Solving the neutron transport equation within a diffusive regime

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The problem

An important problem in nuclear physics is that of efficiently solving the neutron transport equation. This equation governs the behaviour of neutrons within a nuclear fission reactor. It is used to model reactors for testing and simulation purposes.

$$\Omega \cdot \nabla \psi(x,\Omega) + \sigma_T \psi(x,\Omega) = \frac{\sigma_S}{2} \int_{-1}^1 \psi(x,\Omega') d\Omega' + q(x,\Omega)$$

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Boundary conditions:
$$\begin{array}{ll} \psi(a,\mu)=0, & \mu>0 \\ \psi(b,\mu)=0, & \mu<0 \end{array}$$
 , with $\begin{array}{ll} x\in[a,b] \\ \mu\in[-1,1] \end{array}$

Note: $\sigma_T = \sigma_S + \sigma_A$

Operator notation and source iteration Introduce

$$\mathcal{T}(\cdot) \equiv \mu \frac{\partial}{\partial x}(\cdot) + \sigma_{\mathcal{T}}(\cdot)$$
 and define $\phi(x) = \frac{1}{2} \int_{-1}^{1} \psi(x, \mu) d\mu$.

 ϕ is called the scalar flux. The transport equation is then

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Define source iteration as follows

$$\mathcal{T}\psi^{(i+1)} = \sigma_S \phi^{(i)} + q$$
$$\phi^{(i+1)} = \frac{1}{2} \int_{-1}^{1} \psi^{(i+1)} d\mu$$

This basic iterative method is a contraction, with the difference between successive iterations bounded as (F. Scheben, 2011)

$$\left\|\phi^{(i+1)} - \phi^{(i)}\right\|_{2} \leq \frac{\sigma_{S}}{\sigma_{T}} \left\|\phi^{(i)} - \phi^{(i-1)}\right\|_{2}.$$

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$$\sigma_T \equiv \frac{1}{\epsilon}, \quad \sigma_A \equiv \epsilon, \quad \sigma_S \equiv \left[\frac{1}{\epsilon} - \epsilon\right], \quad \text{and} \quad q(x) \equiv \epsilon Q(x)$$

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SPOILER:

$$-\frac{1}{3}\frac{\mathrm{d}^2}{\mathrm{d}x^2}\phi + \phi = Q + O(\epsilon^2).$$

(boundary conditions: ?)

Implications for source iteration

Looking back to the contraction inequality, we note

$$\|\phi^{(i+1)} - \phi^{(i)}\|_{2} \le \frac{\sigma_{S}}{\sigma_{T}} \|\phi^{(i)} - \phi^{(i-1)}\|_{2}$$
$$= [1 - \epsilon^{2}] \|\phi^{(i)} - \phi^{(i-1)}\|_{2}$$

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We want to work with a pair of coupled equations for ψ and ϕ in block operator form, and for this the following will be convenient:

$$\mathcal{P}(\cdot) \equiv \frac{1}{2} \int_{-1}^{1} (\cdot) d\mu, \quad \Sigma(\cdot) \equiv \left[\frac{1}{\epsilon} - \epsilon \right] (\cdot).$$

Then
$$\phi(x) \equiv \mathcal{P}\psi(x,\mu)$$

Working from the transport equation in block operator form

$$\left(\begin{array}{cc} \mathcal{T} & -\Sigma \\ -\mathcal{P} & \mathcal{I} \end{array} \right) \left(\begin{array}{c} \psi \\ \phi \end{array} \right) = \left(\begin{array}{c} \epsilon Q \\ 0 \end{array} \right)$$

Working from the transport equation in block operator form

$$\left(\begin{array}{cc} \mathcal{I} & 0 \\ \mathcal{P}\mathcal{T}^{-1} & \mathcal{I} \end{array} \right) \left(\begin{array}{cc} \mathcal{T} & -\Sigma \\ -\mathcal{P} & \mathcal{I} \end{array} \right) \left(\begin{array}{c} \psi \\ \phi \end{array} \right) = \left(\begin{array}{cc} \mathcal{I} & 0 \\ \mathcal{P}\mathcal{T}^{-1} & \mathcal{I} \end{array} \right) \left(\begin{array}{c} \epsilon Q \\ 0 \end{array} \right)$$

where \mathcal{T}^{-1} is defined in F.Scheben, 2011, for the 1D case.

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$$\begin{pmatrix} \mathcal{T} & -\Sigma \\ 0 & \mathcal{I} - \mathcal{P}\mathcal{T}^{-1}\Sigma \end{pmatrix} \begin{pmatrix} \psi \\ \phi \end{pmatrix} = \begin{pmatrix} \epsilon Q \\ \epsilon \mathcal{P}\mathcal{T}^{-1}Q \end{pmatrix}$$

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Working from the transport equation in block operator form

$$\left(\begin{array}{cc} \mathcal{I} & \mathbf{0} \\ \mathcal{P}\mathcal{T}^{-1} & \mathcal{I} \end{array} \right) \left(\begin{array}{cc} \mathcal{T} & -\Sigma \\ -\mathcal{P} & \mathcal{I} \end{array} \right) \left(\begin{array}{c} \psi \\ \phi \end{array} \right) = \left(\begin{array}{cc} \mathcal{I} & \mathbf{0} \\ \mathcal{P}\mathcal{T}^{-1} & \mathcal{I} \end{array} \right) \left(\begin{array}{c} \epsilon Q \\ \mathbf{0} \end{array} \right)$$

$$\Rightarrow \qquad \left(\begin{array}{cc} \mathcal{T} & -\Sigma \\ 0 & \mathcal{I} - \mathcal{P}\mathcal{T}^{-1}\Sigma \end{array}\right) \left(\begin{array}{c} \psi \\ \phi \end{array}\right) = \left(\begin{array}{c} \epsilon Q \\ \epsilon \mathcal{P}\mathcal{T}^{-1}Q \end{array}\right)$$

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Diffusion Approximation Theorem:

$$\frac{1}{\epsilon^2} (\mathcal{I} - \mathcal{P} \mathcal{T}^{-1} \Sigma) \phi = -\frac{\mathrm{d}}{\mathrm{d}x} \left(\frac{1}{3} \frac{\mathrm{d}}{\mathrm{d}x} \phi \right) + \phi + O(\epsilon^2)$$

Heuristic proof

We know $\mathcal{T}=\frac{1}{\epsilon}\left(\mathcal{I}+\epsilon\mu\frac{\partial}{\partial x}\right)$, where \mathcal{I} is the identity operator. Expand:

$$\mathcal{T}^{-1} = \epsilon \left(\mathcal{I} - \epsilon \mu \frac{\partial}{\partial x} + \epsilon^2 \mu^2 \frac{\partial^2}{\partial x^2} - \dots \right)$$

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Using this expansion

$$\mathcal{P}\mathcal{T}^{-1}\Sigma = \epsilon \mathcal{P}\left(\mathcal{I} - \epsilon \mu \frac{\partial}{\partial x} + \epsilon^2 \mu^2 \frac{\partial^2}{\partial x^2} - \dots\right) \left[\frac{1}{\epsilon} - \epsilon\right]$$
$$= \left[1 - \epsilon^2\right] \left(\mathcal{P} - \epsilon \frac{\partial}{\partial x} \mathcal{P}\mu + \epsilon^2 \frac{\partial^2}{\partial x^2} \mathcal{P}\mu^2 - \dots\right).$$

Note that $\mathcal{P}\mu=$ 0, $\mathcal{P}\mu^2=\frac{1}{3}$ and $\mathcal{P}\phi=\phi$,

Heuristic proof

We know $\mathcal{T} = \frac{1}{\epsilon} \left(\mathcal{I} + \epsilon \mu \frac{\partial}{\partial x} \right)$, where \mathcal{I} is the identity operator. Expand:

$$\mathcal{T}^{-1} = \epsilon \left(\mathcal{I} - \epsilon \mu \frac{\partial}{\partial x} + \epsilon^2 \mu^2 \frac{\partial^2}{\partial x^2} - \dots \right)$$

Using this expansion

$$\mathcal{P}\mathcal{T}^{-1}\Sigma = \epsilon \mathcal{P}\left(\mathcal{I} - \epsilon \mu \frac{\partial}{\partial x} + \epsilon^2 \mu^2 \frac{\partial^2}{\partial x^2} - \dots\right) \left[\frac{1}{\epsilon} - \epsilon\right]$$
$$= \left[1 - \epsilon^2\right] \left(\mathcal{P} - \epsilon \frac{\partial}{\partial x} \mathcal{P}\mu + \epsilon^2 \frac{\partial^2}{\partial x^2} \mathcal{P}\mu^2 - \dots\right).$$

Note that $\mathcal{P}\mu=0$, $\mathcal{P}\mu^2=\frac{1}{3}$ and $\mathcal{P}\phi=\phi$, so

$$\frac{1}{\epsilon^2} (\mathcal{I} - \mathcal{P} \mathcal{T}^{-1} \Sigma) \phi = \frac{1}{\epsilon^2} \left(\mathcal{I} - \mathcal{P} - \frac{\epsilon^2}{3} \frac{\partial^2}{\partial x^2} \mathcal{P} + \epsilon^2 \mathcal{P} + O(\epsilon^4) \right) \phi$$

$$= -\frac{1}{3} \frac{d^2}{dx^2} \phi + \phi + O(\epsilon^2).$$



Discretise

Transport equation in discrete block matrix form

$$\left(\begin{array}{cc} T & -\Sigma \\ -P & I \end{array}\right) \left(\begin{array}{c} \psi \\ \phi \end{array}\right) = \left(\begin{array}{c} \epsilon Q \\ 0 \end{array}\right)$$

Discretise

Transport equation in discrete block matrix form

$$\Rightarrow \qquad \left(\begin{array}{cc} T & -\Sigma \\ 0 & I - PT^{-1}\Sigma \end{array}\right) \left(\begin{array}{c} \psi \\ \phi \end{array}\right) = \left(\begin{array}{c} \epsilon Q \\ \epsilon PT^{-1}Q \end{array}\right)$$

Discretise

Transport equation in discrete block matrix form

$$\Rightarrow \qquad \left(\begin{array}{cc} T & -\Sigma \\ 0 & I - PT^{-1}\Sigma \end{array}\right) \left(\begin{array}{c} \psi \\ \phi \end{array}\right) = \left(\begin{array}{c} \epsilon Q \\ \epsilon PT^{-1}Q \end{array}\right)$$

- ► Form source iteration in block matrix form
- Subtract this from block matrix transport equation
- Conduct the same Gaussian elimination

$$\begin{pmatrix} T & -\Sigma \\ 0 & I - PT^{-1}\Sigma \end{pmatrix} \begin{pmatrix} \psi - \psi^{(i+1)} \\ \phi - \phi^{(i+1)} \end{pmatrix} = \begin{pmatrix} \Sigma(\phi^{(i+1)} - \phi^{(i)}) \\ PT^{-1}\Sigma(\phi^{(i+1)} - \phi^{(i)}) \end{pmatrix}$$

A quick numerical experiment

Iteration	$\ \phi - \phi^{(i)}\ _2$
	for $\epsilon=10^{-3}$
1	1.0e-004
2	1.8e-004
3	9.6e-005
4	4.9e-005
:	:
14	5.1e-008
15	2.6e-008
16	1.3e-008
17	6.4e-009

Source iteration would take roughly 5 million iterations for the same calculation.



Iterative Methods for Criticality Computations in Neutron Transport Theory.

PhD thesis, University of Bath, 2011.

Edward W Larsen, J.E Morel, and Warren F Miller Jr.
Asymptotic solutions of numerical transport problems in optically thick, diffusive regimes.

Journal of Computational Physics, 69(2):283 – 324, 1987.