### Conservation of stuff

#### 1. Continous equations

For example: momentum conservation in volume of

We can also derive conservation laws for other quantities

Kinetic energy: 
$$\vec{u} \cdot \left(\frac{3\vec{y}\vec{u}}{3t} + \vec{\nabla} \cdot \vec{y}\vec{u}\vec{v} = -\vec{\nabla} \cdot \vec{\tau}\right)$$
conservation
$$= q_{\text{uaction}}$$

$$\frac{3}{3t} \frac{1}{2} \vec{y} \vec{u} \cdot \vec{u} + \cdots$$

Enstrophy:  $\nabla_{x}\vec{u} \cdot (x \pmod{\text{momentum equation}})$ 

Key Point: in the continuous "world" I can write conservation laws of any quantity I can think of (some one more useful than others)

- Only specially constructed methods guarantee conservation
   these methods or schemes are called conservative
- · Conservative methods conserve only the quantities we solve for.

e.g. if a method numerically integrates

mass and momentum, a conservative method with

conserve only mass and momentum to machine precision

and nothing more (not kinetic energy, not

enstrophy, etc...)

· If a CFDer wants to conserve additional quantities then a special method is required

#### Remember:

• No conservation means a quantity "leaks" from the interior of the computational domain e.g. somehow mass or momentum is added or depleted in the interior of the computational domain.

Conservation properties are important in CFD

## Discretization forms of convection term

$$\frac{\partial x_j}{\partial x_j}$$
 u;u;

· Advective

· Skew-symmetric

$$\frac{1}{2} \frac{\partial x_{j}}{\partial x_{j}} (u_{j}u_{i}) + \frac{1}{2}u_{j} \frac{\partial u_{i}}{\partial x_{j}}$$

Rotational

$$n^{2}\left(\frac{3x^{2}}{3n^{2}}-\frac{3x^{2}}{3n^{2}}\right)+\frac{5}{1}\frac{3x^{2}}{3n^{2}n^{2}}$$

- · All forms are equivalent in differential form
- . But, have different properties in discrete form
- Please see details in Morinish et al. (1998)
   in "Additional Material" on Husky CT
- · Check out Table 7

# Verification of Euler equations time integration (Part 2 of Project)

We will exploid that for our discretization

- Kinetic energy remains constant\* (not really)
- Scalar variance remains constant\* (not really)
- => these imply that the simulation cannot "blow up" because energy in always bounded by the initial condition

# Verification Test 1 (easy)

- Initialize u,v, & with random numbers but apply u,v BC u and v periodic left/right & v=0 top/bottom walls
- · Perform 100 time steps (any Dt as long the integration is stable)
- . Success if a. Flow is divergence-free
  - b. u, v, & remain finite

## Verification Test 2 (advanced)

- 1. a. Initialize 2, y, & with random numbers
  - b. Apply BCs
  - c. Mate flow divergence free
- 2. Calculate initial

  Kinetic energy  $K_0 = \sum_{i=1}^{\infty} \left(u_{i,j}^2 + v_{i,j}^2\right)$ Temperature "energy"  $V_0 = \sum_{i=1}^{\infty} \theta_{i,j}^2$
- 3. Integrate to time Tend with constant Dt pick a time
- 4. Calculate final Kinetic energy  $K_e = \sum_{i=1}^{2} (u_{ij}^2 + v_{ij}^2)$ Temperature "energy"  $V_e = \sum_{i=1}^{2} \theta_{ij}^2$

Ideally Ke = Ko and Ye = Vo since the method conserves K and V by construction

- However, Runge-Kutta is not exact time integration
   and introduces a Dt<sup>3</sup> error
- The error is dissipative, because of numerical stability
  Thus: Ke< ko and Ve< Yo

To verify the method:

- · Fix (Do NOT change) Dx, Dy and Te
- · Do steps 1-4 with different Dt
- . For each Dt calculate errors:

$$E_{\kappa}(\Delta t) = K_0 - K_e$$
 $E_{\nu}(\Delta t) = Y_0 - Y_e$ 
Value here

. Show that  $E_k \sim \Delta t^3$  and  $E_{\gamma} \sim \Delta t^3$