

B23

a) Grand Canonical ensemble:

$$Q(r) = \sum_{N=0}^{\infty} e^{\beta \mu N} \cdot Z_N(r)$$

$$Z_N(r) \stackrel{(\text{spin})}{=} \frac{2^N}{N!} \left(\int d^3p e^{-\beta \frac{p^2}{2m}} \int dA \cdot e^{+\frac{\beta G_{MM}}{r} \Delta h} \right)^N$$

define $V = A \cdot \Delta h$

$$= \frac{2^N}{N!} \left(\frac{V}{\lambda^3} e^{\frac{\beta G_{MM}}{r}} \right)^N$$

$$Q(r) = \sum_{N=0}^{\infty} \frac{1}{N!} \left(\frac{2V}{\lambda^3} e^{\beta(\mu + \frac{G_{MM}}{r})} \right)^N$$

$$Q(r) = \exp \left(\frac{2V}{\lambda^3} e^{\beta(\mu + \frac{G_{MM}}{r})} \right)$$

$$\Omega(r) = -kT \ln Q(r) = -2kT \frac{V}{\lambda^3} e^{\beta(\mu + \frac{G_{MM}}{r})}$$

$$N(r) = -\frac{\partial \Omega}{\partial \mu} = \frac{2V}{\lambda^3} e^{\beta \mu + \frac{\beta G_{MM}}{r}}$$

$$n(r) = \frac{2}{\lambda^3} e^{\beta \mu + \frac{\beta G_{MM}}{r}}$$

B23 Neutron star, $U = -\frac{GMm}{r}$

a) $n(r) = \frac{2}{V} \sum_p e^{-\beta(\frac{p^2}{2m} - \mu - \frac{mMG}{r})}$
 $= \frac{2}{\lambda^3} \cdot e^{\beta\mu + \beta \frac{mMG}{r}}$

$n(r_0) = \frac{2}{\lambda^3} \cdot e^{\beta\mu + \beta \frac{mMG}{r_0}}$ $\mu(r) = \mu = \text{const}$

$n(r) = n(r_0) \cdot e^{\beta GMm(\frac{1}{r} - \frac{1}{r_0})}$

$n(r \rightarrow \infty) > 0$ ie no confinement.

b) $n(r) = \frac{2}{V} \sum_p \frac{1}{e^{\beta(\frac{p^2}{2m} - \mu - \frac{GMm}{r})} + 1}$
 $= \frac{2}{\lambda^3} f_{3/2} \left(e^{\beta\mu + \beta \frac{GMm}{r}} \right)$

at $T \rightarrow 0$ $f_{3/2}(z^*) = \frac{4}{3\sqrt{\pi}} \left[(\ln z^*)^{3/2} + \frac{\pi^2}{8} (\ln z^*)^{1/2} \right]$
 (pathria 9.1.30)

$\rightarrow n(r) \Big|_{T \rightarrow 0} = 2 \left(\frac{2\pi m}{\beta h^2} \right)^{3/2} \cdot \frac{4}{3\sqrt{\pi}} \beta^{3/2} \left(\mu + \frac{GMm}{r} \right)^{3/2}$
 $= \frac{8\pi}{3h^3} \cdot (2m)^{3/2} \left(\mu + \frac{GMm}{r} \right)^{3/2}$

$n(r) > 0 \rightarrow -\frac{GMm}{r_0} < \mu$, for $\mu < 0$ we have

$n(r) > 0 \rightarrow -\frac{GMm}{r} < \mu \rightarrow \boxed{r_1 < r < \frac{GMm}{|\mu|}}$

c) for $T \neq 0$ the expansion needs $T \ll \mu + \frac{GMm}{r}$
 at $r \rightarrow \infty \rightarrow T \ll \mu \rightarrow \mu \gg 0$

$$n(\mu + \frac{GMm}{r}) \Big|_{r \rightarrow \infty} \rightarrow n(\mu) \text{ always positive}$$

for $\mu > 0$, i.e. no confinement.

for $\mu < 0$ again $r < \frac{GMm}{|\mu|} \rightarrow n(r) \approx r^{-3/2}$
 No confinement!

$$\text{For } e^{\left(\frac{\mu + \frac{GMm}{r}}{kT} \right)} \ll 1$$

i.e. for $|\mu| \gg kT$, $\mu < 0$

i.e. for $e^{-\frac{|\mu|}{kT}} \ll 1$ we have

$$f_{3/2} \left(\frac{\mu + \frac{GMm}{r}}{kT} \right) \approx e^{\frac{\mu + \frac{GMm}{r}}{kT}}$$

giving us the classical sol.