## Output\_BBH\_Properties

February 4, 2025

#### 1 Output Binary Black Hole Properties

In paper Zhu et al. 2025, we investigate how the binary companion influences the superradiance gravitational atom (GA) in reality. By obtaining the SEVN output for the Binary Black Hole (BBH) remnant properties, we are able to calculate both the superradiance rate and the corrections to the superradiance rate. Details of SEVN simulation package can be found in

```
Zhu et al. 2025: https://doi.org/10.48550/arXiv.2409.14159
Spera et al. 2018: https://doi.org/10.1093/mnras/stz359
```

Begin by reviewing the output data. We have evolved a total of around 3.0\*10^7 stellar binaries and selected all the remnants where both remnants are black holes. The simulation output includes the mass of the primary (Mass\_0, in solar mass) and secondary BHs (Mass\_1, in solar masses), the semimajor axis (Semimajor) for the orbit, and the eccentricity (Eccentricity) of the orbit. The remnant type equals to 6 indicates that the remnant is a BH.

```
[1]: # Constants for unit transformation

Msun=1.9884099* 10**30  # to kg
eVtokg=1.7826619216278975e-36 # to kg
Rsun=6.957*10**8  # to meter
Myr=3.1536e13  # to second
```

```
[2]: import warnings warnings.filterwarnings("ignore")
```

```
[3]: import pandas as pd

# Read the data and save it in ddf

ddf=pd.read_csv('/home/hzhuav/SEVN_Result/Output1.dat',delimiter=' ',header=0)
ddf
```

```
[3]:
                        ID
                                          name
                                                    {\tt Mass\_0}
                                                              RemnantType_0
                                                                                    \mathtt{Mass}_{\mathtt{1}}
     0
                         6
                                                 13.704040
                                                                                16.971260
                            402748869422646
                                                                            6
     1
                        12
                            843618165177474
                                                 18.667630
                                                                                19.054640
     2
                            999940604697875
                                                 45.106340
                                                                            6
                                                                                13.726090
     3
                     1004
                           844142907982993
                                                  7.081709
                                                                                 5.750017
     4
                     3005
                            402270096778797
                                                28.347130
                                                                            6
                                                                                 6.560349
```

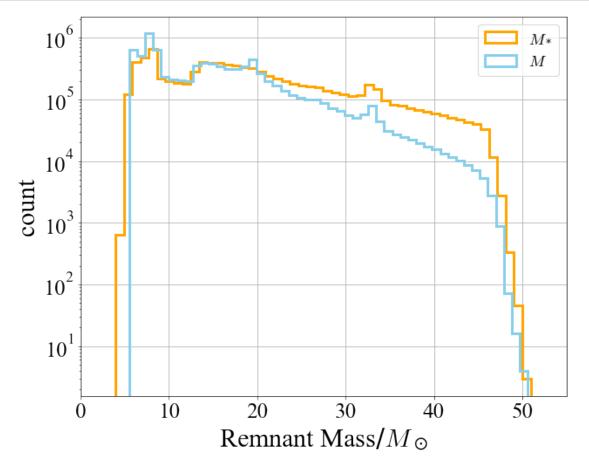
```
8602942 28983152 997087945223631 27.932890
                                                               5.810605
8602943 28983162 190633857428840
                                    7.485959
                                                               9.539038
                                                           6 26.930120
8602944 28984146 407767383028673 33.464130
8602945 28984161 982373687732211
                                     6.162466
                                                               7.684832
8602946 28985074 143475573419797 34.138350
                                                           6 30.229370
        RemnantType_1
                        Semimajor Eccentricity
                       87494.2300
0
                                        0.023631
1
                       25329.4300
                                        0.235629
2
                     6
                         761.9675
                                        0.008499
3
                     6
                          152.7133
                                        0.353820
4
                        2531.5220
                                        0.206282
8602942
                     6
                        2024.3770
                                       0.651537
8602943
                     6
                         225.9608
                                        0.129645
8602944
                     6 15676.7800
                                        0.415756
8602945
                     6
                          21.9911
                                        0.036108
                     6 76240.8400
8602946
                                        0.289955
```

[8602947 rows x 8 columns]

The distribution plot for semimajor axis, BH mass and eccentricity for BBHs obtained from SEVN simulation are as follows:

```
[5]: | # Plot of mass distribution of primary and secondary remnant
     import numpy as np
     import matplotlib
     import matplotlib.pyplot as plt
     from matplotlib.font manager import FontProperties
     # Font setting part, you can delete it.
     font1=FontProperties(fname=r'/usr/share/fonts/truetype/Times/times.ttf',size=30)
     font2=FontProperties(fname=r'/usr/share/fonts/truetype/Times/times.ttf',size=25)
     plt.rcParams["font.family"] = 'DejaVu Sans'
     matplotlib.rcParams['mathtext.fontset'] = 'custom'
     matplotlib.rcParams['mathtext.rm'] = 'STIXSizeFourSym'
     matplotlib.rcParams['mathtext.it'] = 'STIXSizeFourSym'
     #Primary and seconday BH mass in solar mass
     M1=ddf.Mass O.to numpy()
     M2=ddf.Mass_1.to_numpy()
     plt.figure(figsize=(10,8))
     plt.
     →hist(M1,label='M1',lw=3,bins=50,color='orange',edgecolor='orange',histtype='step')
     plt.
      →hist(M2,label='M2',lw=3,bins=50,color='skyblue',edgecolor='skyblue',histtype='step')
     plt.legend(['$M *$','$M$'],fontsize=18)
```

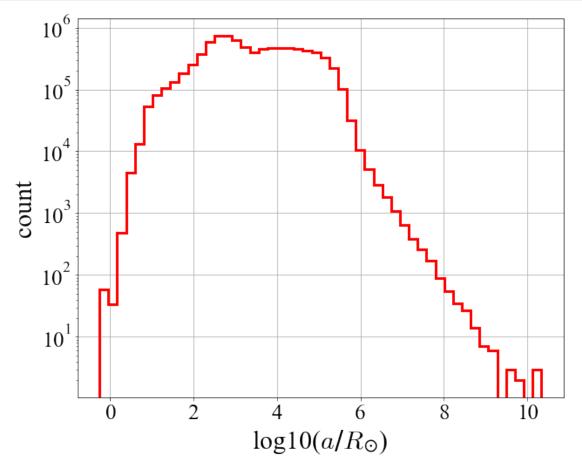
```
plt.grid(True)
plt.xlabel(r'Remnant Mass$/M_\odot$',fontproperties=font1)
plt.ylabel('count',fontproperties=font1)
plt.yscale("log")
plt.xticks(fontproperties=font2)
plt.yticks(fontproperties=font2)
plt.xlim((0,55))
plt.show()
```



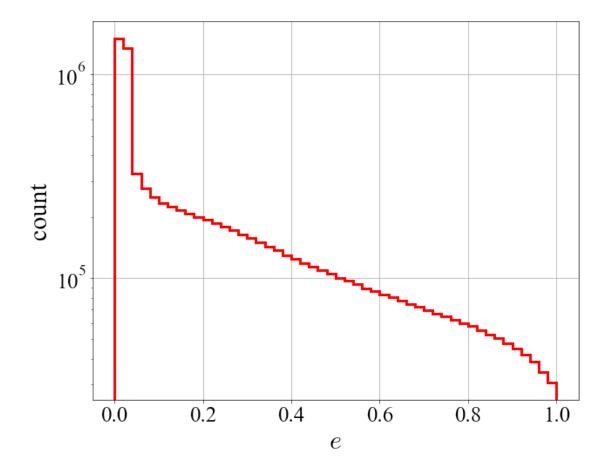
```
[6]: # Semi major axis for BBH

a = ddf.Semimajor * 696340*10**3 # in meters
a_solar =ddf.Semimajor # in solar radius
plt.figure(figsize=(10,8))
plt.hist(np.log10(a_solar),lw=3,bins=50,edgecolor='red',histtype='step')
plt.yscale('log')
plt.grid(True)
plt.xlabel(r'log10'+'('+'$a/R_{\odot}$'+')',fontproperties=font1)
plt.ylabel('count',fontproperties=font1)
```

```
plt.xticks(fontproperties=font2)
plt.yticks(fontproperties=font2)
plt.show()
```



```
[7]: # Eccentricity distribution
    e =ddf.Eccentricity
    plt.figure(figsize=(10,8))
    plt.hist(e,lw=3,bins=50,edgecolor='red',histtype='step')
    plt.yscale('log')
    plt.grid(True)
    plt.xlabel(r'$e$',fontproperties=font1)
    plt.ylabel('count',fontproperties=font1)
    plt.xticks(fontproperties=font2)
    plt.yticks(fontproperties=font2)
    plt.show()
```



# 2 Spin Model

To compute the effective superradiance rate as introduced in our paper, we still require the BH spin. While SEVN includes several spin models, the actual BH spin model remains highly uncertain. Generally, BBHs with small separations tend to possess significant spin due to efficient accretion. However, superradiance termination effect greatly depends on both spin and the binary separation. Consequently, we initially explore a pre-existing Wolf-Rayet model and then investigate our Plateau (PT) model, inspired by the WR model but designed with adjustable parameters to encompass the uncertainties in spin models.

Details of WR spin model can be found in

Bavera et al. 2021: https://iopscience.iop.org/article/10.3847/2515-5172/ac053c

From the paper, we found that the spin of the secondary BH  $\tilde{a}$  is well approximated by

$$\tilde{a} = \begin{cases} f^{\alpha} \log_{10}^{2} \left(\frac{P}{\text{day}}\right) + f^{\beta} \log_{10} \left(\frac{P}{\text{day}}\right) &, \ 0.1 \leq \frac{P}{\text{day}} \leq 1 \\ 0 &, \ \frac{P}{\text{day}} > 1 \\ \tilde{a}|_{P=0.1 \, \text{day}} &, \ \frac{P}{\text{day}} < 0.1 \end{cases}$$
 where  $f^{(\alpha,\beta)} = -c_{1}^{(\alpha,\beta)} \left[c_{2}^{(\alpha,\beta)} + \exp\left(-c_{3}^{(\alpha,\beta)} M/M_{\odot}\right)\right]$ , with coefficients  $c_{1}^{(\alpha,\beta)}$ ,  $c_{2}^{(\alpha,\beta)}$ , and  $c_{3}^{(\alpha,\beta)}$  determined from both some proposition (see Proposition 4.2021 for every data; is)

determined from least-square regression (see Bavera et al. 2021 for more details).

```
[8]: # WR Spin Model
     # MODELLING THE BH SPIN #
     ###########################
     # The WR spin model code we generally use based on the equation above
     def model(m_WR,T,state='c_depletion'):
         '''modelling the spin of the black hole assuming that it is forming from a_{\sqcup}
      \hookrightarrow WR star (Equation 5)
         INPUT:
             m_WR = mass of the Wolf Rayet star (can be approximanted with the BH)
             T = period in days
             state = c_deplation/He_deplation change slightly the parameters
         OUTPUT:
             spin of the black hole'''
         if state == 'c_depletion':
             c1a = 0.051237
             c2a = 0.029928
             c3a = 0.282998
             c1b = 0.027090
             c2b = 0.010905
             c3b = 0.422213
         elif state == 'he_depletion':
             c1a = 0.059305
             c2a = 0.035552
             c3a = 0.270245
             c1b = 0.026960
             c2b = 0.011001
             c3b = 0.420739
         else:
             raise ValueError('state not supported!')
         \#a_BH2(T >= 1.) = 0
         a_BH2 = np.zeros(len(T))
         def constant(m_WR, c1, c2, c3):
```

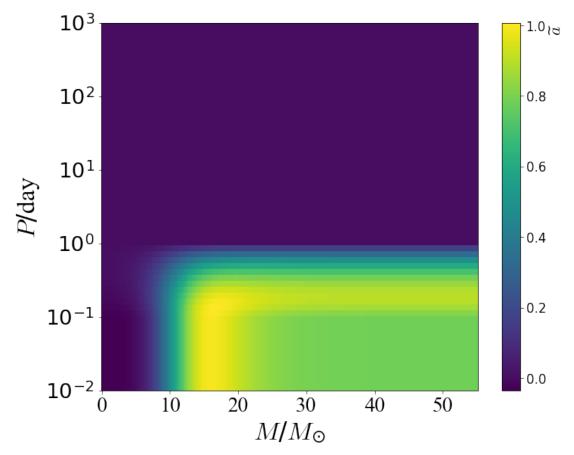
```
return -c1/(c2+np.exp(-c3*m_WR))
    alpha = constant(m_WR[(0.1 \le T)\&(T \le 1.)], c1a, c2a, c3a)
    beta = constant(m_WR[(0.1 \le T)\&(T \le 1.)], c1b, c2b, c3b)
    a_BH2[(0.1<=T)\&(T<1.)] = alpha*np.log10(T[(0.1<=T)\&(T<1.)])**2+beta*np.
 \rightarrow \log 10 (T[(0.1 \le T) \& (T \le 1.)])
    a_BH2[T<0.1] = constant(m_WR[T<0.1]), c1a, c2a, c3a)-constant(m_WR[T<0.1], __
 \rightarrowc1b, c2b, c3b)
    return a_BH2
# This code is defined only for specific plot
def modelnew(m_WR,T):
    c1a = 0.059305
    c2a = 0.035552
    c3a = 0.270245
    c1b = 0.026960
    c2b = 0.011001
    c3b = 0.420739
    def constant(m_WR, c1, c2, c3):
        return -c1/(c2+np.exp(-c3*m_WR))
    alpha = constant(m_WR, c1a, c2a, c3a)
    beta = constant(m_WR, c1b, c2b, c3b)
    a BH2= alpha*T**2+beta*T
    return a_BH2
```

The relation between the mass, spin of the secondary BH, and the binary period is

```
[9]: # Plot of the relation between secondary spin, BH mass and binary period based
import matplotlib.cm as cm
matplotlib.rc('xtick', labelsize=15)
matplotlib.rc(ytick', labelsize=15)
matplotlib.rcParams.update({'font.size': 20})
matplotlib.rcParams['text.usetex'] = False

x = np.linspace(0,55,200)
y = np.linspace(-2, 14,200)
X,Y=np.meshgrid(x,y)
Z=modelnew(X,Y)*(Y<0)*(Y>-1)+modelnew(X,-1)*(Y<=-1)

plt.figure(figsize=(10,8))
```



Inspired by the WR model, we choose to design our own spin model with different parameters, which are to be varied in the later analysis of superradiance termination. Consider

$$\tilde{a} = \frac{1}{4} \operatorname{erfc} \left( \frac{\ln P / P_C}{\sqrt{2} w_P} \right) \operatorname{erfc} \left( \frac{M_C - M}{\sqrt{2} w_M} \right) ,$$

where  $\operatorname{erfc}(x)$  is the complementary error function. This PT model is an emulation of the WR model

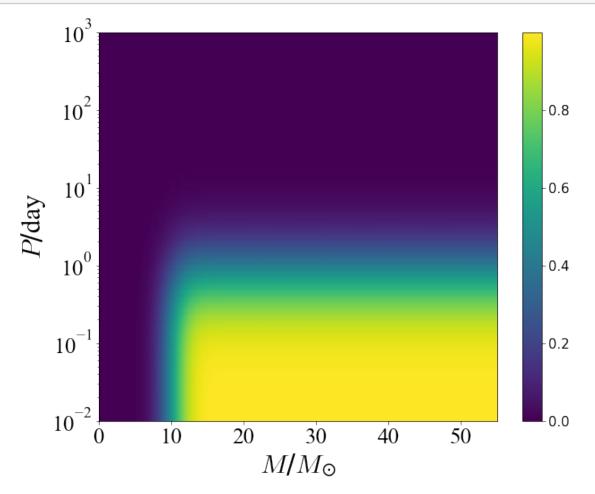
with its essential features preserved. The PT model is characterized by a plateau region with two cutoffs ( $P_C$  and  $M_C$ ), along with two width parameters ( $w_P$  and  $w_M$ ) to adjust the sharpness of the cutoffs.

```
[10]: # PT spin model code based on the formula we gave above
      ################################
      # MODELLING THE BH SPIN #
      ############################
      def ourspin(period, BHmass, p0, M0, wp, wm):
          '''Cook up our own spin model
          INPUT:
              period = period in days
              BHmass = mass of the superradiant BH in solar mass
              p0 = period cutoff
              MO = BH mass cutoff
              wp,wm = two parameters that we can adjust
          OUTPUT:
              spin of the black hole'''
          atilde = (1/4)* special.erfc(np.log(period/p0)/(np.sqrt(2)*wp)) * special.
       →erfc((MO-BHmass)/(np.sqrt(2)*wm))
          return atilde
```

An example illustrating the relationship between the mass, spin of the secondary black hole, and the binary period using the PT model with  $P_C = 0.8$ ,  $w_P = 0.5$ ,  $M_C = 10$ , and  $w_M = 2$  is as follows

```
[11]: # Plot of the relation between secondary spin, BH mass and binary period based
       \rightarrow on PT spin model with P_C=0.8, w_P=0.5, M_C=10, w_M=2
      from scipy import special
      # If you want to modify the PT model parameter to get other plot, directly \Box
      →modify in the line below
      Z1=(1/4)* special.erfc(np.log10(10**Y/0.8)/(np.sqrt(2)*0.5)) * special.
       \rightarrowerfc((10-X)/(np.sqrt(2)*2))
      plt.figure(figsize=(10,8))
      plt.xlabel('$M/M_{\odot}$',size=30)
      plt.ylabel('$P/$'+'day',size=30,fontproperties=font2)
      ax=plt.pcolormesh(X, 10**Y, Z1,cmap='viridis',shading='gouraud')
      plt.yscale('log')
      plt.ylim((0,1000))
      cb1=plt.colorbar(ticks=np.linspace(0, 1,6))
      cb1.ax.tick_params(labelsize=15)
      plt.xticks(fontproperties=font2)
      plt.yticks(fontproperties=font2)
```

matplotlib.pyplot.grid(False)
plt.show()



Therefore, we can plot the distribution of spin using our SEVN simulation output and observe how it behaves with different spin models (such as the WR spin model and the PT spin model with varying parameter choices).

```
[12]: #To calculate the spin, the period of the orbital motion and the superradiant

→BH mass Mb are needed

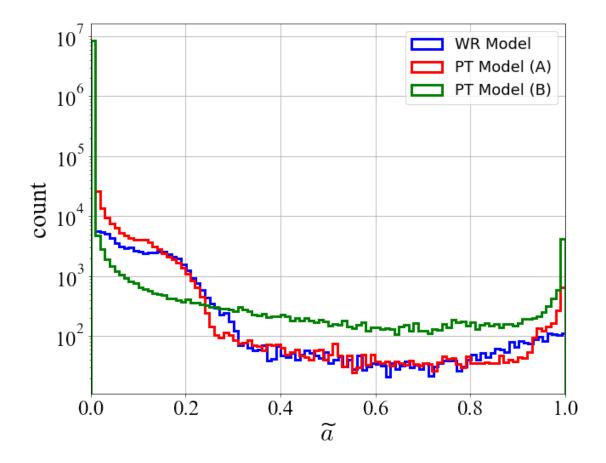
from scipy.constants import *

DAY=3600 * 24

# Mass ratio: primary BH mass over secondary BH mass
q = M1/M2

#Superradiant BH mass
Mb = M2
```

```
Mb_kg = M2*Msun
# Semi latus rectum
p = a * (1 - e**2)
# Period\ T = 2*np.pi*np.sqrt(p**3/(G * Mb*(1+q)*(1-e**2)**3))
T = 2*np.pi*np.sqrt(p**3/(G * Mb_kg*(1+q)*(1-e**2)**3)) # in seconds
T = T / DAY
                                                          # in days
T = np.array(T)
#plt.title('BH Spin for PT Model')
at=model(Mb,T)
                                  # Spin distribution with WR spin model
atPT=ourspin(T,Mb,0.8,10,0.5,2)
                                     # Spin distribution with PT spin model with
\rightarrow P_C=0.8, w_P=0.5, M_C=10, w_M=2
atPT1=ourspin(T,Mb,0.4,5,0.25,1)
                                    # Spin distribution with PT spin model with
\rightarrow P_C=0.4, w_P=0.25, M_C=5, w_M=1
plt.figure(figsize=(10,8))
plt.xlim((0,1))
plt.yscale('log')
plt.hist(at,lw=3,bins=100,edgecolor='blue',histtype='step',label='WR Model')
plt.hist(atPT,lw=3,bins=100,edgecolor='red',histtype='step',label='PT Modelu
\hookrightarrow (A) ')
plt.hist(atPT1,lw=3,bins=100,edgecolor='green',histtype='step',label='PT Model_
plt.grid(True)
plt.xticks(fontproperties=font2)
plt.yticks(fontproperties=font2)
plt.xlabel('$\widetilde{a}$',fontproperties=font1)
plt.ylabel('count',fontproperties=font1)
plt.legend(fontsize=18)
plt.show()
```



### 3 Effective Superradiant Rate

Considering ultralight bosonic cloud around the BH, the cloud will grow or decay exponentially with time by extracting BH spin. The rate of growth/decay can be calculated as

$$\Gamma_{n00} = -\frac{4}{n^3} \left( 1 + \sqrt{1 - \tilde{a}^2} \right) \mu \alpha^5 ,$$
  
$$\Gamma_{nlm} = 2\tilde{r}_+ C_{nl} g_{lm}(\tilde{a}, \alpha, \omega_{nlm}) (m\Omega_H - \omega_{nlm}) \alpha^{4l+5} .$$

Detail form of superradiance rate can be found in

Detweiler 1980: https://doi.org/10.1103/PhysRevD.22.2323

We consider the most dominant mode  $\psi_{211}$  and  $\psi_{322}$ . Those superradiance rate can be calculated with the following codes according to the formula we mentioned above

```
def GAMMA_211(at,Mb,mu = 1.34e-12,verbose=False):
    INPUT:
        at = spin of the black hole
        Mb = mass of the black hole in solar masses
        mu = mass of the boson in electron volts
    OUTPUT:
        GAMMA = superradiance rate for 211 state
    mu = mu * eVtokg # conversion to from eV to kg
    Mb = Mb * Msun # conversion from solar masses to kg
    alpha = (G * Mb * mu) / (s.c * s.hbar)
    gamma211 = c**3 * (alpha**9 * (np.sqrt(1 - at**2) + 1) * ((at - 2 * alpha *_{\sqcup}
 \hookrightarrow (np.sqrt(1 - at**2) + 1))**2 - at**2 + 1) * ((17 * alpha**5) / 64 + alpha**3
 \rightarrow/ 4 - 2 * alpha + at * (1 / (np.sqrt(1 - at**2) + 1) - (50000000000 *_\subseteq
 \rightarrowalpha**6) / 30000000003))) / (48 * s.G * Mb)
    if verbose:
        print('Boson mass:\t\t %.3e kg (%.3e eV)'%(mu,mu/1.78e-36))
        print('Black Hole mass:\t %.3e kg (%.0f Msun)'%(Mb, Mb/2e30))
        print('Alpha:\t\t %.3e '%alpha )
        print("Gamma:\t\t\t %.3e"%gamma211)
    return gamma211
####################################
# SUPERRADIANCE RATE 322 STATE #
######################################
def GAMMA_322(at,Mb,mu = 1.34e-12,verbose=False):
    111
    INPUT:
        at = spin of the black hole
        Mb = mass of the black hole in solar masses
        mu = mass of the boson in electron volts
    OUTPUT:
        GAMMA = superradiance rate'''
    mu = mu * eVtokg # conversion to from eV to kg
    Mb = Mb * Msun # conversion from solar masses to kg
    alpha = (G * Mb * mu) / (c * hbar)
```

```
gamma = c**3 * (8 * alpha**13 * (1 + np.sqrt(1 - at**2)) * (-alpha + □ → alpha**3 / 18. + (23 * alpha**5)/1080. - (1.6e10 * alpha**6 * at) / 1. → 620000000081e12 + at / (1 + np.sqrt(1 - at**2))) * (4 - 4 * at**2 + (2 * at □ → 2 * alpha * (1 + np.sqrt(1 - at**2)))**2) * (1 - at**2 + (2 * at - 2 * □ → alpha * (1 + np.sqrt(1 - at**2)))**2)) / (885735. * s.G * Mb)

if verbose:

print('Boson mass:\t\t %.3e kg (%.3e eV)'%(mu,mu/1.78e-36))
print('Black Hole mass:\t %.3e kg (%.0f Msun)'%(Mb, Mb/2e30))
print('Alpha:\t\t\t %.3e '%alpha )
print("Gamma:\t\t\t\t %.3e"%gamma)
return gamma
```

With the existence of the binary companion, the superradiance eigenmode can be couple with absorptive eigenmode, leading to a correction to the superradiance rate, which is

$$\Delta\Gamma_{nlm}^{(ACR)} \approx -\sum_{n'l'm'} \frac{(\Gamma_{nlm} - \Gamma_{n'l'm'})|\langle \psi_{n'l'm'}|V_*(t)|\psi_{nlm}\rangle|^2}{(E_{nlm} - E_{n'l'm'})^2 + \frac{M(1+q)}{p^3}(m-m')^2} ,$$

where  $V_*$  is the gravitational potential given by the binary companion.

Details of the correction on superradiant rate caused by the binary companion can be found in Tong et al. 2022: https://doi.org/10.1103/PhysRevD.106.043002 Fan et al. 2024: https://doi.org/10.1103/PhysRevD.109.024059

The correction to the superradiant rate can be calculated by the following code based on the formula we mentioned above.

```
# SUPERRADIANCE RATE VARIATION ACR 211 #
     def DELTA_GAMMA_ACR_211(e,Mb,q,a,at,mu = 1.34e-12,verbose=False):
        111
        INPUT:
            e = eccentricity
            Mb = mass of the black hole in solar masses
            q = ratio between primary BH and secondary BH
            a = semimajor axis in solar radii
            mu = mass of the boson in electron volts
            at = spin of the black hole
        OUTPUT:
            DELTA GAMMA = superradiance rate variation ACR 211 for 211 to 21-1
        mu = mu * eVtokg # conversion to from eV to kg
        Mb = Mb * Msun # conversion from solar masses to kg
        a = a * Rsun # conversion from solar radii to meters
```

```
#Semi latus rectum
   p = a * (1 - e**2) # same units of a
   alpha = (G * Mb * mu) / (c * hbar)
   #ACR for 211 to 21-1 (the only coupling allowed by the selection rules)
   numm1=(810000000081 * alpha**3 * at * (3 * (e**2 + 8) * e**2 + 8) * G**5 *_{L}
\rightarrow (Mb**3)**(3/2) * q**2)
   numm21=(-2 * (6400000000000 * alpha**6 + 5500000000051 * alpha**4 + ___
480000000048 * alpha**2 - 5760000000576) * (np.sqrt(1 - at**2) + 1))
   numm22=at**2 * (510000000051 * alpha**4 + 480000000048 * alpha**2 +
→3200000000000 * alpha**6 * (np.sqrt(1 - at**2) + 3) - 5760000000576)
   denn=4096 * p**3 * np.sqrt(-1 / ((e**2 - 1)**3 * Mb)) *_{\sqcup}
\rightarrow90000000180000000009 * G**3 * Mb**3 * (q + 1) * c**3)
   delta gamma acr211=-(numm1*( 480000000048 + alpha**2 * (numm21+numm22)))/
\rightarrow (denn)
   if verbose:
       print('Boson mass:\t\t %.3e kg (%.3e eV)'%(mu,mu/1.78e-36))
       print('Black Hole mass:\t %.3e kg (%.0f Msun)'%(Mb, Mb/2e30))
       print('Alpha:\t\t\t %.3e '%alpha )
       print('eccentricity:\t\t %.3e '%e )
       print('Mass ratio:\t\t %.3f '%q )
       print('Semi-major axis:\t %.3e meters'%a )
       print('Semi-latus rectum:\t %.3e meters'%p )
       print("Delta Gamma:\t\t %.3e"%delta_gamma_acr211)
   return delta_gamma_acr211
# SUPERRADIANCE RATE VARIATION ACR 322 #
def DELTA_GAMMA_ACR_322(e,Mb,q,a,at,mu = 1.34e-12,verbose=False):
   INPUT:
       e = eccentricity
       Mb = mass of the black hole in solar masses
       q = ratio between primary BH and secondary BH
       a = semimajor axis in solar radii
       mu = mass of the boson in electron volts
       at = spin of the black hole
   OUTPUT:
       DELTA GAMMA = superradiance rate variation ACR
```

```
111
      mu = mu * eVtokg # conversion to from eV to kg
      Mb = Mb * Msun # conversion from solar masses to kg
      a = a * Rsun
                                              # conversion from solar radii to meters
       #Semi latus rectum
      p = a * (1 - e**2) # same units of a
      alpha = (G * Mb * mu) / (c * hbar)
       # numerator components
      num1 = 5 * (1 + np.sqrt(1 - at**2)) * (8 + 3 * e**2 * (8 + e**2)) * G**3 *_{11}
\rightarrownp.sqrt(Mb**3) * q**2
      num21 = 98415. / (np.sqrt(9 - alpha**2))
      →at**2)) * alpha)**2))
      num222 = (1 - at**2 + (2 * at-2 * (1 +np.sqrt(1 - at**2)) * alpha)**2)
      num223 = (-alpha + alpha**3 / 18. + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + at * (1 / (1 + 1) + (23 * alpha**5) / 1080 + (23 * al
→np.sqrt(1 - at**2)) - (16e9 * alpha**6) / (1620000000081) ) )
       #denumerator component
      den=5832* np.sqrt(1/(Mb-e**2*Mb))**3*p**6*(((c**3*4*G*Mb*_
\rightarrow (1+q))/p**3)+((s.c**9 * 64 * alpha**10 * (18000000009+10000000000 * at *<sub>11</sub>
\rightarrowalpha)**2)/(65610000006561000000164025 * G**2 * Mb**2)))
       #DELTA GAMMA ACR = -num1*(num21+num221*num222*num223)/(den * c**5)
      delta_gamma_acr = -num1*(num21+num221*num222*num223)/den
      if verbose:
               print('Boson mass:\t\t %.3e kg (%.3e eV)'%(mu,mu/1.78e-36))
                print('Black Hole mass:\t %.3e kg (%.0f Msun)'%(Mb, Mb/2e30))
               print('Alpha:\t\t\t %.3e '%alpha )
               print('eccentricity:\t\t %.3e '%e )
               print('Mass ratio:\t\t %.3f '%q )
                print('Semi-major axis:\t %.3e meters'%a )
               print('Semi-latus rectum:\t %.3e meters'%p )
               print("Delta Gamma:\t\t %.3e"%delta_gamma_acr)
      return delta_gamma_acr
```

Thus we can try to plot the example distribution of superradiance rate and the change of the superradiance rate with our simulation output data with the boson mass choice  $\mu = 1.34 \times 10^{-12}$  eV.

```
[15]: import dask.dataframe as dd import dask.array as da

DaskMb=da.from_array(Mb,chunks=1e5)
```

```
Dasksemi=da.from_array(a_solar.to_numpy(),chunks=1e5)
Daskq=da.from_array(q,chunks=1e5)
Daskat=da.from_array(at,chunks=1e5)
Daske=da.from_array(ddf.Eccentricity.to_numpy(),chunks=1e5)
# Get the full list of 211 and 322 Superradiance Rate (Gamma211List,
\hookrightarrow Gamma322List)
# If you want to choose the boson mass, directly change the mu in the code,
→below
Gamma211List=GAMMA 211(Daskat,DaskMb,mu = 1.34e-12,verbose=False).compute()
Gamma322List=GAMMA 322(Daskat,DaskMb,mu = 1.34e-12,verbose=False).compute()
# Get the full list of 211 and 322 Correction to Superradiance Rate
→ (Gamma211List, Gamma322List)
# If you want to choose the boson mass, directly change the mu in the code,
→below
DeltaGamma322List=DELTA_GAMMA_ACR_322(Daske,DaskMb,Daskq,Dasksemi,Daskat,mu = 1.
→34e-12, verbose=False).compute()
DeltaGamma211List=DELTA_GAMMA_ACR_211(Daske,DaskMb,Daskq,Dasksemi,Daskat,mu = 1.
 →34e-12, verbose=False).compute()
```

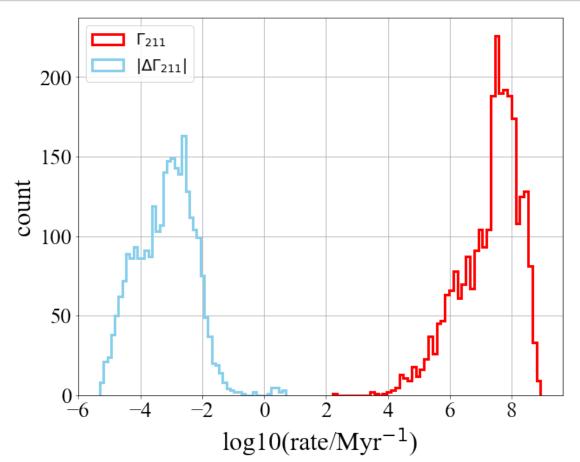
We require the superradiance rate to be greater than both the inverse of the age of the universe and the merger time of the BBH to guarantee the presence of the cloud. Additionally, the gravitational fine structure constant  $\alpha = GM_B\mu/(c\hbar)$  must be less than 0.4 to ensure the validity of the non-relativistic approximation employed in obtaining the analytical solution. Data points meeting these criteria are referred to as "participants."

Thus we can plot the distribution of  $\psi_{211}$  mode superradiant rate for participants

```
[17]: # Filter the 211 participants

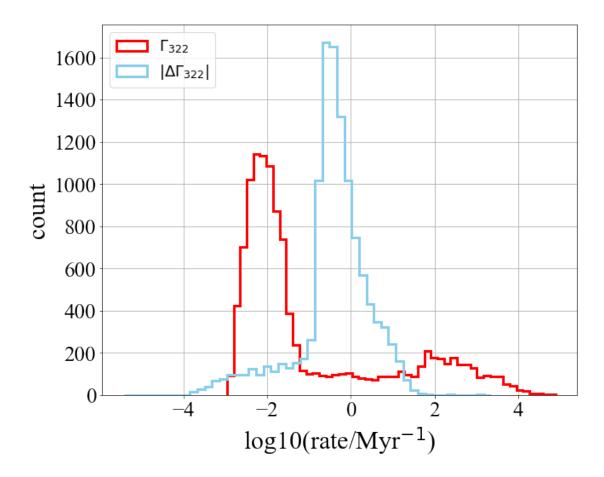
test3=Gamma211List-cuttimed>0
test2=G*Mb * 2e30 * 1.34e-12 *1.78e-36 /(c * hbar )<0.4
test211=test3 * test2
```

```
# Plot the distribution of 211 superradiant rate and the correction of \Box
\hookrightarrow superradiant rate
testGamma211=test211 * Gamma211List
testDelta211=test211 * DeltaGamma211List
plt.figure(figsize=(10,8))
plt.hist(np.
→log10(testGamma211[testGamma211>0]*Myr),bins=50,lw=3,edgecolor='red',histtype='step')
plt.hist(np.
→log10(-testDelta211[testDelta211<0]*Myr),bins=50,lw=3,edgecolor='skyblue',histtype='step')</pre>
plt.legend(['$\Gamma_{211}$','$|\Delta\Gamma_{211}|$'],fontsize=18)
plt.xlabel('log10(rate/Myr'+'${}^{-1}$'+')',fontproperties=font1)
plt.grid(True)
plt.ylabel('count',fontproperties=font1)
plt.xticks(fontproperties=font2)
plt.yticks(fontproperties=font2)
plt.show()
```



Also for  $\psi_{322}$  mode participants

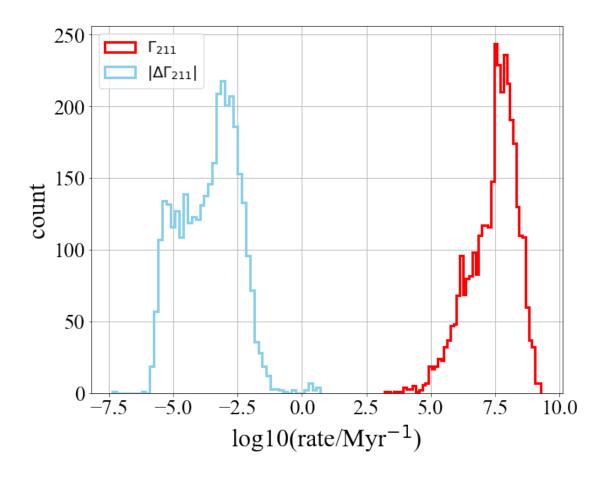
```
[18]: # Filter the 322 data
      test1=Gamma322List-cuttimed>0
      test2=G*Mb * 2e30 * 1.34e-12 *1.78e-36 /(c * hbar)<0.4
      test=test1 * test2
      \# Plot the distribution of 322 superradiant rate and the correction of
      \rightarrow superradiant rate
      testGamma=test * Gamma322List
      testDelta=test * DeltaGamma322List
      plt.figure(figsize=(10,8))
      plt.hist(np.
      →log10(testGamma[testGamma>0]*Myr),bins=50,lw=3,edgecolor='red',histtype='step')
      plt.hist(np.
      →log10(-testDelta[testDelta<0]*Myr),bins=50,lw=3,edgecolor='skyblue',histtype='step')
      plt.legend(['$\Gamma_{322}$','$|\Delta\Gamma_{322}|$'],fontsize=18,loc='upper_
      ⇔left')
      plt.xlabel('log10(rate/Myr'+'${}^{-1}$'+')',fontproperties=font1)
      plt.grid(True)
      plt.ylabel('count',fontproperties=font1)
      plt.xticks(fontproperties=font2)
      plt.yticks(fontproperties=font2)
      plt.show()
```



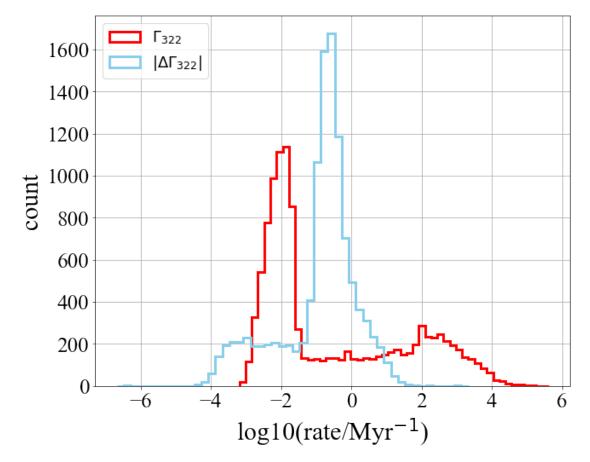
Similar estimation can be done for PT model

 $\label{eq:deltaGamma322ListPT=DELTA_GAMMA_ACR_322(Daske,DaskMb,Daskq,Dasksemi,Daskat2,mu_ \Rightarrow 1.34e-12,verbose=False).compute()$ 

```
[20]: # Filter the 211 data
      test3PT=Gamma211ListPT-cuttimed>0
      test2=G*Mb * 2e30 * 1.34e-12 *1.78e-36 /(c * hbar)<0.4
      test211PT=test3PT * test2
      # Plot the distribution of 211 superradiant rate and the correction of \Box
      \rightarrow superradiant rate
      testGamma211PT=test211PT * Gamma211ListPT
      testDelta211PT=test211PT * DeltaGamma211ListPT
      plt.figure(figsize=(10,8))
      plt.hist(np.
      →log10(testGamma211PT[testGamma211PT>0]*Myr),bins=50,lw=3,edgecolor='red',histtype='step')
      plt.hist(np.
       →log10(-testDelta211PT[testDelta211PT<0]*Myr),bins=50,lw=3,edgecolor='skyblue',histtype='ste
      plt.legend(['$\Gamma_{211}\$', '$|\Delta\Gamma_{211}\$'],fontsize=18)
      plt.xlabel('log10(rate/Myr'+'${}^{-1}$'+')',fontproperties=font1)
      plt.grid(True)
      plt.ylabel('count',fontproperties=font1)
      plt.xticks(fontproperties=font2)
      plt.yticks(fontproperties=font2)
      plt.show()
```



```
plt.xlabel('log10(rate/Myr'+'${}^{-1}$'+')',fontproperties=font1)
plt.grid(True)
plt.ylabel('count',fontproperties=font1)
plt.xticks(fontproperties=font2)
plt.yticks(fontproperties=font2)
plt.show()
```



### 4 Survival Rate From Superradiance Termination

In our previous work (Tong et al. 2022), we mentioned that if the binary separation is small enough, the absolute value of correction of superradiant rate can be bigger than the superradiant rate, terminating the superradiant effect. To investigate how this termination effect influence the superradiant BH in reality statistically, we define those with  $\Gamma_{nlm} + \Delta \Gamma_{nlm} > 0$  as survivors, indicating that they can survive the termination effect and produce an abundantly populated boson cloud.

To characterize the likelihood of a superradiant cloud surviving the termination from a binary companion, we define the survival rate as the ratio of the number of survivors and the number of participants

$$R_{\rm surv} \equiv \frac{N_{\rm surv}}{N_{\rm part}} \ ,$$

with  $R_{\rm surv} \approx 0$  standing for complete cloud termination and  $R_{\rm surv} \approx 1$  standing for assured survival.

```
[22]: # Code to calculate survival rate based on the definition above
      def RateOfSurvive_322d(BHspin,BHmass,mumass,Massratio,ecc,semimajor,cut):
          #mumass should be in eV
          #superradiance rate bigger than one over the threashold time (Return True)
       \rightarrow or False list)
          sel_cand_1 = GAMMA_322(at=BHspin,Mb=BHmass,mu=mumass,verbose=False)-cut>0
          #alpha<0.4 (Return another True or False list)
          sel\_cand\_2 = G*BHmass * Msun * mumass *eVtokg /(c * hbar )<0.4
          #Only two condition True together can be True (Return True or False list)
          sel_part_322 = sel_cand_1 * sel_cand_2
          #Delta Gamma (Return True or False list)
       ⇒sel surv 322=-DELTA GAMMA ACR 322(e=ecc,Mb=BHmass,q=Massratio,a=semimajor,at=BHspin,mu,
       →= mumass, verbose=False)/
       →GAMMA_322(at=BHspin,Mb=BHmass,mu=mumass,verbose=False)<1
          #Survival for 322 (Need Delta Gamma < Gamma and it should be a participant)
          SurviveOrNot322 = sel_part_322 * sel_surv_322
          return da.sum(SurviveOrNot322).compute()/da.sum(sel_part_322).compute()
```

To reflect the statistical error in a sample of finite dimensionality, we also introduce an uncertainty estimator

$$\sigma(R_{\rm surv}) \equiv \frac{1}{\sqrt{N_{\rm part}}} \; ,$$

which is based on a Poisson distribution.

```
[23]: # Code to calculate error bar based on the definition above

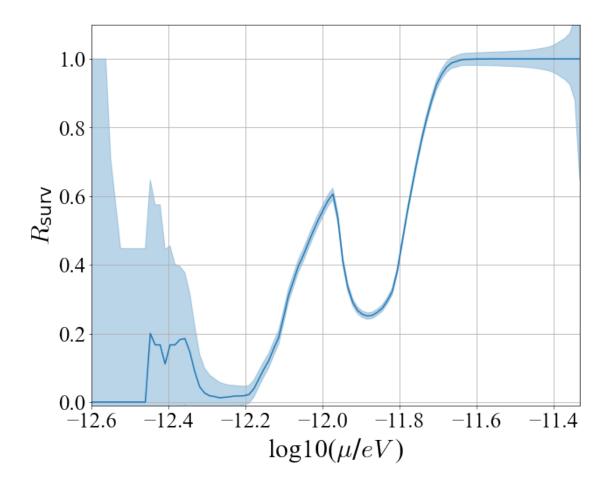
def ErrorBar_322d(BHspin,BHmass,mumass,Massratio,ecc,semimajor,cut):
    #mumass should be in eV
    sel_cand_1 = GAMMA_322(at=BHspin,Mb=BHmass,mu=mumass,verbose=False)-cut>0
    sel_cand_2 = G*BHmass * Msun * mumass *eVtokg /(c * hbar )<0.4

    sel_part_322 = sel_cand_1 * sel_cand_2
    errorbar=1/(np.sqrt(da.sum(sel_part_322).compute()))
    return errorbar</pre>
```

It will be useful to estimate how the survival rate  $R_{\text{surv}}$  change with the boson mass  $\mu$ .

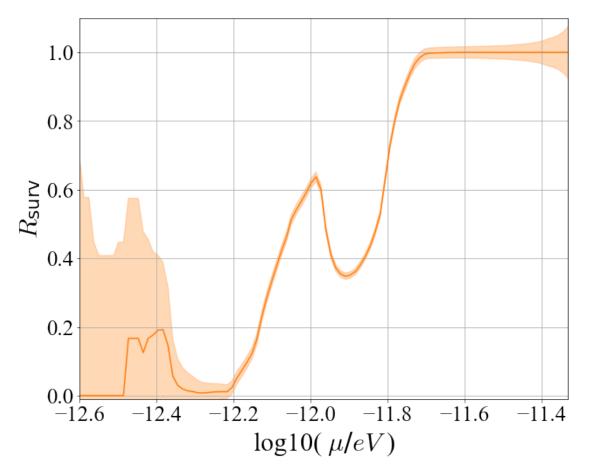
```
[24]: # Creating a list of mu, and then calculate the survival rate for each mu
      mulist=1e-13*np.array([1.03**i for i in range(200)])
      Daskmulist=da.from_array(mulist,chunks=1e5)
[25]: # Survival rate -mu realtion with WR spin model
      # Survival rate list and the errorbar list for each mu
      Ratelistd=[RateOfSurvive_322d(BHspin=Daskat,BHmass=DaskMb,mumass=mui,Massratio=Daskq,ecc=Daske
       →for mui in mulist ]
      Errord=[ErrorBar_322d(BHspin=Daskat,BHmass=DaskMb,mumass=mui,Massratio=Daskq,ecc=Daske,semimageneral)
       →for mui in mulist]
      error1=np.array(Ratelistd)+(np.array(Errord))
      error2=np.array(Ratelistd)-(np.array(Errord))
[26]: # Plot for the survival rate with the shaded region representing the error bar
      plt.figure(figsize=(10,8))
      plt.rcParams['axes.grid'] = True
      plt.rcParams['image.cmap'] = 'gray'
      plt.figure(figsize=(10,8))
      plt.plot(np.log10(mulist), Ratelistd, color='C0')
      plt.fill_between( np.log10(mulist),y1=error1, y2=error2, color='C0', alpha=0.3)
      plt.ylim((-0.01,1.1))
      plt.xlim((-12.6,-11.33))
      plt.ylabel('$R_{\mathrm{surv}}$',fontproperties=font1)
      plt.xlabel('log10('+'$\mu/eV$'+' )',fontproperties=font1)
      plt.xticks(fontproperties=font2)
      plt.yticks(fontproperties=font2)
      plt.show()
```

<Figure size 720x576 with 0 Axes>



We can also investigate the behavior of survival rate with our own PT model with different parameter choice. Here we give one example of

<Figure size 720x576 with 0 Axes>



[]:[