

# PROBLEM SHEET 12

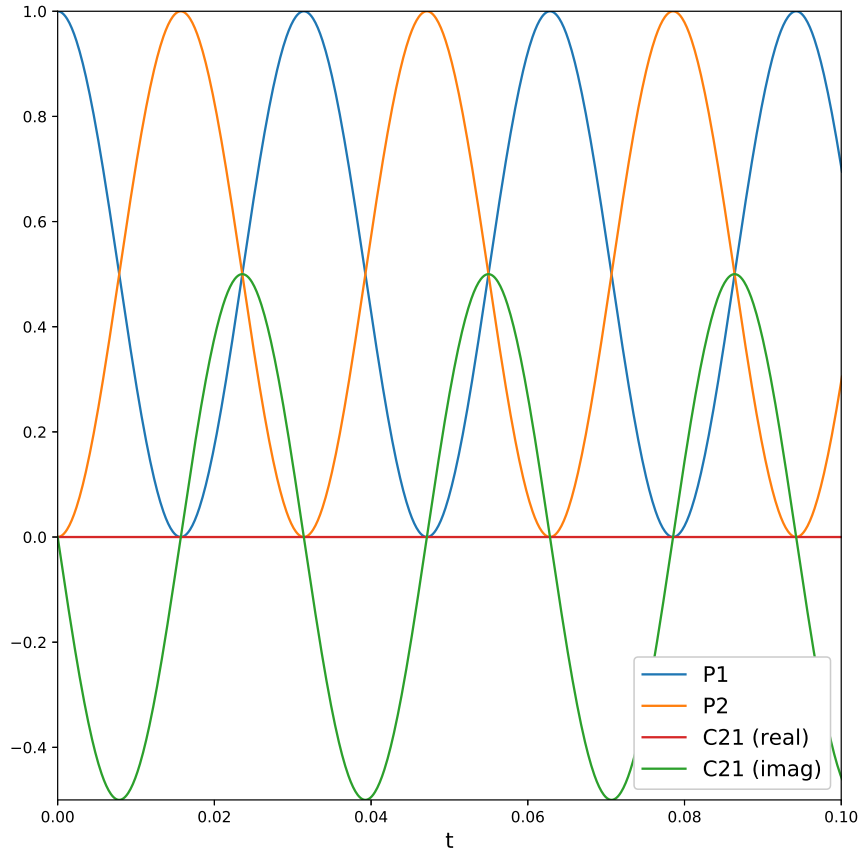
## Open Quantum Systems WS18-19

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Code available here: [https://github.com/iyanmv/OQS\\_ws18-19](https://github.com/iyanmv/OQS_ws18-19).

### Exercise 26

a) Simulation of a two-site closed system with energy difference  $\Delta\Omega = \Omega_1 - \Omega_2 = 0$  and electronic coupling  $J = 100$ . Perfect coherent energy transfer is observed since there is no interaction with the thermal bath. See Figure 1. Values are plotted from  $t = 0$  to  $t = 0.1$  so it is possible to observe the fast oscillations.

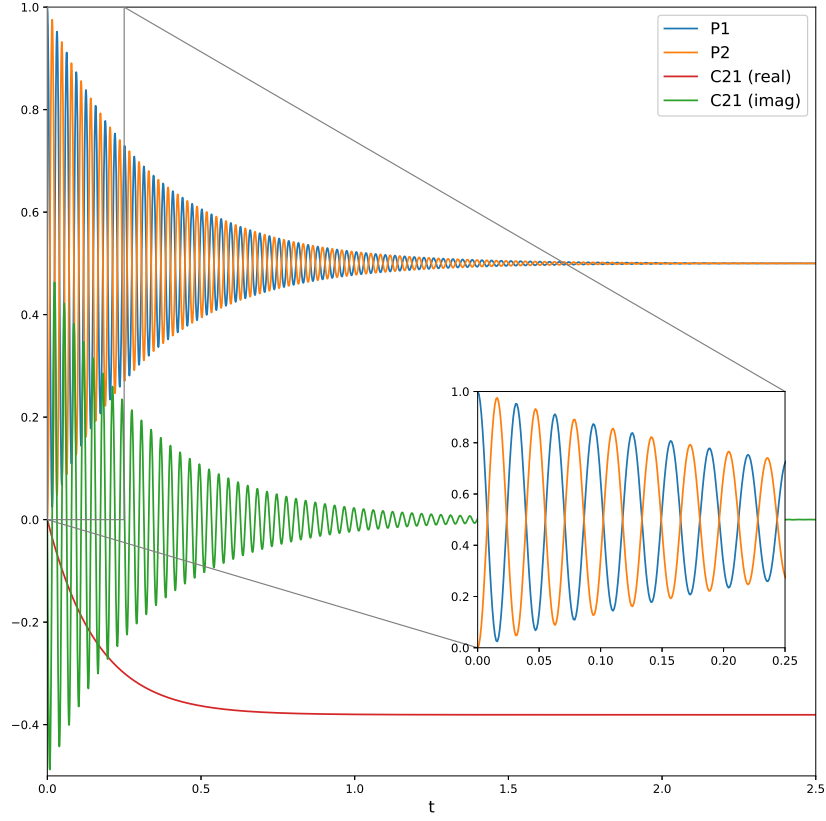


**Figure 1:** Obtained populations  $P_1(t) = \langle 1 | \rho_s(t) | 1 \rangle$  and  $P_2(t) = \langle 2 | \rho_s(t) | 2 \rangle$ , and inter-site coherence  $C_{21}(t) = \langle 2 | \rho_s(t) | 1 \rangle$ , using Runge-Kutta method from  $t = 0$  to  $t = 0.1$ .

b) Same parameters as before but now with each site coupled to an independent thermal bath modelled with quantum harmonic oscillators for temperature  $(k_B T / \hbar) = 100$ . An Ohmic spectral density

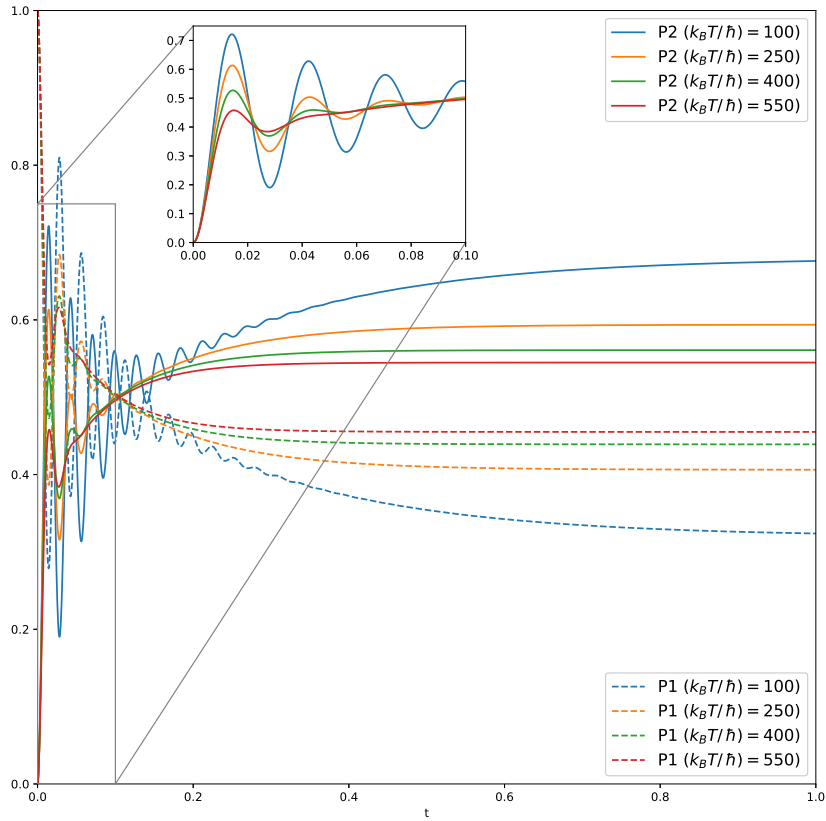
$$\mathcal{T}_k(\omega) = \frac{2\lambda_k}{\pi} \frac{\omega \gamma_k}{\omega^2 + \gamma_k^2} \quad (1)$$

was used with  $\lambda_k = 10$  and  $\gamma_k = 50$ . See Figure 2.



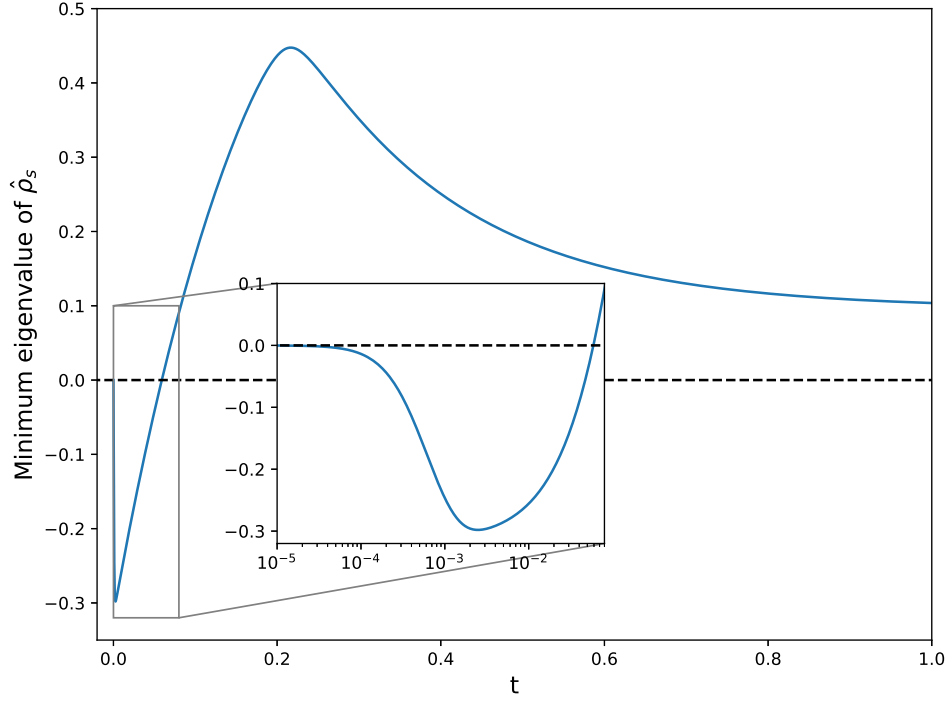
**Figure 2:** Obtained populations  $P_1(t) = \langle 1 | \rho_s(t) | 1 \rangle$  and  $P_2(t) = \langle 2 | \rho_s(t) | 2 \rangle$ , and inter-site coherence  $C_{21}(t) = \langle 2 | \rho_s(t) | 1 \rangle$ , using Runge-Kutta method from  $t = 0$  to  $t = 2.5$ . It can be seen that  $P_1(t) = P_2(t) = 0.5$  for long enough times because of  $\Omega_1 = \Omega_2$ .

c) Same parameters from b) but now with  $\Delta\Omega = \Omega_1 - \Omega_2 = 100$ . See Figure 3.



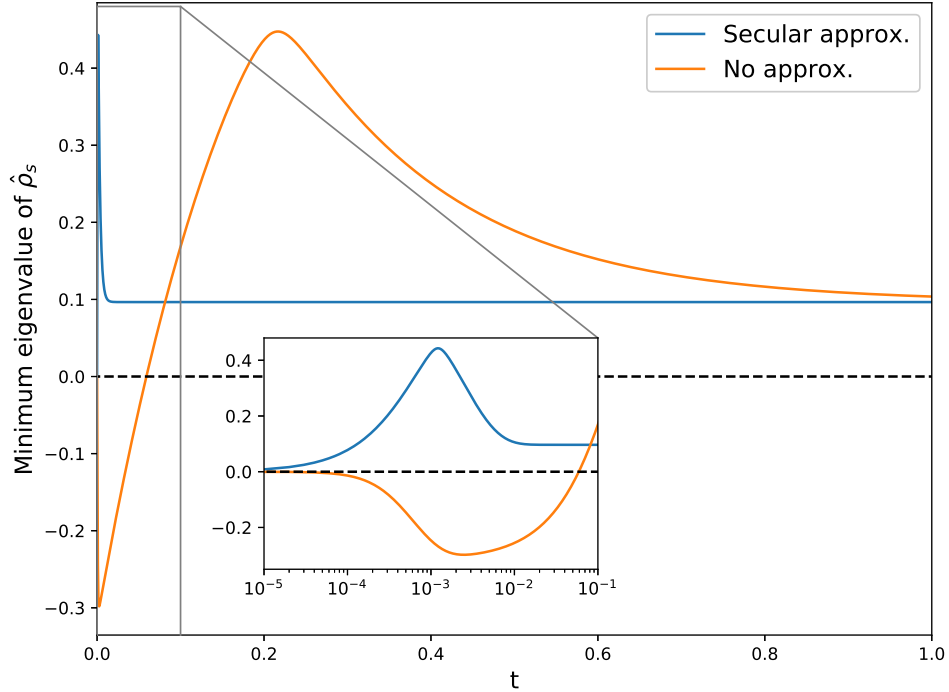
**Figure 3:** Populations  $P_1(t) = \langle 1 | \rho_s(t) | 1 \rangle$  and  $P_2(t) = \langle 2 | \rho_s(t) | 2 \rangle$  from  $t = 0$  to  $t = 1$ . When saturated,  $P_2^{(s)}$ , becomes higher than  $P_1^{(s)}$  because  $\Omega_1 > \Omega_2$ . This difference,  $P_2^{(s)} - P_1^{(s)}$  is reduced when temperature increases.

**d)** The time evolution of the negativity of  $\hat{\rho}_s(t)$  is simulated with the same parameters from **c)**, but now with a stronger environmental coupling  $\lambda_k = 1000$ . See Figure 4.

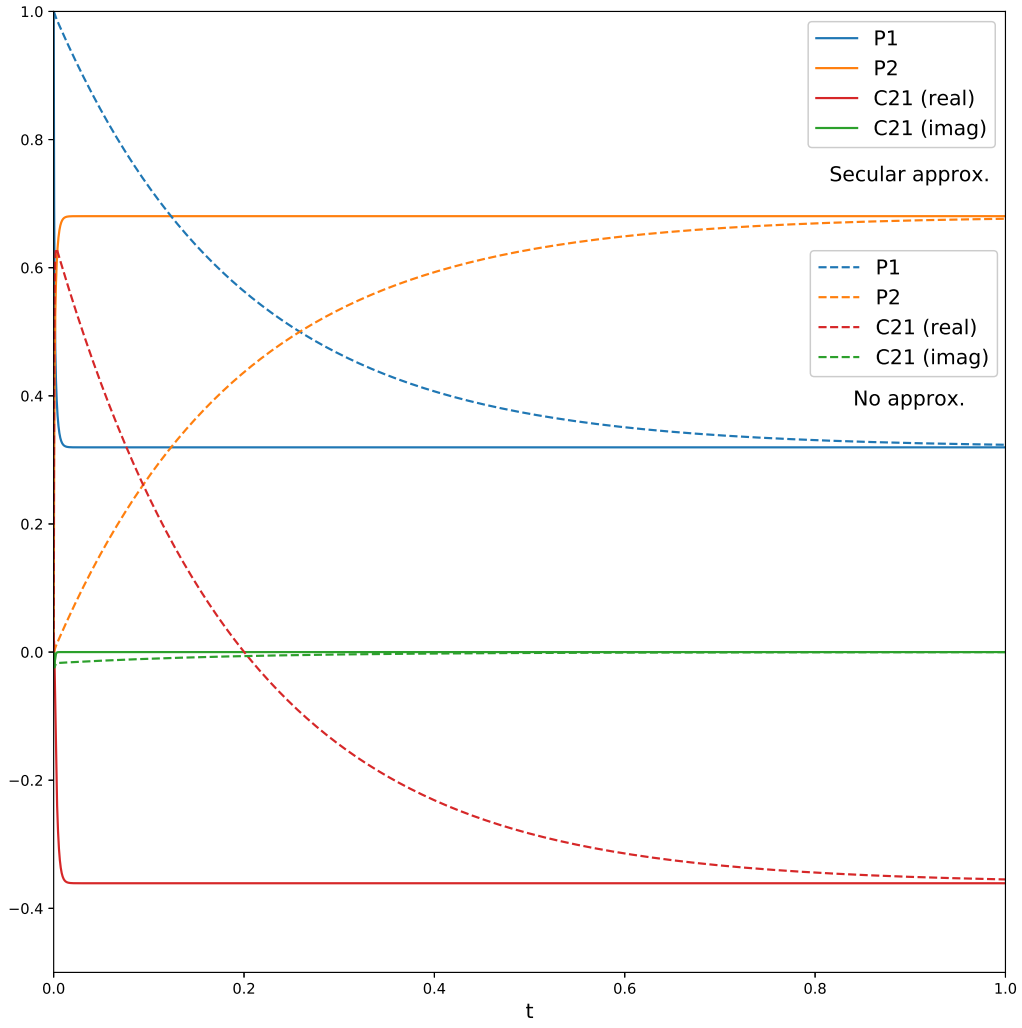


**Figure 4:** The minimum eigenvalue of  $\hat{\rho}_s$  is plotted from  $t = 0$  to  $t = 1$ . Negativity of  $\hat{\rho}_s(t)$  can be observed at first with a minimum of  $\approx -0.3$  which is reached at  $t \approx 2.5$  ms. After  $t = 0.1$ , the system  $\hat{\rho}_s$  is positive semidefinite again.

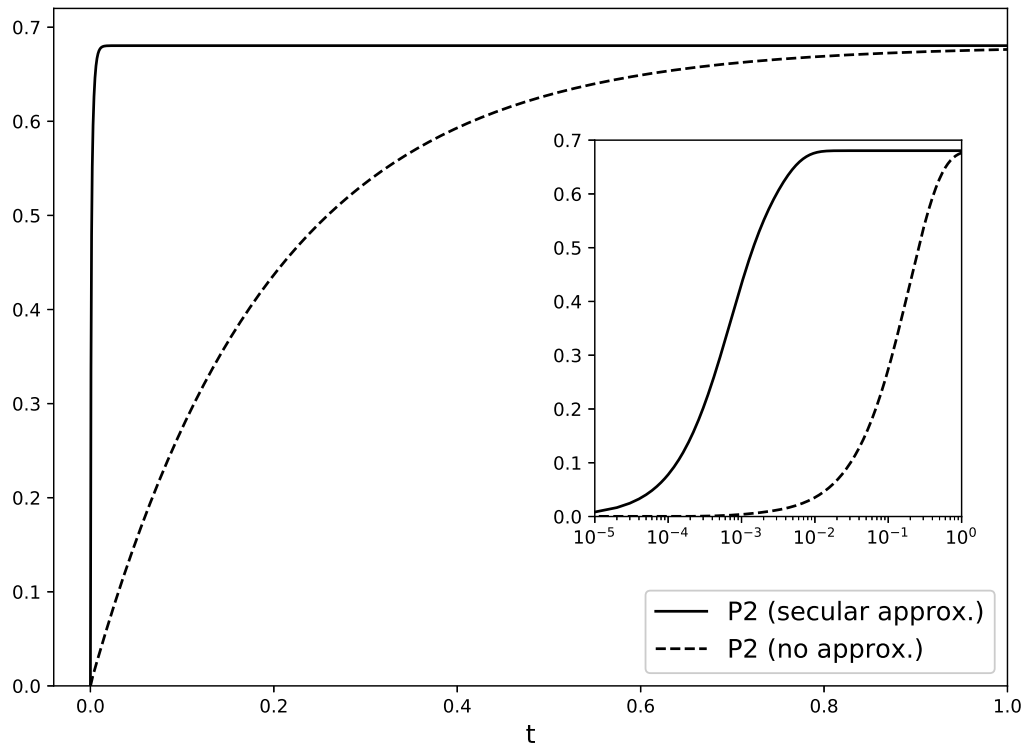
**e)** With the same parameters from **d)**, the secular approximation is used to avoid the negativity. See Figures 5, 6 and 7



**Figure 5:** Negativity of  $\hat{\rho}_s(t)$  from  $t = 0$  to  $t = 1$ . When using the secular approximation (blue) the negativity disappears.

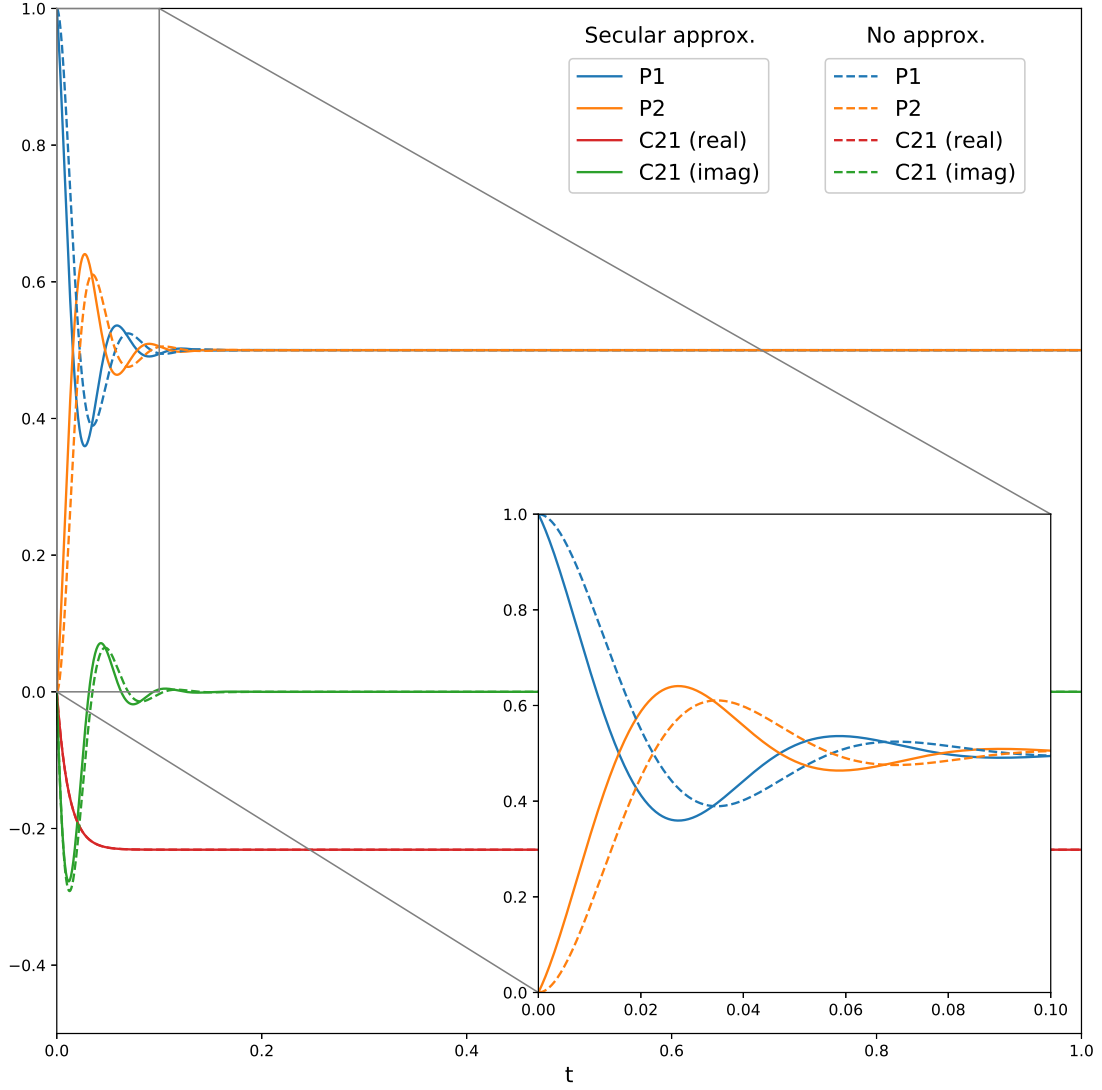


**Figure 6:** Time evolution of the populations and inter-site coherence using the secular approximation (solid lines) and without (dashed lines) from  $t = 0$  to  $t = 1$ .



**Figure 7:** Time evolution of just the population of site 2 from  $t = 0$  to  $t = 1$ .

f) For the last part, a moderate environmental coupling strength is considered ( $J \approx \lambda_k$ ). This does not cause the negativity effect as before, but induces non-secular effects modulating system dynamics. See Figure 8.



**Figure 8:** Time evolution of the site populations from  $t = 0$  to  $t = 1$ . Non-secular effects can be observed before the saturation time ( $t < 0.1$ ) in the zoomed area.