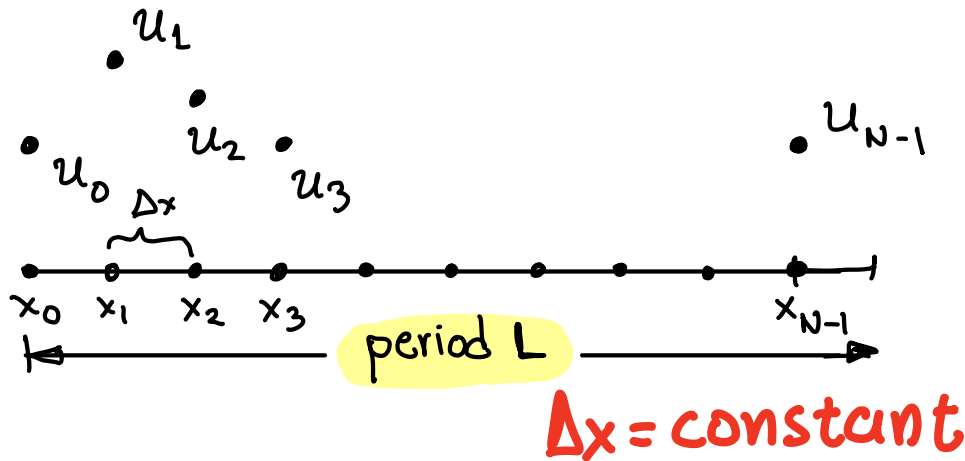


Trigonometric Interpolation + Differentiation

- periodic discrete function



Discrete Fourier Transform:

$$\hat{u}_n = \sum_{i=0}^{N-1} u_i e^{-i k_n x_i} \quad \text{where: } k_n = n \frac{2\pi}{L}$$

Forward $u_i \rightarrow \hat{u}_n$

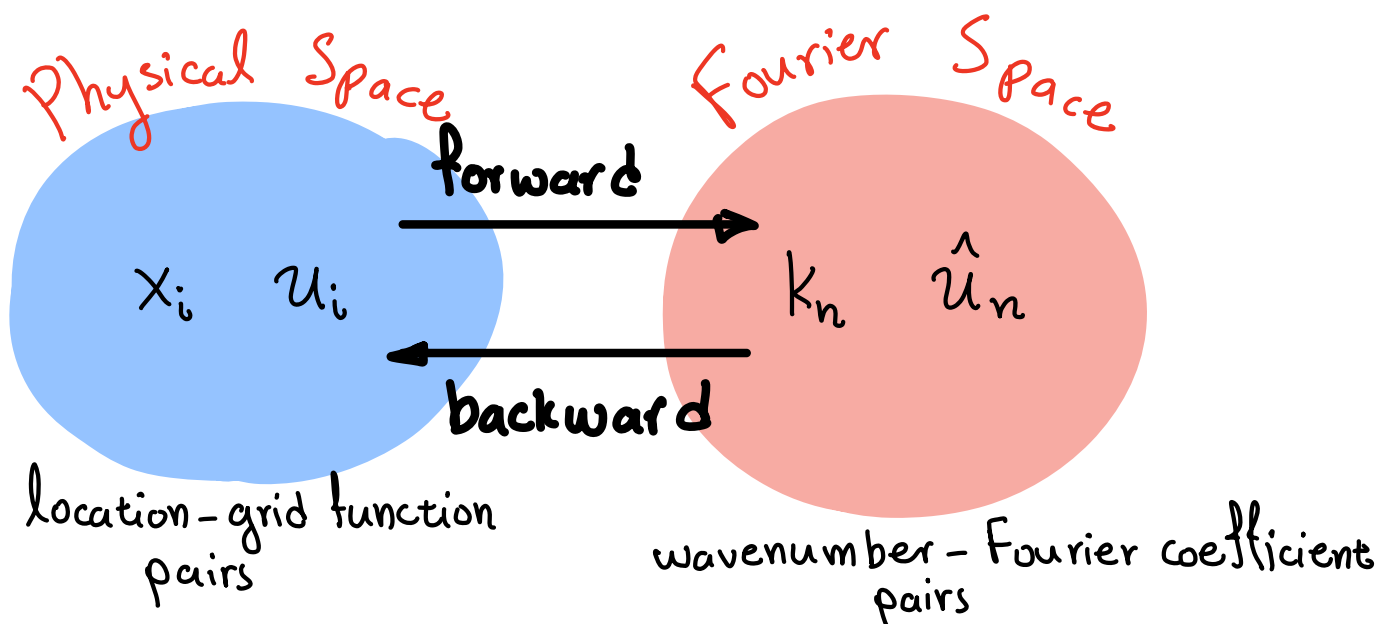
Important:

- $\underline{i} = \sqrt{-1}$ imaginary unit
- $e^{-\underline{i} k_n x_j} = \cos(k_n x_j) + \underline{i} \sin(k_n x_j)$
- \hat{u}_n are complex
- When u_i is real $\hat{u}_n = \hat{u}_{n-n}^*$ complex conjugate
- Aliasing: \hat{u}_n are the same for $n \rightarrow n \pm mN$ integer
- **CHECK IMPLEMENTATION on ordering of \hat{u}_n, k_n**

Backward DFT:

$$u_i = \frac{1}{N} \sum_{n=0}^{N-1} \hat{u}_n e^{\underline{i} k_n x_i}$$

This is the expression to find u_i . However, in applications the transform will take as input all Fourier coefficients and return all values of u_i



Spectral Derivative

We will utilize the discrete Fourier transform to estimate the derivatives of u_i at locations x_i :

u_i written using expression for inverse Fourier Transform:

$$u_i = \frac{1}{N} \sum_{n=0}^N \hat{u}_n e^{i k_n x_i}$$

differentiate in x both sides:

$$\frac{du_i}{dx} = \frac{d}{dx} \left[\frac{1}{N} \sum_{n=0}^N \hat{u}_n e^{i k_n x_i} \right]$$

$$\frac{du_i}{dx} = \frac{1}{N} \sum_{n=0}^N \left[\hat{u}_n \frac{d}{dx} (e^{i k_n x_i}) \right]$$

$$\frac{du_i}{dx} = \frac{1}{N} \sum_{n=0}^N \left[\hat{u}_n i k_n e^{i k_n x_i} \right]$$

Remember:

This is a sum over n

This is now in the form:

$$\frac{du_i}{dx} = \frac{1}{N} \sum_{n=0}^N \left[(\text{Fourier Coefficients}) e^{i k_n x_i} \right]$$

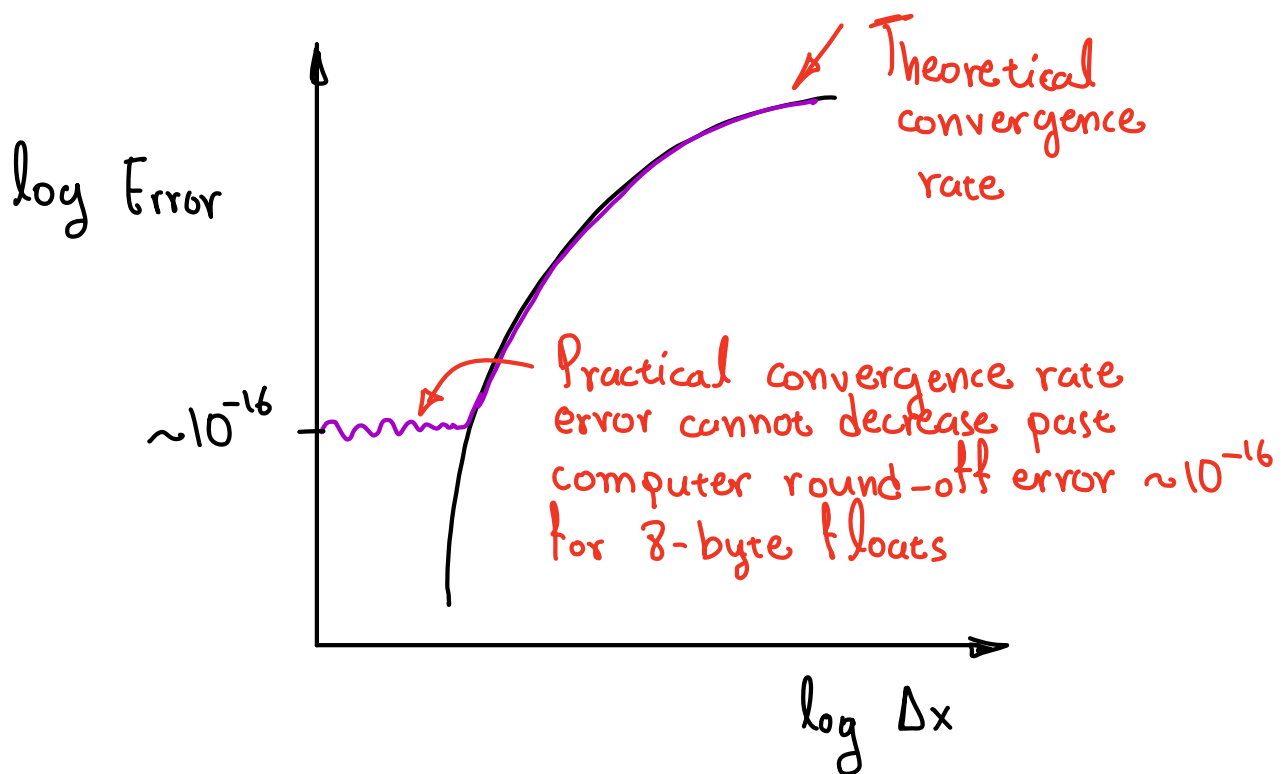
By the uniqueness property of the Fourier Transform we have shown that $i k_n \hat{u}_n$ are the coefficients of the Fourier transform of $\frac{du_i}{dx}$ ← The derivative of u_i at x_i

Similarly: $(ik_n)^m \hat{u}_n$ are the Fourier coefficients of the m -th derivative

For example: second derivative $m=2$: $(ik_n)^2 \hat{u}_n = -k_n^2 \hat{u}_n$

Important: • ONLY periodic functions
• ONLY constant Δx

Order of accuracy of the spectral derivative is "infinite"
This means error decreases faster than any power of Δx^2
if $u(x)$ is infinitely differentiable



Recipe:

How to apply this in practice:

We have u_i at x_i locations and we want to find $\frac{du}{dx}$ at x_i

- Steps:
1. Compute wavenumbers k_n
 2. $\hat{u}_n = \text{DFT}(u_i)$ compute Fourier coefficients \hat{u}_n
 3. Form Fourier coefficients of $\frac{du}{dx}$

$$\widehat{\frac{du}{dx}}_n = i k_n \hat{u}_n$$

multiply the corresponding k_n and u_n times i (imaginary unit)

4. $\frac{du_i}{dx} = \text{inverse DFT}(i k_n \hat{u}_n)$

take the inverse transform to get back to physical space and find $\frac{du}{dx}$ at x_i

Finite Difference Approximation (Lomax 3.2-3.4)

Taylor series $u_{i+j} = u_i + (j\Delta x) \frac{\partial u}{\partial x} \Big|_i + \frac{1}{2} (j\Delta x)^2 \frac{\partial^2 u}{\partial x^2} \Big|_i + \dots$

expanding at x_i to distance $x_i - x_j$

$$+ \frac{1}{n!} (j\Delta x)^n \frac{\partial^n u}{\partial x^n} \Big|_i + \dots$$

Example $j=1$ $u_{i+1} = u_i + \Delta x \frac{\partial u}{\partial x} \Big|_i + \frac{1}{2} \Delta x^2 \frac{\partial^2 u}{\partial x^2} \Big|_i + \dots$

$$\frac{u_{i+1} - u_i}{\Delta x} - \frac{\partial u}{\partial x} \Big|_i = \frac{1}{2} \Delta x \frac{\partial^2 u}{\partial x^2} \Big|_i + \dots$$

infinite terms

approximating $\frac{\partial u}{\partial x} \Big|_i$ with error

leading error term

$$u_{i-1} = u_i - \Delta x \frac{\partial u}{\partial x} \Big|_i + \frac{1}{2} \Delta x^2 \frac{\partial^2 u}{\partial x^2} \Big|_i - \frac{1}{6} \Delta x^3 \frac{\partial^3 u}{\partial x^3} \Big|_i + \dots$$

$$\frac{u_{i+1} - u_{i-1}}{2\Delta x} - \frac{\partial u}{\partial x} \Big|_i = \frac{1}{6} \Delta x^2 \frac{\partial^3 u}{\partial x^3} \Big|_i + \dots$$

point operator

leading error term

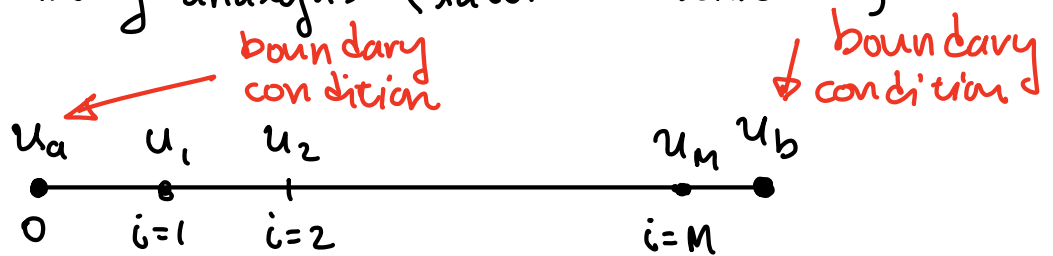
ORDER OF ACCURACY = 2

Move from point operators to matrix operators

Why? 1. Boundary conditions

2. Stability analysis (later this semester)

Example



$\frac{\partial^2 u}{\partial x^2}$? approximation is $(\delta_{xx} u)_i = \frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta x^2} + O(\Delta x^2)$

$$\delta_{xx} u_1 = \frac{1}{\Delta x^2} (u_a - 2u_1 + u_2)$$

$$\delta_{xx} u_2 = \frac{1}{\Delta x^2} (u_1 - 2u_2 + u_3)$$

\vdots

$$\delta_{xx} u_M = \frac{1}{\Delta x^2} (u_{M-1} - 2u_M + u_b)$$

$$\delta_{xx} \begin{bmatrix} u_1 \\ u_2 \\ \vdots \\ u_M \end{bmatrix} = \frac{1}{\Delta x^2} \begin{bmatrix} -2 & 1 & & & \\ 1 & -2 & 1 & & \\ & \ddots & \ddots & \ddots & \\ & & 1 & -2 & 1 \\ & & & 1 & -2 \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ \vdots \\ u_M \end{bmatrix} + \begin{bmatrix} u_a \\ 0 \\ \vdots \\ 0 \\ u_b \end{bmatrix}$$

diagonal
boundary conditions

$$\delta_{xx} \vec{u} = A \vec{u} + \vec{b}$$

↑ vector \vec{u} is only a function of time

$$\delta_{xx} \vec{u} = \frac{1}{\Delta x^2} B(M; 1, -2, 1) \vec{u} + \vec{b}$$

General FD approximation:

$$\sum_{j=-r}^s b_j \left. \frac{\partial^m u}{\partial x^m} \right|_{i+j} - \sum_{j=-p}^q a_j u_{i+j} = \text{error}$$

\uparrow \uparrow
 $r+s+1$ $p+q+1$
 points points
 implicit explicit
 Hermitian difference
 Padé
 compact

Stencil width = $\max(s, q) + \max(r, p) + 1$

Example: Implicit Finite difference:

$$\left. \frac{\partial u}{\partial x} \right|_{i-1} + 4 \left. \frac{\partial u}{\partial x} \right|_i + \left. \frac{\partial u}{\partial x} \right|_{i+1} - \frac{3}{\Delta x} (u_{i+1} - u_{i-1}) = \frac{\Delta x^4}{120} \frac{\partial^5 u}{\partial x^5} + \dots$$

Order = 4 stencil width = 3

Matrix operator:

$$\frac{1}{6} B(1, 4, 1) \delta_x \vec{u} = \frac{1}{2\Delta x} B(-1, 0, 1) \vec{u} + \vec{b}$$

$$\delta_x \vec{u} = 6 B(1, 4, 1)^{-1} \left[\frac{1}{2\Delta x} B(-1, 0, 1) \vec{u} + \vec{b} \right]$$