

Introduce: what is \mathbb{R}^n ?

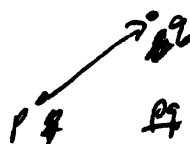
• metric topological space



$d(p, o(p)) < r$ open ball
open sets \equiv unions of open balls } used for manifolds

• vector space (linear space), $x, y \in \mathbb{R}^n \rightarrow \alpha x + \beta y \in \mathbb{R}^n$ } tangent space $T_p M, p \in M$

• affine space: rot + translations



• inner product space, $x, y \rightarrow x \cdot y \in \mathbb{R}, \|x\|^2 = x \cdot x$ } Riemann metric, geometric embeddings

2.2 Extremizing proper time



$p, q \in M$ connected by a time-like curve γ
Small deformations of γ also time-like, ∞ many time-like curves
Which curve is extremizing

$$\tau(\gamma) = \int_0^1 \sqrt{-g_{\mu\nu} \dot{x}^\mu \dot{x}^\nu} du = \int_0^1 L(x, \dot{x}) du$$

$$\gamma(0) = p$$

$$\gamma(1) = q$$

$$\gamma: [0, 1] \rightarrow M$$

$$E-L \text{ eq (4 ODEs)}$$

$$\frac{d}{du} \left(\frac{\partial L}{\partial \dot{x}^\mu} \right) = \frac{\partial L}{\partial x^\mu} \Rightarrow \frac{d}{du} \left(-\frac{1}{c} g_{\mu\nu} \dot{x}^\nu \right) - \left(-\frac{1}{2c} g_{\mu\rho, \sigma} \dot{x}^\nu \dot{x}^\sigma \right) = 0$$

Use proper time as a parameter along γ .

$$\text{where } g_{\mu\rho, \sigma} \equiv \frac{\partial g_{\mu\rho}}{\partial x^\sigma}$$

$$\frac{d\tau}{du} = L, \quad \frac{d}{du} \rightarrow L \frac{d}{d\tau}$$

$$\text{Exercise (*) } \left[\frac{1^2 x^\mu}{L^2 \tau^2} + \Gamma_{\nu\rho}^\mu \frac{dx^\nu}{d\tau} \frac{dx^\rho}{d\tau} = 0 \right]$$

geodesic equations

$$\Gamma_{\nu\rho}^\mu = \frac{1}{2} g^{\mu\sigma} (g_{\sigma\nu, \rho} + g_{\sigma\rho, \nu} - g_{\nu\rho, \sigma})$$

Christoffel symbols

[Not components of a tensor]

$$\Gamma_{\nu\rho}^{\mu} = \Gamma_{\rho\nu}^{\mu}$$

• Minkowski space, inertial frame $g = -dx^0{}^2 + dx^2$ $\rightarrow \Gamma_{\nu\rho}^{\mu} = 0$
 Change coordinates $\tilde{x}^{\mu} = \tilde{x}^{\mu}(x^{\nu}) \rightarrow \tilde{\Gamma}_{\nu\rho}^{\mu} \neq 0$.

Postulate of GR Massive test bodies in GR follow curves of extremal proper time

Exercise The same geodesic equations arise from $L = G^2 = -g_{\mu\nu} \dot{x}^{\mu} \dot{x}^{\nu}$.
 (Exercise to obtain E-L eq, and to compute $\Gamma_{\nu\rho}^{\mu}$)

Note $\frac{\partial L}{\partial t} = 0 \rightarrow$ conservation law, $H = -\left(\dot{x}^{\mu} \frac{\partial L}{\partial \dot{x}^{\mu}} - L\right) = g_{\mu\nu} \dot{x}^{\mu} \dot{x}^{\nu}$
 (constant, RHS = -1)

Example The Schwarzschild metric

$$f = 1 - \frac{2m}{r}, m = \text{const}, x^{\mu} = (t, r, \theta, \phi)$$

$$g = -f dt^2 + f^{-1} dr^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2)$$

$$t: \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{t}} \right) = \frac{d}{dt} \left(2 + \frac{dt}{dt} \right) = 0$$

$$\ddot{t} + f^{-1} f' \dot{t} \dot{r} = 0 \quad \text{compare with (*)} \quad \Gamma_{01}^0 = \Gamma_{10}^0 = \frac{1}{2} f'/f, \quad \Gamma_{\mu\nu}^0 = 0 \quad \text{otherwise}$$

2.3 Covariant derivative and Levi-Civita connection

How to differentiate tensors? E.g. $V^a =$ vector field, $\frac{\partial V^a}{\partial x^b}$ not components of a tensor

Def^h A connection (covariant derivative) is a map $\nabla: (X, Y) \in T_p M \times T_p M \rightarrow \nabla_X Y \in T_p M$ such that

$$(i) \nabla_{fX+gY} Z = f \nabla_X Z + g \nabla_Y Z \quad (f, g \text{ any functions})$$

$$(ii) \nabla_X (Y+Z) = \nabla_X Y + \nabla_X Z$$

$$(iii) \text{ (Leibniz) } \nabla_X (fY) = X(f) Y + f \nabla_X Y, \quad \nabla_X f = X(f)$$

} $\forall X, Y, Z$ vector fields

Defⁱ A covariant derivative of a vector field Y is a $(1,1)$ tensor ∇Y such that

$$(\nabla Y)^a_b \equiv \nabla_b Y^a \equiv Y^a_{;b}$$