

1. (10 points) Fix some  $x_0 \in \mathbb{R}^n$  and  $t_0 > 0$ . For each  $s \in [0, t_0]$  denote by

$$\Omega_s = \{(x, t) : 0 \leq t \leq s, |x - x_0| \leq c(t_0 - t)\}$$

as well as its side boundary

$$\mathcal{S}_{side} = \{(x, t) : 0 \leq t \leq s, |x - x_0| = c(t_0 - t)\}.$$

Let  $\nu = (\nu_1, \dots, \nu_n, \nu_{n+1})$  be the unit normal vector on  $\mathcal{S}_{side}$ . Show that  $\nu_{n+1}^2 = c^2(\nu_1^2 + \dots + \nu_n^2)$ . (This is an exercise left during the lecture.)

Consider  $F(x, t) = |x - x_0| - c(t_0 - t)$

Let  $\hat{\nu}_x = \{\nu_i\}_{i=1}^n = \left\{ \frac{\partial F}{\partial x^i} \right\}_{i=1}^n$  note  $\left\| \frac{\partial F}{\partial x^{(i)}} \right\| = \left\| \frac{x^i - x_0^i}{|x - x_0|} \right\| = 1$

so  $\|\hat{\nu}_x\|^2 = n$

Let  $\nu_x = \frac{1}{\sqrt{n}} \hat{\nu}_x$  so that  $\|\nu_x\| = 1$ . It follows

$$\nu = \frac{(\nu_x, \nu_{n+1})}{\|(\nu_x, \nu_{n+1})\|} = \frac{1}{\sqrt{1+c^2}} \left( \left\{ \frac{x^i - x_0^i}{|x - x_0|} \right\}_{i=1}^n, c \right)$$

$$\nu_{n+1}^2 = \frac{c^2}{1+c^2} \quad \sum_{i=1}^n \nu_i^2 = \frac{1}{1+c^2} \|\nu_x\|^2 = \frac{1}{1+c^2}$$

$$\text{Thus } c^2 \sum_{i=1}^n \nu_i^2 = \frac{c^2}{1+c^2} = \nu_{n+1}^2$$

2. (15 points) Let  $u$  solve the initial-value problem for the wave equation in one dimension:

$$\begin{cases} u_{tt} - u_{xx} = 0 & \text{in } \mathbb{R} \times (0, \infty), \\ u = g, \quad u_t = h & \text{on } \mathbb{R} \times \{t = 0\}. \end{cases} \quad (1)$$

Suppose  $g, h$  have compact support. The kinetic energy is  $k(t) := \frac{1}{2} \int_{\mathbb{R}} u_t^2(x, t) dx$  and the potential energy is  $p(t) := \frac{1}{2} \int_{\mathbb{R}} u_x^2(x, t) dx$ . Prove  $k(t) = p(t)$  for all large enough time  $t$ .

$$u(x, t) = \frac{1}{2}(g(x+t) + g(x-t)) + \frac{1}{2} \int_{x-t}^{x+t} h(y) dy$$

$$u_x(x, t) = \frac{1}{2}(g'(x+t) + g'(x-t)) + \frac{1}{2} \frac{d}{dx} \int_{x-t}^{x+t} h(y) dy$$

$$= \frac{1}{2}(g'(x+t) + g'(x-t)) + \frac{1}{2}(h(x+t) - h(x-t))$$

$$u_t(x, t) = \frac{1}{2}(g'(x+t) - g'(x-t)) + \frac{1}{2} \frac{d}{dt} \int_{x-t}^{x+t} h(y) dy$$

$$= \frac{1}{2}(g'(x+t) - g'(x-t)) + \frac{1}{2}(h(x+t) + h(x-t))$$

Since  $g, h$  have compact support  $\exists G, H \subset \mathbb{R} \ni g, h = 0$  in  $G^c, H^c$  and  $G, H$  are compact. Also  $\text{supp}(g') \subseteq G$  which is compact.

$$\begin{aligned} \frac{d}{dt} E(t) &= \frac{d}{dt} \int u_x^2 + u_t^2 dx = \int u_x u_{x,t} + u_t u_{t,t} dx \\ &= \dots \end{aligned}$$

$$\frac{d}{dt} \int_{-\infty}^{\infty} (u_t^2 + u_x^2 + u^3) dx = \int_{-\infty}^{\infty} (2u_t u_{tt} + 2u_x u_{xt} + 3u^2 u_t) dx$$

$$= \int_{-\infty}^{\infty} (u_x u_{xt} + u_t u_{xx}) dx$$

$$= \int_{-\infty}^{\infty} u_x u_{xt} dx + [u_t u_x]_{-\infty}^{\infty} - \int_{-\infty}^{\infty} u_{tx} u_x dx$$

$$= [u_t u_x]_{-\infty}^{\infty}$$

$$= 0 \quad \text{since } u \text{ and } g \text{ have compact support}$$

$$\therefore E(0) = E(t) \quad (\text{constant in time})$$

Since energy is conserved we can choose a large  $t$  so that  $u(x,t)$  is separated into two waves centered around  $x-t$  and  $x+t$ .

$$\text{For the } x-t \text{ we see } u_x = \frac{1}{2}(g'(x-t) - h(x-t))$$

$$u_t = \frac{1}{2}(-g'(x-t) + h(x-t))$$

$$\Rightarrow u_x = -u_t$$

$$\text{For the } x+t \text{ we see } u_x = \frac{1}{2}(g'(x+t) + h(x+t))$$

$$u_t = \frac{1}{2}(g'(x+t) + h(x+t))$$

$$\Rightarrow u_x = u_t$$

In both cases  $u_x^2 = u_t^2 \Rightarrow \rho(t) = k(t)$  for large enough  $t$

3. (15 points) Use energy method to show the following (nonlinear) wave equation only has a zero solution

$$\begin{cases} u_{tt} - \Delta u + u + u^3 = 0 & x \in \Omega, 0 < t < T; \\ u(x, 0) = u_t(x, 0) = 0 & x \in \Omega; \\ u(x, t) = 0 & x \in \partial\Omega, 0 < t < T \end{cases}$$

where  $\Omega \subset \mathbb{R}^n$  is an open smooth bounded domain. (Hint: Recall that for our original wave equation  $u_{tt} - \Delta u = 0$ , the energy is defined as  $\frac{1}{2} \int_{\Omega} u_t^2 + |\nabla u|^2 dx$ . Think about why the energy is defined in that way, and here you need to construct a new suitable energy  $E(t)$  for our nonlinear wave equation and prove it is zero at  $t = 0$  and  $\frac{d}{dt} E(t) = 0$ .)

$$E(t) = \frac{1}{2} \int_{\Omega} u_t^2 + |\nabla u|^2 + u^2 + \frac{1}{2} u^4 dx$$

$$\frac{d}{dt} E(t) = \int_{\Omega} u_t u_{tt} + \nabla u \cdot \nabla u_t + u u_t + u^3 u_t dx$$

$$= \int_{\Omega} u_t (\Delta u - u - u^3) + \nabla u \cdot \nabla u_t + u u_t + u^3 u_t dx$$

$$= \int_{\Omega} u_t \Delta u + \nabla u \cdot \nabla u_t dx$$

$$\begin{aligned}
&= \int_{\Omega} u_t + \Delta u + \nabla u \cdot \nabla u + d\Omega \\
&= \int_{\partial\Omega} u_t \cdot \hat{n} dS - \int_{\Omega} \nabla u_t \cdot \nabla u d\Omega + \int_{\Omega} \nabla u \cdot \nabla u d\Omega \\
&= 0
\end{aligned}$$

Thus  $E(t) = E(0) = \frac{1}{2} \int_{\Omega} 0 + 0 + 0 + 0 d\Omega = 0$

$\therefore u(x, t) = 0$

4. (20 points) Let  $Q = \Omega \times (0, T]$  and  $u \in C^{2,1}(Q) \cap C(\bar{Q})$  satisfy

$$u_t - \Delta u + cu \leq 0 \quad \text{in } Q$$

where  $c \geq 0$  is a constant.

- If  $u \geq 0$ , show that the weak maximum principle holds for  $u$ , i.e.,  $u$  achieves its maximum at the parabolic boundary of  $Q$ . (Hint: the proof is similar to the one we did during the lecture.)
- Give a counterexample to show the weak maximum principle may not hold without the condition  $u \geq 0$ . (Hint: consider  $\Omega = (-1, 1)$  and consider  $u = -x^2 + at - b$  with some suitable  $a$  and  $b$ .)

a) Assume at some point in the interior of  $Q$ ,  $(x_0, t_0)$  a maximum is achieved, say  $u(x_0, t_0) = M$ . Also assume  $u_t - \Delta u + cu < 0$  in  $Q$ . At  $(x_0, t_0)$  since it is a max,  $u_t(x_0, t_0) \geq 0$  and  $\Delta u(x_0, t_0) \leq 0$ . It follows

$$0 \leq u_t(x_0, t_0) - \Delta u(x_0, t_0) < -cM$$

$$\Rightarrow cM < 0$$

\* Since  $u \geq 0$  and  $c \geq 0 \therefore (x_0, t_0) \in \Gamma$

Next consider  $u_t - \Delta u + cu \leq 0$  in  $Q$ . It follows  $\forall \varepsilon > 0$

$$v_\varepsilon(x, t) = u(x, t) - \varepsilon t$$

$$v_\varepsilon \in C^{2,1}(Q) \cap C(\bar{Q}) \quad \partial_t v_\varepsilon - \Delta v_\varepsilon + c v_\varepsilon = u_t - \Delta u - \varepsilon + cu - c\varepsilon t < 0$$

$$\max_Q u = \max_Q (v_\varepsilon + \varepsilon t)$$

$$\leq \max_Q v_\varepsilon + \varepsilon T$$

$$\leq \max_\Gamma v_\varepsilon + \varepsilon T$$

$$\leq \max_\Gamma u + \varepsilon T$$

By applying the previous case

as  $\varepsilon \rightarrow 0 \quad \max_Q u \leq \max_\Gamma u$

Also, as  $\Gamma \subseteq \bar{Q}$   $\max_{\Gamma} u \leq \max_{\bar{Q}} u$

$$\therefore \max_{\Gamma} u = \max_{\bar{Q}} u$$

b)  $-x^2 + at - b < 0$   $\Omega = (-1, 1)$

$$u_t = a \quad \Delta u = -2 \quad c = 1$$

$$a + 2 - x^2 + at - b \leq 0$$

$$-x^2 + (1+t)a - b \leq 0 \quad b - at + a + at - b < 0$$

$$-x^2 + at - b < 0 \quad a < 0$$

$$-x^2 + at - b < 0$$

$$t > \frac{b}{a}$$

$$b > 0$$

$$-x^2 - t - 1 < 0 \quad \forall x \in \mathbb{R} \quad t > 0$$

max is at  $(0, t)$   $0 \in (-1, 1)$   
 $0 \notin \{-1, 1\}$

I tried to draw a picture

