

1. (15 points) For some fixed $a > 0$, compute the following.

- (a) The Fourier transform of $f(x) = e^{-a|x|}$ where $x \in \mathbb{R}$.
- (b) The Fourier transform of $f(x) = e^{-a|x|}$ where $x \in \mathbb{R}^n$.

(Hint: For part (b), you can use the result from part (a), the trivial identity $\int_0^\infty e^{-(a^2+\xi^2)s} ds = \frac{1}{a^2+\xi^2}$, and the Fourier transform of $e^{-b|x|^2}$ we discussed during lectures to derive an integral representation of $e^{-a|x|}$ in terms of $e^{-c|x|^2}$ for some c , and then continue from there.)

a) Note $f(x) = f(-x)$ so

$$\hat{f}(\xi) = \frac{1}{\sqrt{2\pi}} \int_0^\infty e^{-ax} (\cos \xi x) dx$$

$$\begin{aligned} u &\quad \checkmark \\ (\cos \xi x) &+ e^{-ax} \\ -\xi \sin \xi x &- \frac{1}{a} e^{-ax} \\ -\xi^2 \cos \xi x &- \frac{1}{a^2} e^{-ax} \end{aligned}$$

$$\int u dv = uv - \int v du$$

$$\int_0^\infty e^{-ax} \cos \xi x dx = -\frac{1}{a} \cos \xi x e^{-ax} \Big|_0^\infty + \frac{\xi}{a} \sin \xi x e^{-ax} \Big|_0^\infty + \frac{\xi^2}{a^2} \int_0^\infty \cos \xi x e^{-ax} dx$$

$$(1 + \frac{\xi^2}{a^2}) \int_0^\infty e^{-ax} \cos \xi x dx = \frac{1}{a} + 0$$

$$\int_0^\infty e^{-ax} \cos \xi x dx = \frac{1}{a} \frac{a^2}{\xi^2 + a^2} = \frac{a}{\xi^2 + a^2}$$

$$\hat{f}(\xi) = \sqrt{\frac{a}{\xi^2 + a^2}}$$

b) $\widehat{e^{-c|x|^2}} = (2c)^{-n/2} e^{-\frac{|\xi|^2}{4c}}$ $\hat{f}(\xi) = (2\pi)^{-n/2} \int_{\mathbb{R}^n} f(x) e^{-ix^T \xi} dx$

$$\hat{f}(\xi) = (2\pi)^{-n/2} \int_{\mathbb{R}^n} e^{-(ax^2 + i x^T \xi)} dx$$

$$\text{Let } r = |x| \Rightarrow dx = r^{n-1} dr dS^{n-1} \quad x^T \xi = |x| |\xi| \cos \theta$$

$$\text{Let } \xi = |\xi| \nu$$

$$\hat{f}(\xi) = (2\pi)^{-n/2} \int_0^\infty \int_{S^{n-1}} e^{-(ar + i x^T \xi \cos \theta)} r^{n-1} dr dS^{n-1}$$

$$= (2\pi)^{-n/2} \int_0^\infty (2\pi)^{n/2} \frac{J_{\frac{n-2}{2}}(r|\xi|)}{(r|\xi|)^{\frac{n-2}{2}}} e^{-ar} r^{n-1} dr$$

$$= \int_0^\infty e^{-ar} \frac{J_{\frac{n-2}{2}}(r|\xi|)}{|\xi|^{\frac{n-2}{2}}} r^{\frac{n-2}{2}+1} dr$$

$$= \frac{2^{\frac{n+1}{2}}}{\pi(a^2 + |\xi|^2)^{\frac{n+1}{2}}} \Gamma\left(\frac{n+1}{2}\right)$$

J is the Bessel function

These are known integrals from integral tables

$$\sqrt{(\alpha^2 + |\zeta|^2)^{\frac{n+1}{2}}} \cdot (-2)$$

These are known integrals from integral tables I found.

2. (15 points) Let $K(x, y, t)$ be the heat kernel given as

$$K(x, y, t) = \frac{1}{(4\pi t)^{\frac{n}{2}}} e^{-\frac{|x-y|^2}{4t}}, \quad x, y \in \mathbb{R}^n, t > 0$$

- (a) Show that $K(x, y, t) \in C^\infty(\mathbb{R}^n \times \mathbb{R}^n \times (0, \infty))$, and show it satisfies the heat equation

$$\frac{\partial K(x, y, t)}{\partial t} - \Delta_x K(x, y, t) = 0.$$

- (b) Show that for every $x \in \mathbb{R}^n$ and $t > 0$,

$$\int_{\mathbb{R}^n} K(x, y, t) dy = 1.$$

$$a) K(x, y, t) = (4\pi t)^{-n/2} e^{-\frac{|x-y|^2}{4t}} \quad \Phi(z, t) = (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}}$$

$$x-y = z.$$

$$\text{Let } 1 \leq i \leq j \leq n \quad \frac{\partial \Phi}{\partial z_i} = \frac{z_i}{2t} (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}} \in C(\mathbb{R}^n)$$

$$\frac{\partial^2 \Phi}{\partial z_i \partial z_j} = \frac{z_i z_j}{4t^2} (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}} \in C(\mathbb{R}^n)$$

In general $D^\alpha \Phi = \alpha_\alpha(z) e^{-\frac{|z|^2}{4t}}$ where α_α is a function of polynomials in z . As the product of continuous functions, $D^\alpha \Phi$ is continuous.

$$\frac{\partial}{\partial z_i} D^\alpha \Phi = -\frac{z_i}{2t} \alpha_\alpha(z) e^{-\frac{|z|^2}{4t}} + \alpha'_\alpha(z) e^{-\frac{|z|^2}{4t}} \in C(\mathbb{R}^n).$$

Thus by induction $\Phi \in C^\infty(\mathbb{R}^n)$ in space.

$$\text{Let } C = (4\pi)^{-n/2}$$

$$\frac{\partial \Phi}{\partial t} = -\frac{n}{2} C \frac{1}{t^{\frac{n}{2}+1}} e^{-\frac{|z|^2}{4t}} + \frac{C |z|^2}{4t^{n+2}} e^{-\frac{|z|^2}{4t}}$$

$$-\frac{n}{2t} \Phi + \frac{|z|^2}{4t^2} \Phi \in C(\mathbb{R}_+ \setminus \{0\})$$

Similarly to the spacial variables, we will have a polynomial in $1/t$ for higher derivatives. Thus by the same space argument, $\Phi \in C^\infty(\mathbb{R}_+ \setminus \{0\})$ in time.

$$\frac{\partial \Phi}{\partial z_i} = \frac{z_i}{2t} (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}} = -\frac{z_i}{2t} \Phi$$

$$\frac{\partial^2 \Phi}{\partial z_i \partial z_j} = -\frac{z_i z_j}{2t^2} (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}} + \frac{1}{2t} \Phi$$

$$\begin{aligned}\frac{\partial \Phi}{\partial z_i} &= -\frac{z_i}{2t} \left(-\frac{z_i}{2t} \Phi \right) + -\frac{1}{2t} \Phi \\ &= \left(\frac{|z|^2}{4t^2} - \frac{1}{2t} \right) \Phi\end{aligned}$$

$$\begin{aligned}\frac{\partial \Phi}{\partial t} - \Delta \times \Phi &= \left(\frac{n}{2t} - \frac{|z|^2}{4t^2} \right) \Phi - \sum_i \frac{\partial^2 \Phi}{\partial z_i^2} \\ &= \left(\frac{n}{2t} + \frac{|z|^2}{4t^2} \right) \Phi - \sum_i \left(\frac{|z|^2}{4t^2} - \frac{1}{2t} \right) \Phi \\ &= 0\end{aligned}$$

b) WTS $\int_{\mathbb{R}^n} \Phi(z, t) dz = 1 \quad \forall z \in \mathbb{R}^n$

$$(4\pi t)^{-n/2} \int_{\mathbb{R}^n} e^{-|z|^2/(4t)} dz = (4\pi t)^{-n/2} \int_{-\infty}^{\infty} e^{-z_1^2/(4t)} dz_1 \int_{\mathbb{R}^{n-1}} e^{-|z'|^2/(4t)} dz'$$

$$(z' = (z_2, \dots, z_n)) \rightarrow = (4\pi t)^{-(n-1)/2} \int_{\mathbb{R}^{n-1}} e^{-|z'|^2/(4t)} dz'$$

Note $\int_{-\infty}^{\infty} e^{-x_1^2/(4t)} dx_1 = \sqrt{4\pi t}$...

Integral table 5 = 1

3. (20 points) Suppose $g \in L^1(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$ and set

$$u(x, t) = \int_{\mathbb{R}^n} K(x, y, t) g(y) dy$$

where $K(x, y, t)$ denotes the heat kernel as defined in problem 2. Show that for all $t > 0$:

- (a) $\|u(\cdot, t)\|_{L^\infty(\mathbb{R}^n)} \leq Ct^{-n/2} \|g\|_{L^1(\mathbb{R}^n)}$
- (b) $\|u(\cdot, t)\|_{L^1(\mathbb{R}^n)} \leq \|g\|_{L^1(\mathbb{R}^n)}$
- (c) $\|u_{x_i}(\cdot, t)\|_{L^\infty(\mathbb{R}^n)} \leq Ct^{-1/2} \|g\|_{L^\infty(\mathbb{R}^n)}$.
- (d) $\|u_{x_i}(\cdot, t)\|_{L^\infty(\mathbb{R}^n)} \leq Ct^{-\frac{n+1}{2}} \|g\|_{L^1(\mathbb{R}^n)}$

a) $|u(x, t)| = |\int_{\mathbb{R}^n} K(x, y, t) g(y) dy| \leq (4\pi t)^{-n/2} \int_{\mathbb{R}^n} e^{-|x-y|^2/(4t)} |g(y)| dy$

$$\sup e^{-|x-y|^2/(4t)} = 1 \text{ when } x = y$$

$$\leq C t^{-n/2} \|g\|_{L^1}$$

Thus $\|u(x, t)\|_{L^\infty} \leq C t^{-n/2} \|g\|_{L^1}$

b) $\|u(x, t)\|_{L^1} = \int_{\mathbb{R}^n} |u(x, t)| dx = \int_{\mathbb{R}^n} \left| \int_{\mathbb{R}^n} K(x, y, t) g(y) dy \right| dx$

$$\leq \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} K(x, y, t) |g(y)| dy dx$$

$$= \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} K(x, y, t) |g(y)| dy dx$$

$$\begin{aligned} &\leq \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} K(x, y, t) |g(y)| dy dx \\ &= \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \tilde{\Phi}(x, t) dz |g(y)| dy \\ &= \|g\|_{L^1} \end{aligned}$$

c) $\frac{\partial u}{\partial x_i} = (4\pi t)^{-n/2} \int_{\mathbb{R}^n} \frac{\partial}{\partial x_i} e^{-\frac{|x-y|^2}{4t}} g(y) dy$

$$= \frac{1}{2t} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |x_i - y_i| e^{-\frac{|x_i - y_i|^2}{4t}} g(y) \int_{\mathbb{R}^{n-1}} e^{-\frac{|x' - y'|^2}{4t}} dy dy$$

$$\|u_{x_i}(x, t)\|_{L^\infty} \leq \frac{1}{2t} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |x_i - y_i| e^{-\frac{|x_i - y_i|^2}{4t}} \|g\|_\infty \int_{\mathbb{R}^{n-1}} e^{-\frac{|x' - y'|^2}{4t}} dy dy$$

$$= (2t)^{-1} (4\pi t)^{-n/2} \int_{\mathbb{R}^{n-1}} e^{-\frac{|z'|^2}{4t}} dz \int_{\mathbb{R}^n} |z_i| e^{-\frac{|z_i|^2}{4t}} dz \|g\|_\infty$$

Note $\int_{\mathbb{R}^1} 1 \times e^{-\frac{|x|^2}{4t}} dx = 2 \int_0^\infty e^{-\frac{x^2}{4t}} dx = 4t$
 By symmetry & Integral tables

$$\begin{aligned} \|u_{x_i}(x, t)\|_\infty &\leq (2t)^{-1} (4\pi t)^{-n/2} 4t (4\pi t)^{n/2} \|g\|_\infty \\ &= C t^{-1/2} \|g\|_\infty \end{aligned}$$

d) $|u_{x_i}(x, t)| \leq \frac{1}{2t} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |x_i - y_i| e^{-\frac{|x_i - y_i|^2}{4t}} |g(y)| dy \quad |x_i - y_i| \leq |x - y|$

$$\leq \frac{1}{2t} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |z| e^{-\frac{|z|^2}{4t}} |g(x - z)| dz$$

Consider $h(r) = r e^{-\frac{r^2}{4t}} \quad r = |z|$

$$h'(r) = e^{-\frac{r^2}{4t}} - r \cdot \frac{2r}{4t} e^{-\frac{r^2}{4t}} = e^{-\frac{r^2}{4t}} \left(1 - \frac{r^2}{2t}\right)$$

$$h'(r) = 0 \quad @ r = \sqrt{2t}$$

$$h(\sqrt{2t}) = \sqrt{2t} e^{-\frac{1}{2}}$$

$$\begin{aligned} |u_{x_i}(x, t)| &\leq \frac{\sqrt{2t}}{2t} e^{-\frac{1}{2}} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |g(x - z)| dz \\ &= C t^{-\frac{n+1}{2}} \|g\|_{L^1} \end{aligned}$$

4. (10 points) For a fixed constant $c \in \mathbb{R}$, find an explicit solution for

$$\begin{aligned} u_t - \Delta u + cu &= f \quad (x, t) \in \mathbb{R}^n \times (0, \infty), \\ u(x, 0) &= g(x) \quad x \in \mathbb{R}^n. \end{aligned}$$

Let $u = u_1 + u_2 \quad \exists$

$$\underline{u_1} \left\{ \begin{array}{l} u_t - \Delta u + cu = 0 \\ u(x, 0) = g(x) \end{array} \right. \quad x \in \mathbb{R}^n, t > 0 \quad \underline{u_2} \left\{ \begin{array}{l} u_t - \Delta u + cu = f \\ u(x, 0) = 0 \end{array} \right. \quad x \in \mathbb{R}^n, t > 0$$

$$\boxed{u_1} \quad u(x,0) = g(x) \quad x \in \mathbb{R} \quad t > 0 \quad \boxed{u_2} \quad u_t - \Delta u + cu = 0 \quad x \in \mathbb{R} \quad t > 0$$

Let $\hat{u} = F(u)$, $\hat{g} = F(g)$

$$\boxed{u_1} \quad \frac{\partial}{\partial t} \hat{u} - |\xi|^2 \hat{u} + c \hat{u} = 0$$

$$\frac{\partial}{\partial t} \hat{u} + (|\xi|^2 + c) \hat{u} = 0$$

$$\Rightarrow \hat{u}(\xi, t) = C e^{-(|\xi|^2 + c)t}$$

$$\hat{u}(\xi, 0) = \hat{g}(\xi) = C$$

$$\hat{u}(\xi, t) = \hat{g}(\xi) e^{-(|\xi|^2 + c)t}$$

$$F^{-1}(\hat{u})(x, t) = e^{-ct} \mathcal{F}(x, t) * g(x)$$

$$= e^{-ct} \int_{\mathbb{R}^n} (4\pi t)^{-\frac{n}{2}} e^{-\frac{|x-y|^2}{4t}} g(y) dy$$

u_2 By the Duhamel principle

$$u_2 = \int_0^t w(x, t-s) ds \text{ where}$$

$$w \text{ solves } \begin{cases} w_t - \Delta w + cw = 0 & x \in \mathbb{R}^n \quad t > 0 \\ w(x, 0) = f(x, 0) & x \in \mathbb{R}^n \end{cases}$$

$$\text{Thus } u_2 = \int_0^t e^{-c(t-s)} \int_{\mathbb{R}^n} (4\pi(t-s))^{-\frac{n}{2}} e^{-\frac{|x-y|^2}{4(t-s)}} f(y, 0) dy ds$$

$$\therefore u(x, t) = e^{-ct} \int_{\mathbb{R}^n} (4\pi t)^{-\frac{n}{2}} e^{-\frac{|x-y|^2}{4t}} g(y) dy +$$

$$\int_0^t e^{-c(t-s)} \int_{\mathbb{R}^n} (4\pi(t-s))^{-\frac{n}{2}} e^{-\frac{|x-y|^2}{4(t-s)}} f(y, 0) dy ds$$