

Inverse problems for linear and non-linear elliptic equations

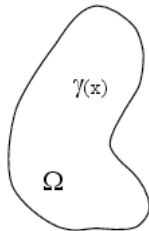
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Outline of talk

- 1 Introduction
- 2 Cloaking in electrostatics model
- 3 Polyharmonic first order perturbation

Calderón problem



Consider a body $\Omega \subset \mathbb{R}^N$, $N \geq 2$ with conductivity $\gamma(x)$ where $x \in \Omega$. An electrical potential $u(x)$ causes the current $I(x) = \gamma(x)\nabla u(x)$.

The conductivity $\gamma(x)$ can be isotropic (scalar valued), or anisotropic (matrix valued).

If the current has no sources or sinks, we have

$$-\operatorname{div}(\gamma(x)\nabla u(x)) = 0 \text{ in } \Omega.$$

Dirichlet-to-Neumann map

- Consider the following boundary value problem

$$\begin{aligned} -\operatorname{div}(\gamma(x)\nabla u(x)) &= 0 \text{ in } \Omega, \\ u &= f \text{ on } \partial\Omega. \end{aligned}$$

$\gamma(x)$ = conductivity,
 f = voltage potential at $\partial\Omega$.

- Current flux at $\partial\Omega = (\nu \cdot \gamma \nabla u)|_{\partial\Omega}$ where ν is the unit outer normal.

Information is encoded in map $\Lambda_\gamma(f) = \nu \cdot \gamma \nabla u|_{\partial\Omega}$

- Inverse problem

Does Λ_γ determine γ ?

Λ_γ = Dirichlet-to-Neumann map

- Different aspects of an inverse problem

- Uniqueness
- Stability
- Reconstruction techniques
- Numerical implementation
- Partial data problems

- Interior uniqueness: Does $\Lambda_{\gamma_1} = \Lambda_{\gamma_2}$ imply $\gamma_1 = \gamma_2$ in Ω ?

- Yes, if γ_i are scalar valued. First answered in 1987 for $N \geq 3$ by Sylvester and Uhlmann for C^2 conductivities.
- Mathematics of the problem for $N = 2$ is quite different.

Past work: Scalar case

- Uniqueness result was generalized for $N \geq 3$ by **Brown and Torres** in 1996 for $\gamma \in W^{\frac{3}{2},p}$ with $p > 2N$, **Haberman and Tataru** in 2013 for $\gamma \in C^1$ and **Caro and Rogers** in 2016 for $\gamma \in C^{0,1}$.
- For $N = 2$, the global uniqueness problem was first solved by **Nachman** in 1996 for $\gamma \in W^{2,p}$ for $p > 1$.
- This result was completely generalized by **Astala and Päivärinta** in 2005 for $\gamma \in L^\infty$.

Diffeomorphism: Anisotropic conductivities

- No for anisotropic conductivities. Choose a smooth diffeomorphism $\Phi : \Omega \rightarrow \Omega$ such that $\Phi(x) = x$ on $\partial\Omega$ and define

$$\Phi_*\gamma(y) = \frac{D\Phi(x)^T \gamma(x) D\Phi(x)}{|D\Phi|} \circ \Phi^{-1}(y),$$

where $\Phi(x) = y$.

- $\Phi^*\gamma(y)$ is the push-forward of the conductivity γ by Φ . Since $\Phi = Id$ on $\partial\Omega$, we get

$$\Lambda_\gamma = \Lambda_{\Phi_*\gamma}.$$

- Such a restricted uniqueness was first shown in 1989 by Lee and Uhlmann for $N \geq 3$ for real analytic conductivities, and in 2006 by Astala, Lassas and Päiväranta for $N = 2$ for $\gamma \in L^\infty$.
- Uniqueness up to diffeomorphism holds only when we have upper and lower bounds for γ .

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Three key ideas

There are 3 key ideas associated with such a *non-uniqueness* result:

- 1 Cloaking
- 2 Approximate cloaking
- 3 Isotropic approximate cloaking

Cloaking

- Let $E \subset \Omega$ be fixed and let σ_c be a non negative matrix valued function defined on $\Omega \setminus E$. We say σ_c *cloaks* E if any extension of σ_c across E of following form,

$$\sigma_A(x) = \begin{cases} A(x) & x \in E, \\ \sigma_c(x) & x \in \Omega \setminus E, \end{cases}$$

produces the same DN map as a uniform isotropic region, irrespective of the choice of $A(x)$.

- Existence of such cloaks first shown by **Greenleaf, Lassas and Uhlmann** in 2003.
- A change-of-variable scheme is employed which essentially blows up a point to a region in space and is highly singular.

Approximate cloaking

- Consider a regularized change-of-variable scheme which blows a small ball to the region being cloaked.
- One can look at the asymptotic behavior as radius of the ball goes to 0 and recover the singular transform.
- For simplicity, let $\Omega = B_2$ and restrict our attention the case when $E = B_1$ needs to be nearly cloaked.
- Fix a small parameter $r > 0$.

Regular change of variables

- Consider $F^r : B_2 \rightarrow B_2$ such that

$$F^r(x) = \begin{cases} \frac{x}{r} & |x| \leq r, \\ \left(\frac{2-2r}{2-r} + \frac{1}{2-r}|x|\right) \frac{x}{|x|} & r \leq |x| \leq 2. \end{cases}$$

- Consider

$$\sigma_A^r(x) = \begin{cases} A(x) & x \in B_1, \\ F_*^r 1 & x \in B_2 \setminus B_1. \end{cases}$$

- By approximate cloaking, we mean

$$|\langle \Lambda_{\sigma_A^r} f, g \rangle - \langle \Lambda_1 f, g \rangle| = o(1) \|f\|_{H^{\frac{1}{2}}(\partial B_2)} \|g\|_{H^{\frac{1}{2}}(\partial B_2)},$$

where the $o(1)$ term is independent of f and g .

Isotropic approximate cloaking

- The approximate cloaks using regular change of variables are anisotropic.
- In 2008, **Greenleaf, Kurylev, Lassas and Uhlmann** constructed isotropic and nonsingular parameters that give *approximate* cloaking to any desired degree of accuracy.
- They used the notion of H-convergence in linear settings to construct such isotropic cloaks.
- The H-limit of a sequence of isotropic cloaks need not be isotropic, and this key property allowed them to construct isotropic approximate cloaks.

Our work

- ① We extend previously known results in electrostatics on cloaking, approximate cloaking and isotropic approximate cloaking to a *quasi-linear* operator.
- ② This is achieved using the same techniques:
 - singular change of variables for cloaking,
 - regular change of variables for approximate cloaking &
 - H-convergence for isotropic approximate cloaking.

Basic Set up: quasi-linear equation

- Consider

$$\begin{aligned} -\operatorname{div}(A(x, u(x))\nabla u(x)) &= 0 \text{ in } \Omega, \\ u &= f \text{ on } \partial\Omega, \end{aligned}$$

where $A(x, t) \in \mathcal{M}(\alpha, \beta, L; B_2 \times \mathbb{R})$.

- The above equation arises in modeling of thermal conductivity of the Earth's crust and heat conduction in composite materials.
- More specifically, steady state heat conduction in an inhomogeneous anisotropic nonlinear medium with Dirichlet boundary conditions is governed by the above equation.

Quasi-linear approximate cloaking: [Ghosh, I '18]

- Consider

$$\sigma_A^r(x, t) = \begin{cases} A(x, t) & (x, t) \in B_1 \times \mathbb{R}, \\ F_*^r 1 & (x, t) \in B_2 \setminus B_1 \times \mathbb{R}, \end{cases}$$

where $A(x, t) \in \mathcal{M}(\alpha, \beta, L; B_2 \times \mathbb{R})$.

- We show

$$|\langle \Lambda_{(F^r)^{-1} \sigma_A^r} f, g \rangle - \langle \Lambda_1 f, g \rangle| = o(1) \|f\|_{H^{\frac{1}{2}}(\partial B_2)} \|g\|_{H^{\frac{1}{2}}(\partial B_2)},$$

where

$$(F^r)^{-1} \sigma_A^r = \begin{cases} (F^r)^{-1} A = \tilde{A}^r(x, t) & (x, t) \in B_r \times \mathbb{R}, \\ 1 & (x, t) \in (B_2 \setminus B_r) \times \mathbb{R}. \end{cases}$$

Sketch of proof

- If $A \in \mathcal{M}(\alpha, \beta, L; B_1 \times \mathbb{R})$, then $\tilde{A}^r(x, t) \in \mathcal{M}(\frac{\alpha}{r^{N-2}}, \frac{\beta}{r^{N-2}}, L; B_r \times \mathbb{R})$.
- This implies that for $r \ll 1$, $(F^r)^{-1}_* \sigma_A^r \in \mathcal{M}(1, \frac{\beta}{r^{N-2}}, L; B_r \times \mathbb{R})$.
- This gives us

$$|\langle \Lambda_{(F^r)^{-1}_* \sigma_A^r} f, g \rangle - \langle \Lambda_1 f, g \rangle| \leq C r^{-\frac{N}{2}+2} \|f\|_{H^{\frac{1}{2}}(\partial B_2)} \|g\|_{H^{\frac{1}{2}}(\partial B_2)}.$$

- Pass to the limit as $r \rightarrow 0$ for perfect cloaking.

Homogenization set up

- The approximate cloaks discussed earlier are anisotropic. We will now construct *isotropic* approximate cloaks.
- **Set up:** Let $A(x, y, t) = [a_{ij}(x, y, t)] \in \mathcal{M}(\alpha, \beta, L; \Omega \times Y \times \mathbb{R})$ be such that $y \mapsto a_{ij}(x, y, t)$ are $Y = [0, 1]^N$ -periodic functions for a.e. $(x, t) \in \Omega \times \mathbb{R}$.

Let

$$A^\epsilon(x, t) = \left[a_{ij}\left(x, \frac{x}{\epsilon}, t\right) \right], \quad (x, t) \in \Omega \times \mathbb{R}.$$

Quasi-linear H-convergence: [Ghosh, I '18]

Theorem

Let A^ϵ and A^* belong to $\mathcal{M}(\alpha, \beta, L; \Omega \times \mathbb{R})$. We say $A^\epsilon \xrightarrow{H} A^*$, if up to a subsequence, the corresponding solutions $\{u^\epsilon\}$ to

$$\begin{aligned} -\operatorname{div}(A^\epsilon(x, u^\epsilon) \nabla u^\epsilon(x)) &= 0 \text{ in } \Omega, \\ u^\epsilon &= f \in H^{\frac{1}{2}}(\partial\Omega), \end{aligned}$$

are such that

$$\begin{aligned} u^\epsilon &\rightharpoonup u \text{ weakly in } H^1(\Omega) \text{ and} \\ A^\epsilon(x, u^\epsilon) \nabla u^\epsilon &\rightharpoonup A^*(x, u) \nabla u \text{ weakly in } L^2(\Omega)^N, \end{aligned}$$

where $u \in H^1(\Omega)$ uniquely solves

$$\begin{aligned} -\operatorname{div}(A^*(x, u(x)) \nabla u(x)) &= 0 \text{ in } \Omega, \\ u &= f \text{ on } \partial\Omega. \end{aligned}$$

H-limit

The homogenized conductivity $A^*(x, t)$

$$A^*(x, t) = [a_{ij}^*(x, t)] \in \mathcal{M}(\tilde{\alpha}, \tilde{\beta}, \tilde{L}; \Omega \times \mathbb{R}),$$

can be defined by its entries as

$$a_{kl}^*(x, t) = \int_Y a_{ij}(x, y, t) \frac{\partial}{\partial y_i} (\chi_k(x, y, t) + y_k) \frac{\partial}{\partial y_j} (\chi_l(x, y, t) + y_l) dy,$$

where for each canonical basis vector $e_k \in \mathbb{R}^N$, $\chi_k(x, y, t)$ satisfy the following cell-problem for almost every $(x, t) \in \Omega \times \mathbb{R}$:

$$\begin{aligned} -\operatorname{div}_y A(x, y, t)(\nabla_y \chi_k(x, y, t) + e_k) &= 0 \quad \text{in } \mathbb{R}^N, \\ y \rightarrow \chi_k(x, y, t) &\text{ is } Y\text{-periodic for all } (x, t) \in \Omega \times \mathbb{R}. \end{aligned}$$

Isotropic approximate cloaking

- Our isotropic cloaks take the form

$$A^\epsilon(x, t) = \sigma\left(x, \frac{|x|}{\epsilon}, t\right) I_{N \times N}, \quad (x, t) \in \Omega \times \mathbb{R}.$$

- Temporarily fix $R > 1$ and introduce a new parameter $\eta > 0$.
- More specifically,

$$\sigma_n(x, t) := \left[1 + a_{R_n, \eta_n}^1(x, t) \zeta_1\left(\frac{r}{\varepsilon_n}\right) - a_{R_n, \eta_n}^2(x, t) \zeta_2\left(\frac{r}{\varepsilon_n}\right) \right]^2, \quad r = |x|,$$

where $a_1, a_2, \zeta_1, \zeta_2$ are chosen to satisfy Lipschitz condition in t variable and periodicity in r variable.

Concluding steps

- We first let $\epsilon_n \rightarrow 0$ (approximate isotropic \rightarrow approximate anisotropic).
- Then let $\eta_n \rightarrow 0$ and finally $R_n \searrow 1$ (approximate anisotropic \rightarrow cloaking).
- We get the desired strong convergence of the DN maps,

$$\|(\Lambda_{\sigma_n} - \Lambda_1)(f)\|_{H^{\frac{1}{2}}(\partial\Omega)} \rightarrow 0 \text{ as } n \rightarrow \infty.$$

- This finishes the proof.

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Preliminaries

- Let $\Omega \subset \mathbb{R}^N$, $N \geq 3$ be a bounded open set with C^∞ boundary.
- Consider the operator $\mathcal{H}_{A,q} = -\Delta + A \cdot D + q$, where A is a vector-valued potential and q is a scalar valued potential.
- If 0 is not a Dirichlet eigenvalue of $\mathcal{H}_{A,q}$, then we can define the DN map $\Lambda_{A,q}^{\mathcal{H}}$.
- **Research question** Are A and q uniquely determined by $\Lambda_{A,q}^{\mathcal{H}}$?
- No. When $A \neq 0$, there is a natural obstruction to uniqueness.

Counter example

- Consider the magnetic Schrödinger operator

$$\mathcal{M}_{A,q} = \sum_{j=1}^N \left(-i \frac{\partial}{\partial x_j} + A_j(x) \right)^2 + q(x).$$

- Let $A' = A + \nabla g$, where $g = \frac{\partial g}{\partial N} = 0$ on $\partial\Omega$.
- $\Lambda_{A',q}^{\mathcal{M}} = \Lambda_{A,q}^{\mathcal{M}}$
- Research question:** Is this the only obstruction to uniqueness?
- Yes. In 1993, **Sun** proved such a restricted uniqueness result for $A \in W^{2,\infty}(\mathbb{R}^N) \cap \mathcal{E}'(\bar{\Omega})$ and $q \in L^\infty(\Omega)$.

Polyharmonic operator

- Surprisingly, for the polyharmonic operator

$$\mathcal{L}_{A,q} = (-\Delta)^m + A \cdot D + q, \quad m \geq 2,$$

one can uniquely recover both A and q .

- This result was first proved by **Krupchyk, Lassas and Uhlmann** in 2014, for $A \in W^{1,\infty}(\mathbb{R}^N) \cap \mathcal{E}'(\bar{\Omega})$ and $q \in L^\infty(\Omega)$.
- Higher order polyharmonic operators occur in areas of physics and geometry such as
 - Kirchoff plate equation in the theory of elasticity.
 - Paneitz-Branson operator in conformal geometry.

Definition of A and q

- Let first order perturbation A be in $W^{-\frac{m}{2}+1,p'}(\mathbb{R}^N) \cap \mathcal{E}'(\bar{\Omega})$, where

$$\begin{cases} p' \in [2N/m, \infty) & \text{if } m < N, \\ p' \in (2, \infty) & \text{if } m = N \text{ or } m = N + 2, \\ p' \in [2, \infty) & \text{otherwise.} \end{cases} \quad (\text{A})$$

- For a fixed δ with $0 < \delta < \frac{1}{2}$, let the zeroth order perturbation q be in $W^{-\frac{m}{2}+\delta,r'}(\mathbb{R}^N) \cap \mathcal{E}'(\bar{\Omega})$, where

$$\begin{cases} r' \in [2N/(m - 2\delta), \infty), & \text{if } m < N, \\ r' \in [2N/(m - 2\delta), \infty), & \text{if } m = N, \\ r' \in [2, \infty), & \text{if } m \geq N + 1. \end{cases} \quad (\text{q})$$

Main Result

Theorem (Assylbekov-I '17)

Let $\Omega \subset \mathbb{R}^N$, $N \geq 3$ be a bounded open set with C^∞ boundary, and let $m \geq 2$ be an integer. Let $0 < \delta < 1/2$. Suppose that A_1, A_2 satisfy (A) and q_1, q_2 satisfy (q) and 0 is not in the spectrum of \mathcal{L}_{A_1, q_1} and \mathcal{L}_{A_2, q_2} .

If $\Lambda_{A_1, q_1}^{\mathcal{L}} = \Lambda_{A_2, q_2}^{\mathcal{L}}$, then $A_1 = A_2$ and $q_1 = q_2$.

Key steps

We use the following two key steps in the proof:

- Carleman estimates
- Construction of Complex Geometric Optics (CGO) solutions to $\mathcal{L}_{A,q}$.

Carleman estimates

- **Carleman estimates** : For $0 < h \ll 1$, we have

$$\|u\|_{H_{scl}^{m/2}(\mathbb{R}^N)} \lesssim \frac{1}{h^m} \|e^{\phi/h} (h^{2m} \mathcal{L}_{A,q}) e^{-\phi/h} u\|_{H_{scl}^{-3m/2}(\mathbb{R}^N)}$$

for all $u \in C_0^\infty(\Omega)$.

- Here, ϕ is the so-called Limiting Carleman Weight for $-h^2\Delta$ and $H_{scl}^s(\mathbb{R}^N)$ are weighted Sobolev spaces.
- The two key ingredients in the proof of the Carleman estimates are:
 - Estimates for the Laplacian due to **Salo and Tzou** with a gain of two derivatives and
 - Continuity of multiplication between two Sobolev spaces.

Proposition

Let $\zeta \in \mathbb{C}^N$ and $h > 0$ be such that $\zeta \cdot \zeta = 0$, $\zeta = \zeta_0 + \zeta_1$ with ζ_0 independent of h and $\zeta_1 = \mathcal{O}(h)$ as $h \rightarrow 0$. For all $h > 0$ small enough, there exists $u(x, \zeta; h) \in H^{m/2}(\Omega)$ solving $\mathcal{L}_{A,q}u = 0$, and of the form

$$u(x, \zeta; h) = e^{\frac{ix \cdot \zeta}{h}} (a(x, \zeta_0) + h^{m/2} r(x, \zeta; h)),$$

where $a(\cdot, \zeta_0) \in C^\infty(\overline{\Omega})$ satisfies

$$(\zeta_0 \cdot \nabla)^2 a = 0 \quad \text{in } \Omega,$$

and the correction term r is such that $\|r\|_{H_{scl}^{m/2}(\Omega)} = \mathcal{O}(1)$ as $h \rightarrow 0$.

Concluding steps

- Construct $u_2(\cdot, \zeta_2; h)$ and $v(\cdot, \zeta_1; h)$ in $H^{\frac{m}{2}}(\Omega)$ solving $\mathcal{L}_{A_2, q_2} u_2 = 0$ and $\mathcal{L}_{A_1, q_1}^* v = 0$ in Ω and plug them in to the following integral identity,

$$\int_{\Omega} ((A_2 - A_1) \cdot Du_2) \bar{v} \, dx + \int_{\Omega} (q_2 - q_1) u_2 \bar{v} \, dx = 0.$$

- First let $h \rightarrow 0$ and obtain $A_1 = A_2$.
- Now plug in $A_1 = A_2$ and $a_1 = a_2 = 1$ into the integral identity, and let $h \rightarrow 0$ to get $q_1 = q_2$, which finishes the proof.
- Remark:** Obtain conditions on A and q by working backwards from construction of CGO solutions.

Summary

- We looked at the injectivity of the DN map for two different PDEs, one quasi-linear and one linear.
- For the quasi-linear PDE, the DN map is not injective. Uniqueness up to diffeomorphism fails in certain cases and this gives a recipe for *cloaking*.
- We used the same ideas as in the linear settings, to give a scheme for cloaking, approximate cloaking and isotropic approximate cloaking in quasi-linear settings.
- For the linear poly-harmonic PDE, the DN map is injective and goal of our work was to prove injectivity of the DN map for rough A and q .
- Our main contribution is Carleman estimates and construction of CGO solutions with proper decay of remainder term.