

HW3

Tuesday, February 18, 2025 11:16 AM

1. (10 points) Fix some $x_0 \in \mathbb{R}^n$ and $t_0 > 0$. For each $s \in [0, t_0]$ denote by

$$\Omega_s = \{(x, t) : 0 \leq t \leq s, |x - x_0| \leq c(t_0 - t)\}$$

as well as its side boundary

$$\mathcal{S}_{\text{side}} = \{(x, t) : 0 \leq t \leq s, |x - x_0| = c(t_0 - t)\}.$$

Let $\nu = (\nu_1, \dots, \nu_n, \nu_{n+1})$ be the unit normal vector on $\mathcal{S}_{\text{side}}$. Show that $\nu_{n+1}^2 = c^2(\nu_1^2 + \dots + \nu_n^2)$. (This is an exercise left during the lecture.)

Consider $F(x, t) = |x - x_0| - c(t_0 - t)$
 Let $\hat{\nu}_x = \{\nu_i\}_{i=1}^n = \left\{ \frac{\partial F}{\partial x^i} \right\}_{i=1}^n$ note $\left\| \frac{\partial F}{\partial x^{n+1}} \right\| = \left\| \frac{x^i - x_0^i}{|x - x_0|} \right\| = 1$
 so $\|\hat{\nu}_x\|^2 = n$
 Let $\nu_x = \frac{1}{\sqrt{n}} \hat{\nu}_x$ so that $\|\nu_x\| = 1$. It follows
 $\nu = \frac{(\nu_x, \hat{\nu}_{n+1})}{\|(\nu_x, \hat{\nu}_{n+1})\|} = \frac{1}{\sqrt{1+c^2}} \left(\left\{ \frac{x^i - x_0^i}{|x - x_0|} \right\}_{i=1}^n, c \right)$
 $\nu_{n+1}^2 = \frac{c^2}{1+c^2} \quad \sum_{i=1}^n \nu_i^2 = \frac{1}{1+c^2} \|\nu_x\|^2 = \frac{1}{1+c^2}$
 Thus $c^2 \sum_{i=1}^n \nu_i^2 = \frac{c^2}{1+c^2} = \nu_{n+1}^2$

2. (15 points) Let u solve the initial-value problem for the wave equation in one dimension:

$$\begin{cases} u_{tt} - u_{xx} = 0 & \text{in } \mathbb{R} \times (0, \infty), \\ u = g, \quad u_t = h & \text{on } \mathbb{R} \times \{t = 0\}. \end{cases} \quad (1)$$

Suppose g, h have compact support. The kinetic energy is $k(t) := \frac{1}{2} \int_{\mathbb{R}} u_t^2(x, t) dx$ and the potential energy is $p(t) := \frac{1}{2} \int_{\mathbb{R}} u_x^2(x, t) dx$. Prove $k(t) = p(t)$ for all large enough time t .

$$\begin{aligned} u(x, t) &= \frac{1}{2}(g(x+t) + g(x-t)) + \frac{1}{2} \int_{x-t}^{x+t} h(y) dy \\ u_x(x, t) &= \frac{1}{2}(g'(x+t) + g'(x-t)) + \frac{1}{2} \frac{d}{dx} \int_{x-t}^{x+t} h(y) dy \\ &= \frac{1}{2}(g'(x+t) + g'(x-t)) + \frac{1}{2}(h(x+t) - h(x-t)) \\ u_+(x, t) &= \frac{1}{2}(g'(x+t) - g'(x-t)) + \frac{1}{2} \frac{d}{dt} \int_{x-t}^{x+t} h(y) dy \\ &= \frac{1}{2}(g'(x+t) - g'(x-t)) + \frac{1}{2}(h(x+t) + h(x-t)) \end{aligned}$$

Since g, h have compact support $\exists G, H \subset \mathbb{R}$ such that $g, h = 0$ in $\mathbb{R} \setminus G, H$ and G, H are compact.
 Also $\text{supp}(g') \subseteq G$ which is compact.

$$\begin{aligned} \frac{d}{dt} E(t) &= \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}} (u_x^2 + u_t^2) dx = \int (u_x u_{xt} + u_t u_{tt}) dx \\ &= \int_{\mathbb{R}} u_x u_{xt} dx \end{aligned}$$

$$\frac{d}{dt} \int_{\Omega} (U^2) - 2 \int_{\Omega} U_x \cdot U_t + U^2 - J \int_{\Omega} U_x U_t + U^2 =$$

$$= \int_{\Omega} U_x U_{xt} + U_t U_{xx} dx$$

$$= \int_{\Omega} U_x U_{xt} dx + [U + U_x]_{-\infty}^{\infty} - \int_{\Omega} U_{tx} U_x dx$$

$$= [U + U_x]_{-\infty}^{\infty}$$

$= 0$ since h and g have compact support

$$\therefore E(0) = E(t) \quad (\text{constant in time})$$

Since energy is conserved we can choose a large t so that $U(x,t)$ is separated into two waves centered around $x-t$ and $x+t$.

$$\text{For the } x-t \text{ we see } U_x = \frac{1}{2}(g'(x-t) - h(x-t))$$

$$U_t = \frac{1}{2}(-g'(x-t) + h(x-t))$$

$$\Rightarrow U_x = -U_t$$

$$\text{For the } x+t \text{ we see } U_x = \frac{1}{2}(g'(x+t) + h(x+t))$$

$$U_t = \frac{1}{2}(g'(x+t) + h(x+t))$$

$$\Rightarrow U_x = U_t$$

In both cases $U_x^2 = U_t^2 \Rightarrow \rho(t) = k(t)$ for large enough t

3. (15 points) Use energy method to show the following (nonlinear) wave equation only has a zero solution

$$\begin{cases} u_{tt} - \Delta u + u + u^3 = 0 & x \in \Omega, 0 < t < T; \\ u(x, 0) = u_t(x, 0) = 0 & x \in \Omega; \\ u(x, t) = 0 & x \in \partial\Omega, 0 < t < T \end{cases}$$

where $\Omega \subset \mathbb{R}^n$ is an open smooth bounded domain. (Hint: Recall that for our original wave equation $u_{tt} - \Delta u = 0$, the energy is defined as $\frac{1}{2} \int_{\Omega} u_t^2 + |\nabla u|^2 dx$. Think about why the energy is defined in that way, and here you need to construct a new suitable energy $E(t)$ for our nonlinear wave equation and prove it is zero at $t = 0$ and $\frac{d}{dt} E(t) = 0$.)

$$E(t) = \frac{1}{2} \int_{\Omega} U_t^2 + |\nabla U|^2 + U^2 + \frac{1}{2} U^4 d\Omega$$

$$\frac{d}{dt} E(t) = \int_{\Omega} U_t U_{tt} + \nabla U \cdot \nabla U_t + U U_{tt} + U^3 U_t d\Omega$$

$$= \int_{\Omega} U_t (\Delta U - U - U^3) + \nabla U \cdot \nabla U_t + U U_{tt} + U^3 U_t d\Omega$$

$$= \int_{\Omega} U_t \Delta U + \nabla U \cdot \nabla U_t + U U_{tt} + U^3 U_t d\Omega$$

$$\begin{aligned}
 &= \int_{\Omega} u_t + \Delta u + \nabla u \cdot \nabla u + d u \\
 &= \cancel{\int_{\partial\Omega} u_t + \nabla u \cdot \hat{n} dS} - \int_{\Omega} \nabla u_t \cdot \nabla u d\Omega + \int_{\Omega} \nabla u_t \cdot \nabla u d\Omega \\
 &= 0
 \end{aligned}$$

Thus $E(t) = E(0) = \frac{1}{2} \int_{\Omega} 0 + 0 + 0 + 0 d\Omega = 0$

$$\therefore u(x, t) = 0$$

4. (20 points) Let $Q = \Omega \times (0, T]$ and $u \in C^{2,1}(Q) \cap C(\bar{Q})$ satisfy

$$u_t - \Delta u + cu \leq 0 \quad \text{in } Q$$

where $c \geq 0$ is a constant.

- (a) If $u \geq 0$, show that the weak maximum principle holds for u , i.e., u achieves its maximum at the parabolic boundary of Q . (Hint: the proof is similar to the one we did during the lecture.)
- (b) Give a counterexample to show the weak maximum principle may not hold without the condition $u \geq 0$. (Hint: consider $\Omega = (-1, 1)$ and consider $u = -x^2 + at - b$ with some suitable a and b .)

a) Assume at some point in the interior of Q , (x_0, t_0) a maximum is achieved, say $u(x_0, t_0) = M$. Also assume $u_t - \Delta u + cu < 0$ in Q . At (x_0, t_0) since it is a max, $u_t(x_0, t_0) > 0$ and $\Delta u(x_0, t_0) < 0$. It follows

$$0 \leq u_t(x_0, t_0) - \Delta u(x_0, t_0) < -cM$$

$$\Rightarrow cM < 0$$

* since $u \geq 0$ and $c \geq 0$ $\therefore (x_0, t_0) \in \Gamma$

Next consider $u_t - \Delta u + cu \leq 0$ in Q . It follows $\forall \varepsilon > 0$

$$V_\varepsilon(x, t) = u(x, t) - \varepsilon t$$

$$V_\varepsilon \in C^{2,1}(\bar{Q}) \cap C(\bar{Q}) \quad \partial_t V_\varepsilon - \Delta V_\varepsilon + cV_\varepsilon = u_t - \Delta u^2 - \varepsilon + cu - c\varepsilon t < 0$$

$$\max_Q u = \max_Q (V_\varepsilon + \varepsilon t)$$

$$\leq \max_Q V_\varepsilon + \varepsilon T$$

$$\leq \max_{\Gamma} V_\varepsilon + \varepsilon T$$

$$\leq \max_{\Gamma} u + \varepsilon T$$

By applying the previous case

$$\text{as } \varepsilon \rightarrow 0 \quad \max_Q u \leq \max_{\Gamma} u$$

Also, as $\Gamma \subseteq \bar{Q}$ $\max_{\Gamma} u \leq \max_{\bar{Q}} u$

$$\therefore \max_{\Gamma} u = \max_{\bar{Q}} u$$

b) $-x^2 + at - b < 0 \quad \Gamma = (-1, 1)$

$$u_t = a \quad a+2-x^2+at-b \leq 0 \quad c=1$$

$$-x^2 + (1+t)a - b \leq 0 \quad \frac{b-a+at-a}{a} < 0$$

$$-x^2 + at - b \leq 0$$

$$t > \frac{b}{a}$$

$$b > 0$$

$$-x^2 - t - 1 < 0 \quad \forall x \in \mathbb{R} \quad t > 0$$

$$\max \text{ is at } (0, t) \quad \begin{array}{l} 0 \in (-1, 1) \\ 0 \notin \{-1, 1\} \end{array}$$

I tried to draw
a picture

