x \end{array} = -\frac{g}{\hbar} \int d^d x$$

3. Divide by the symmetry factor.

Remarks: • The vacuum subdiagrams are cancelled

by the denominator, as can be seen

from our perturbative expansion with $N=1$.

• The symmetry factors come from the

factor $\frac{1}{3!^n n!}$, the relabelling of the internal

vertices z_1, \dots, z_n and the permutation of

the wick contractions within each group

$\phi(z_i)^3$.

b) We have $\langle \phi(x)\phi(y) \rangle$ is,

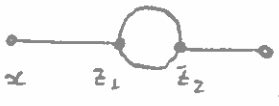
$O(g^0)$:

$$* \quad \text{---} \quad = i G_0(x-y),$$

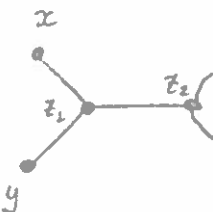
$O(g^1)$: None, since $\langle \phi(x)\phi(y)\phi(z_1)^3 \rangle$ has an

odd number of fields,

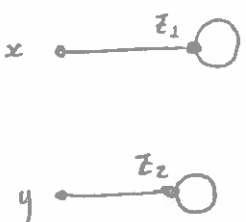
$\mathcal{O}(g^2)$

*  $= \frac{1}{2} \left(-\frac{g}{\hbar} \right)^2 \int d^d z_1 d^d z_2 \hbar G(x-z_1) \hbar G(z_1-z_2) \times$
 $\hbar G(z_1-z_2) \hbar G(z_2-y)$

$$= \frac{1}{2} g^2 \hbar^2 \int d^d z_1 d^d z_2 G(x-z_1) G(z_1-z_2)^2 G(z_2-y),$$

*  $= \frac{1}{2} \left(-\frac{g}{\hbar} \right)^2 \int d^d z_1 d^d z_2 \hbar G(x-z_1) \hbar G(y-z_1) \times$
 $\hbar G(z_1-z_2) \hbar G(z_2-z_2)$

$$= \frac{1}{2} g^2 \hbar^2 \int d^d z_1 d^d z_2 G(x-z_1) G(y-z_1) G(z_1-z_2) G(0),$$

*  $= \frac{1}{4} \left(-\frac{g}{\hbar} \right)^2 \int d^d z_1 d^d z_2 \hbar G(x-z_1) \hbar G(z_1-z_1) \times$
 $\hbar G(y-z_2) \hbar G(z_2-z_2)$

$$= \frac{1}{4} g^2 \hbar^2 \int d^d z_1 G(x-z_1) \int d^d z_2 G(y-z_2) G(0)^2.$$