

[12.8] Let $\alpha = \sum_{r=1}^n \alpha_r dx^r, \dots, \gamma = \sum_{u=1}^n \gamma_u dx^u, \lambda = \sum_{j=1}^n \lambda_j dx^j, \dots, v = \sum_{m=1}^n v_m dx^m$ be independent

1-forms in \mathbb{R}^n . Let $\phi = \alpha \wedge \dots \wedge \gamma$ be a simple p -form and $\chi = \lambda \wedge \dots \wedge v$ be a simple q -form. Show that $\phi \wedge \chi = \alpha \wedge \dots \wedge \gamma \wedge \lambda \wedge \dots \wedge v$.

Proof: Juergen Beckmann's proof uses the antisymmetrized-coefficients expression for the wedge product and is well presented and very insightful. This problem can also be solved using the un-antisymmetrized-coefficients expression.

Let $M = \{1, 2, \dots, n\}$, $\mathcal{P}_{r \dots u}$ be the set of permutations of the p -tuple (r, \dots, u) and $\mathcal{P}_{j \dots m}$ the set of permutations of the q -tuple (j, \dots, m) . As Juergen reminds us in his proof, the un-antisymmetrized expressions for ϕ and χ are

$$\phi = \alpha \wedge \dots \wedge \gamma = \sum_{(r, \dots, u) \in M^p} \alpha_{[r} \dots \gamma_{u]} dx^r \wedge \dots \wedge dx^u \quad (1)$$

where

$$\alpha_{[r} \dots \gamma_{u]} = \frac{1}{p!} \sum_{\pi \in \mathcal{P}_{r \dots u}} \text{sign}(\pi) \alpha_{\pi(r)} \dots \gamma_{\pi(u)}$$

and

$$\chi = \lambda \wedge \dots \wedge v = \sum_{(j, \dots, m) \in M^q} \lambda_{[j} \dots v_{m]} dx^j \wedge \dots \wedge dx^m \quad (2)$$

where

$$\lambda_{[j} \dots v_{m]} = \frac{1}{q!} \sum_{\pi \in \mathcal{P}_{j \dots m}} \text{sign}(\pi) \lambda_{\pi(j)} \dots v_{\pi(m)}.$$

The un-antisymmetrized expression for $\phi \wedge \chi$ is

$$\begin{aligned} \phi \wedge \chi &= \sum_{(r, \dots, u) \in M^p} \sum_{(j, \dots, m) \in M^q} (\alpha_r \dots \gamma_u) (\lambda_j \dots v_m) (dx^r \wedge \dots \wedge dx^u) \wedge (dx^j \wedge \dots \wedge dx^m) \\ &= \sum_{(r, \dots, u) \in M^p} \sum_{(j, \dots, m) \in M^q} \alpha_r \dots \gamma_u \lambda_j \dots v_m dx^r \wedge \dots \wedge dx^u \wedge dx^j \wedge \dots \wedge dx^m. \end{aligned} \quad (3)$$

The un-antisymmetrized expression for $\alpha \wedge \dots \wedge \gamma \wedge \lambda \wedge \dots \wedge v$ is

$$\alpha \wedge \dots \wedge \gamma \wedge \lambda \wedge \dots \wedge v = \sum_{(r, \dots, u, j, \dots, m) \in M^{p+q}} \alpha_r \dots \gamma_u \lambda_j \dots v_m dx^r \wedge \dots \wedge dx^u \wedge dx^j \wedge \dots \wedge dx^m. \quad (4)$$

Both expressions (3) and (4) have n^{p+q} terms, and all terms are identical. Therefore

$\phi \wedge \chi = \alpha \wedge \cdots \wedge \gamma \wedge \lambda \wedge \cdots \wedge v$, which finishes the proof. ■

Note: Juergen's proof has several nice features and he clarifies issues only hinted at by Penrose.

First, he introduces an elegant notation to represent the antisymmetrization of two antisymmetrizations:

$$\phi \wedge \chi = \sum_{(r, \dots, u, j, \dots, m) \in M^{p+q}} \alpha_{[r} \cdots \gamma_{u]} \lambda_{[j} \cdots v_{m]} dx^r \wedge \cdots \wedge dx^u \wedge dx^j \wedge \cdots \wedge dx^m. \quad (5)$$

I, for one, had a hard time figuring out where to put the outside brackets in order to express the outer antisymmetrization.

He also clarified the meaning of the RHS of (5). I considered expanding the RHS using a $2!$ permutation because I thought that the outer antisymmetrization was operating on the 2 quantities $\alpha_{[r} \cdots \delta_{u]}$ and $\lambda_{[j} \cdots v_{m]}$, but that was incorrect. Juergen's explanation uses a permutation expression involving $(p+q)!$ and I believe it translates to

$$\begin{aligned} & \phi \wedge \chi \\ &= \sum_{(r, \dots, u, j, \dots, m) \in M^{p+q}} \frac{1}{(p+q)!} \sum_{\pi \in \mathcal{P}_{r \dots u j \dots m}} \text{sign}(\pi) \alpha_{\pi[r} \cdots \gamma_{u]} \lambda_{\pi[j} \cdots v_{m]} dx^r \wedge \cdots \wedge dx^u \wedge dx^j \wedge \cdots \wedge dx^m. \end{aligned}$$

Next he expresses the inner antisymmetrizations in terms of permutations:

$$\begin{aligned} \phi \wedge \chi &= \sum_{(r, \dots, u, j, \dots, m) \in M^{p+q}} \frac{1}{(p+q)!} \sum_{\pi \in \mathcal{P}_{r \dots u j \dots m}} \text{sign}(\pi) \\ & \quad \frac{1}{p!} \sum_{\pi_1 \in \mathcal{P}_{r \dots u}} \text{sign}(\pi_1) \alpha_{\pi(\pi_1(r)} \cdots \gamma_{\pi_1(u))} dx^r \wedge \cdots \wedge dx^u \\ & \quad \frac{1}{q!} \sum_{\pi_2 \in \mathcal{P}_{j \dots m}} \text{sign}(\pi_2) \lambda_{\pi(\pi_2(j)} \cdots v_{\pi_2(m))} \wedge dx^j \wedge \cdots \wedge dx^m \\ &= \sum_{(r, \dots, u, j, \dots, m) \in M^{p+q}} \frac{1}{(p+q)!} \frac{1}{p!} \frac{1}{q!} \sum_{\pi \in \mathcal{P}_{r \dots u j \dots m}} \sum_{\pi_1 \in \mathcal{P}_{r \dots u}} \sum_{\pi_2 \in \mathcal{P}_{j \dots m}} \text{sign}(\pi) \text{sign}(\pi_1) \text{sign}(\pi_2) \\ & \quad \alpha_{\pi \circ \pi_1(r)} \cdots \gamma_{\pi \circ \pi_1(u)} \lambda_{\pi \circ \pi_2(j)} \cdots v_{\pi \circ \pi_2(m)} dx^r \wedge \cdots \wedge dx^u \wedge dx^j \wedge \cdots \wedge dx^m. \end{aligned}$$

Juergen then finished his proof with some clever re-writing of the permutations to eliminate the two inner antisymmetrizations (i.e., the 2^{nd} and 3^{rd} summation signs), leading to the (desired) expression for $\alpha \wedge \cdots \wedge \gamma \wedge \lambda \wedge \cdots \wedge v$. My only purpose here was to attempt to translate his notation without his reference to a generic antisymmetric operator.

We provide two examples, one illustrating the permutation notation and the other illustrating the [12.8] problem statement.

Example: Permutation Notation

Let $\alpha = \alpha_1 dx^1 + \alpha_2 dx^2$ and $\beta = \beta_1 dx^1 + \beta_2 dx^2$ be 1-forms in \mathbb{R}^2 . Then $n = p = 2$, $M = \{1, 2\}$, and $M^p = \{1, 2\}^2 = \{1, 2\} \times \{1, 2\} = \{(1,1), (1,2), (2,1), (2,2)\}$. So $\mathcal{P}_{12} = \{\pi_1, \pi_2\}$ where

$$\pi_1: \begin{cases} 1 \rightarrow 1 \\ 2 \rightarrow 2 \end{cases}, \quad \pi_2: \begin{cases} 1 \rightarrow 2 \\ 2 \rightarrow 1 \end{cases}, \quad \text{sign}(\pi_1) = +1, \text{ and } \text{sign}(\pi_2) = -1.$$

So, for example,

$$\begin{aligned} \alpha_1 \beta_2 &= \frac{1}{2} (\alpha_1 \beta_2 - \alpha_2 \beta_1) = \frac{1}{2!} (\text{sign}(\pi_1) \alpha_{\pi_1(1)} \beta_{\pi_1(2)} + \text{sign}(\pi_2) \alpha_{\pi_2(1)} \beta_{\pi_2(2)}) \\ &= \frac{1}{p!} \sum_{\pi \in \mathcal{P}_{12}} \text{sign}(\pi) \alpha_{\pi(1)\pi(2)}. \end{aligned} \quad \blacksquare$$

Example: $\alpha \wedge (\lambda \wedge v)$

Let $\alpha = \sum_{r=1}^n \alpha_r dx^r$, $\lambda = \sum_{s=1}^n \lambda_s dx^s$, and $v = \sum_{t=1}^n v_t dx^t$. Set $\sigma = \lambda \wedge v = \sum_{s,t} \sigma_{st} dx^s dx^t$ where

$$\sigma_{st} = \frac{1}{2} \sum_{\pi_1 \in \mathcal{P}_{st}} \text{sign}(\pi_1) \lambda_{\pi_1(s)} v_{\pi_1(t)} = \frac{1}{2} (\lambda_s v_t - \lambda_t v_s). \text{ So } \sigma = \sum_{s,t} \frac{1}{2} (\lambda_s v_t - \lambda_t v_s) dx^s \wedge dx^t.$$

Next, let $\tau = \alpha \wedge \sigma = \alpha \wedge (\lambda \wedge v) = \sum_{r,s,t} \tau_{rst} dx^r \wedge (dx^s \wedge dx^t) = \sum_{r,s,t} \tau_{rst} dx^r \wedge dx^s \wedge dx^t$ where

$$\begin{aligned} \tau_{rst} &= \alpha_{[r} \sigma_{st]} = \frac{1}{3!} \sum_{\pi \in \mathcal{P}_{rst}} \text{sign}(\pi) \alpha_{\pi(r)} \sigma_{\pi(st)} = \frac{1}{3!} \sum_{\pi \in \mathcal{P}_{rst}} \text{sign}(\pi) \alpha_{\pi(r)} \lambda_{[\pi(s)} v_{\pi(t)]} \\ &= \frac{1}{3!} \sum_{\pi \in \mathcal{P}_{rst}} \text{sign}(\pi) \alpha_{\pi(r)} \frac{1}{2!} \sum_{\pi_1 \in \mathcal{P}_{st}} \text{sign}(\pi_1) \lambda_{\pi \circ \pi_1(s)} v_{\pi \circ \pi_1(t)} \\ &= \frac{1}{3!} \frac{1}{2!} \sum_{\pi \in \mathcal{P}_{rst}} \sum_{\pi_1 \in \mathcal{P}_{st}} \text{sign}(\pi) \text{sign}(\pi_1) \alpha_{\pi(r)} \lambda_{\pi \circ \pi_1(s)} v_{\pi \circ \pi_1(t)}. \end{aligned}$$