

AstroStatistics and Cosmology Homework

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Exercise 4

After being given a probability distribution $\mathbb{P}(x)$, we define the *characteristic function* ϕ as its Fourier transform, which can also be expressed as the expectation value of $\exp(-i\vec{k} \cdot \vec{x})$:

$$\phi(\vec{k}) = \int d^n x \exp(-i\vec{k} \cdot \vec{x}) \mathbb{P}(x) = \mathbb{E} \left[\exp(-i\vec{k} \cdot \vec{x}) \right]. \quad (1.1)$$

Claim 1.1. *A multivariate normal distribution*

$$\mathcal{N}(\vec{x}|\vec{\mu}, C) = \frac{1}{(2\pi)^{n/2} \sqrt{\det C}} \exp\left(-\frac{1}{2} \vec{y}^\top C^{-1} \vec{y}\right) \Big|_{\vec{y}=\vec{x}-\vec{\mu}}, \quad (1.2)$$

has a characteristic function equal to

$$\phi(\vec{k}) = \exp\left(-i\vec{\mu} \cdot \vec{k} - \frac{1}{2} \vec{k}^\top C \vec{k}\right). \quad (1.3)$$

Proof: completing the square. The integral we need to compute is given, absorbing the normalization into a factor N , by

$$\phi(\vec{k}) = N \int d^n x \exp\left(-i\vec{k} \cdot \vec{x} - \frac{1}{2} \vec{y}^\top C^{-1} \vec{y}\right) \Big|_{\vec{y}=\vec{x}-\vec{\mu}}. \quad (1.4)$$

The only integrals we really know how to do are Gaussian ones, so we want to rewrite the argument of the exponential so that it is a quadratic form. The manipulation goes as follows, considering the opposite of the argument the exponential in order to have less minus signs and defining the symmetric matrix $V = C^{-1}$:

$$i\vec{k} \cdot \vec{x} + \frac{1}{2} \vec{y}^\top V \vec{y} = i\vec{k} \cdot \vec{x} + \frac{1}{2} \vec{x}^\top V \vec{x} - \vec{x}^\top V \vec{\mu} + \frac{1}{2} \vec{\mu}^\top V \vec{\mu} \quad (1.5)$$

$$= \frac{1}{2} \vec{x}^\top V \vec{x} + \vec{x}^\top (i\vec{k} - V\vec{\mu}) + \frac{1}{2} \vec{\mu}^\top V \vec{\mu} \quad (1.6)$$

$$\begin{aligned}
&= \underbrace{\frac{1}{2} \left(\vec{x} + V^{-1}(i\vec{k} - V\vec{\mu}) \right)^\top V \left(\vec{x} + V^{-1}(i\vec{k} - V\vec{\mu}) \right)}_{\textcircled{1}} + \\
&\quad \underbrace{-\frac{1}{2} \left(i\vec{k} - V\vec{\mu} \right)^\top V^{-1} \left(i\vec{k} - V\vec{\mu} \right) + \frac{1}{2} \vec{\mu}^\top V \vec{\mu}}_{\textcircled{2}}, \tag{1.7}
\end{aligned}$$

which we can now integrate, since it is now a quadratic form in terms of a shifted variable, $\vec{x} + \vec{p}$, where the constant (with respect to \vec{x}) vector \vec{p} is given by $V^{-1}(i\vec{k} - V\vec{\mu})$.¹

Now, shifting the integral from one in $d^n x$ to one in $d^n(x + p)$ does not change the measure, since the Jacobian of a shift is the identity. Then, we have

$$\phi(\vec{k}) = N \int d^n(x + p) \exp(-\textcircled{1} - \textcircled{2}) \tag{1.12}$$

$$= N \sqrt{\frac{(2\pi)^n}{\det V}} \exp(-\textcircled{2}) \tag{1.13}$$

$$= \underbrace{\frac{1}{\sqrt{\det V \det C}}}_{=1} \exp(-\textcircled{2}), \tag{1.14}$$

since the determinant of the inverse is the inverse of the determinant.

Now, we only need to simplify $\textcircled{2}$:

$$\textcircled{2} = -\frac{1}{2} \left[-\vec{k}^\top V^{-1} \vec{k} - 2i\vec{\mu}^\top V V^{-1} \vec{k} + \vec{\mu}^\top V V^{-1} V \vec{\mu} \right] + \frac{1}{2} \vec{\mu}^\top V \vec{\mu} \tag{1.15}$$

$$= \frac{1}{2} \vec{k}^\top C \vec{k} + i\vec{\mu}^\top \vec{k}, \tag{1.16}$$

inserting which into the exponent yields the desired result. \square

Proof: by diagonalization. We now follow a different approach: the covariance matrix C is symmetric, so we will always be able to find an orthogonal matrix O (satisfying $O^\top = O^{-1}$) such that $C = O^\top D O$, where D is diagonal. We will then also have $V = C^{-1} = O^\top D^{-1} O$. Let us denote the eigenvalues of D as λ_i , and the eigenvalues of D^{-1} as $d_i = \lambda_i^{-1}$.

Defining $\vec{z} = O\vec{x}$, $\vec{m} = O\vec{\mu}$, $\vec{u} = O\vec{k}$ the negative of the argument of the integral becomes:

$$i\vec{k} \cdot \vec{x} + \frac{1}{2} (\vec{x} - \vec{\mu})^\top C^{-1} (\vec{x} - \vec{\mu}) = i\vec{u} \cdot \vec{z} + \frac{1}{2} (\vec{z} - \vec{m})^\top D^{-1} (\vec{z} - \vec{m}) \tag{1.17}$$

¹ In the last step we applied the matrix square completion formula: for a symmetric matrix A and vectors \vec{x} , \vec{b} we have

$$\frac{1}{2} (\vec{x} + A^{-1} \vec{b})^\top A (\vec{x} + A^{-1} \vec{b}) - \frac{1}{2} \vec{b}^\top A^{-1} \vec{b} = \tag{1.8}$$

$$= \frac{1}{2} \left[\vec{x}^\top A \vec{x} + \vec{x}^\top A A^{-1} \vec{b} + (A^{-1} \vec{b})^\top A \vec{x} + (A^{-1} \vec{b})^\top A A^{-1} \vec{b} - \vec{b}^\top A^{-1} \vec{b} \right] \tag{1.9}$$

$$= \frac{1}{2} \left[\vec{x}^\top A \vec{x} + \vec{x}^\top \vec{b} + \vec{b}^\top (A^{-1})^\top A \vec{x} + \vec{b}^\top (A^{-1})^\top \vec{b} - \vec{b}^\top A^{-1} \vec{b} \right] \tag{1.10}$$

$$= \frac{1}{2} \vec{x}^\top A \vec{x} + \vec{b}^\top \vec{x}, \tag{1.11}$$

which we used with $\vec{b} = i\vec{k} - V\vec{\mu}$.

$$= i\vec{u} \cdot \vec{z} + \frac{1}{2} \sum_i d_i (z_i - m_i)^2 \quad (1.18)$$

$$= \sum_i \left[iu_i z_i + \frac{d_i}{2} (z_i^2 + m_i^2 - 2m_i z_i) \right] \quad (1.19)$$

$$= \sum_i \left[z_i^2 \frac{d_i}{2} + z_i (iu_i - m_i d_i) + \frac{d_i}{2} m_i^2 \right]. \quad (1.20)$$

With this, and since by $\det O = 1$ we have $d^n z = d^n x$, we can decompose our Gaussian integral into a product of Gaussian integrals:

$$\phi(\vec{k}) = N \int d^n x \exp \left(-i\vec{k} \cdot \vec{x} - \frac{1}{2} (\vec{x} - \vec{\mu})^\top C^{-1} (\vec{x} - \vec{\mu}) \right) \quad (1.21)$$

$$= N \int d^n z \exp \left(- \sum_i \left[z_i^2 \frac{d_i}{2} + z_i (iu_i - m_i d_i) + \frac{d_i}{2} m_i^2 \right] \right) \quad (1.22)$$

$$= N \prod_i \int dz_i \exp \left(-z_i^2 \frac{d_i}{2} - z_i (iu_i - m_i d_i) - \frac{d_i}{2} m_i^2 \right) \quad (1.23)$$

$$= N \prod_i \sqrt{\frac{2\pi}{d_i}} \exp \left(\frac{(iu_i - m_i d_i)^2}{2d_i} - \frac{d_i m_i^2}{2} \right) \quad (1.24)$$

$$= \frac{1}{\sqrt{\det C \det V}} \prod_i \exp \left(\frac{-u_i^2 + m_i^2 d_i^2 - 2iu_i m_i d_i}{2d_i} - \frac{d_i m_i^2}{2} \right) \quad (1.25)$$

$$= \exp \left(\sum_i \left[-\frac{u_i^2}{2d_i} - iu_i m_i \right] \right) \quad (1.26)$$

$$= \exp \left(-\frac{1}{2} \vec{u}^\top C \vec{u} - i\vec{u} \cdot \vec{m} \right) \quad (1.27)$$

$$= \exp \left(-\frac{1}{2} \vec{k}^\top C \vec{k} - i\vec{k} \cdot \vec{\mu} \right), \quad (1.28)$$

where we have used the expression for the single-variable Gaussian integral:

$$\int dz \exp(-az^2 + bz + c) = \sqrt{\frac{\pi}{a}} \exp \left(\frac{b^2}{4a} + c \right), \quad (1.29)$$

which comes from the one-variable completion of the square:

$$-az^2 + bz + c = -a \left(z - \frac{b}{2a} \right)^2 + \frac{b^2}{4a} + c. \quad (1.30)$$

Also, we used the fact that orthogonal transformation do not change fully-contracted objects, such as scalar products or bilinear forms. \square

Exercise 5

We can calculate the moments of a distribution through its characteristic function:

$$\mathbb{E}[x_\alpha^{n_\alpha} \dots x_\beta^{n_\beta}] = \frac{\partial^{n_\alpha \dots n_\beta} \phi(\vec{k})}{\partial(-ik_\alpha)^{n_\alpha} \dots \partial(-ik_\beta)^{n_\beta}} \Big|_{\vec{k}=0}. \quad (1.31)$$

In the multivariate Gaussian case we can then calculate the mean (component by component) as

$$\mathbb{E}(x_\alpha) = \frac{\partial \phi(\vec{k})}{\partial(-ik_\alpha)} \Big|_{\vec{k}=0} \quad (1.32)$$

$$= \frac{\partial}{\partial(-ik_\alpha)} \Big|_{\vec{k}=0} \exp\left(-\frac{1}{2} \vec{k}^\top C \vec{k} - i \vec{k} \cdot \vec{\mu}\right) \quad (1.33)$$

$$= \left[-i \sum_{\beta} k_\beta C_{\beta\alpha} + \mu_\alpha \right] \exp\left(-\frac{1}{2} \vec{k}^\top C \vec{k} - i \vec{k} \cdot \vec{\mu}\right) \Big|_{\vec{k}=0} \quad (1.34)$$

$$= \mu_\alpha, \quad (1.35)$$

where we used the fact that the differentiation of a symmetric bilinear form is as follows:

$$\frac{\partial}{\partial k_\alpha} \left(\sum_{\beta\gamma} k_\beta k_\gamma C_{\beta\gamma} \right) = 2 \sum_{\beta\gamma} \delta_{\beta\alpha} k_\gamma C_{\beta\gamma} = 2 \sum_{\gamma} k_\gamma C_{\alpha\gamma}. \quad (1.36)$$

The covariance matrix can be computed by linearity as

$$\tilde{C}_{\alpha\beta} = \mathbb{E}\left[(x_\alpha - \mathbb{E}(x_\alpha))(x_\beta - \mathbb{E}(x_\beta))\right] = \mathbb{E}[x_\alpha x_\beta] - \mu_\alpha \mu_\beta, \quad (1.37)$$

the first term of which reads as follows:

$$\mathbb{E}[x_\alpha x_\beta] = \frac{\partial^2 \phi(\vec{k})}{\partial(-ik_\beta) \partial(-ik_\alpha)} \Big|_{\vec{k}=0} \quad (1.38)$$

$$= \frac{\partial}{\partial(-ik_\beta)} \Big|_{\vec{k}=0} \left[-i \sum_{\beta} k_\beta C_{\beta\alpha} + \mu_\alpha \right] \exp\left(-\frac{1}{2} \vec{k}^\top C \vec{k} - i \vec{k} \cdot \vec{\mu}\right) \quad (1.39)$$

$$= C_{\alpha\beta} + \mu_\alpha \mu_\beta, \quad (1.40)$$

therefore, as expected, $\tilde{C}_{\alpha\beta}$ is indeed $C_{\alpha\beta}$.

Exercise 6

Claim 1.2. *The characteristic function of a multivariate Gaussian is, up to normalization, a multivariate Gaussian.*

Proof. The characteristic function is the exponential of (minus)

$$\frac{1}{2}\vec{k}^\top C\vec{k} + i\vec{k} \cdot \vec{\mu} = \frac{1}{2}(\vec{k} + iC^{-1}\vec{\mu})^\top C(\vec{k} + iC^{-1}\vec{\mu}) + \frac{1}{2}\vec{\mu}^\top C^{-1}\vec{\mu}, \quad (1.41)$$

which means that the characteristic function is in the form

$$\phi(\vec{k}) = \text{const} \times \exp\left(-\frac{1}{2}(\vec{k} - \vec{m})^\top C(\vec{k} - \vec{m})\right), \quad (1.42)$$

a multivariate normal with mean $\vec{m} = -iC^{-1}\vec{\mu}$ and covariance matrix C^{-1} , the inverse of the covariance matrix of the corresponding MVN. \square

Exercise 8

For clarity, we denote with Greek indices those ranging from 1 to N , the size of the vector of data; and with Latin indices those ranging from 1 to M , the number of templates.

We are assuming that the data have a Gaussian distribution with a covariance matrix C , and we are modelling their mean μ_α as a sum of templates $t_{i\alpha}$ with coefficients A_i :

$$\mu_\alpha = t_{i\alpha}A_i, \quad (1.43)$$

where the Einstein summation convention has been used. Therefore, the likelihood is proportional to

$$\mathcal{L}(d_\alpha|A_i) \propto \exp\left(-\frac{1}{2}(d_\alpha - A_it_{i\alpha})C_{\alpha\beta}^{-1}(d_\beta - A_jt_{j\beta})\right). \quad (1.44)$$

The normalization only depends on the covariance matrix $C_{\alpha\beta}$, which we assume is fixed. Therefore, maximizing the likelihood² is equivalent to minimizing the χ^2 , which reads

$$\chi^2 = (d_\alpha - A_it_{i\alpha})C_{\alpha\beta}^{-1}(d_\beta - A_jt_{j\beta}). \quad (1.45)$$

We want to minimize this as the amplitudes vary: therefore, we set the derivative with respect to A_k to zero,³

$$\frac{\partial\chi^2}{\partial A_k} = -2t_{k\alpha}C_{\alpha\beta}^{-1}(d_\beta - A_jt_{j\beta}) = 0, \quad (1.47)$$

which means that

$$t_{k\alpha}C_{\alpha\beta}^{-1}d_\beta = (t_{k\alpha}C_{\alpha\beta}^{-1}t_{j\beta})A_j, \quad (1.48)$$

² Which is equivalent to maximizing the posterior if we are using a flat prior.

³ The fact that the stationary point we will find is indeed a minimum can be checked by looking at the second derivative of χ^2 :

$$\frac{\partial^2\chi^2}{\partial A_k\partial A_m} = 2t_{k\alpha}C_{\alpha\beta}^{-1}t_{m\beta}, \quad (1.46)$$

and recalling that the inverse of the covariance matrix is positive definite.

a linear system of M equations (indexed by k) in the M variables A_j . If we denote the evaluations of bilinear forms in the data (N -dimensional) space with brackets, as $a_\alpha C_{\alpha\beta} b_\beta \stackrel{\text{def}}{=} (a|C|b)$, this reads

$$(t|C^{-1}|d)_k = (t|C^{-1}|t)_{kj} A_j \quad (1.49)$$

$$\left[(t|C^{-1}|t)^{-1} \right]_{mk} (t|C^{-1}|d)_k = \underbrace{\left[(t|C^{-1}|t)^{-1} \right]_{mk} (t|C^{-1}|t)_{kj} A_j}_{=\delta_{mj}} = A_m \quad (1.50)$$

$$A_m = \left[(t|C^{-1}|t)^{-1} \right]_{mk} (t|C^{-1}|d)_k, \quad (1.51)$$

where the inverse of $(t|C^{-1}|t)$ is to be computed in the M -dimensional vector space.

Exercise 9

Our model for the mean value is in the form $\mu(\Theta, A) = A\bar{x}(\Theta)$, where \bar{x} is a generic function of Θ , while A is our scale parameter.⁴ Our likelihood then reads

$$\mathcal{L}(x|\Theta, A) = \underbrace{\frac{1}{(2\pi)^{N/2} \sqrt{\det C}}}_{B_1} \exp\left(-\frac{1}{2}(x - A\bar{x}(\Theta))^\top C^{-1}(x - A\bar{x}(\Theta))\right). \quad (1.52)$$

If the priors for both A and Θ are flat, this corresponds to the joint posterior $P(\Theta, A|x)$. We want to marginalize over A , which amounts to integrating over it: dropping the dependence on Θ of \bar{x} and defining $V = C^{-1}$ we find

$$P(\Theta|x) = B_1 \int \exp\left(-\frac{1}{2}(x - A\bar{x})^\top V(x - A\bar{x})\right) dA \quad (1.53)$$

$$= B_1 \int \exp\left(-\frac{1}{2}\left(x^\top Vx - 2A\bar{x}^\top Vx + A^2\bar{x}^\top V\bar{x}\right)\right) dA. \quad (1.54)$$

Used the symmetry of V .

The amplitude being negative makes little sense in a typical physical context, however the Gaussian integral can be done analytically only over the whole of \mathbb{R} .

In order to get analytical results, here we will marginalize by integrating over negative amplitudes as well ($A \in \mathbb{R}$); the last figure (1) will show how only integrating over positive amplitudes only would have looked (by numerical calculation) in a simple case. In general if one wishes to perform the integral over $A \in (0, +\infty)$ the tabulated values of the error function may be used.

Applying the formula for the single-variable Gaussian integral (1.29) (the bilinear forms are all evaluated to yield scalars, we are only integrating over the scalar A !) we then get

$$P(\Theta|x) = \underbrace{B_1 \exp\left(-\frac{1}{2}x^\top Vx\right)}_{B_2} \exp\left(\frac{1}{2} \frac{(\bar{x}^\top Vx)^2}{(\bar{x}^\top V\bar{x})}\right) \sqrt{\frac{2\pi}{\bar{x}^\top V\bar{x}}} \quad (1.55)$$

⁴ This is not specified in the problem, but it seems natural to think that $|\bar{x}(\Theta)|$ is a constant for varying Θ .

$$= B_2 \sqrt{\frac{2\pi}{\bar{x}^\top V \bar{x}}} \exp\left(\frac{1}{2} \frac{\bar{x}^\top \Omega \bar{x}}{\bar{x}^\top V \bar{x}}\right), \quad (1.56)$$

where we defined the bilinear form $\Omega = V x x^\top V^\top$.⁵

An application of posterior marginalization in this fashion

Let us consider a simple example of this as a sanity check: suppose that x is two-dimensional, and $\bar{x}(\Theta) = (\cos \Theta, \sin \Theta)^\top$; further, suppose that V is diagonal, so that

$$V = \begin{bmatrix} \sigma_x^{-2} & 0 \\ 0 & \sigma_y^{-2} \end{bmatrix}. \quad (1.57)$$

Also, suppose that the observed data parameter is

$$x = A_x \begin{bmatrix} \cos \varphi \\ \sin \varphi \end{bmatrix}. \quad (1.58)$$

Then, the multiplicative constant in front of the marginalized posterior reads

$$B_2 = B_1 \exp\left(-\frac{1}{2} A_x^2 \left(\frac{\cos^2 \varphi}{\sigma_x^2} + \frac{\sin^2 \varphi}{\sigma_y^2}\right)\right); \quad (1.59)$$

while the bilinear form Ω is

$$\Omega = A_x^2 \begin{bmatrix} \sigma_x^{-2} & 0 \\ 0 & \sigma_y^{-2} \end{bmatrix} \begin{bmatrix} \cos^2 \varphi & \cos \varphi \sin \varphi \\ \cos \varphi \sin \varphi & \sin^2 \varphi \end{bmatrix} \begin{bmatrix} \sigma_x^{-2} & 0 \\ 0 & \sigma_y^{-2} \end{bmatrix} \quad (1.60)$$

$$= A_x^2 \begin{bmatrix} \cos^2 \varphi / \sigma_x^4 & \cos \varphi \sin \varphi / \sigma_x^2 \sigma_y^2 \\ \cos \varphi \sin \varphi / \sigma_x^2 \sigma_y^2 & \sin^2 \varphi / \sigma_y^4 \end{bmatrix}. \quad (1.61)$$

Then, when we evaluate the marginalized posterior we will find something in the form

$$P(\Theta|x) = B_1 \sqrt{2\pi} \left(\frac{\cos^2 \Theta}{\sigma_x^2} + \frac{\sin^2 \Theta}{\sigma_y^2}\right)^{-1/2} \exp\left(A_x^2 F(\Theta, \varphi)\right), \quad (1.62)$$

where $F(\Theta, \varphi)$ is some function whose specific form does not really matter.⁶

The amplitude of the observed data vector, A_x , appears in a rather simple way, as a multiplicative prefactor in the exponent: it can affect the shape of the distribution, but not its mean. Specifically, we can see that scaling A_x is equivalent to scaling σ_x and σ_y simultaneously in the opposite direction — this is rather intuitive, since the angular size of the distribution as seen from the origin is smaller if it is further away.

⁵ With explicit indices, $\Omega_{im} = V_{ij} x_j x_k V_{km}$.

⁶ For completeness, here is the full expression:

$$F(\Theta, \varphi) = -\frac{1}{2} \left(\frac{\cos^2 \varphi}{\sigma_x^2} + \frac{\sin^2 \varphi}{\sigma_y^2}\right) + \left(\frac{\cos^2 \Theta}{\sigma_x^2} + \frac{\sin^2 \Theta}{\sigma_y^2}\right)^{-1} \left[\frac{\cos^2 \Theta \cos^2 \varphi}{\sigma_x^4} + 2 \frac{\cos \Theta \sin \Theta \cos \varphi \sin \varphi}{\sigma_x^2 \sigma_y^2} + \frac{\sin^2 \Theta \sin^2 \varphi}{\sigma_y^4} \right]. \quad (1.63)$$

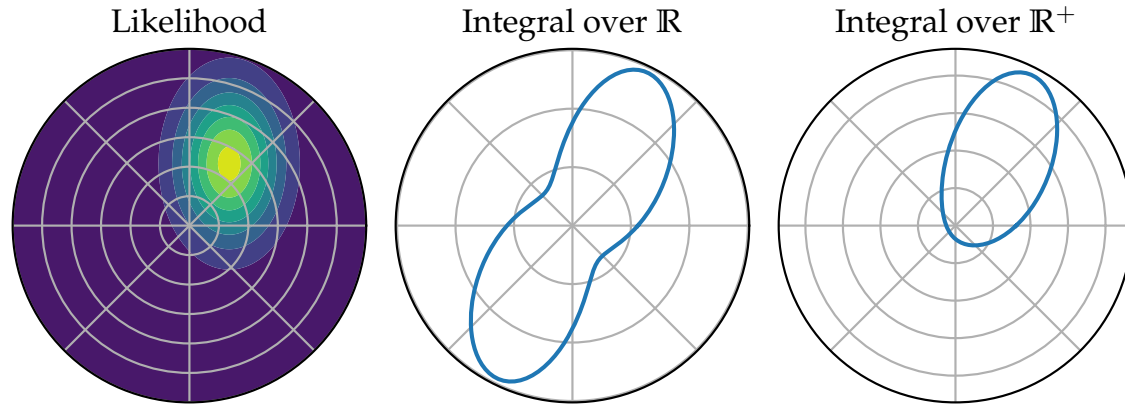


Figure 1: Marginalization: the left plot shows the full likelihood in terms of A and Θ ; the middle plot shows the result of marginalization as shown in the previous calculation (the posterior as a function of Θ); the right plot shows the result of the more physically meaningful marginalization over $A \in (0, +\infty)$ only. Here the likelihood is a diagonal Gaussian with $\sigma_x = 1.2$ and $\sigma_y = 1.8$, centered in $A_x = 2.5$ and $\varphi = 1$ rad.

Likelihood marginalization

So far we have considered the posterior $P(\Theta|x)$, the marginalized posterior, a function of the parameter(s) Θ ; however we may also be interested in the marginalized likelihood $\mathcal{L}(x|\Theta)$, whose expression is the same as the one we found for $P(\Theta|x)$. Let us write it in a way which makes the dependence on x more explicit:

$$\mathcal{L}(x|\Theta) = \underbrace{B_1 \sqrt{\frac{2\pi}{\bar{x}^\top V \bar{x}}}}_{B_3} \exp\left(-\frac{1}{2}x^\top V x + \frac{1}{2} \frac{(\bar{x}^\top V x)^2}{\bar{x}^\top V \bar{x}}\right), \quad (1.64)$$

which can be simplified by making use of the fact that the best-fit template amplitude we found in the last exercise (equation (1.51)) can be applied here, with the single template $t = \bar{x}$, the single amplitude A , the data $d = x$, and the inverse covariance matrix $C^{-1} = V$:

the fitting value for A is

$$\hat{A} = \frac{\bar{x}^\top V x}{\bar{x}^\top V \bar{x}}; \quad (1.65)$$

therefore the likelihood is

$$\mathcal{L}(x|\Theta) = B_3 \exp\left(-\frac{1}{2}x^\top V x + \frac{1}{2}\hat{A}\bar{x}^\top V x\right). \quad (1.66)$$

This can be rewritten in the canonical MVN form by making use of the matrix square completion formula (1.8), with $A = -V$ and $\bar{b}^\top = \hat{A}\bar{x}^\top V$:

$$\begin{aligned} -\frac{1}{2}x^\top V x + \frac{1}{2}\hat{A}\bar{x}^\top V x &= -\frac{1}{2}\left(x - \frac{1}{2}V^{-1}\hat{A}(\bar{x}^\top V)^\top\right)^\top V \left(x - \frac{1}{2}V^{-1}\hat{A}(\bar{x}^\top V)^\top\right) \\ &\quad + \frac{1}{8}\hat{A}^2(\bar{x}^\top V)V^{-1}(\bar{x}^\top V)^\top \end{aligned} \quad (1.67)$$

$$= -\frac{1}{2}\left(x - \frac{1}{2}\hat{A}\bar{x}\right)^\top V \left(x - \frac{1}{2}\hat{A}\bar{x}\right) + \frac{1}{8}\hat{A}^2\bar{x}^\top V \bar{x}. \quad (1.68)$$

Therefore, the marginalized likelihood reads

$$\mathcal{L}(x|\Theta) = B_3 \exp\left(\frac{1}{8}\hat{A}^2\bar{x}^\top V \bar{x}\right) \exp\left(-\frac{1}{2}\left(x - \frac{1}{2}\hat{A}\bar{x}\right)^\top V \left(\frac{1}{2}x - \hat{A}\bar{x}\right)\right). \quad (1.69)$$

We must be careful with this expression: it looks like a multivariate normal in \bar{x} , however \hat{A} is not in general independent of it. If neither the estimate for \hat{A} nor the expression $\bar{x}^\top V \bar{x}$ vary significantly in the range of \bar{x} we are interested in, then we can consider this expression for the likelihood as Gaussian in x , or as a marginalized posterior which is Gaussian in \bar{x} .

2 December exercises

Exercise 10

We have a time series of N data points, $D = \{d_i\}$, corresponding to the times t_i , which are separated by the constant spacing Δ .

We model them as

$$d_i = \underbrace{B_1 \cos(\omega t_i) + B_2 \sin(\omega t_i)}_{f(t_i)} + n_i, \quad (2.1)$$

where $f(t)$ the signal we want to characterize, which depends on the unknown amplitudes B_1 and B_2 and the unknown frequency ω ; while n_i is the noise: each n_i is i.i.d. as a zero-mean Gaussian with known variance σ^2 .

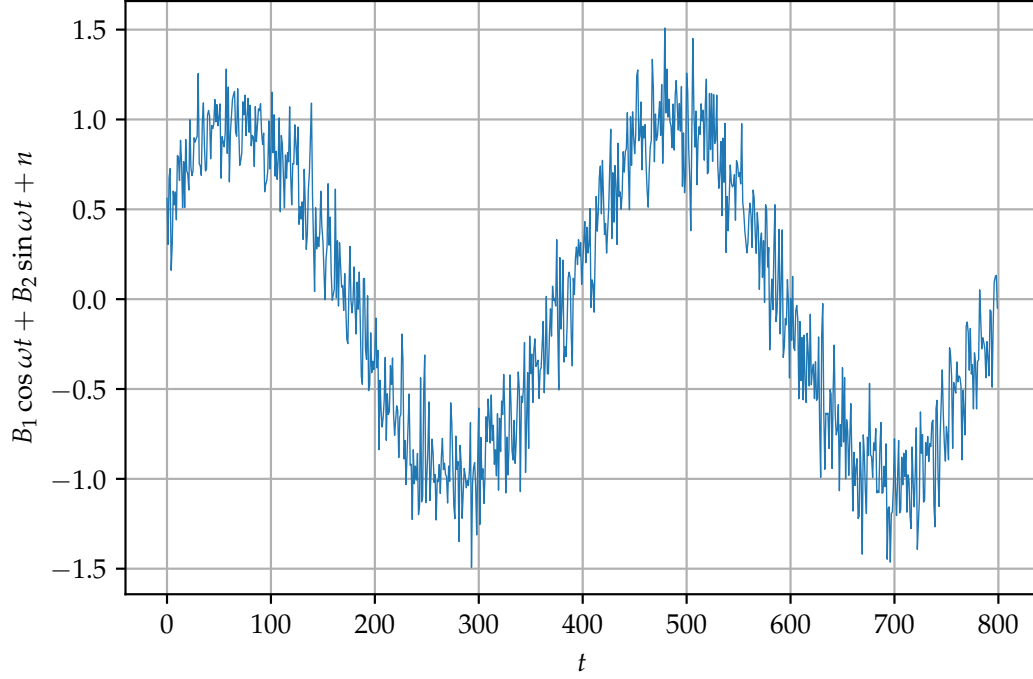


Figure 2: Example of a portion of the signal the model assumes; the true signal should be longer than a few periods, and it is taken to be 20 times as long as this in the numerical examples which follow. The parameters are $B_1 = 1/2$, $B_2 = \sqrt{3}/2$, $\sigma = 0.2$, $\Delta = 1$, $\omega = 0.015$.

The full likelihood

The likelihood of a single datum of index i attaining the value d_i is given⁷ by

$$\mathcal{L}(d_i|\omega, B_1, B_2) = \frac{1}{\sqrt{2\pi}\sigma} \exp\left(-\frac{1}{2\sigma^2}(d_i - f(t_i))^2\right). \quad (2.2)$$

Now, since the noise at each point is independent, the full likelihood is the product of the likelihoods of each datum:

$$\mathcal{L}(D|\omega, B_1, B_2) = \frac{1}{(\sqrt{2\pi}\sigma)^N} \prod_{i=1}^N \exp\left(-\frac{1}{2\sigma^2}(d_i - f(t_i))^2\right) \quad (2.3)$$

$$= \frac{1}{(\sqrt{2\pi}\sigma)^N} \exp\left(-\frac{1}{2\sigma^2} \sum_{i=1}^N (d_i - f(t_i))^2\right) \quad (2.4)$$

⁷ Omitting the dependence on previous information for simplicity.

$$= \frac{1}{(\sqrt{2\pi}\sigma)^N} \exp \left(-\frac{1}{2\sigma^2} \underbrace{\sum_{i=1}^N (d_i - B_1 \cos(\omega t_i) - B_2 \sin(\omega t_i))^2}_Q \right). \quad (2.5)$$

Let us manipulate the sum in the exponent, which we denote as Q :

$$Q = \sum_i d_i^2 - 2 \sum_i d_i (B_1 \cos(\omega t_i) + B_2 \sin(\omega t_i)) + \sum_i (B_1 \cos(\omega t_i) + B_2 \sin(\omega t_i))^2 \quad (2.6)$$

$$\begin{aligned} &= N\bar{d}^2 - 2B_1 \underbrace{\sum_i d_i \cos(\omega t_i)}_{R_1(\omega)} - 2B_2 \underbrace{\sum_i d_i \sin(\omega t_i)}_{R_2(\omega)} + \\ &\quad + B_1^2 \underbrace{\sum_i \cos^2(\omega t_i)}_c + B_2^2 \underbrace{\sum_i \sin^2(\omega t_i)}_s + 2B_1 B_2 \sum_i \cos(\omega t_i) \sin(\omega t_i) \end{aligned} \quad (2.7)$$

$$= N\bar{d}^2 - 2B_1 R_1(\omega) - 2B_2 R_2(\omega) + B_1^2 c + B_2^2 s + B_1 B_2 \underbrace{\sum_i \sin(2\omega t_i)}_h. \quad (2.8)$$

Large pulsation limit

The condition we ask of our signal is to have many data points for each period ($\Delta \ll \omega^{-1}$) and many sampled periods ($N\Delta \gg \omega^{-1}$), which is equivalent to $(N\Delta)^{-1} \ll \omega \ll \Delta^{-1}$.

The values of c , s and h for various values of ω are shown in figure 3.

Let us then assume that we are working in that region. Then, denoting the argument of the functions as $x_i = \omega t_i \bmod 2\pi$, we will have:

$$c = \sum_{i=1}^N \cos^2 x_i \approx N \langle \cos^2 x \rangle_{\text{period}} = \frac{N}{2} \quad (2.9)$$

$$s = \sum_{i=1}^N \sin^2 x_i \approx N \langle \sin^2 x \rangle_{\text{period}} = \frac{N}{2} \quad (2.10)$$

$$h = \sum_{i=1}^N \sin(2x_i) \approx N \langle \sin(2x) \rangle_{\text{period}} = 0, \quad (2.11)$$

since x will be approximately uniformly distributed in the $[0, 2\pi)$ interval. The figure also shows this; the small deviations from these values are due to non-integer amounts of periods being included in the sample, but this is an edge effect, which becomes negligible as N becomes very large.

Marginalization

With these simplifications, the likelihood looks like

$$\mathcal{L}(D|\omega, B_1, B_2) = \frac{1}{(\sqrt{2\pi}\sigma)^N} \exp \left(-\frac{Q}{2\sigma^2} \right) \quad (2.12)$$

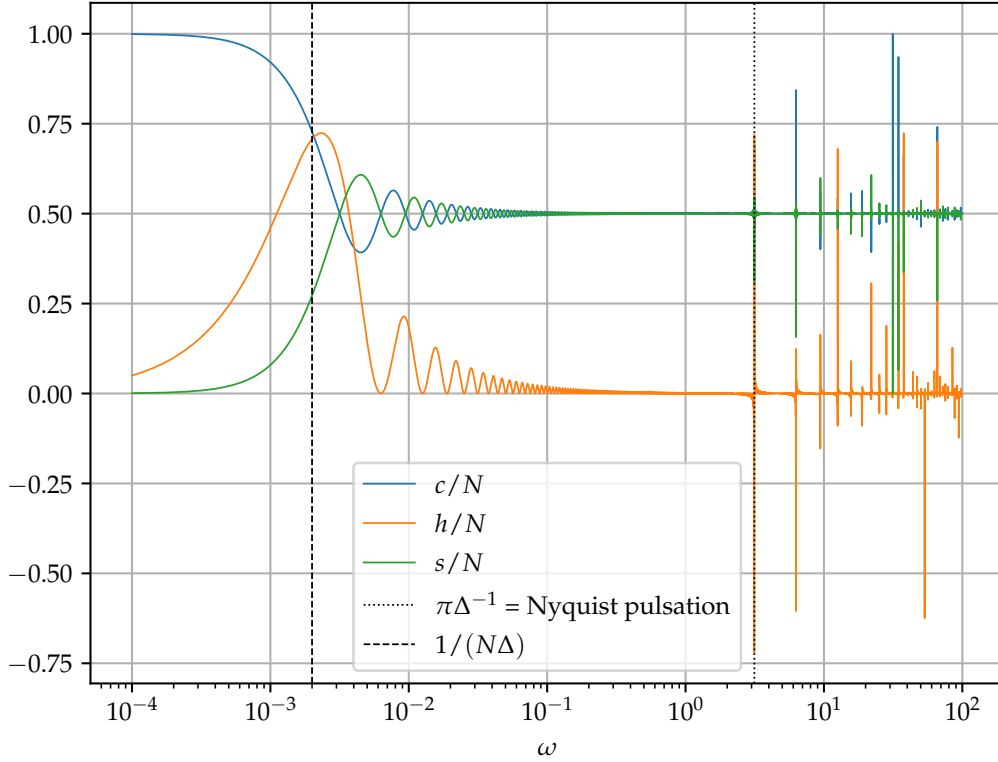


Figure 3: Values of c , s and h for different orders of magnitude of ω . The sampling frequency is fixed at $\Delta^{-1} = 1$, and $N = 500$ points are always sampled. The conditions we ask, in the context of this plot, would mean that we put ourselves roughly in the middle of the two marked vertical lines, in the region $\omega \sim 10^{-1}$.

$$Q = N\bar{d}^2 - 2B_1R_1(\omega) - 2B_2R_2(\omega) + B_1^2\frac{N}{2} + B_2^2\frac{N}{2} \quad (2.13)$$

$$= N\left(\bar{d}^2 + \frac{B_1^2 + B_2^2}{2}\right) - 2B_1R_1(\omega) - 2B_2R_2(\omega). \quad (2.14)$$

The posterior is proportional to the likelihood, since we are assuming the priors on ω and B_i are uniform. We wish to marginalize it over the parameters $B_i \in \mathbb{R}$, for $i = 1, 2$. This amounts to solving the integral

$$P(\omega|D) \propto \int_{\mathbb{R}^2} dB_1 dB_2 P(\omega, B_1, B_2|D) \quad (2.15)$$

$$\propto \int_{\mathbb{R}^2} dB_1 dB_2 \exp\left(-\frac{N}{2\sigma^2} \left(\underbrace{\bar{d}^2}_{\text{constant}} + \frac{B_1^2 + B_2^2}{2} - 2B_1R_1(\omega) - 2B_2R_2(\omega)\right)\right) \quad (2.16)$$

$$\propto \int_{\mathbb{R}^2} dB_1 dB_2 \exp \left(-\frac{1}{2\sigma^2} \left(\sum_i \frac{NB_i^2}{2} - 2B_i R_i \right) \right) \quad (2.17)$$

$$\propto \prod_i \int_{\mathbb{R}} dB_i \exp \left(-\frac{NB_i^2}{4\sigma^2} + \frac{B_i R_i}{\sigma^2} \right) \quad (2.18)$$

$$\propto \prod_i \sqrt{\frac{\pi}{N/(4\sigma^2)}} \exp \left(\frac{R_i^2}{\sigma^4} \frac{1}{4} \frac{4\sigma^2}{N} \right) \quad (2.19)$$

$$\propto N^{-1} \prod_i \exp \left(\frac{R_i^2}{\sigma^2 N} \right) \quad (2.20)$$

$$\propto N^{-1} \exp \left(\frac{R_1^2(\omega) + R_2^2(\omega)}{\sigma^2 N} \right). \quad (2.21)$$

In the last step we have used the usual expression for a univariate Gaussian integral (1.29).

Since the exponential is monotonic and we are keeping σ and N constant, the Maximum A-Posteriori (MAP) estimate is given by the maximum of $R_1^2(\omega) + R_2^2(\omega)$.

The periodogram

The periodogram C is defined as

$$C(\omega) = \frac{2}{N} \left| \sum_{k=1}^N d_k \exp(-i\omega t_k) \right|^2, \quad (2.22)$$

and while this definition could be applied for an arbitrary set of times t_k , we will only consider it for evenly spaced times $t_k = k\Delta + t_0$ for some t_0 : a discrete-time Fourier transform.

We can rewrite the periodogram as

$$C(\omega) = \frac{2}{N} \left| \sum_{k=1}^N d_k (\cos(\omega t_k) - i \sin(\omega t_k)) \right|^2 \quad (2.23)$$

$$= \frac{2}{N} \left[\left(\sum_{k=1}^N d_k \cos(\omega t_k) \right)^2 + \left(\sum_{k=1}^N d_k \sin(\omega t_k) \right)^2 \right] \quad (2.24)$$

$$= \frac{2}{N} [R_1^2(\omega) + R_2^2(\omega)]. \quad (2.25)$$

Therefore, the value of ω which maximizes $C(\omega)$ is the same which maximizes $R_1^2(\omega) + R_2^2(\omega)$, which is the MAP estimate.

Least-squares fitting

Least-squares fitting the sinusoid with the same model means we minimize $\chi^2 = Q/\sigma^2$. This is precisely equivalent to the MAP estimate for the full likelihood, which under the aforementioned conditions can be estimated through the maximum of $R_1^2(\omega) + R_2^2(\omega)$.

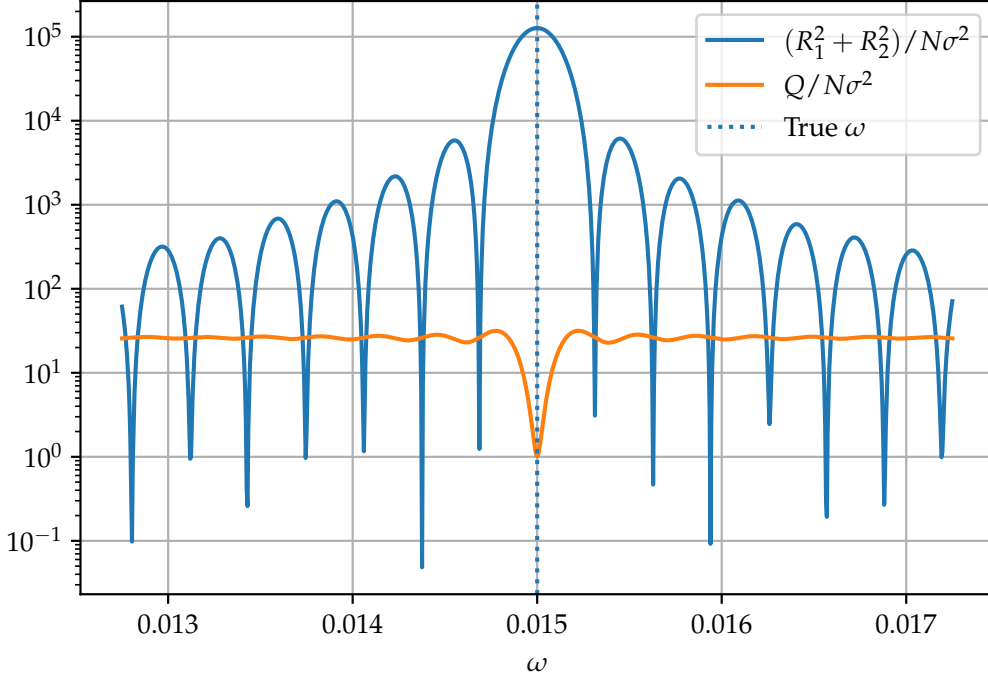


Figure 4: N data points d_i are generated with the same distribution as the theoretical model (see figure 2). The value of Q is computed by fixing B_1 and B_2 to their true values, which is why its extremal point is sharper: it assumes more knowledge than the alternative, which is what we must maximize after we marginalize over all possible amplitudes $B_{1,2}$.

As is shown in figure 4, the maximum of $R_1^2 + R_2^2$ and the minimum of Q do indeed coincide. In fact, Q can also be computed for different values of B_1 and B_2 , and it attains its global minimum near the true values of the whole triple (ω, B_1, B_2) .

This procedure would yield a Gaussian likelihood for ω under the following (sufficient) conditions:

1. i.i.d. Gaussian noise on each data point;
2. linear dependence of the model $f(t)$ on its parameter ω .

The first condition is satisfied under our hypotheses, the second is not unless the entire data range lies in a very small region: $N\Delta \ll \omega^{-1}$, in which case the model can be approximated to linear order; for example, if the region is near the origin we can approximate it as $f(t) \sim B_1 + B_2\omega t$.

In the general (and typical) case, instead, the likelihood $\mathcal{L} \propto \exp(-Q/\sigma^2)$ is not Gaussian, and the model is not linear; however the minimum of Q is still rather peaked, and as long as we start near it we can numerically find the MAP, or equivalently least-squares, estimate of ω .