DALBEATTIE.

Nov. 13, 1867.

Dear Tait

If you have any spare copies of your translation of Helmholtz on "Water Twists" I should be obliged if you could send me

GLENLAIR.

I set [sic] the Helmholtz dogma to the Senate House in '66, and got it very nearly done by some men, completely as to the calculation, nearly as to the interpretation.

Thomson has set himself to spin the chains of destiny out of a fluid plenum as M. Scott set an eminent person to spin ropes from the sea sand, and I saw you had put your calculus in it too. May you both prosper and disentangle your formulae in proportion as you entangle your worbles. But I fear the simplest indivisible whirl is either two embracing worbles or a worble embracing itself.

For a simple closed worble may be easily split and the parts separated



but two embracing worbles preserve each others solidarity thus



though each may split into many, every one of the one set must embrace every one of the other. So does a knotted one.



yours truly J. Clerk Maxwell

 \mathbf{c}

VECTOR-POTENTIAL OF A CLOSED CURVE. 422.

to be intertwined alternately in opposite directions, so that they are inseparably linked together though the value of the integral is zero. See Fig. 4.

It was the discovery by Gauss of this very integral, expressing the work done on a magnetic pole while de-

scribing a closed curve in presence of a closed electric current, and indicating the geometrical connexion between the two closed curves, that led him to lament the small progress made in the Geometry of Position since the time of Leibnitz, Euler and Vandermonde. We have now, however, some progress to report, chiefly due to Riemann, Helmholtz

and Listing. 422.] Let us now investigate the result of integrating with

respect to s round the closed curve.

One of the terms of Π in equation (7) is

$$\frac{\xi - x}{r^3} \frac{d\eta}{d\sigma} \frac{dz}{ds} = \frac{d\eta}{d\sigma} \frac{d}{d\xi} \left(\frac{1}{r} \frac{dz}{ds} \right). \tag{8}$$

If we now write for brevity

$$F = \int \frac{1}{r} \frac{dx}{ds} ds$$
, $G = \int \frac{1}{r} \frac{dy}{ds} ds$, $H = \int \frac{1}{r} \frac{dz}{ds} ds$,

the integrals being taken once round the closed curve s, this term of Π may be written $\frac{d\eta}{d\sigma} \frac{d^2H}{d\xi ds}$,

and the corresponding term of $\int \Pi ds$ will be

$$\frac{d\eta}{d\sigma} \frac{dH}{d\xi}$$

Collecting all the terms of II, we may now write

$$\begin{split} &-\frac{d\omega}{d\sigma} = -\int \Pi \; ds \\ &= \left(\frac{dH}{d\eta} - \frac{dG}{d\zeta}\right) \frac{d\xi}{d\sigma} + \left(\frac{dF}{d\xi} - \frac{dH}{d\eta}\right) \frac{d\eta}{d\sigma} + \left(\frac{dG}{d\xi} - \frac{dF}{d\eta}\right) \frac{d\zeta}{d\sigma} \cdot (10) \end{split}$$

This quantity is evidently the rate of decrement of ω, the magnetic potential, in passing along the curve σ, or in other words, it is the magnetic force in the direction of $d\sigma$.

By assuming $d\sigma$ successively in the direction of the axes of x, y and z, we obtain for the values of the components of the magnetic force

h two dused curves and or the lecture between Diper, and LMN are direction comment of do do de v respection State I M N = 1 ds do (- #)(- # the internation being extended revend hall seeme and on being the alephanest murcher of times that one came embases they other in the sume direction. If the curve are not linked together no 0 but if n = 0 the curvey are not necessarily independent by I the time closed curves are unseparable but n = 0. In hig 2 the 3 closed curves a inseparable but n = 0 for every pour of them Fig 3 is the simplest sind hud on curve. The simplest equation I can 1 = 3 + a cos 3 8 2= c sm 3 A when c is - was in the figure the best is righthered, where c is + we it is left healed. I righthered house - ve as in the cannot be changed into a left hunded an

