

# Week 8 Worksheet Solutions

## (Nondegenerate) Perturbation Theory

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**Exercise 1.** Let  $H = H^0 + H^1$  be a perturbed hamiltonian. Suppose we know

$$H^0 |n^0\rangle = E_n^0 |n^0\rangle ,$$

where the  $|n^0\rangle$  are the unperturbed, orthonormal, nondegenerate eigenstates.

- Expand the exact solutions for  $H$ ,  $|n\rangle$  and  $E_n$ , in perturbation expansions.
- Write the Schrödinger equation for  $H$  in terms of the above expansions.
- Study the first order part of the equation from (b), and derive the first order corrections to the energies. You should get

$$E_n^1 = \langle n^0 | H^1 | n^0 \rangle .$$

- Along the way to solving (c), you should have come up with the equation

$$H^0 |n^1\rangle + H^1 |n^0\rangle = E_n^1 |n^0\rangle + E_n^0 |n^1\rangle .$$

Using this equation, find the expansion of  $|n^1\rangle$  in the eigenbasis of  $H^0$ .

*Hints:* In order to find the component of  $|n^1\rangle$  which is parallel to  $|n^0\rangle$ , enforce normalization of  $|n\rangle$  to first order, i.e.  $|n^0\rangle + |n^1\rangle$  should have norm 1. It will be helpful to write  $|n^1\rangle = |n_{\parallel}\rangle + |n_{\perp}\rangle$ , where  $\langle n^0 | n_{\perp} \rangle = 0$ ; also, use the fact that—to first order— $e^{ia} = 1 + ia$ .

- Derive the second order corrections to the energies,  $E_n^2$ .

- We write

$$|n\rangle = \sum_{k=0}^{\infty} \lambda^k |n^k\rangle$$
$$E_n = \sum_{k=0}^{\infty} \lambda^k E_n^k .$$

b) Recall that  $H = H^0 + \lambda H^1$ . Then we have

$$H |n\rangle = E_n |n\rangle$$

$$(H^0 + \lambda H^1) \sum \lambda^k |n^k\rangle = \sum \lambda^{k+j} E_n^k |n^j\rangle.$$

c) The order  $\lambda$  part of the equation from (b) is

$$H^0 |n^1\rangle + H^1 |n^0\rangle = E_n^1 |n^0\rangle + E_n^0 |n^1\rangle.$$

Bracket this equation with  $\langle n^0|$  to obtain

$$E_n^0 \langle n^0 | n^1 \rangle + \langle n^0 | H^1 | n^0 \rangle = E_n^1 + E_n^0 \langle n^0 | n^1 \rangle,$$

from which we immediately obtain the desired result.

d) The idea is that we want to write  $|n^1\rangle$  in the eigenbasis of  $H^0$ . So we want to write

$$|n^1\rangle = \sum c_m |m^0\rangle,$$

and we know that the coefficients  $c_m$  are given by

$$c_m = \langle m^0 | n^1 \rangle.$$

Thus, we want to bracket the equation from (d) with  $\langle m^0|$  with  $m$  not necessarily equal to  $n$ , and try to write down an equation for  $c_m$ . Doing this bracket, we have

$$E_m^0 \langle m^0 | n^1 \rangle + \langle m^0 | H^1 | n^0 \rangle = E_n^1 \delta_{nm} + E_n^0 \langle m^0 | n^1 \rangle.$$

Now, notice that if  $m = n$ , the equation we obtained doesn't tell us anything. But, if  $m \neq n$ , we immediately have

$$c_m = \frac{\langle m^0 | H^1 | n^0 \rangle}{E_n^0 - E_m^0}.$$

For the case  $m = n$ , we follow the hint. Normalization of  $|n\rangle$  to first order gives

$$1 = \langle n^0 | n^0 \rangle + \langle n^0 | n_{\parallel}^1 \rangle + \langle n_{\parallel}^1 | n^0 \rangle,$$

where we note that terms like  $\langle n^1 | n^1 \rangle$  are second order and can be ignored. Since  $|n^0\rangle$  is also assumed normalized, we have

$$0 = 2\text{Re} [\langle n^0 | n_{\parallel}^1 \rangle],$$

so that we can set

$$|n_{\parallel}^1\rangle = i a |n^0\rangle,$$

with  $a \in \mathbb{R}$ . Thus,

$$|n\rangle = (1 + i a) |n^0\rangle + |n_{\perp}^1\rangle$$

$$= e^{ia} |n^0\rangle + |n_{\perp}^1\rangle.$$

Since we can absorb the phase into the definition of  $|n^0\rangle$ , we find

$$|n^1\rangle = |n_{\perp}^1\rangle,$$

i.e. we can set  $c_n = 0$ .

e) We write down the second order equation obtained from the result of (b):

$$H^0 |n^2\rangle + H^1 |n^1\rangle = E_n^2 |n^0\rangle + E_n^1 |n^1\rangle + E_n^0 |n^2\rangle.$$

Then we bracket with  $\langle n^0|$  to find

$$\langle n^0|H^1|n^1\rangle = E_n^2,$$

so

$$E_n^2 = \sum_{m \neq n} \frac{|\langle m|H^1|n\rangle|^2}{E_n^0 - E_m^0}.$$

**Exercise 2.** Suppose you want to calculate the expectation value of some observable  $A$  in the  $n^{\text{th}}$  energy eigenstate of a system perturbed by  $H^1$ ,

$$\langle A \rangle = \langle n|A|n \rangle.$$

Suppose further that all eigenstates are nondegenerate.

- Replace  $|n\rangle$  by its perturbation expansion, and write down the formula for the first order correction to  $\langle A \rangle$ ,  $\langle A \rangle^1$ .
- Use the first order corrections to the states,

$$|n^1\rangle = \sum_{m \neq n} \frac{\langle m^0|H^1|n^0\rangle}{E_n^0 - E_m^0} |m^0\rangle,$$

to rewrite  $\langle A \rangle^1$  in terms of the unperturbed eigenstates.

- If  $A = H^1$ , what does the result of (b) tell you? Explain why this is consistent with the result of Exercise 1(e),

$$E_n^2 = \sum_{m \neq n} \frac{|\langle m|H^1|n\rangle|^2}{E_n^0 - E_m^0}.$$

- We write

$$|n\rangle = |n^0\rangle + \lambda |n^1\rangle + \lambda^2 |n^2\rangle + \dots.$$

Thus,

$$\langle A \rangle = \langle n|A|n \rangle + 2\text{Re} \langle n|A|n^1 \rangle + \dots,$$

so  $\langle A \rangle^1 = 2\text{Re} \langle n|A|n^1 \rangle$ .

- Plugging in the expression for  $|n\rangle$  given above, we get

$$\langle A \rangle^1 = 2\text{Re} \sum_{m \neq n} \frac{\langle m|H^1|n\rangle}{E_n^0 - E_m^0} \langle n|A|m \rangle.$$

c) If  $A = H'$ , then we get that the first order correction to the expectation value of  $H'$  is given by

$$2 \sum_{m \neq n} \frac{|H'_{mn}|^2}{E_n^0 - E_m^0},$$

where  $H'_{mn} = \langle m | H' | n \rangle$ . This is consistent with the equation from 1(d), since we are looking for the expectation value *in the  $n^{\text{th}}$  energy eigenstate*.