

HW4

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1. (15 points) For some fixed $a > 0$, compute the following.

(a) The Fourier transform of $f(x) = e^{-a|x|}$ where $x \in \mathbb{R}$.

(b) The Fourier transform of $f(x) = e^{-a|x|}$ where $x \in \mathbb{R}^n$.

(Hint: For part (b), you can use the result from part (a), the trivial identity $\int_0^\infty e^{-(a^2+\xi^2)s} ds = \frac{1}{a^2+\xi^2}$, and the Fourier transform of $e^{-b|x|^2}$ we discussed during lectures to derive an integral representation of $e^{-a|x|}$ in terms of $e^{-c|x|^2}$ for some c , and then continue from there.)

a) Note $f(x) = f(-x)$ so

$$\hat{f}(\xi) = \frac{2}{\sqrt{2\pi}} \int_0^\infty e^{-ax} \cos \xi x \, dx$$

$$\begin{aligned} & \cos \xi x + e^{-ax} \\ & -\xi \sin \xi x - \frac{1}{a} e^{-ax} \quad \int u dv = uv - \int v du \\ & -\xi \cos \xi x + \frac{1}{a} e^{-ax} \end{aligned}$$

$$\int_0^\infty e^{-ax} \cos \xi x \, dx = -\frac{1}{a} \cos \xi x e^{-ax} \Big|_0^\infty + \frac{\xi}{a^2} \sin \xi x e^{-ax} \Big|_0^\infty + \frac{\xi^2}{a^2} \int_0^\infty e^{-ax} \cos \xi x \, dx$$

$$(1 + \frac{\xi^2}{a^2}) \int_0^\infty e^{-ax} \cos \xi x \, dx = \frac{1}{a} + 0$$

$$\int_0^\infty e^{-ax} \cos \xi x \, dx = \frac{1}{a} \frac{a^2}{\xi^2 + a^2} = \frac{a}{\xi^2 + a^2}$$

$$\hat{f}(\xi) = \sqrt{\frac{2}{\pi}} \frac{a}{\xi^2 + a^2}$$

b) $\widehat{e^{-c|x|^2}} = (2c)^{-n/2} e^{-\frac{|\xi|^2}{4c}} \quad \hat{f}(\xi) = (2\pi)^{-n/2} \int_{\mathbb{R}^n} f(x) e^{-ix \cdot \xi} \, dx$ } Hints

$$\hat{f}(\xi) = (2\pi)^{-n/2} \int_{\mathbb{R}^n} e^{-(a|x| + i x \cdot \xi)} \, dx$$

Let $r = |x| \Rightarrow dx = r^{n-1} dr dS^{n-1} \quad x \cdot \xi = |x| |\xi| \cos \theta$

Let $\xi = |\xi| \nu$

$$\hat{f}(\xi) = (2\pi)^{-n/2} \int_0^\infty \int_{S^{n-1}} e^{-(ar + i r |\xi| \cos \theta)} r^{n-1} dr dS^{n-1}$$

$$= (2\pi)^{-n/2} \int_0^\infty (2\pi)^{n/2} \frac{J_{\frac{n-2}{2}}(r|\xi|)}{(r|\xi|)^{\frac{n-2}{2}}} e^{-ar} r^{n-1} dr$$

$$= \int_0^\infty e^{-ar} \frac{J_{\frac{n-2}{2}}(r|\xi|)}{|\xi|^{\frac{n-2}{2}}} r^{\frac{n-2}{2}+1} dr$$

$$= \frac{2^{\frac{n-1}{2}}}{\sqrt{\pi} (a^2 + |\xi|^2)^{\frac{n-1}{2}}} \Gamma\left(\frac{n+1}{2}\right)$$

J is the Bessel function

These are known integrals from integral tables

$$\sqrt{(a^2 + |z|^2)^{n/2} - (z \cdot z)}$$

These are known integrals from integral tables I found

2. (15 points) Let $K(x, y, t)$ be the heat kernel given as

$$K(x, y, t) = \frac{1}{(4\pi t)^{n/2}} e^{-\frac{|x-y|^2}{4t}}, \quad x, y \in \mathbb{R}^n, t > 0$$

(a) Show that $K(x, y, t) \in C^\infty(\mathbb{R}^n \times \mathbb{R}^n \times (0, \infty))$, and show it satisfies the heat equation

$$\frac{\partial K(x, y, t)}{\partial t} - \Delta_x K(x, y, t) = 0.$$

(b) Show that for every $x \in \mathbb{R}^n$ and $t > 0$,

$$\int_{\mathbb{R}^n} K(x, y, t) dy = 1.$$

$$a) K(x, y, t) = (4\pi t)^{-n/2} e^{-\frac{|x-y|^2}{4t}} \quad \Phi(z, t) = (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}}$$

$$x - y = z.$$

$$\text{Let } 1 \leq i \leq j \leq n \quad \frac{\partial \Phi}{\partial z_i} = -\frac{z_i}{2t} (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}} \in C(\mathbb{R}^n)$$

$$\frac{\partial^2 \Phi}{\partial z_i \partial z_j} = \frac{z_i z_j}{4t^2} (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}} \in C(\mathbb{R}^n)$$

In general $D^\alpha \Phi = \alpha_\alpha(z) e^{-\frac{|z|^2}{4t}}$ where α_α is a function of polynomials in z . As the product of continuous functions, $D^\alpha \Phi$ is continuous.

$$\frac{\partial}{\partial z_i} D^\alpha \Phi = -\frac{z_i}{2t} \alpha_\alpha(z) e^{-\frac{|z|^2}{4t}} + \alpha'_\alpha(z) e^{-\frac{|z|^2}{4t}} \in C(\mathbb{R}^n).$$

Thus by induction $\Phi \in C^\infty(\mathbb{R}^n)$ in space.

$$\text{Let } C = (4\pi)^{-n/2}$$

$$\frac{\partial \Phi}{\partial t} = -\frac{n}{2} C \frac{1}{t^{n/2+1}} e^{-\frac{|z|^2}{4t}} + \frac{C|z|^2}{4t^{n/2+2}} e^{-\frac{|z|^2}{4t}} \\ = -\frac{n}{2t} \Phi + \frac{|z|^2}{4t^2} \Phi \in C(\mathbb{R}_+ \setminus \{0\})$$

Similarly to the spacial variables, we will have a polynomial in $1/t$ for higher derivatives.

Thus by the same space argument, $\Phi \in C^\infty(\mathbb{R}_+ \setminus \{0\})$ in time.

$$\frac{\partial \Phi}{\partial z_i} = -\frac{z_i}{2t} (4\pi t)^{-n/2} e^{-\frac{|z|^2}{4t}} = -\frac{z_i}{2t} \Phi$$

$$\frac{\partial^2 \Phi}{\partial z_i^2} = -\frac{z_i}{2t} \left(-\frac{z_i}{2t} \Phi \right) + \frac{1}{t} \Phi$$

$$\frac{\partial^2 \Phi}{\partial z_i^2} = -\frac{z_i}{2t} \left(-\frac{z_i}{2t} \Phi \right) + -\frac{1}{2t} \Phi$$

$$= \left(\frac{|z_i|^2}{4t^2} - \frac{1}{2t} \right) \Phi$$

$$\frac{\partial \Phi}{\partial t} - \Delta_x \Phi = \left(\frac{n}{2t} - \frac{|z|^2}{4t^2} \right) \Phi - \sum_i \frac{\partial^2 \Phi}{\partial z_i^2}$$

$$= \left(\frac{n}{2t} + \frac{|z|^2}{4t^2} \right) \Phi - \sum_i \left(\frac{|z_i|^2}{4t^2} - \frac{1}{2t} \right) \Phi$$

$$= 0$$

b) WTS $\int_{\mathbb{R}^n} \Phi(z, t) dz = 1 \quad \forall t > 0, z \in \mathbb{R}^n$

$$(4\pi t)^{-n/2} \int_{\mathbb{R}^n} e^{-\frac{|z|^2}{4t}} dz = (4\pi t)^{-n/2} \int_{-\infty}^{\infty} e^{-\frac{z_1^2}{4t}} dz_1 \int_{\mathbb{R}^{n-1}} e^{-\frac{|z'|^2}{4t}} dz'$$

$$(z' = (z_2, \dots, z_n)) \rightarrow = (4\pi t)^{-(n-1)/2} \int_{\mathbb{R}^{n-1}} e^{-\frac{|z'|^2}{4t}} dz'$$

note $\int_{-\infty}^{\infty} e^{-\frac{x^2}{4t}} dx = \sqrt{4\pi t}$...

Integral table 5 = 1

3. (20 points) Suppose $g \in L^1(\mathbb{R}^n) \cap L^\infty(\mathbb{R}^n)$ and set

$$u(x, t) = \int_{\mathbb{R}^n} K(x, y, t) g(y) dy$$

where $K(x, y, t)$ denotes the heat kernel as defined in problem 2. Show that for all $t > 0$:

(a) $\|u(\cdot, t)\|_{L^\infty(\mathbb{R}^n)} \leq C t^{-n/2} \|g\|_{L^1(\mathbb{R}^n)}$

(b) $\|u(\cdot, t)\|_{L^1(\mathbb{R}^n)} \leq \|g\|_{L^1(\mathbb{R}^n)}$

(c) $\|u_{x_i}(\cdot, t)\|_{L^\infty(\mathbb{R}^n)} \leq C t^{-1/2} \|g\|_{L^\infty(\mathbb{R}^n)}$

(d) $\|u_{x_i}(\cdot, t)\|_{L^\infty(\mathbb{R}^n)} \leq C t^{-\frac{n+1}{2}} \|g\|_{L^1(\mathbb{R}^n)}$

a) $|u(x, t)| = \left| \int K g dy \right| \leq (4\pi t)^{-n/2} \int_{\mathbb{R}^n} e^{-\frac{|x-y|^2}{4t}} |g| dy$

$\sup e^{-\frac{|x-y|^2}{4t}} = 1$ when $x = y$

$$\leq C t^{-n/2} \|g\|_{L^1}$$

Thus $\|u(x, t)\|_{L^\infty} \leq C t^{-n/2} \|g\|_{L^1}$

b) $\|u(x, t)\|_{L^1} = \int_{\mathbb{R}^n} |u(x, t)| dx = \int_{\mathbb{R}^n} \left| \int_{\mathbb{R}^n} K(x, y, t) g(y) dy \right| dx$

$$\leq \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} K(x, y, t) |g(y)| dy dx$$

$$= \int_{\mathbb{R}^n} |g(y)| \left(\int_{\mathbb{R}^n} K(x, y, t) dx \right) dy$$

$$\leq \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} K(x, y, t) |g(y)| dy dx \\ = \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \Phi(z, t) dz |g(y)| dy \\ = \|g\|_{L^1}$$

$$c) \frac{\partial u}{\partial x_i} = (4\pi t)^{-n/2} \int_{\mathbb{R}^n} \frac{\partial}{\partial x_i} e^{-\frac{|x-y|^2}{4t}} g(y) dy \\ = \frac{1}{2t} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} (x_i - y_i) e^{-\frac{|x-y|^2}{4t}} g(y) \int_{\mathbb{R}^{n-1}} e^{-\frac{|x'-y'|^2}{4t}} dy' dy \\ \|u_{x_i}(x, t)\|_{L^\infty} \leq \frac{1}{2t} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |x_i - y_i| e^{-\frac{|x-y|^2}{4t}} \|g\|_{L^\infty} \int_{\mathbb{R}^{n-1}} e^{-\frac{|x'-y'|^2}{4t}} dy' dy \\ = (2t)^{-1} (4\pi t)^{-n/2} \int_{\mathbb{R}^{n-1}} e^{-\frac{|z'|^2}{4t}} dz' \int_{\mathbb{R}} |z_i| e^{-\frac{|z|^2}{4t}} dz \|g\|_{L^\infty}$$

Note $\int_{\mathbb{R}} |x| e^{-\frac{x^2}{2t}} dx = 2 \int_0^\infty x e^{-\frac{x^2}{2t}} dx = 4t$
By symmetry & Integral tables

$$\|u_{x_i}(x, t)\|_{L^\infty} \leq (2t)^{-1} (4\pi t)^{-n/2} 4t (4\pi t)^{n/2} \|g\|_{L^\infty} \\ = C t^{-1/2} \|g\|_{L^\infty}$$

$$d) |u_{x_i}(x, t)| \leq \frac{1}{2t} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |x - y| e^{-\frac{|x-y|^2}{4t}} |g(y)| dy \quad |x_i - y_i| \leq |x - y| \\ \leq \frac{1}{2t} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |z| e^{-\frac{|z|^2}{4t}} |g(x-z)| dz$$

Consider $h(r) = r e^{-\frac{r^2}{4t}} \quad r = |z|$
 $h'(r) = e^{-\frac{r^2}{4t}} - r \cdot \frac{2r}{4t} e^{-\frac{r^2}{4t}} = e^{-\frac{r^2}{4t}} (1 - \frac{r^2}{2t})$
 $h'(r) = 0 \quad @ \quad r = \sqrt{2t}$

$$h(\sqrt{2t}) = \sqrt{2t} e^{-1/2}$$

$$|u_{x_i}(x, t)| \leq \frac{\sqrt{2t}}{2t} e^{-1/2} (4\pi t)^{-n/2} \int_{\mathbb{R}^n} |g(x-z)| dz \\ = C t^{-\frac{n+1}{2}} \|g\|_{L^1}$$

4. (10 points) For a fixed constant $c \in \mathbb{R}$, find an explicit solution for

$$u_t - \Delta u + cu = f \quad (x, t) \in \mathbb{R}^n \times (0, \infty), \\ u(x, 0) = g(x) \quad x \in \mathbb{R}^n.$$

Let $u = u_1 + u_2 \quad \exists$

$$u_1: \begin{cases} u_t - \Delta u + cu = 0 \\ u(x, 0) = g(x) \end{cases} \quad x \in \mathbb{R}^n, t > 0 \quad u_2: \begin{cases} u_t - \Delta u + cu = f \\ u(x, 0) = 0 \end{cases} \quad x \in \mathbb{R}^n, t > 0$$

$$\underline{u_1)} \begin{cases} u_t - \Delta u = g(x) & x \in \mathbb{R}^n, t > 0 \\ u(x, 0) = 0 \end{cases} \quad \underline{u_2)} \begin{cases} u_t - \Delta u + cu = f & x \in \mathbb{R}^n, t > 0 \\ u(x, 0) = 0 \end{cases}$$

Let $\hat{u} = F(u)$. $\hat{g} = F(g)$

$$\underline{u_1)} \quad \frac{d}{dt} \hat{u} - \Delta \hat{u} + c \hat{u} = 0$$

$$\frac{d}{dt} \hat{u} + (|\xi|^2 + c) \hat{u} = 0$$

$$\Rightarrow \hat{u}(\xi, t) = C e^{-(|\xi|^2 + c)t}$$

$$\hat{u}(\xi, 0) = \hat{g}(\xi) = C$$

$$\hat{u}(\xi, t) = \hat{g}(\xi) e^{-(|\xi|^2 + c)t}$$

$$\begin{aligned} F^{-1}(\hat{u})(x, t) &= e^{-ct} \mathcal{F}^{-1}(\hat{u})(x, t) * g(x) \\ &= e^{-ct} \int_{\mathbb{R}^n} (4\pi t)^{-n/2} e^{-\frac{|x-y|^2}{4t}} g(y) dy \end{aligned}$$

u₂) By the Duhamel principle

$$u_2 = \int_0^t w(x, t-s) ds \quad \text{where}$$

$$w \text{ solves } \begin{cases} w_t - \Delta w + cw = 0 & x \in \mathbb{R}^n, t > s \\ w(x, 0) = f(x, 0) & x \in \mathbb{R}^n \end{cases}$$

$$\text{Thus } u_2 = \int_0^t e^{-c(t-s)} \int_{\mathbb{R}^n} (4\pi(t-s))^{-n/2} e^{-\frac{|x-y|^2}{4(t-s)}} f(y, 0) dy ds$$

$$\begin{aligned} \therefore u(x, t) &= e^{-ct} \int_{\mathbb{R}^n} (4\pi t)^{-n/2} e^{-\frac{|x-y|^2}{4t}} g(y) dy + \\ &\quad \int_0^t e^{-c(t-s)} \int_{\mathbb{R}^n} (4\pi(t-s))^{-n/2} e^{-\frac{|x-y|^2}{4(t-s)}} f(y, 0) dy ds \end{aligned}$$