

Week 13 Worksheet Solutions

Scattering

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Exercise 1. Warm up.

- a) How do the phase shifts δ_ℓ appear in partial wave scattering, and what is their physical significance?
- b) What is the fundamental assumption on the form of the wavefunctions in the Born approximation?
Hint: If the scattering potential is weak, what approximation can we make?
- c) Starting from the Lippmann-Schwinger equation,

$$\psi(\mathbf{x}) = \varphi_{\mathbf{k}}(\mathbf{x}) + \int d^3x' G_0(\mathbf{x}, \mathbf{x}', E) V(\mathbf{x}') \psi(\mathbf{x}'),$$

where G_0 is the free particle, time-independent Green's function and

$$\begin{aligned}\varphi_{\mathbf{k}}(\mathbf{x}) &= \langle \mathbf{x} | \mathbf{k} \rangle \\ &= \frac{e^{i\mathbf{k} \cdot \mathbf{x}}}{(2\pi)^{3/2}},\end{aligned}$$

explain how you would derive the Born approximation.

Hints: The Green's function is (note that $E = \hbar^2 k^2 / 2m$)

$$G_0(\mathbf{x}, \mathbf{x}', E) = \frac{e^{ik|\mathbf{x}-\mathbf{x}'|}}{|\mathbf{x} - \mathbf{x}'|}.$$

This is done in Griffiths. See also the final review.

Exercise 2. Spin-spin Interaction. Consider two spin-1/2 particles that interact in a potential of the form

$$V(r) = V_0(r) + V_s(r) \boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)}.$$

Suppose that both the orbital and spin interactions are short range in the interparticle separation r (i.e. vanish faster than $1/r$ as $r \rightarrow \infty$).

- a) The first Born approximation for the scattering amplitude is given by

$$f(\mathbf{k}, \mathbf{k}') = -\frac{4\pi^2 m}{\hbar^2} \langle \mathbf{k}' | V | \mathbf{k} \rangle.$$

Use a Fourier transform to express the scattering amplitude in terms of

$$\int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V_0(r_0) d^3r_0,$$

and a similar expression for $V_s(r_0)$.

Hint: Don't forget to account for the incoming and outgoing spins!

- b) You computed on midterm 1 that the eigenvectors of $\boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)}$ are the singlet and triplet states, with eigenvalues -3 and 1 , respectively. If the incoming particles have parallel spins, is a spin flip possible? Why or why not? Explain why the scattering is elastic or inelastic in this case, and then calculate it.
- c) Calculate the scattering amplitude for incident particles with opposite spins. You should be able to split it into two channels: an elastic one and an inelastic one.
- a) The Fourier transform of

$$\begin{aligned} \langle \mathbf{k}' | V | \mathbf{k} \rangle &= \int d^3r_0 \langle \mathbf{k}' | V | \mathbf{r}_0 \rangle \langle \mathbf{r}_0 | \mathbf{k} \rangle \\ &= \int d^3r_0 V(r_0) \langle \mathbf{k}' | \mathbf{r}_0 \rangle \langle \mathbf{r}_0 | \mathbf{k} \rangle \\ &= \left(\frac{1}{(2\pi)^{3/2}} \right)^2 \int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V(r_0) d^3r_0. \end{aligned}$$

Thus, the scattering amplitude is

$$f(\mathbf{k}, \mathbf{k}') = -\frac{m}{2\pi\hbar^2} \left(\int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V_0(r_0) d^3r_0 \langle f | i \rangle + \int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V_s(r_0) d^3r_0 \langle f | \boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)} | i \rangle \right),$$

where $|i\rangle$ ($|f\rangle$) denote the spin states of the incoming (resp. outgoing) particles. Note that the incoming spin state space is *four*-dimensional, as we should account for the spins of *both* particles (each of which has a two-dimensional state space). Indeed, the spins of both particles can change from their initial configurations to some different final configurations.

- b) Note that parallel spins are eigenstates of $\boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)}$; thus, it is impossible for a scattered particle to change spin, so the scattering will be purely elastic. Since these have eigenvalue 1 , we get

$$f(\theta) = -\frac{m}{2\pi\hbar^2} \left(\int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V_0(r_0) d^3r_0 + \int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V_s(r_0) d^3r_0 \right).$$

- c) If the spins are not parallel, then the scattered wave can have either the same spins or opposite spins (spin-flip). This is because the singlet and mixed triplet states are superpositions of the antiparallel

configurations. It follows that we have an elastic channel (no spin-flip), as well as an inelastic one (spin-flip). We compute

$$\begin{aligned}\langle \uparrow\downarrow | \boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)} | \uparrow\downarrow \rangle &= -1 \\ \langle \uparrow\downarrow | \boldsymbol{\sigma}^{(1)} \cdot \boldsymbol{\sigma}^{(2)} | \downarrow\uparrow \rangle &= 2,\end{aligned}$$

This is most easily seen by writing e.g. $|\uparrow\downarrow\rangle$ as a linear combination of triplet and singlet states. We then end up with two amplitudes.

$$\begin{aligned}f_{\uparrow\downarrow,\uparrow\downarrow}(\theta) &= -\frac{m}{2\pi\hbar^2} \left(\int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V_0(r_0) d^3r_0 - \int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V_s(r_0) d^3r_0 \right) \\ f_{\uparrow\downarrow,\downarrow\uparrow}(\theta) &= -2\frac{m}{2\pi\hbar^2} \int e^{-i(\mathbf{k}-\mathbf{k}')\cdot\mathbf{r}_0} V_s(r_0) d^3r_0.\end{aligned}$$

Note that the *kinetic* energy of the flipped spin states is still the same as the original states; only potential energy can change via the potentials V_0, V_s .