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Boundary Conditions

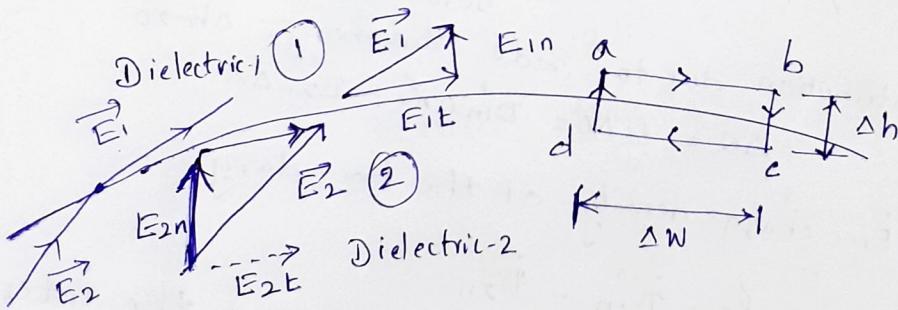
So far, we have considered the existence of the electric field in a homogeneous medium. If the field exist in a region consisting of two different media, the conditions that the field must satisfy at the interface separating the media are called boundary conditions. There are three types

* Dielectric (ϵ_{r1}) and Dielectric (ϵ_{r2}) Boundary Conditions

* Conductor - Dielectric Boundary Conditions

* Conductor - Free Space Boundary Conditions

Dielectric - Dielectric Boundary Conditions



$$\vec{E}_1 = \vec{E}_{1t} + \vec{E}_{1n, in}$$

$$\vec{E}_2 = \vec{E}_{2t} + \vec{E}_{2n}$$

Apply $\oint \vec{E} \cdot d\vec{r} = 0$ to the loop abcd a

$$E_{1t} \Delta W - E_{1n} \frac{\Delta h}{2} - E_{2n} \frac{\Delta h}{2} - E_{2t} \Delta W$$

$$+ E_{2n} \frac{\Delta h}{2} + E_{1n} \frac{\Delta h}{2} = 0$$

$$\vec{E}_{1t} \Delta W = \vec{E}_{2t} \Delta W$$

$$\boxed{\vec{E}_{1t} = \vec{E}_{2t}}$$

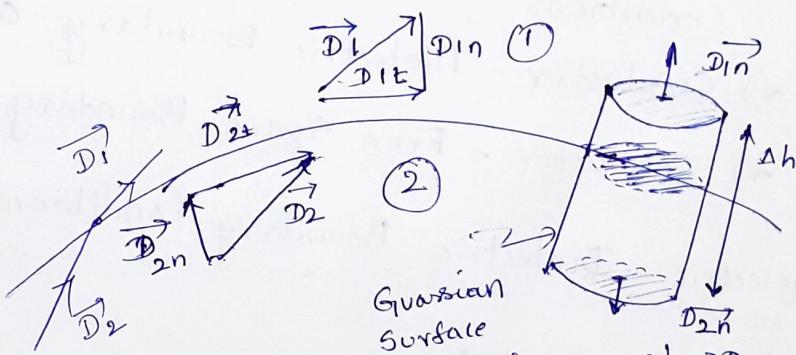
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Experiment

$$\frac{\vec{D}_{1t}}{\epsilon_0 \epsilon_{r1}} = \frac{\vec{D}_{2t}}{\epsilon_0 \epsilon_{r2}}$$

or

$$\frac{\vec{D}_{1t}}{\epsilon_1} = \frac{\vec{D}_{2t}}{\epsilon_2}$$



Contribution due to side vanishes $\therefore \Delta h \rightarrow 0$

$$\Delta Q = \rho_s \Delta A = \vec{D}_{in} \Delta A - \vec{D}_{2n} \Delta A$$

ρ_s = Free charge density at the boundary

$$\rho_s = \vec{D}_{in} - \vec{D}_{2n}$$

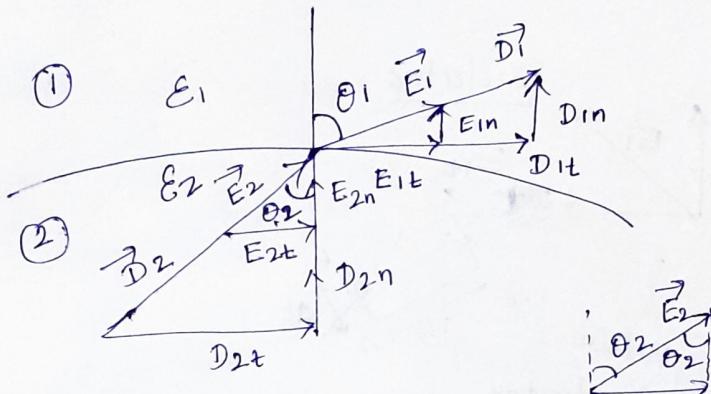
If no ~~charge~~ free charges exist at the interface

then

$$0 = \vec{D}_{in} - \vec{D}_{2n}$$

$$\vec{D}_{in} = \vec{D}_{2n}$$

$$\epsilon_1 \vec{E}_{in} = \epsilon_2 \vec{E}_{2n}$$



$$E_{1t} = E_1 \sin \theta_1$$

$$E_{1t} = E_{2t}$$

$$E_{2t} = E_2 \sin \theta_2$$

$$E_1 \sin \theta_1 = E_2 \sin \theta_2$$

$$D_{in} = D_{2n}$$

$$E_1 E_1 \cos \theta_1 = E_2 E_2 \cos \theta_2$$

$$\frac{E_1 \sin \theta_1}{E_1 E_1 \cos \theta_1} = \frac{E_2 \sin \theta_2}{E_2 E_2 \cos \theta_2}$$

$$\frac{\tan \theta_1}{E_1} = \frac{\tan \theta_2}{E_2}$$

$$\frac{\tan \theta_1}{\tan \theta_2} = \frac{E_1}{E_2}$$

[Law of Refraction
of Electric field
at a boundary
free of charge.]

assumed at the interface

Conductor - Dielectric Boundary conditions

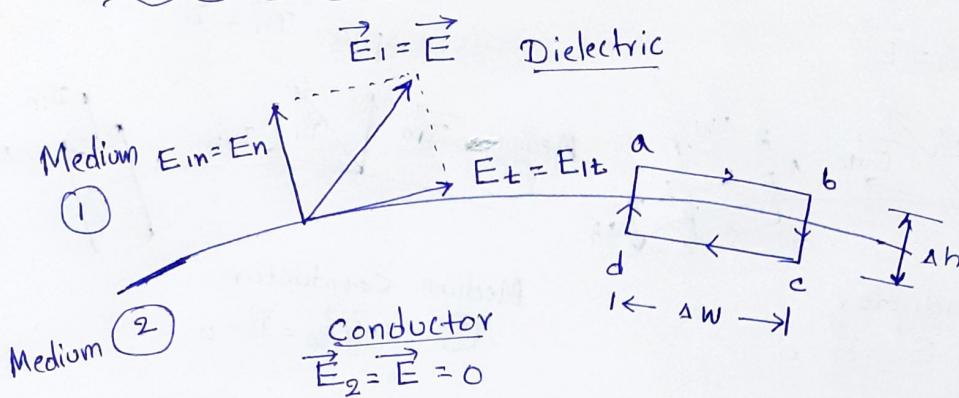


Figure shows the case of conductor-dielectric boundary conditions. The conductor is assumed to be a perfect conductor (i.e. $\omega \rightarrow \infty$ or $\sigma_c = \infty$). To find the boundary conditions for a conductor-dielectric interface.

In (2), Let the electric field inside the conductor $\vec{E}_2 = \vec{E} = 0$, $E_{2t} = 0$, $E_{2n} = 0$

In (1), the electric field outside the conductor i.e. in dielectric $\vec{E}_1 = \vec{E}$

In (1), The tangential component of $\vec{E}_1 = \vec{E}_{1t} = \vec{E}_t$

In (1) The normal component of $\vec{E}_1 = \vec{E}_{1n} = \vec{E}_n$

Applying $\oint \vec{E} \cdot d\vec{l} = 0$ to abcd

$$+ E_{1t} \Delta w - E_{1n} \frac{\Delta h}{2} + 0 + E_{1n} \frac{\Delta h}{2} = 0$$

$$\therefore \boxed{E_t = 0} \quad \boxed{E_{1t} = 0}$$

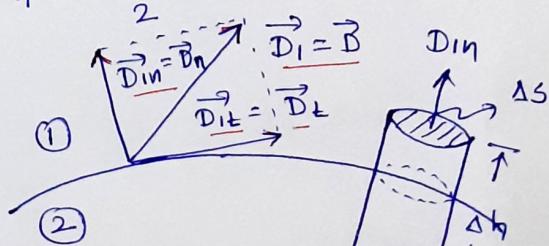
Let, inside the conductor $\vec{D}_2 = \vec{D} = 0$, $D_{2t} = 0$, $D_{2n} = 0$

In the dielectric $\vec{D}_1 = \vec{D}$

$$\vec{D}_{1t} = \vec{D}_t \quad \vec{D}_{1n} = \vec{D}_n$$

$$\therefore D_{1n} \Delta s - 0 \cdot \Delta s = \Delta Q \Rightarrow D_{1n} = \frac{\Delta Q}{\Delta s}$$

$$\therefore \boxed{D_{1n} = D_n = \sigma_s}$$



Conductor - free space Boundary conditions

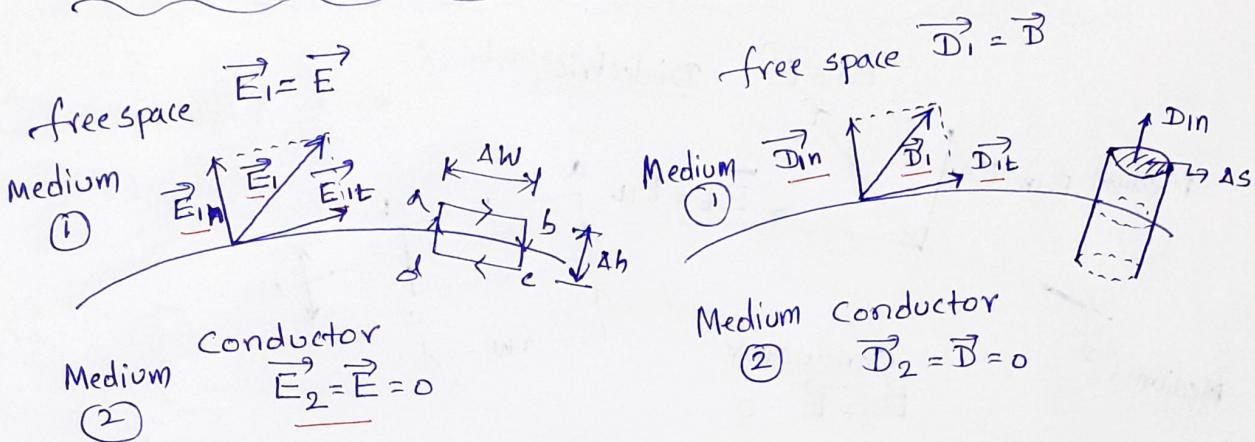


Figure shows the conductor-free space Boundary conditions

In medium ① $\vec{E}_1 = \vec{E}$, $\vec{E}_{1t} = \vec{E}_t$, $\vec{E}_{1n} = \vec{E}_n$

In medium ② $\vec{E}_2 = 0$, $\vec{E}_{2t} = 0$, $\vec{E}_{2n} = 0$

Applying $\oint \vec{E} \cdot d\vec{l} = 0$ to abcd a

$$E_{1t} \Delta W - E_{1n} \frac{\Delta h}{2} + E_{1n} \frac{\Delta h}{2} = 0 \Rightarrow \boxed{E_{1t} = 0} \quad \therefore \boxed{E_t = 0}$$

In medium ① $\vec{D}_1 = \vec{B}$

In medium ② $\vec{D}_2 = 0$

$$\vec{D}_{1t} = \vec{B}_t$$

$$\vec{D}_{2n} = 0$$

$$\vec{D}_{2t} = 0$$

$$\vec{D}_{1n} = \vec{B}_n$$

$$\Delta Q = \Delta Q$$

$$\therefore D_{1n} = \frac{\Delta Q}{AS} = D_n = S_s$$

$$\therefore D_{1n} = D_n = S_s$$

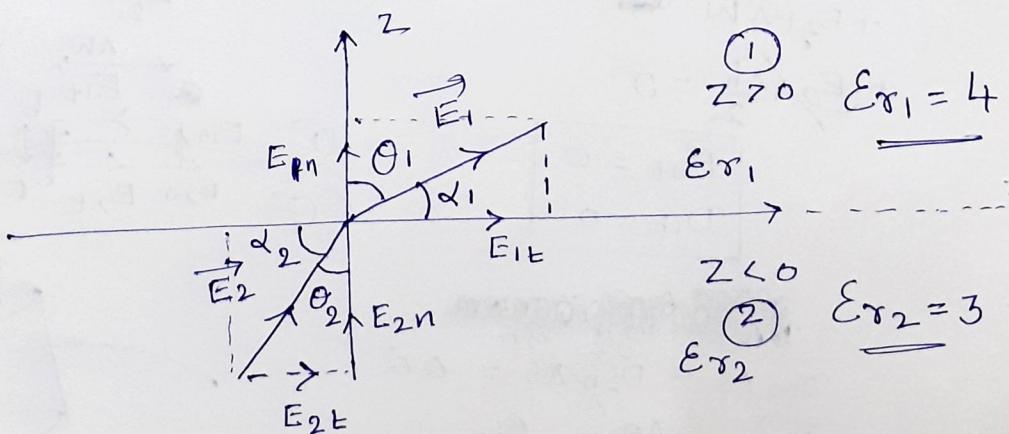
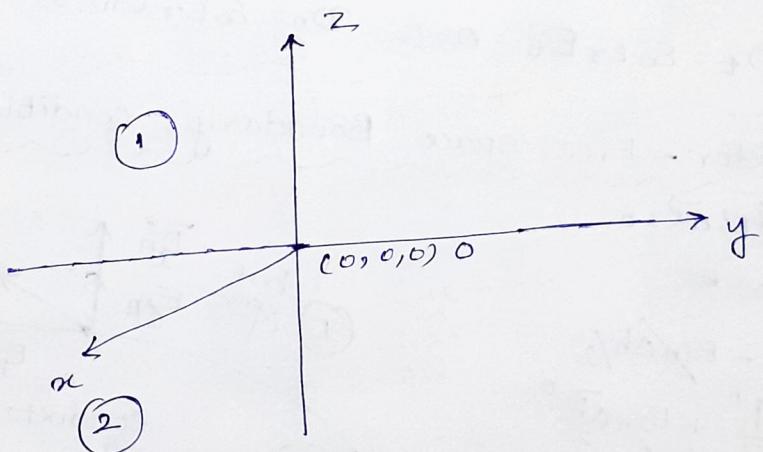
① Two extensive homogeneous isotropic dielectrics meet on plane $z=0$. For $z>0$, $\epsilon_{r1}=4$ and for $z<0$, $\epsilon_{r2}=3$. A uniform electric field $\vec{E}_1 = 5\hat{a}_x - 2\hat{a}_y + 3\hat{a}_z \text{ kV/m}$ exist for $z>0$. Find

(8+)

- (a) \vec{E}_2 for $z \leq 0$
- (b) The angles \vec{E}_1 and \vec{E}_2 make with interface
- (c) The energy densities (in J/m^3) in both dielectrics
- (d) The energy within a cube of side 2m centered at $(3, 4, -5)$.

Solution:

(a)



It is a case of dielectric to dielectric

$$\begin{aligned} \vec{E}_{1t} &= \vec{E}_{2t} & \left\{ \begin{array}{l} \vec{D}_{in} = \vec{D}_{2n} \Rightarrow \epsilon_0 \epsilon_{r1} \vec{E}_{in} \\ = \epsilon_0 \epsilon_{r2} \vec{E}_{2n} \end{array} \right. \\ \vec{E}_1 &= \vec{E}_{1t} + \vec{E}_{in} \\ \vec{E}_1 &= 5\hat{a}_x - 2\hat{a}_y \\ &\quad + 3\hat{a}_z \text{ KV/m} \end{aligned}$$

$$\begin{aligned} \vec{E}_{in} &= 3\hat{a}_z \\ \vec{E}_{1t} &= 5\hat{a}_x - 2\hat{a}_y \\ \therefore \vec{E}_{2t} &= 5\hat{a}_x - 2\hat{a}_y \rightarrow ③ \end{aligned}$$

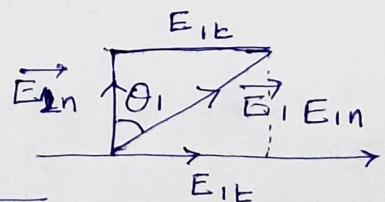
$$\begin{aligned} \cancel{\epsilon_0 \epsilon_{r2} \vec{E}_{2n}} &= \cancel{\epsilon_0 \epsilon_{r1} \vec{E}_{in}} \\ \vec{E}_{2n} &= \frac{\epsilon_{r1} \vec{E}_{in}}{\epsilon_{r2}} = \frac{4}{3} \times \cancel{3\hat{a}_z} = 4\hat{a}_z \text{ KV/m } ④ \\ \text{But } \vec{E}_2 &= \vec{E}_{2t} + \vec{E}_{2n} = 5\hat{a}_x - 2\hat{a}_y + 4\hat{a}_z \rightarrow ⑤ \\ \boxed{\vec{E}_2 = 5\hat{a}_x - 2\hat{a}_y + 4\hat{a}_z} &\text{ for } z \leq 0 \end{aligned}$$

(b) We know for dielectric - dielectric boundary condition

$$\frac{\tan \theta_1}{\tan \theta_2} = \frac{\epsilon_1}{\epsilon_2}$$

$$\frac{\tan \theta_1}{\tan \theta_2} = \frac{\epsilon_0 \epsilon_{r1}}{\epsilon_0 \epsilon_{r2}}$$

$$\tan \theta_1 = \frac{\vec{E}_{1t}}{\vec{E}_{in}}$$



$$\tan \theta_1 = \frac{|5\hat{a}_x - 2\hat{a}_y|}{|3\hat{a}_z|} = \frac{\sqrt{5^2 + 2^2}}{\sqrt{3^2}}$$

$$\tan \theta_1 = \frac{\sqrt{29}}{3} \Rightarrow \theta_1 = \tan^{-1} \left(\frac{\sqrt{29}}{3} \right)$$

$$\theta_1 = 60.9^\circ$$

$$\alpha_1 = 90 - \theta_1$$

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$$\alpha_1 = 90 - 60.9 = 22.1^\circ$$

$$\frac{\tan \theta_1}{\tan \theta_2} = \frac{\epsilon r_1}{\epsilon r_2}$$

$$\tan \theta_1 \times \frac{\epsilon r_2}{\epsilon r_1} = \tan \theta_2$$

$$\tan \theta_2 = \frac{\epsilon r_2}{\epsilon r_1} \times \tan \theta_1$$

$$= \frac{3}{4} \times \frac{\sqrt{29}}{3} = 1.3462$$

$$\theta_2 = \tan^{-1}(1.3462) = 53.4^\circ$$

$$\alpha_2 = 90 - \theta_2 = 90 - 53.4 = 36.6^\circ$$

(c) The energy densities are given by

$$W_{E1} = \frac{d|W_{E1}|}{dV} = \frac{1}{2} \epsilon_1 |E_1|^2 = \frac{1}{2} \epsilon_0 \epsilon_{r1} |E_1|^2$$

$$\frac{d W_{E1}}{dV} = \frac{1}{2} \times \frac{10^9}{36\pi} \times 4 \times |5\hat{x} - 2\hat{y} + 3\hat{z}|^2 \rightarrow \text{it is in KV}$$

$$\frac{d W_{E1}}{dV} = \frac{1}{2} \times \frac{10^9}{36\pi} \times 4 \times (\sqrt{(5^2 + 2^2 + 3^2) \times 10^6})^2$$

$$\begin{aligned} \frac{d W_{E1}}{dV} &= \frac{1}{2} \times \frac{10^9}{36\pi} \times 4 \times [25 + 4 + 9] \times 10^6 \\ &= 0.672 \times 10^{-3} \text{ J/m}^3 = 672 \mu\text{J/m}^3 \end{aligned}$$

$$W_{E2} = \frac{d|W_{E2}|}{dV} = \frac{1}{2} \epsilon_2 |E_2|^2 = \frac{1}{2} \epsilon_0 \epsilon_{r2} |E_2|^2$$

$$= \frac{1}{2} \times \frac{10^9}{36\pi} \times 3 [5^2 + 2^2 + 4^2] \times 10^6$$

$$= 597 \mu\text{J/m}^3$$

$$\therefore (d) W_E = \int_{x=2}^4 \int_{y=3}^5 \int_{z=-6}^{-4} W_{E2} dz dy dx = 597 \times 10^6 \times 2 \times 2 \times 2 = 4.776 \text{ mJ}$$

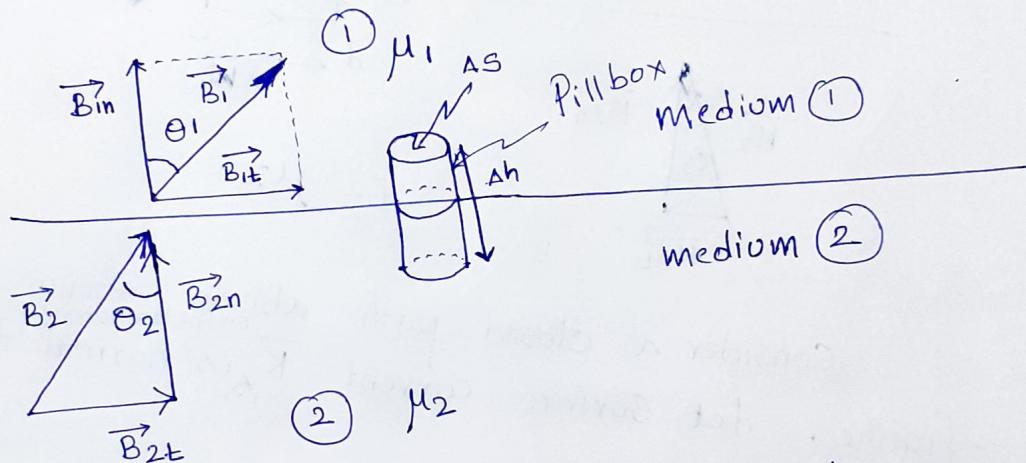
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8.7 Magnetic Boundary Conditions

We define magnetic boundary conditions as the conditions that \vec{H} (or \vec{B}) fields must satisfy at the boundary between two different media.

Here we will make use of

Magnetic Gauss's law for magnetic fields $\oint \vec{B} \cdot d\vec{s} = 0$ and Ampere's law or Ampere's circuit law $\oint \vec{H} \cdot d\vec{l} = I$



Consider the boundary between two magnetic media ① and ②, characterized by μ_1 and μ_2 respectively

Consider Gauss's law for magnetic fields

$$\oint \vec{B} \cdot d\vec{s} = 0 \rightarrow ①$$

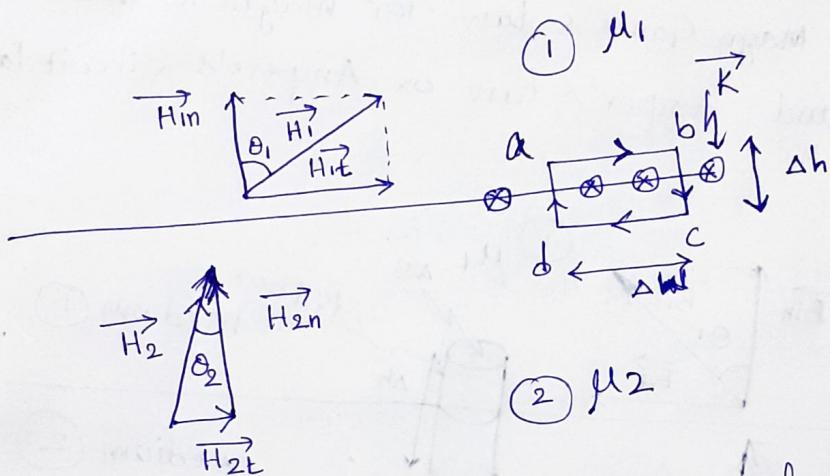
Applying equation ① to the pill box (Gaussian surface) and allowing $\Delta h \rightarrow 0$

$$B_{1n} \Delta S - B_{2n} \Delta S = 0 \Rightarrow \boxed{B_{1n} = B_{2n}} \rightarrow ②$$

$$\mu_1 H_{1n} = \mu_2 H_{2n} \rightarrow ③$$

$$\boxed{\vec{B}_{1n} = \vec{B}_{2n}} \quad \text{or} \quad \boxed{\mu_1 \vec{H}_{1n} = \mu_2 \vec{H}_{2n}} \rightarrow (3) \quad (91)$$

Equation (3) shows that the normal component of \vec{B} is continuous at the boundary. It also shows that the normal component of \vec{H} is discontinuous at the boundary. \vec{H} undergoes some changes at the interface.



Consider a closed path $abcd$ shown on the boundary. Let surface current K is normal to the surface

$$K = K \hat{n} A / m$$

$$\therefore K \Delta W = H_{1t} \Delta W + H_{1n} \frac{\Delta h}{2} + H_{2n} \frac{\Delta h}{2} - H_{2t} \Delta W - H_{2n} \frac{\Delta h}{2} - H_{1n} \frac{\Delta h}{2}$$

As $\Delta h \rightarrow 0$

$$K \Delta W = H_{1t} \Delta W - H_{2t} \Delta W$$

$$H_{1t} \Delta W - H_{2t} \Delta W = K \Delta W$$

$$H_{1t} - H_{2t} = K \rightarrow (4) \quad \frac{B_{1t}}{\mu_1} - \frac{B_{2t}}{\mu_2} = K \rightarrow (5)$$

In general case equation ④ becomes

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$$(\vec{H}_1 - \vec{H}_2) \times \hat{a}_{n12} = \vec{K}$$

Where \hat{a}_{n12} is unit vector normal to the interface and is directed from medium ① to medium ②.

If the boundary is free of current or the medium media are not conductors, then

~~$$K = 0$$~~

∴ Equation ④ becomes

$$H_{1t} - H_{2t} = K$$

$$H_{1t} - H_{2t} = 0$$

$$\vec{H}_{1t} = \vec{H}_{2t}$$

$$\frac{\vec{B}_{1t}}{\mu_1} = \frac{\vec{B}_{2t}}{\mu_2} \rightarrow ⑥$$

Thus the tangential component of \vec{H} is continuous while that of \vec{B} is discontinuous at the boundary if the fields make an angle θ with the normal to the interface.

To the interface

$$B_{1n} = B_1 \cos \theta_1 = B_{2n} = B_2 \cos \theta_2$$

$$B_{1n} = B_{2n}$$

$$B_1 \cos \theta_1 = B_2 \cos \theta_2 \rightarrow ⑦$$

$$H_{1t} = H_{2t}$$

~~$$\frac{B_{1n}}{\mu_1} \sin \theta_1 = \frac{B_{2n}}{\mu_2} \sin \theta_2$$~~

$$H_1 \sin \theta_1 = H_2 \sin \theta_2$$

$$\frac{B_1}{\mu_1} \sin \theta_1 = \frac{B_2}{\mu_2} \sin \theta_2 \rightarrow ⑧$$

\Rightarrow

~~⑦~~

~~⑧~~

⑧
⑦

$$\frac{\frac{B_1}{\mu_1} \sin \theta_1}{B_1 \cos \theta_1} = \frac{\frac{B_2}{\mu_2} \sin \theta_2}{B_2 \cos \theta_2}$$

⑨3

$$\frac{\tan \theta_1}{\tan \theta_2} = \frac{\mu_1}{\mu_2}$$

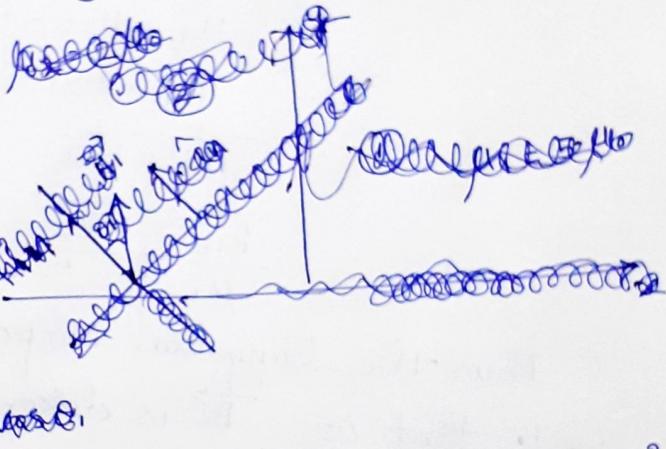
8.8) Example 8.8

Given that $\vec{H}_1 = -2\hat{a}_x + 6\hat{a}_y + 4\hat{a}_z \text{ A/m}$ in region

$y - x - 2 \leq 0$, where $\mu_1 = 5\mu_0$ calculate

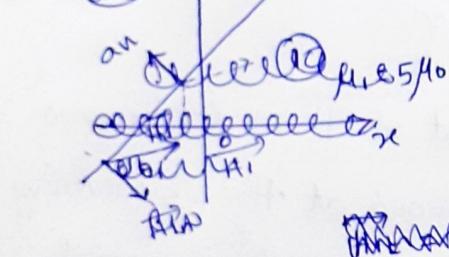
\vec{H}_2 and \vec{B}_2 in region $y - x - 2 \geq 0$, where $\mu_2 = 2\mu_0$

Solution :



Medium 1

Medium 2



$$\begin{aligned} \mu_2 &= 2\mu_0 \\ \text{Medium 2} & \quad \text{Medium 1} \\ a_n & \quad a_n \\ -2 & \quad \mu_1 = 5\mu_0 \end{aligned}$$

$$H_{in} = H_1 \cos \theta_1$$

$$\vec{H}_{in} = (\vec{H}_1 \cdot \hat{a}_n) \hat{a}_n$$

$$\vec{H}_{in} = \left[(-2, 6, 4) \cdot \left(-\frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}}, 0 \right) \right] \left(\frac{-1}{\sqrt{2}}, \frac{1}{\sqrt{2}}, 0 \right)$$

$$\vec{H}_{in} = \left[\frac{+2+6}{\sqrt{2}} \right] \left[\left(\frac{-1}{\sqrt{2}}, \frac{1}{\sqrt{2}}, 0 \right) \right]$$

$$\vec{H}_{in} = -4\hat{a}_x + 4\hat{a}_y$$

$$\vec{H}_1 = \vec{H}_{in} + \vec{H}_{it}$$

$$\vec{H}_{it} = \vec{H}_1 - \vec{H}_{in} = -2\hat{a}_x + 6\hat{a}_y + 4\hat{a}_z - (-4\hat{a}_x + 4\hat{a}_y)$$

$$\vec{H}_{it} = 2\hat{a}_x + 2\hat{a}_y + 4\hat{a}_z \text{ A/m}$$

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$$\vec{B}_{1n} = \vec{B}_{2n}$$

$$\mu_1 \vec{H}_{1n} = \mu_2 \vec{H}_{2n}$$

$$\vec{H}_{2n} = \frac{\mu_1}{\mu_2} \vec{H}_{1n} = \frac{5}{2} (-4\hat{a}_x + 4\hat{a}_y)$$

$$\vec{H}_{2n} = -10\hat{a}_x + 10\hat{a}_y$$

$$\vec{H}_{1t} = \vec{H}_{2t}$$

$$\vec{H}_{2t} = 2\hat{a}_x + 2\hat{a}_y + 4\hat{a}_z \text{ A/m}$$

$$\vec{H}_2 = \vec{H}_{2t} + \vec{H}_{2n}$$

$$\vec{H}_2 = -8\hat{a}_x + 12\hat{a}_y + 4\hat{a}_z \text{ A/m}$$

$$\begin{aligned} \vec{B}_2 &= \mu_2 \vec{H}_2 = \mu_0 \mu_r H_2 \\ \vec{B}_2 &= -20.11\hat{a}_x \\ &\quad + 30.16\hat{a}_y + 10.05\hat{a}_z \end{aligned}$$

8.9) Example 8.9: The xy -plane serves as the interface between two different media. Medium 1 ($z < 0$) is filled with a material whose $\mu_{r1} = 6$ and medium 2 ($z > 0$) is filled with a material whose $\mu_{r2} = 4$. If the interface carries current $(\frac{1}{\mu_0}) \hat{a}_y \text{ mA/m}$, and $\vec{B}_2 = 5\hat{a}_x + 8\hat{a}_z \text{ mWb/m}^2$ find \vec{H}_1 and \vec{B}_1 .

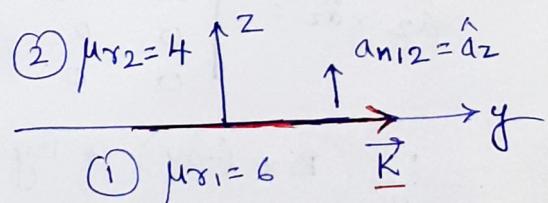
Solution!

$$B_{1n} = B_{2n}$$

$$8\hat{a}_z = 8\hat{a}_z$$

$$\therefore \boxed{B_z = 8}$$

$$\boxed{B_{1n} = 8\hat{a}_z} \text{ (Normal component)}$$



$$\vec{B}_2 = (5\hat{a}_x + 8\hat{a}_z) \text{ mA/m}$$

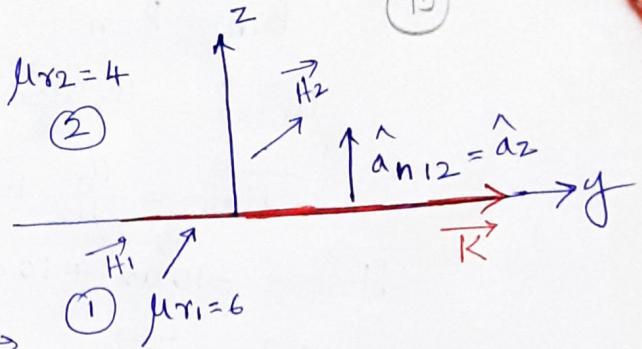
$$\vec{H}_2 = \frac{\vec{B}_2}{\mu_2} = \frac{(5\hat{a}_x + 8\hat{a}_z)}{4\mu_0} \text{ mA/m} = \frac{1}{4\mu_0} (5\hat{a}_x + 8\hat{a}_z) \text{ mA/m}$$

$$\vec{H}_1 = \frac{\vec{B}_1}{\mu_1} = \frac{1}{6\mu_0} (B_x \hat{a}_x + B_y \hat{a}_y + B_z \hat{a}_z) \text{ mA/m}$$

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Tangential components
can be found from the following

$$(\vec{H}_1 - \vec{H}_2) \times \hat{a}_{n12} = \vec{K}$$



$$\vec{H}_1 \times \hat{a}_{n12} = \vec{H}_2 \hat{a}_{n12} + \vec{K}$$

$$\frac{1}{6\mu_0} (B_x \hat{a}_x + B_y \hat{a}_y + B_z \hat{a}_z) \times \hat{a}_z = \frac{1}{4\mu_0} (5 \hat{a}_x + 8 \hat{a}_z) \times \hat{a}_z + \frac{\hat{a}_y}{\mu_0}$$

$$\frac{\mu_r=6}{6\mu_0 \hat{a}_z} \quad \left| \begin{array}{ccc} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{array} \right| \quad \text{or} \quad \left| \begin{array}{ccc} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{array} \right| \quad \text{or} \quad \left| \begin{array}{ccc} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ 0 & 0 & 1 \\ 0 & 0 & 1 \end{array} \right|$$

$$\hat{a}_x \times \hat{a}_z = \left| \begin{array}{ccc} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{array} \right| = -\hat{a}_y$$

$$\hat{a}_y \times \hat{a}_z = \left| \begin{array}{ccc} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{array} \right| = \hat{a}_x$$

$$\hat{a}_z \times \hat{a}_z = \left| \begin{array}{ccc} \hat{a}_x & \hat{a}_y & \hat{a}_z \\ 0 & 0 & 1 \\ 0 & 0 & 1 \end{array} \right| = 0$$

$$\frac{1}{6\mu_0} (B_x (-\hat{a}_y) + B_y \hat{a}_x + B_z (0)) = \frac{1}{4\mu_0} [5(-\hat{a}_y) + 8(0)] + \frac{\hat{a}_y}{\mu_0}$$

Equating the components

$$-\frac{B_x}{6} = -\frac{5}{4} + 1 \Rightarrow -\frac{B_x}{6} = \frac{1}{4}$$

$$B_x = \frac{6}{4} = 1.5$$

$$\boxed{B_y = 0}$$

$$\boxed{B_z = 0}$$

$$\therefore \boxed{\vec{B}_{1T} = 1.5 \hat{a}_x + 0 \hat{a}_y, \vec{B}_{in} = 8 \hat{a}_z \text{ mWb/m}^2}$$

$$\therefore \boxed{\vec{B}_1 = \vec{B}_{1T} + \vec{B}_{in} = (1.5 \hat{a}_x + 8 \hat{a}_z) \text{ mWb/m}^2}$$

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$$\vec{H}_1 = \frac{\vec{B}_1}{\mu_0} = -\frac{1}{\mu_0} \text{ (Ansatz)}$$

$$\vec{H}_1 = \frac{(1.5\hat{a}_x + 8\hat{a}_z)}{6\mu_0} \text{ m A/m}$$

$$\vec{H}_1 = \frac{1}{\mu_0} (0.25\hat{a}_x + 1.33\hat{a}_z) \text{ m A/m}$$

$$\vec{H}_2 = \frac{\vec{B}_2}{\mu_0 \mu_{r2}} = \frac{(5\hat{a}_x + 8\hat{a}_z)}{4\mu_0} \text{ m}$$

$$\vec{H}_2 = \frac{1}{\mu_0} (1.25\hat{a}_x + 2\hat{a}_z) \text{ m A/m}$$

$$\vec{H}_2 = \frac{1}{\mu_0} (1.25\hat{a}_x + 2\hat{a}_z) \text{ m A/m}$$

$$= (994.71 \times 10^3 \hat{a}_x + 1.5915 \times 10^6 \hat{a}_z) \text{ m A/m}$$

$$\vec{H}_2 = (994.71 \hat{a}_x + 1591.5 \hat{a}_z) \text{ A/m}$$

$$\boxed{\vec{H}_1 = 198.94 \hat{a}_x + 1068.8 \hat{a}_z} \quad \text{A/m}$$

8.8) Given that $\vec{H} = -2\hat{a}_x + 6\hat{a}_y + 4\hat{a}_z$ A/m in
 region $y - x - 2 \leq 0$, where $\mu_r = 5/4$
 calculate M_1 and B_1

solution:

$$M_1 = \chi_m H_1$$

$$M_1 = (\mu_r - 1) \vec{H}_1$$

$$M_1 = [5 - 1] \vec{H}_1$$

$$(\chi_m + 1) = \mu_r$$

$$\chi_m = \mu_r - 1$$

$$M_1 = 4 [-2\hat{a}_x + 6\hat{a}_y + 4\hat{a}_z]$$

$$M_1 = -8\hat{a}_x + 24\hat{a}_y + 16\hat{a}_z$$