

# Math 6417 Homework 4

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December 1, 2023

## Question 1.

Define the **Fourier transform operator**  $\mathcal{F} : L^1(\mathbf{R}) \rightarrow L^\infty(\mathbf{R})$  by

$$\mathcal{F}(f)(y) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{iyx} f(x) \, dx. \quad (1)$$

1.1) We note that the function  $x \mapsto e^{iyx} f(x)$  is clearly integrable if  $f$  is, so the integral in (1) exists for all  $y$ . We show that  $\mathcal{F}(f) \in L^\infty(\mathbf{R})$  as claimed, and  $\|\mathcal{F}f\|_{L^\infty} \leq \frac{1}{\sqrt{2\pi}} \|f\|_{L^1}$ . Indeed, for  $y \in \mathbf{R}$ ,

$$|\mathcal{F}(f)(y)| = \frac{1}{\sqrt{2\pi}} \left| \int_{-\infty}^{\infty} e^{iyx} f(x) \, dx \right| \quad (2)$$

$$\leq \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} |e^{iyx} f(x)| \, dx = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} |f(x)| \, dx = \frac{1}{\sqrt{2\pi}} \|f\|_{L^1}. \quad (3)$$

Therefore,  $\|\mathcal{F}f\|_{L^\infty} \leq \frac{1}{\sqrt{2\pi}} \|f\|_{L^1}$ .

1.2) Suppose that  $f \in C^2(\mathbf{R})$ , and  $f, f', f'' \in L^1(\mathbf{R})$ , and  $f(x), f'(x), f''(x) \rightarrow 0$  as  $x \rightarrow \pm\infty$ . Then there exists a constant  $C$  such that  $|y^2 \mathcal{F}(f)(y)| \leq C$  for all  $y \in \mathbf{R}$ . Furthermore,  $\mathcal{F}(f) \in L^1(\mathbf{R})$ .

*Proof.* Since  $f'' \in L^1(\mathbf{R})$ , we can take its Fourier transform, which yields

$$\mathcal{F}(f'')(y) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{iyx} f''(x) \, dx. \quad (4)$$

We can integrate by parts because  $f', f \in L^1(\mathbf{R})$  and are continuous, and  $f(x), f'(x) \rightarrow 0$  as  $x \rightarrow \pm\infty$ . This gives

$$\mathcal{F}(f'')(y) = \frac{1}{\sqrt{2\pi}} \left[ f'(x) e^{iyx} \Big|_{-\infty}^{\infty} - iy \int_{-\infty}^{\infty} e^{iyx} f'(x) \, dx \right] \quad (5)$$

$$= \frac{iy}{\sqrt{2\pi}} \left[ -f(x) e^{iyx} \Big|_{-\infty}^{\infty} + iy \int_{-\infty}^{\infty} e^{iyx} f(x) \, dx \right] \quad (6)$$

$$= -y^2 \mathcal{F}(f)(y). \quad (7)$$

By the reasoning in 1.1), it follows that

$$|y^2 \mathcal{F}(f)(y)| = |\mathcal{F}(f'')(y)| \leq \frac{1}{\sqrt{2\pi}} \|f''\|_{L^1} \quad (8)$$

for all  $y \in \mathbf{R}$ .

Thus, if  $C = \frac{1}{\sqrt{2\pi}} \|f''\|_{L^1}$ , then  $|\mathcal{F}(f)(y)| \leq \frac{C}{y^2}$  for all  $y \in \mathbf{R}$ . On the other hand,  $\mathcal{F}(f) \in L^\infty(\mathbf{R})$  by part 1.1), so  $\mathcal{F}(f)$  is dominated by the integrable function

$$\phi(y) = \begin{cases} \|\mathcal{F}(f)\|_{L^\infty} & y \in [-1, 1], \\ \frac{C}{y^2} & \text{otherwise.} \end{cases} \quad (9)$$

By the integral comparison test,  $\mathcal{F}(f) \in L^1(\mathbf{R})$ . □

1.3) Formally,  $\mathcal{F}^2(f)(y) = f(-y)$ .

*Proof.* We note that if  $f \in C^1 \cap L^1(\mathbf{R})$ , and  $f' \in L^1(\mathbf{R})$ , and  $f(x) \rightarrow 0$  as  $x \rightarrow \pm\infty$ , then we can use integration by parts to show that

$$\mathcal{F}(f')(y) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{iyx} f'(x) dx = \frac{1}{\sqrt{2\pi}} \left[ e^{iyx} f(x) \Big|_{-\infty}^{\infty} - iy \int_{-\infty}^{\infty} e^{iyx} f(x) dx \right] \quad (10)$$

$$= -iy \mathcal{F}(f)(y). \quad (11)$$

On the other hand, let  $f \in L^1(\mathbf{R})$ , and define  $g(x) = ix f(x)$ . If  $g \in L^1(\mathbf{R})$  as well, then

$$\frac{d}{dy} \mathcal{F}(f)(y) = \frac{d}{dy} \int_{-\infty}^{\infty} e^{iyx} f(x) dx = \int_{-\infty}^{\infty} \frac{\partial}{\partial y} [e^{iyx} f(x)] dx \quad (12)$$

$$= \int_{-\infty}^{\infty} e^{iyx} ix f(x) dx = \mathcal{F}(g)(y). \quad (13)$$

If we take  $f(x) = e^{-ax^2}$ , then  $f$  satisfies the above assumptions. Since  $f'(x) = -2ax f(x)$ ,

$$2ai \frac{d}{dy} \mathcal{F}(f)(y) = 2ai \mathcal{F}(i(\cdot) f(\cdot))(y) = \mathcal{F}(-2a(\cdot) f(\cdot))(y) = \mathcal{F}(f')(y) = -iy \mathcal{F}(f)(y). \quad (14)$$

Hence,  $\mathcal{F}(f)(y)$  is the unique solution of the IVP

$$u' = -\frac{y}{2a} u, \quad u(0) = \mathcal{F}(f)(0). \quad (15)$$

The general solution of the differential equation is

$$u(y) = u(0) e^{-\frac{y^2}{4a}}. \quad (16)$$

Since

$$\mathcal{F}(f)(0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-ax^2} dx = \frac{1}{\sqrt{2a}}, \quad (17)$$

it follows that

$$\mathcal{F}(f)(y) = \frac{1}{\sqrt{2a}} e^{-\frac{y^2}{4a}}. \quad (18)$$

Thus, if  $\phi_a(x) = e^{-ax^2}$ , then, formally,

$$\mathcal{F}(1)(y) = \mathcal{F} \left( \lim_{a \rightarrow 0^+} \phi_a \right) (y) = \lim_{a \rightarrow 0^+} \mathcal{F}(\phi_a)(y) = \lim_{a \rightarrow 0^+} \frac{1}{\sqrt{2a}} e^{-\frac{y^2}{4a}}. \quad (19)$$

We would like to interpret the last limit formally as a constant multiple of the Dirac delta function. Clearly,

$$\lim_{a \rightarrow 0^+} \frac{1}{\sqrt{2a}} e^{-\frac{y^2}{4a}} = \begin{cases} 0 & y \neq 0, \\ \infty & y = 0. \end{cases} \quad (20)$$

At the same time, for any  $a > 0$ ,

$$\int_{-\infty}^{\infty} \frac{1}{\sqrt{2a}} e^{-\frac{y^2}{4a}} dy = \frac{1}{\sqrt{2a}} \sqrt{4a\pi} = \sqrt{2\pi}, \quad (21)$$

so it makes sense formally that we should have  $\mathcal{F}(1)(y) = \sqrt{2\pi} \delta(y)$ .

Now, if we consider applying the Fourier transform twice to a function  $f$ , we get

$$\mathcal{F}\mathcal{F}(f)(y) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{ixy} e^{izx} f(z) \, dz \, dx \quad (22)$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(z) \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{ix(y+z)} \, dx \, dz \quad (23)$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(z) \mathcal{F}(1)(y+z) \, dz \quad (24)$$

$$= \int_{-\infty}^{\infty} f(z) \delta(y+z) \, dz \quad (25)$$

$$= \int_{-\infty}^{\infty} f(z-y) \delta(z) \, dz \quad (26)$$

$$= f(-y). \quad (27)$$

□

- 1.4) Define  $g(y) = f(-y)$  for some function  $f$ . Based on the formal result from part 1.3), we see immediately that

$$\mathcal{F}^4(f)(y) = \mathcal{F}^2(\mathcal{F}^2(f))(y) = \mathcal{F}^2(g)(y) = g(-y) = f(y). \quad (28)$$

Since  $f$  was arbitrary, it follows formally that  $\mathcal{F}^4 = I$ , the identity operator.

- 1.5) Let  $p(x) = x^4$ . By the Spectral Mapping Theorem,

$$p(\sigma(\mathcal{F})) = \sigma(p(\mathcal{F})). \quad (29)$$

Since  $p(\mathcal{F}) = \mathcal{F}^4 = I$ , the spectrum of  $p(\mathcal{F})$  is just  $\sigma(I) = \{1\}$ , as the operator  $I - \lambda I = (1 - \lambda)I$  is invertible, with inverse  $\frac{1}{1-\lambda}I$ , if and only if  $\lambda \neq 1$ . Therefore, if  $\lambda \in \sigma(\mathcal{F})$ , then  $p(\lambda) = 1$ , that is,  $\lambda^4 = 1$ . The possible solutions of this equation are  $1, -1, i, -i$ , so  $\sigma(\mathcal{F}) \subseteq \{1, -1, i, -i\}$ .

- 1.6) If we reuse the result in equation (18) with  $a = \frac{1}{2}$ , we see that if  $f(x) = e^{-\frac{1}{2}x^2}$ , then

$$\mathcal{F}(f)(y) = e^{-\frac{1}{2}y^2} \quad (30)$$

as well. Thus,  $\mathcal{F}f = f$ , so  $f$  is an eigenfunction of  $\mathcal{F}$  with corresponding eigenvalue 1.