

WS3: Numerical Integration

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1 Integration via Newton-Cotes Formulae

I compared numerical integration using the Trapezoid Rule and Simpson's Rule, testing both schemes on two integrals: $\sin(x)$ from 0 to π , and $x \sin(x)$ from 0 to π .

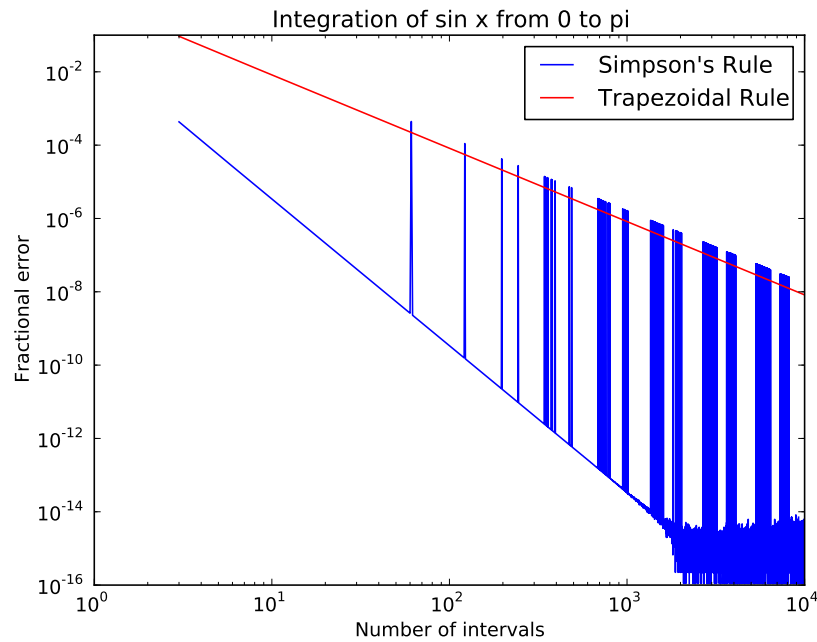


Figure 1: Fractional error on numerical integration of $\sin(x)$ from 0 to π , using the Trapezoid Rule and Simpson's Rule, for a range of step sizes.

Figures 1 and 2 show the fractional errors, $(\text{calculated answer} - \text{real answer})/(\text{real answer})$, of the integrations, which decrease with increasing number

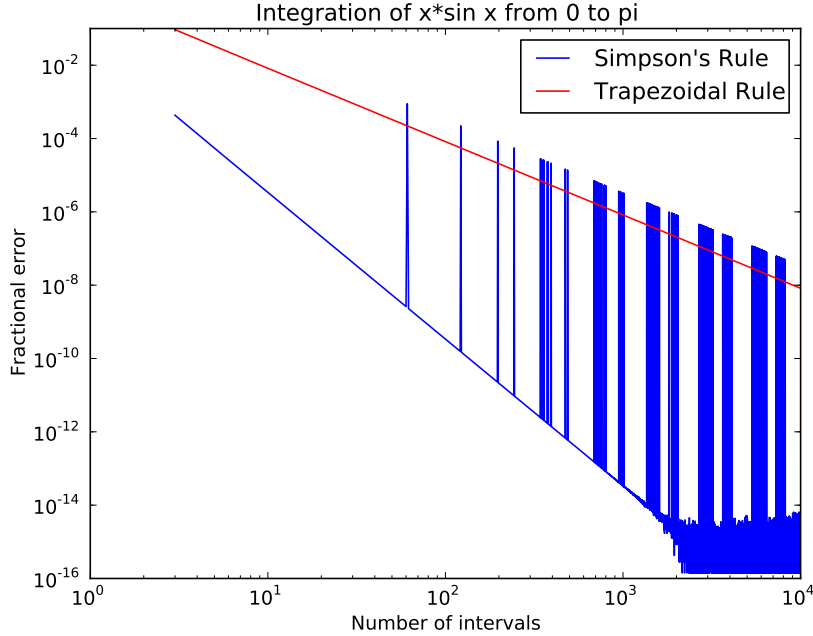


Figure 2: Fractional error on numerical integration of $x \sin(x)$ from 0 to π , using the Trapezoid Rule and Simpson's Rule, for a range of step sizes.

of intervals N . $N = \pi/h$ for step size h . Since the error on the Trapezoid Rule goes as h^2 or N^{-2} , the slope of the error on the Trapezoid Rule is -2 on a log-log plot. Similarly, Simpson's Rule is $O(h^4)$, so the slope of the fractional error is -4 on the log-log plot. The fractional error flattens out around 10^{-14} , limited by floating-point numerical precision.

There are a bunch of weird spikes in the error on Simpson's Rule, indicating that when using Simpson's Rule, not all tiny step sizes are equal. It's interesting to note that the spikes follow the same slope as the Trapezoid Rule, rather than the slope of Simpson's Rule.

2 Gaussian Quadrature

We consider an environment where photons and electron-positron pairs are in equilibrium, so that the electron/positron density is determined by the Fermi distribution. We assume a temperature of $k_B T = 20$ MeV, so that the electrons are highly relativistic. The density of electrons is determined by the integral:

$$n_e = \frac{8\pi(k_B T)^3}{(2\pi\hbar c)^3} \int_0^\infty \frac{x^2 dx}{e^x + 1} \quad (1)$$

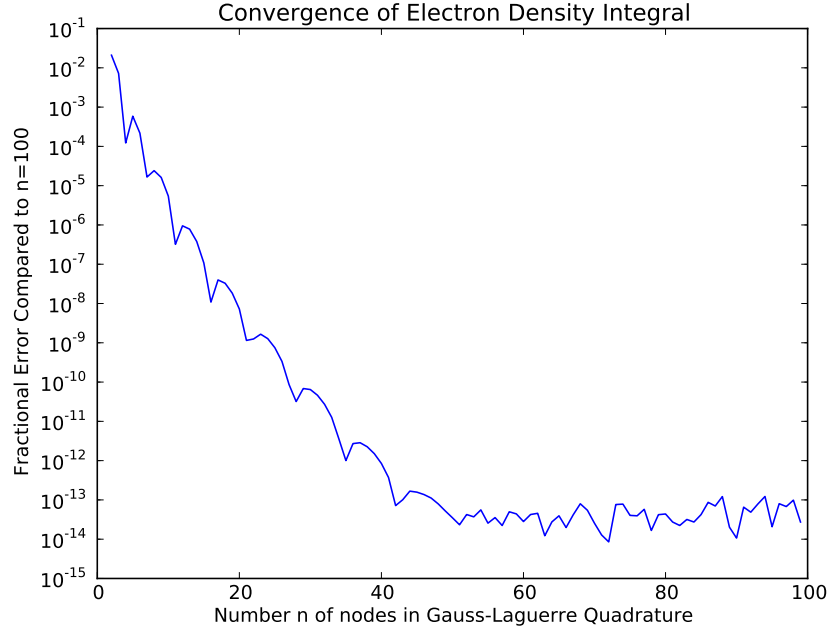


Figure 3: Fractional error on Gauss-Laguerre integration to obtain total electron density at $k_B T = 20$ MeV, for n nodes.

I used Gauss-Laguerre integration, which is useful for integrals that are unbounded on one end, to solve this integral, obtaining a total electron density of $1.902 \times 10^{35} \text{ cm}^{-3}$ for $n=150$ nodes. Figure 3 shows the fractional error on this integral, calculated by considering the $n=100$ solution to be the “true” solution. The fractional error flattens off between 10^{-13} and 10^{-14} due to floating-point rounding error.

Figure 4 shows the electron energy spectrum, dn_e/dE , from 0 to 250 MeV. I used Gauss-Legendre integration over intervals of 5 MeV to calculate the density of electrons in that energy range, and divided by the interval width to get the derivative. When I add up the derivative from 0 to 250 MeV and multiply by the interval width (effectively integrating), I obtain $n_e = 1.901 \times 10^{35}$, consistent with the previous integration.

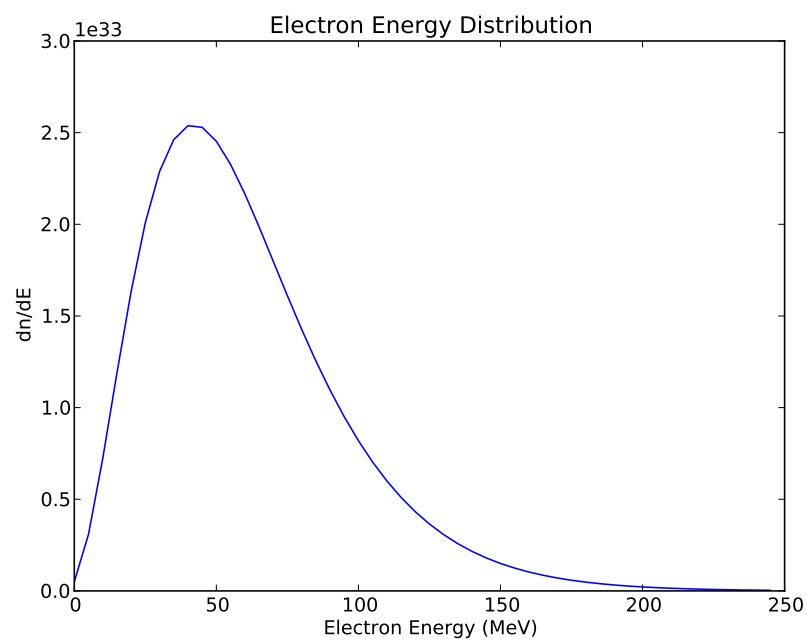


Figure 4: Electron energy distribution, obtained by Gauss-Legendre integration over 5 MeV intervals.