

Lecture 2 Forms & connections

①

Dimension 4 is special in differential geometry:

- 1) 2-forms on 4 manifolds have special properties
- 2) 2-forms ^{naturally} arise to curvature of connections.

X any smooth manifold

$$0 \rightarrow \Omega^0(X) \xrightarrow{d} \Omega^1(X) \xrightarrow{d} \Omega^2(X) \rightarrow \dots \rightarrow \Omega^n(X) \rightarrow 0 \quad \text{de Rham complex.}$$

($d^2=0$).

\rightarrow cohomology is $H^i(X; \mathbb{R})$.

Hodge theory $(X, g) \Rightarrow H^i(X; \mathbb{R})$ comes with natural representatives.

$V, \langle \rangle, \text{orientation} \rangle : * : \overset{\text{wedge}}{\Lambda^k V} \rightarrow \Lambda^{n-k} V.$ (more intuitively, sends a volume form of a plane to

$\hookrightarrow \dim n$ $e_1 \dots e_k$ $\rightarrow e_{k+1} \dots e_n$
 \uparrow ON vectors \uparrow ON vectors
 $+ e_1 \dots e_k \dots e_{k+1} \dots e_n$ is n -vector basis.

example $V = \mathbb{R}^3$. $*e_1 = e_2 \wedge e_3$
 $*e_2 = e_3 \wedge e_1$

Exercise $*^2 = (-1)^{(n-1)k}$

$(X, g, \text{oriented}) : * : \Omega^k(X) \rightarrow \Omega^{n-k}(X).$

$\hookrightarrow \langle \alpha, \beta \rangle = \alpha \wedge * \beta.$ in particular:
 \uparrow inner product ($\|\alpha\|_{L^2(X)}^2 = \int \alpha \wedge * \alpha.$)

$d : \Omega^{k-1}(X) \rightarrow \Omega^k(X)$

$d^* = (-1)^{n(k-1)+1} * d *.$

\hookrightarrow if the formal L^2 -adjoint $\alpha \in \Omega^k, \beta \in \Omega^{k-1}$

$\langle \alpha, d\beta \rangle_{L^2(X)} = \langle d^* \alpha, \beta \rangle_{L^2(X)}.$

Exercise $\overset{\text{Hut}}{0} = \int d(\beta \wedge * \alpha) = \dots$

Hodge theory There is an L^2 orthogonal decomposition

$$\Omega^k = \underbrace{d\Omega^{k-1}}_{\text{exact}} \oplus \underbrace{\mathcal{H}_k}_{\text{ker}(dd^*)} \oplus \underbrace{d^*\Omega^{k+1}}_{\text{coexact}} \rightarrow \text{needs some analytical work}$$

"harmonic" $\mathcal{H}_k \cong \mathcal{H}^k(X; \mathbb{R})$

Remark

$$(d+d^*)^2 = dd^* + d^*d = \Delta \text{ is the Hodge Laplacian.}$$

Δ is the square of an operator!

On a 4-manifold: $\star^2: \Omega^2 \rightarrow \Omega^2$ is the identity $\Rightarrow \Omega^2 = \Omega^+ \oplus \Omega^-$

$\begin{matrix} \text{self-dual} & \swarrow & \text{anti-self-dual} \\ (+1) & \nearrow & (-1) \\ \text{eigenspace} & & \end{matrix}$

$$\Omega^+ = \text{Span} \left\{ \begin{matrix} dx_1 \wedge dx_2 + dx_3 \wedge dx_4 \\ dx_1 \wedge dx_3 + dx_4 \wedge dx_2 \\ dx_1 \wedge dx_4 + dx_2 \wedge dx_3 \end{matrix} \right\}$$

$$\text{Sol}(4) \cong \text{Sol}(3) \oplus \text{Sol}(3)$$

the only indecomposable Soln.
($\cong A_4$ is the only simple)

$$\Rightarrow \mathcal{H}_2 = \mathcal{H}^+ \oplus \mathcal{H}^-$$

$\alpha \mapsto \star \alpha \mapsto \star \star \alpha \mapsto \alpha$ on $H^2(X; \mathbb{Z}) \otimes \mathbb{R}$.

$$\alpha = \star \alpha \Rightarrow \int \alpha \wedge \alpha = \int \alpha \wedge \star \alpha = \|\alpha\|_{L^2(X)}^2 > 0$$

$$\alpha = -\star \alpha \Rightarrow \int \alpha \wedge \alpha = -\int \alpha \wedge \star \alpha = -\|\alpha\|_{L^2(X)}^2 < 0$$

$$d^+ = \pi_+ \circ d$$

Exercise $0 \rightarrow \Omega^0(X) \xrightarrow{d} \Omega^1(X) \xrightarrow{d^+} \Omega^2(X) \rightarrow 0$
has homology $\mathbb{R}, H^1(X), H^2(X)$. Hint (what's the adjoint of d^+ ?)

$\begin{matrix} \mathbb{R} \\ \downarrow \\ E \\ \downarrow \\ X \end{matrix}$ vector bundle. A connection is

$$\nabla: \Omega^0(E) \rightarrow \Omega^1(E) \quad \left| \begin{array}{l} \nabla_{fX} s = f \nabla_X s \\ \nabla_X fs = df s + f \nabla_X s \end{array} \right.$$

Geometrically: $\nabla s = 0$ parallel transport.

If ∇ is a connection, any other is $\nabla + a \mapsto \Omega^1(\mathfrak{g}(E))$. (affine space). (3)

Curvature of a connection A is $F_A \in \Omega^2(\mathfrak{g}(E))$ local invariant

$$F_A(X, Y)s = \nabla_X \nabla_Y s - \nabla_Y \nabla_X s - \nabla_{[X, Y]} s$$

\nearrow vector fields \nwarrow to take turns

measures how parallel transport is non-commutative

$$[X, Y] = 0$$



Rule If E is unitary, we can ask ∇ to preserve norm. $\Rightarrow F_A \in \Omega^2(\mathfrak{u}(E))$.

Chern-Weil theory The curvature can recover global invariants of E .

E complex $\Rightarrow c_i(E) \in H^{2i}(X; \mathbb{Z})$.

$F_A \in \Omega^2(\mathfrak{g}(E))$. If we pick basis of local sections of E
 $\Rightarrow F_A$ is an $n \times n$ matrix of 2-forms (well defined up to conjugacy).
of degree k .

$P_k: \mathfrak{g}(E) \rightarrow \mathbb{C}$. invariant polynomial. ($P(YXY^{-1}) = P(X)$)

$\Rightarrow P_k(F_A)$ is a well defined form (if possibly $\in \Omega^{2k}(X, \mathbb{R})$)

Then $P_k(F_A)$ is a closed form; $[P_k(F_A)]$ is independent of A
 \Rightarrow invariant of E .

$$\text{If } P_k^{(X)} = \left(\frac{i}{2\pi}\right)^k \text{tr}(X^k) \Rightarrow c_k(E).$$

(we will only use: $\frac{i}{2\pi} F_A$ is a closed 2-form representing $c_2(E)$).

Exercise show that $c_1(\text{topological}) = -1 \in H^2(\mathbb{CP}^1; \mathbb{Z})$