

#### 5th NCSP DKIST Data-Training Workshop: He I Diagnostics in the Solar Atmosphere





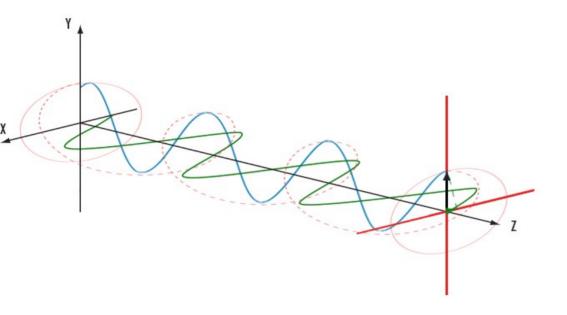
# Scattering Polarization and Hanle effect in He 10830

Ivan Milic (CU Boulder / LASP); ivan.milic (at) colorado.edu

## Scattering and polarization are like pizza and tomato

- The "slabs" (i.e. prominences and filaments), are not very dense, so they purely scatter the light.
- Scattering automatically creates the polarization.
- Polarization is a very powerful tool, it probes geometry, spatial anisotropy (or anisotropies), as well as the magnetic field.
- In the same way we can look at the variation of the intensity with the wavelength (i.e. the spectrum), we can look at the variation of the polarization state of the light with the wavelength (i.e. polarized spectrum, in case of Stokes Q only, second solar spectrum).

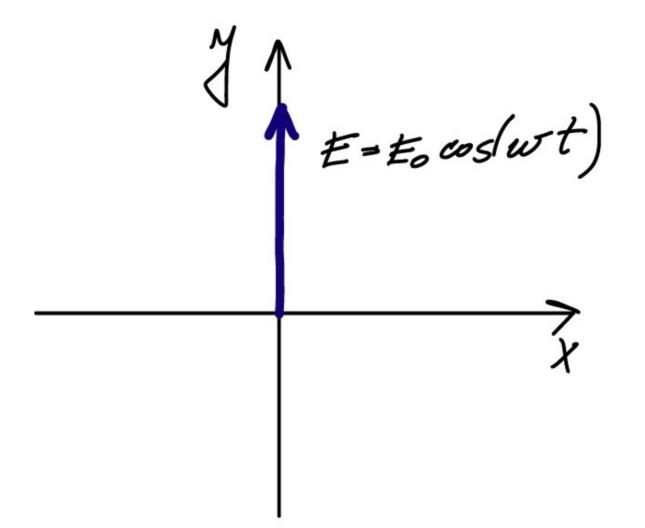
### How do we describe the polarization of the light?



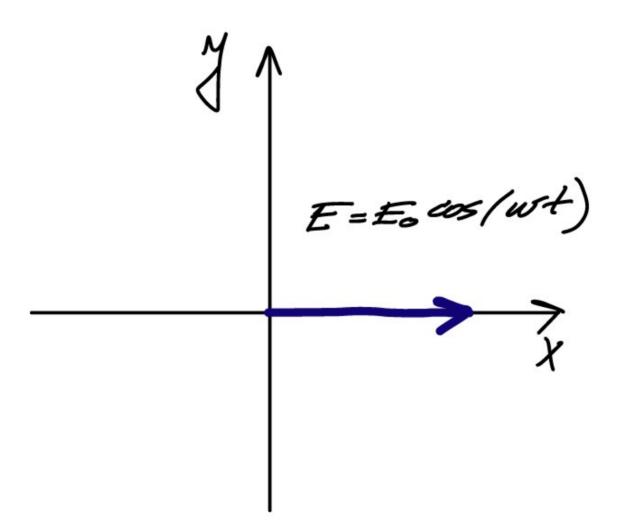
Credits: www.edmundoptics.com

"Instrumental" definition of Stokes parameters:

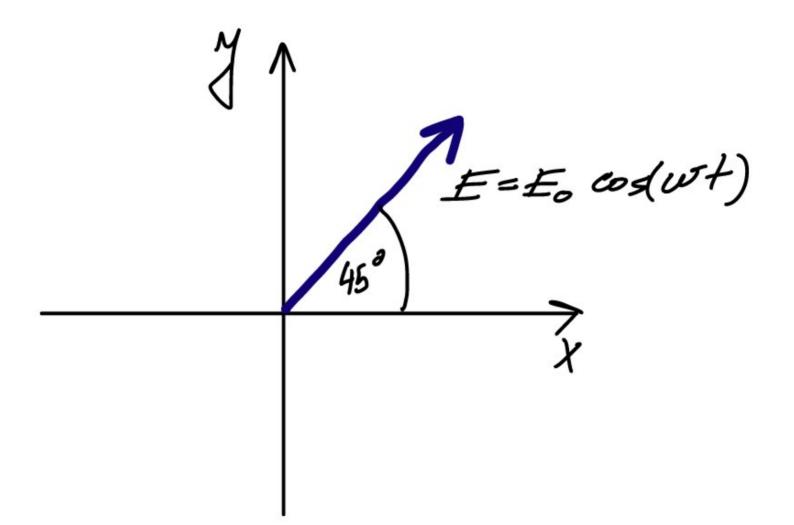
So let's have some questions! Is this positive Q or negative Q?



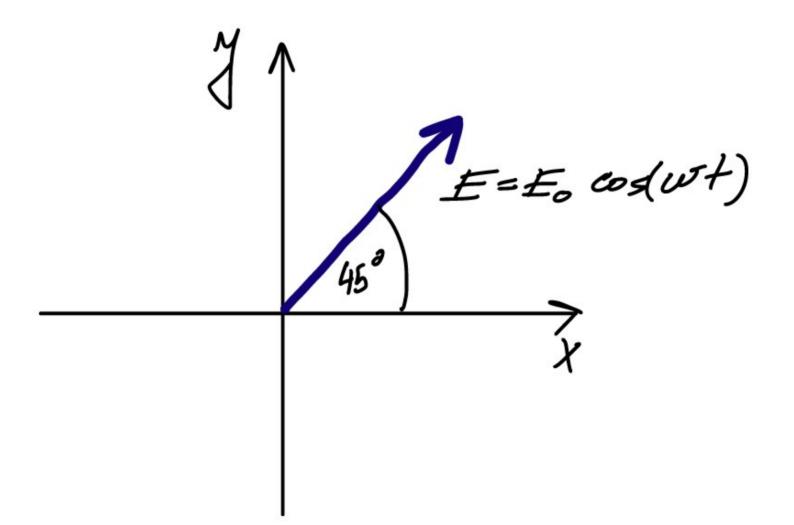
How about this?



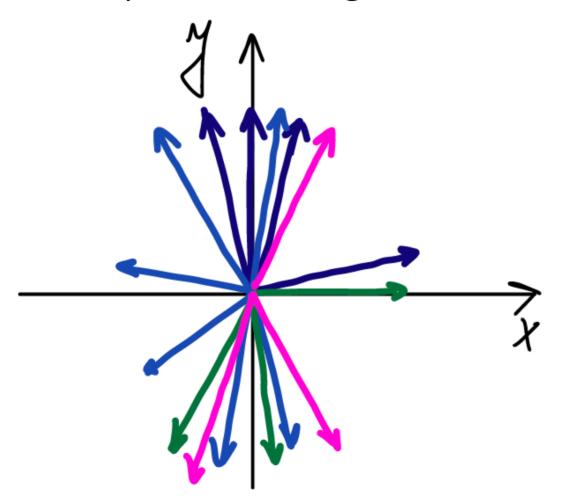
And this?



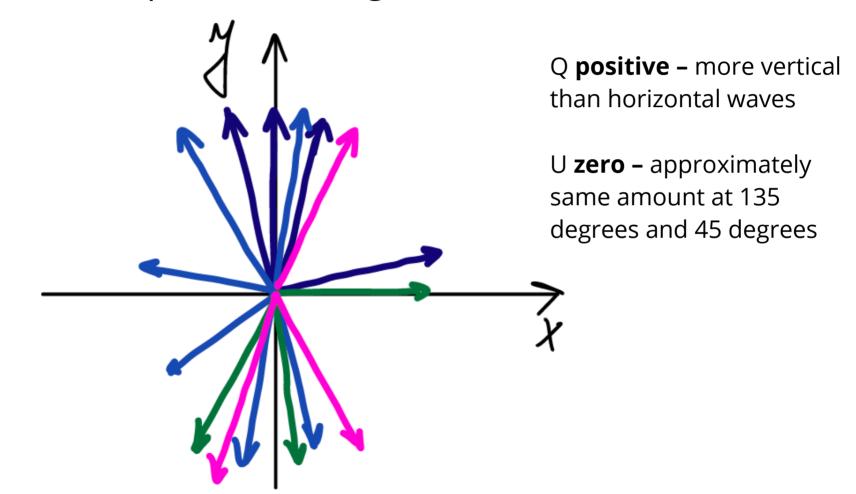
And this? That's right, this is zero Q but negative U



Here we have an incoherent superposition of many wavepackets. Is Q positive or negative? How about U?



Here we have an incoherent superposition of many wavepackets. Is Q positive or negative? How about U?



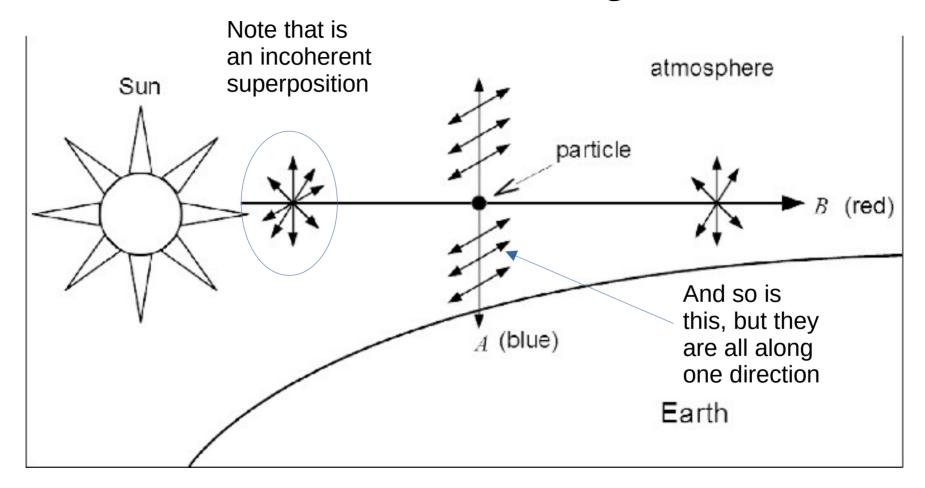
#### Reminders about Stokes formalism:

- Stokes formalism describes intensities
- I is total intensity of the light (counting photons not caring about their orientation)
- Other 3 are differences between number of photons with specific polarizations
- Note that if light was really monochromatic plane wave, sum of of squares of Q,U,V would be equal to square of I.
- It is not (again, the light we see is an incoherent superposition of many short wavepackets)

## Classical coherent (continuum) scattering



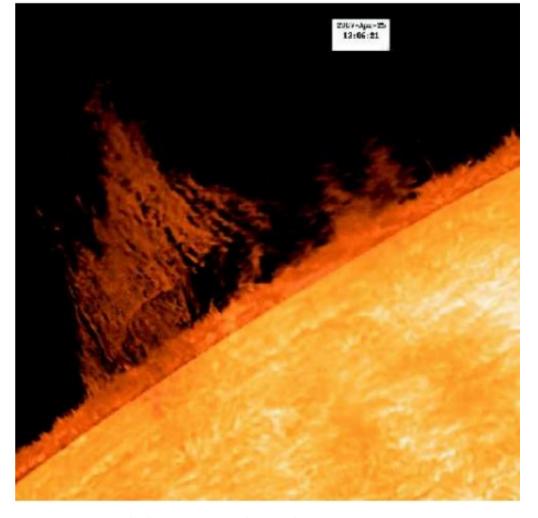
## Classical coherent (continuum) scattering



Credits: Soga et al. 2018

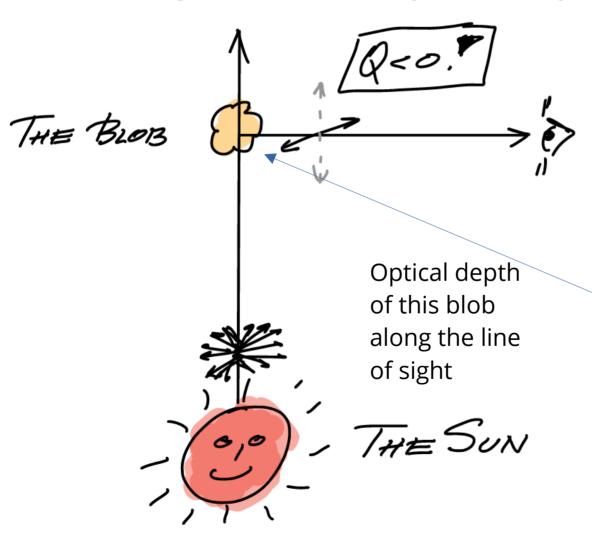
## Let's go back to prominences

- We agreed that prominences scatter the light and that is why we see them.
- They are more dense than the surrounding corona so we see them even out of the eclipse
- Easy way to model brightness of this prominence would be to split it in little 'blobs' and say that each blob scatters light



HINODE NFI observations of a prominence in H alpha, Henzel et al. 2008.

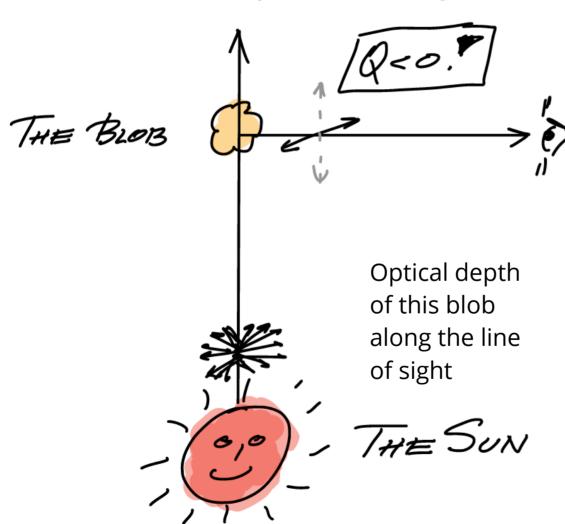
## Let's imagine a Blob very far away from the Sun



- In this (extreme) case, all the light is linearly polarized "parallel to the limb"
- Can we also agree thatI = |Q| in this case?
- And, separately from that, in the optically thin limit:

$$I_{\lambda} = S \tau_{\lambda} = S \tau \phi_{\lambda}$$

## Now let's complicate things a little:



- Since I wrote RTE (simple one but still a RTE), for stokes I, can we all agree to write one for Q too?
- After all Q is some sort of intensity too!

$$I_{\lambda} = S_I au_{\lambda}$$
 $Q_{\lambda} = S_Q au_{\lambda}$ 

## RTE for scattering polarization

• In principle, I could have went in a bit more detail:

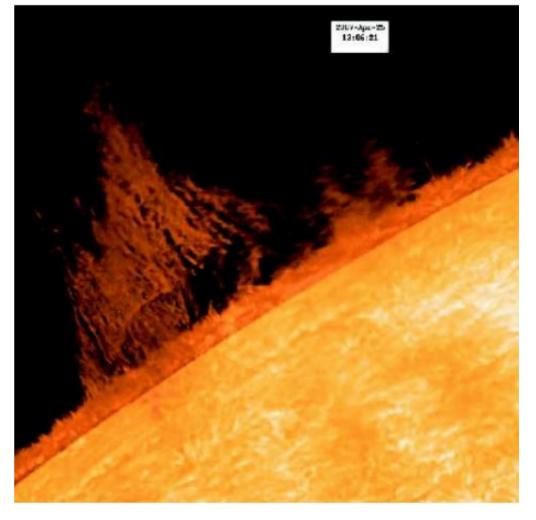
$$\frac{dI_{\lambda}}{d\tau_{\lambda}} = I_{\lambda} - S_{I}$$

$$\frac{dQ_{\lambda}}{d\tau_{\lambda}} = Q_{\lambda} - S_{Q}$$

- Two radiative transfer equations, for two components, each of these has it's own Source function.
- This implies that both Stokes I and Q "experience" the same absorption coefficient
- What would be source function for S in general?

### Prominence again

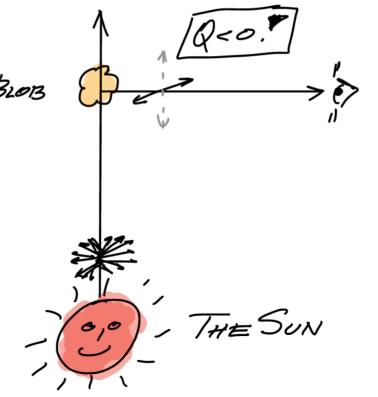
- These prominences are not suspended super super high above the Sun
- That is, they are not illuminated only from one direction
- They are illuminated by a cone of radiation, as we saw yesterday
- How do we proceed now?

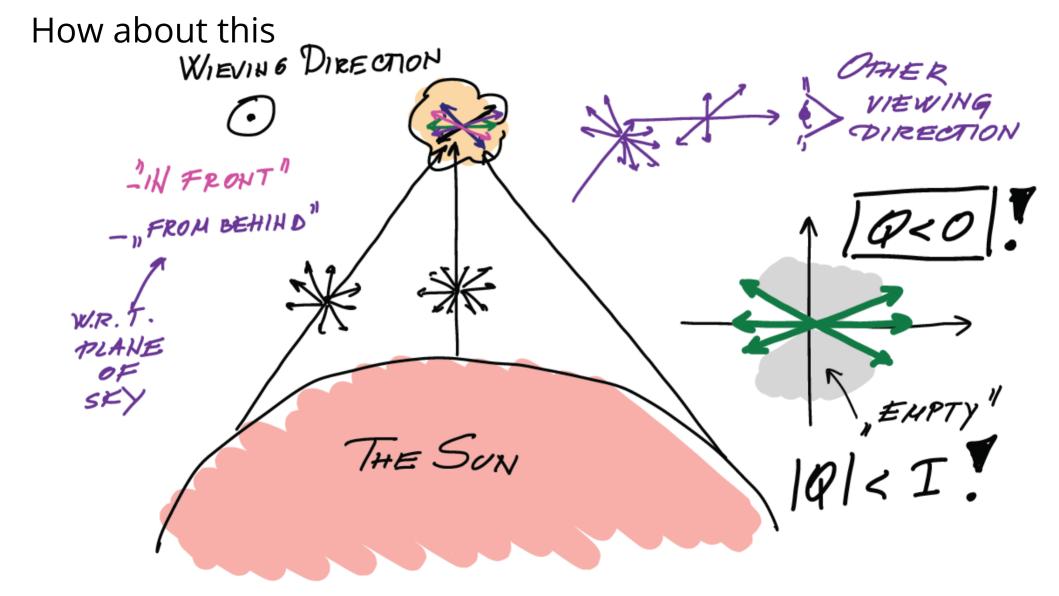


HINODE NFI observations of a prominence in H alpha, Henzel et al. 2008.

Now, take pen, paper and 5 mins and make a sketch similar to this one but for when the prominence is slose to the surface. What is the sign of Q? Is Q still equal to I?

(Hint: You might want to change the "orientation" to make it more obvious. i.e. make the plane of the sky he the plane of you notebook)





It turns out that (continuum, so no wavelength for now)

$$S_I \approx \frac{1}{2} \int_{-1}^{1} I(\mu') d\mu'$$
 $S_Q \approx \frac{3(\mu^2 - 1)}{16} \int_{-1}^{1} I(\mu') (3\mu'^2 - 1) d\mu'$ 

- These follow from Rayleigh's scattering. Note that the Source function for Q is direction-dependent!
- Can someone tell me difference between mu and mu'?
- These directions are with respect to the blob
- Let's use the equation to answer some more questions!

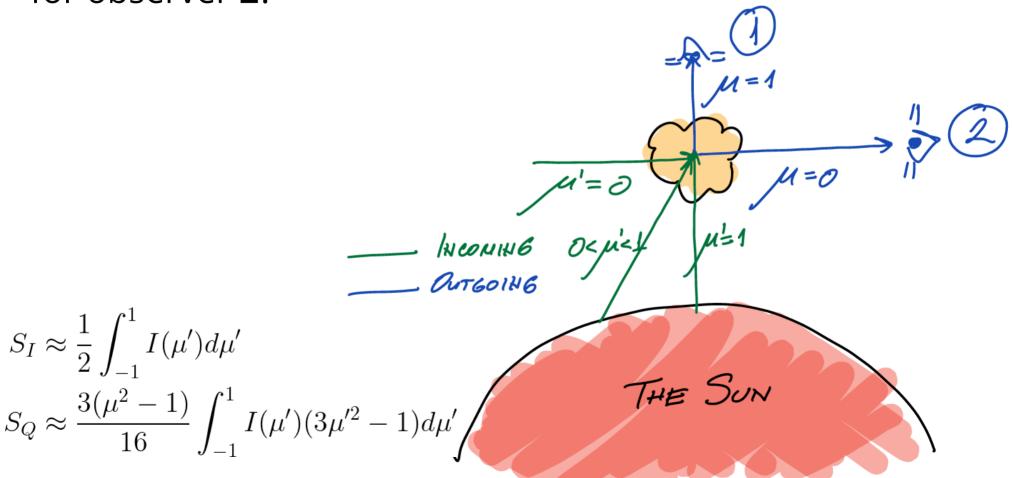
## What is Source function for Q for isotropic radiation?

• 3 mins to try and calculate this one

$$S_I \approx \frac{1}{2} \int_{-1}^{1} I(\mu') d\mu'$$
 $S_Q \approx \frac{3(\mu^2 - 1)}{16} \int_{-1}^{1} I(\mu') (3\mu'^2 - 1) d\mu'$ 

- We very often refer to  $S_{\mathbb{Q}}$  (up to a constant) as a measure of anisotropy of the radiation
- Scattering polarization is a consequence of scattering of light by anisotropically illuminated particles (lot of bold text here)

What is polarization for observer **1?** What sign is polarization for observer **2**?



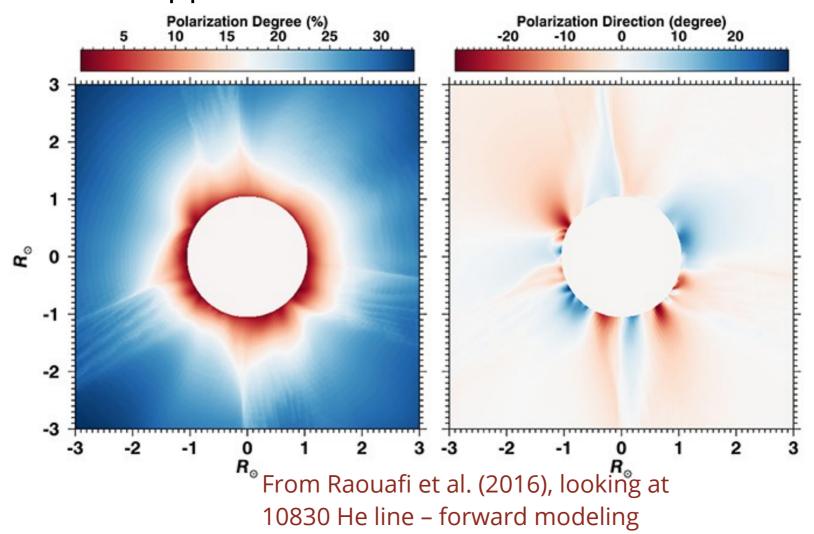
What is polarization for observer **1?** What sign is polarization for observer **2**?

$$S_I \approx \frac{1}{2} \int_{-1}^1 I(\mu') d\mu'$$
 
$$S_Q \approx \frac{3(\mu^2 - 1)}{16} \int_{-1}^1 I(\mu') (3\mu'^2 - 1) d\mu'$$
 
$$\mu' = 0$$
 In the proof of the

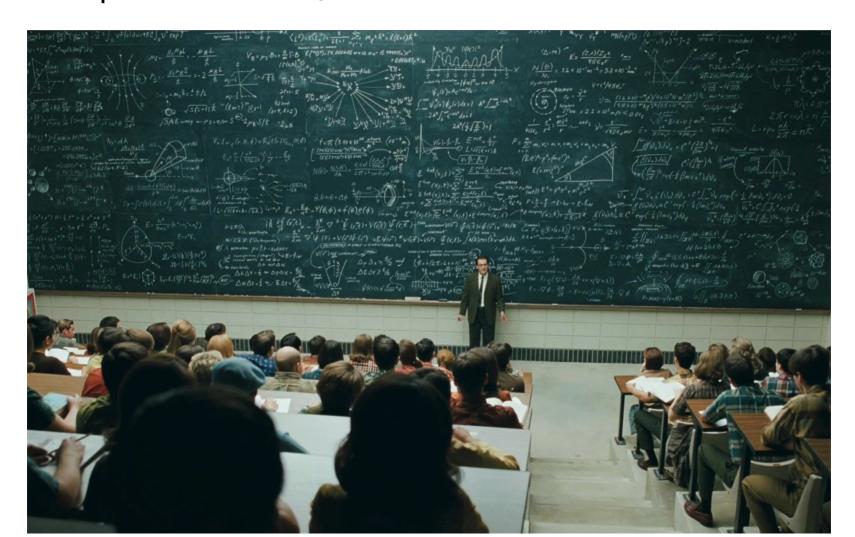
## Summary of scattering polarization

- It appears when, we have, well, scattering
- It is a measure of anisotropy. More anisotropically illuminated the scattering material is , more polarization
- But it is also a function of geometry: We have to be in a convenient viewing direction (no polarization at the disk center)
- There is more to this when you take into account stuff is 3D, but let's leave that to some other workshop (HAO workshop 2022 !?)

### All this is also applicable to corona



#### Now, the spectral lines! (Serious man, Cohen & Cohen, 2009)



## Now, the spectral line

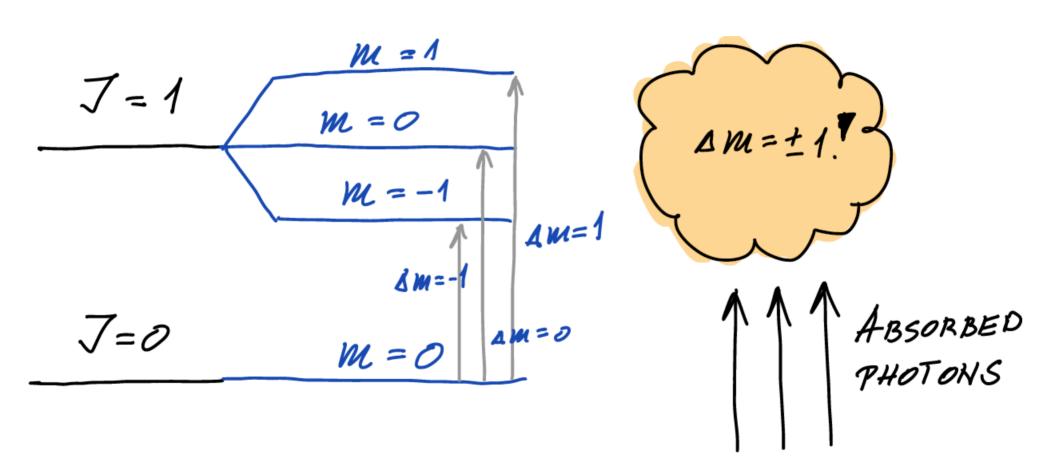
 $S_I \approx \frac{1}{2} \int_{-1}^{1} \int_{-\infty}^{\infty} I(\mu', \lambda) \phi_{\lambda} d\lambda d\mu'$ 

Additional integration over wavelength since the line redistributes the light

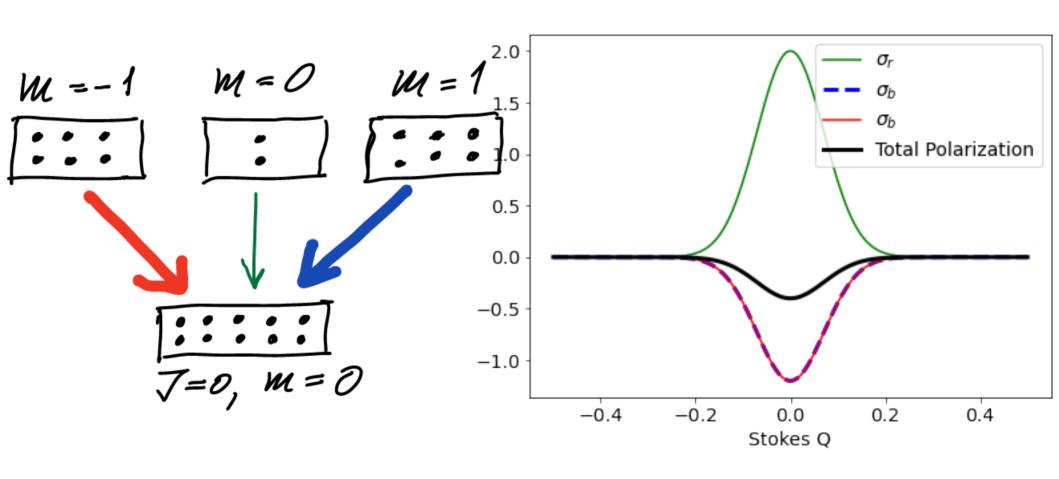
Intrinsic line polarizability.

 $\frac{3(\mu^{2}-1)}{16} \frac{1}{1+\delta/\Gamma_{R}} w_{0}^{2} \int_{-1}^{1} \int_{-\infty}^{\infty} I(\mu',\lambda) (3\mu'^{2}-1) \phi_{\lambda} d\lambda$ 

Why are spectral lines polarized due to scattering?

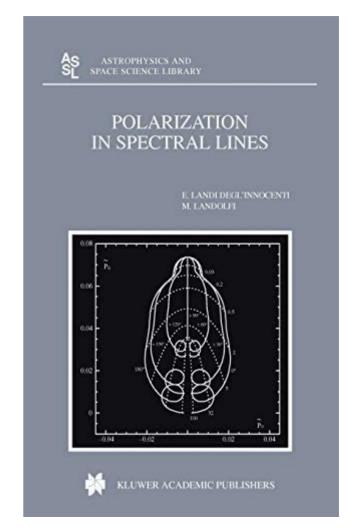


Why are spectral lines polarized due to scattering?



#### The ultimate text: Polarization in Spectral Lines

- A texbook that starts from first principles and derives spectral line polarization via density matrices and proper quantum mechanics
- Not for faint of heart
- Also includes magnetic fields and their effect (i.e. Hanle effect)
- Very little "radiative transfer" in the sense that the goal is to understand the atoms

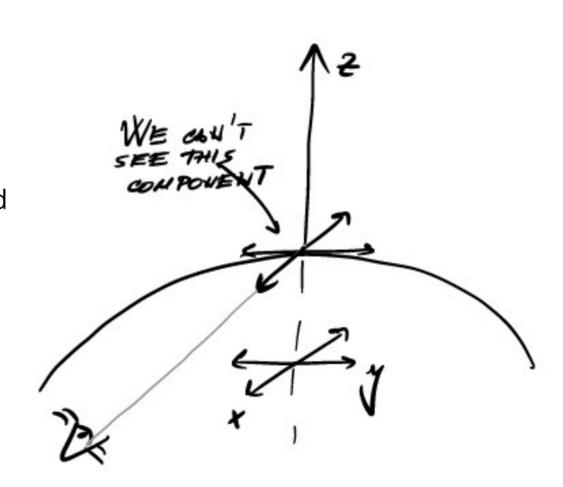


#### What is this "Hanle effect"?

- Hanle effect is a consequence of interaction between the magnetic fields and "polarized" Zeeman sublevels
- Instead of shifting them completely, magnetic field "shuffles" the Zeeman level coherences around / changes the referent axis
- Two important consequences:
- Decrease of polarization (depolarization)
- Rotation of polarization (Some Q goes into U)
- Another way to look at this is classically: The oscillator precesses around the magnetic field axis.

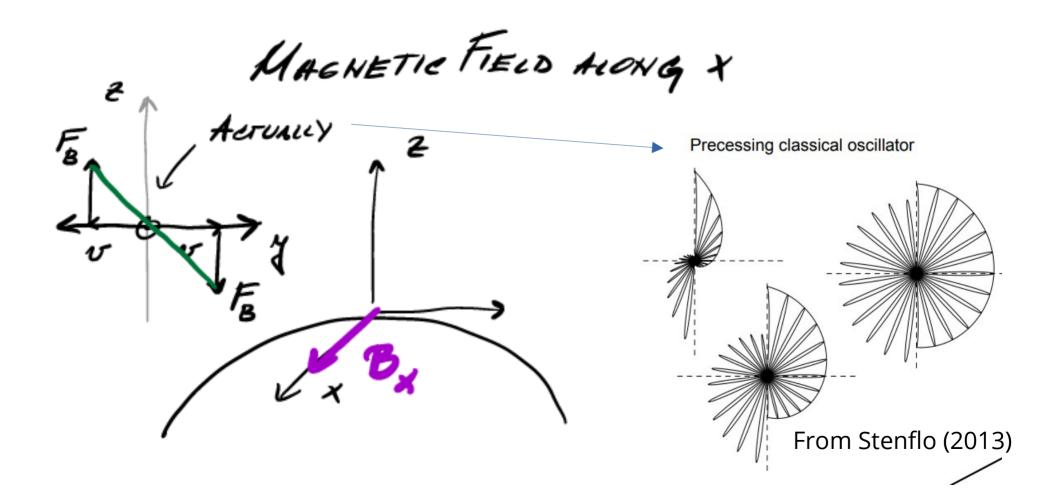
## So let's look at it classicaly:

- To simplify the situation, let's consider scattering at the limb, and the radiation is coming directly from below.
- In that case, E oscillates in x and y only (incoherently), and we only can see the y component because we are looking along y



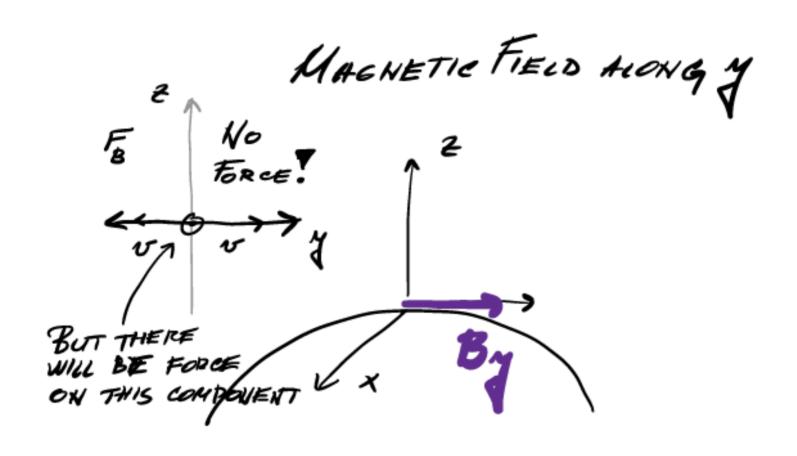
### Now let's add the magnetic field along x

following Trujillo Bueno 2001 ("Atomic polarization and the Hanle effect")



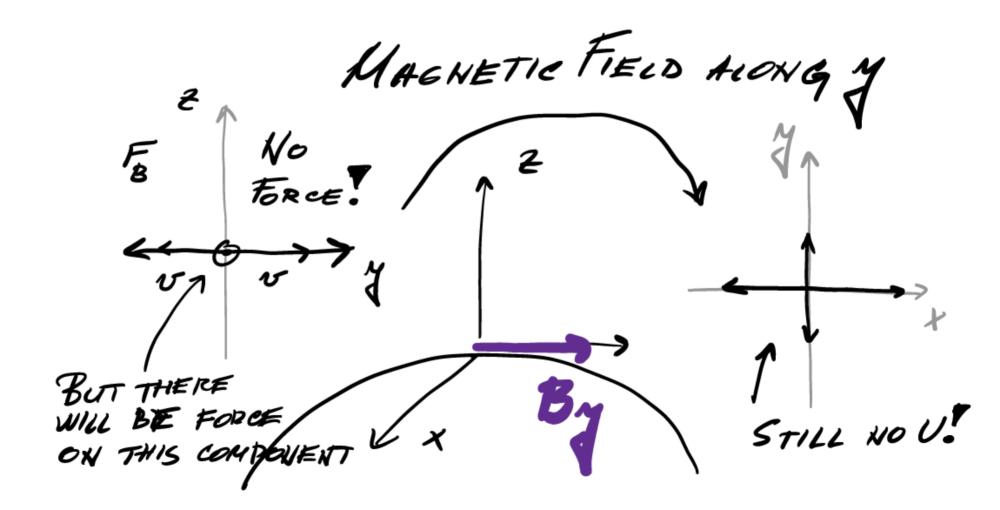
### Now let's add the magnetic field along y

following Trujillo Bueno 2001 ("Atomic polarization and the Hanle effect")



### Now let's add the magnetic field along y

following Trujillo Bueno 2001 ("Atomic polarization and the Hanle effect")



## Formal (but simplified) way to illustrate this would be:

Polarized RTE 
$$\frac{\hat{I}_{\lambda}}{d\tau} = \hat{I}_{\lambda} - \hat{S}_{\lambda} \qquad \text{and decreases polarized sources. Becomes important at:} \\ \hat{S}_{\lambda} = \hat{H}\hat{J}_{\lambda} \qquad \frac{0.88gB \times 10^7}{A_{ul}} \approx 1 \\ \hat{I} = (I,Q,U)^{\dagger} \qquad \text{Zero in the absence of horizontal anisotropy} \\ \hat{J} = (J_I,J_Q,J_U)^{\dagger} \qquad \text{Nean intensity} \qquad \text{Vertical anisotropy} \\ \text{Mean intensity} \qquad \text{Vertical anisotropy} \\ \end{pmatrix}$$

#### Three regimes of spectral line polarization

