

# Partial Differential Equations

## Classification of PDEs

**General Second-Order Linear PDE:**

$$A \frac{\partial^2 u}{\partial x^2} + B \frac{\partial^2 u}{\partial x \partial y} + C \frac{\partial^2 u}{\partial y^2} + D \frac{\partial u}{\partial x} + E \frac{\partial u}{\partial y} + F u = G$$

**Classification:**

- **Hyperbolic:**  $B^2 - 4AC > 0$  (e.g., wave equation)
- **Parabolic:**  $B^2 - 4AC = 0$  (e.g., heat equation)
- **Elliptic:**  $B^2 - 4AC < 0$  (e.g., Laplace equation)

## Three Fundamental PDEs

### Wave Equation (Hyperbolic)

**1D Wave Equation:**

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$$

where  $c$  is the wave speed

**2D Wave Equation:**

$$\frac{\partial^2 u}{\partial t^2} = c^2 \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) = c^2 \nabla^2 u$$

**3D Wave Equation:**

$$\frac{\partial^2 u}{\partial t^2} = c^2 \nabla^2 u$$

### Heat (Diffusion) Equation (Parabolic)

**1D Heat Equation:**

$$\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$$

where  $\alpha = k/(\rho c_p)$  is the thermal diffusivity

**2D Heat Equation:**

$$\frac{\partial u}{\partial t} = \alpha \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) = \alpha \nabla^2 u$$

**3D Heat Equation:**

$$\frac{\partial u}{\partial t} = \alpha \nabla^2 u$$

**Laplace Equation (Elliptic)**

**2D Laplace Equation:**

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$
$$\nabla^2 u = 0$$

**3D Laplace Equation:**

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} = 0$$

**Poisson Equation:**

$$\nabla^2 u = f(x, y, z)$$

Solutions to Laplace equation are called *harmonic functions*.

## Method of Separation of Variables

**General Approach:**

1. Assume solution has the form  $u(x, t) = X(x)T(t)$  (or similar)
2. Substitute into PDE
3. Separate variables to get two ODEs
4. Solve each ODE with appropriate boundary conditions
5. Apply superposition principle to satisfy initial conditions

**Heat Equation on Finite Domain**

For:  $\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$ ,  $0 < x < L$ ,  $t > 0$

**Boundary Conditions (Fixed Ends):**

$$u(0, t) = 0, \quad u(L, t) = 0$$

**Initial Condition:**

$$u(x, 0) = f(x)$$

**Solution:**

$$u(x, t) = \sum_{n=1}^{\infty} B_n e^{-\alpha(n\pi/L)^2 t} \sin\left(\frac{n\pi x}{L}\right)$$

where

$$B_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

**Boundary Conditions (Insulated Ends):**

$$\frac{\partial u}{\partial x}(0, t) = 0, \quad \frac{\partial u}{\partial x}(L, t) = 0$$

**Solution:**

$$u(x, t) = \frac{A_0}{2} + \sum_{n=1}^{\infty} A_n e^{-\alpha(n\pi/L)^2 t} \cos\left(\frac{n\pi x}{L}\right)$$

where

$$A_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx$$

**Wave Equation on Finite Domain**

For:  $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ ,  $0 < x < L$ ,  $t > 0$

**Boundary Conditions:**

$$u(0, t) = 0, \quad u(L, t) = 0$$

**Initial Conditions:**

$$u(x, 0) = f(x), \quad \frac{\partial u}{\partial t}(x, 0) = g(x)$$

**Solution:**

$$u(x, t) = \sum_{n=1}^{\infty} \left[ A_n \cos\left(\frac{n\pi c t}{L}\right) + B_n \sin\left(\frac{n\pi c t}{L}\right) \right] \sin\left(\frac{n\pi x}{L}\right)$$

where

$$A_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

$$B_n = \frac{2}{n\pi c} \int_0^L g(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

## D'Alembert's Solution (Wave Equation)

For infinite domain:  $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ ,  $-\infty < x < \infty$

**Initial Conditions:**

$$u(x, 0) = f(x), \quad \frac{\partial u}{\partial t}(x, 0) = g(x)$$

**D'Alembert's Solution:**

$$u(x, t) = \frac{f(x - ct) + f(x + ct)}{2} + \frac{1}{2c} \int_{x-ct}^{x+ct} g(s) ds$$

For zero initial velocity ( $g(x) = 0$ ):

$$u(x, t) = \frac{f(x - ct) + f(x + ct)}{2}$$

This represents two waves traveling in opposite directions at speed  $c$ .

## Laplace Equation Solutions

### Rectangular Domain

For:  $\nabla^2 u = 0$  on rectangle  $0 < x < a$ ,  $0 < y < b$

**Example Boundary Conditions:**

$$u(0, y) = 0, \quad u(a, y) = 0, \quad u(x, 0) = 0, \quad u(x, b) = f(x)$$

**Solution:**

$$u(x, y) = \sum_{n=1}^{\infty} B_n \frac{\sinh(n\pi y/a)}{\sinh(n\pi b/a)} \sin\left(\frac{n\pi x}{a}\right)$$

where

$$B_n = \frac{2}{a} \int_0^a f(x) \sin\left(\frac{n\pi x}{a}\right) dx$$

### Circular Domain (Polar Coordinates)

For:  $\nabla^2 u = \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = 0$

On disk  $r < R$  with boundary condition  $u(R, \theta) = f(\theta)$ :

**Solution:**

$$u(r, \theta) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left(\frac{r}{R}\right)^n [a_n \cos(n\theta) + b_n \sin(n\theta)]$$

where

$$a_n = \frac{1}{\pi} \int_0^{2\pi} f(\theta) \cos(n\theta) d\theta$$

$$b_n = \frac{1}{\pi} \int_0^{2\pi} f(\theta) \sin(n\theta) d\theta$$

## 2D Wave Equation (Rectangular Membrane)

For:  $\frac{\partial^2 u}{\partial t^2} = c^2 \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$

On rectangle  $0 < x < a$ ,  $0 < y < b$  with fixed boundaries:

**Solution:**

$$u(x, y, t) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right) \\ \times [A_{mn} \cos(\omega_{mn}t) + B_{mn} \sin(\omega_{mn}t)]$$

where the frequencies are:

$$\omega_{mn} = \pi c \sqrt{\frac{m^2}{a^2} + \frac{n^2}{b^2}}$$

## Heat Equation in Infinite Rod

For:  $\frac{\partial u}{\partial t} = \alpha \frac{\partial^2 u}{\partial x^2}$ ,  $-\infty < x < \infty$ ,  $t > 0$

**Initial Condition:**  $u(x, 0) = f(x)$

**Solution (using Fourier Transform):**

$$u(x, t) = \frac{1}{\sqrt{4\pi\alpha t}} \int_{-\infty}^{\infty} f(\xi) e^{-(x-\xi)^2/(4\alpha t)} d\xi$$

This is a convolution with the heat kernel (fundamental solution):

$$K(x, t) = \frac{1}{\sqrt{4\pi\alpha t}} e^{-x^2/(4\alpha t)}$$

## Sturm-Liouville Problems

**General Form:**

$$\frac{d}{dx} \left[ p(x) \frac{dy}{dx} \right] + [q(x) + \lambda r(x)] y = 0$$

with boundary conditions at endpoints

**Properties:**

- Eigenvalues  $\lambda_n$  are real
- Eigenfunctions  $y_n(x)$  are orthogonal with weight function  $r(x)$ :

$$\int_a^b y_m(x)y_n(x)r(x) dx = 0 \quad (m \neq n)$$

- Eigenfunctions form a complete basis

### Common Examples:

*For heat/wave equation with fixed ends:*

$$\frac{d^2 X}{dx^2} + \lambda X = 0, \quad X(0) = X(L) = 0$$

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2, \quad X_n(x) = \sin\left(\frac{n\pi x}{L}\right)$$

*For heat/wave equation with insulated ends:*

$$\frac{d^2 X}{dx^2} + \lambda X = 0, \quad X'(0) = X'(L) = 0$$

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2, \quad X_n(x) = \cos\left(\frac{n\pi x}{L}\right)$$

## Maximum Principle

### For Heat Equation:

The maximum and minimum values of  $u(x, t)$  in a domain occur either:

- At the initial time  $t = 0$ , or
- On the spatial boundary

### For Laplace Equation:

Harmonic functions achieve their maximum and minimum values on the boundary of the domain (no local extrema in the interior).

## Uniqueness and Well-Posedness

### Well-Posed Problem (Hadamard):

1. Solution exists

2. Solution is unique
3. Solution depends continuously on initial/boundary data

**Heat Equation:** Well-posed with initial condition and boundary conditions

**Wave Equation:** Well-posed with initial position, initial velocity, and boundary conditions

**Laplace Equation:** Well-posed with boundary conditions (Dirichlet, Neumann, or mixed)

## Boundary Conditions

**Dirichlet BC:** Specifies value of  $u$  on boundary

$$u|_{\partial\Omega} = f$$

**Neumann BC:** Specifies normal derivative on boundary

$$\left. \frac{\partial u}{\partial n} \right|_{\partial\Omega} = g$$

**Robin (Mixed) BC:** Linear combination

$$\alpha u + \beta \left. \frac{\partial u}{\partial n} \right|_{\partial\Omega} = h$$

**Periodic BC:**

$$u(0, t) = u(L, t), \quad \frac{\partial u}{\partial x}(0, t) = \frac{\partial u}{\partial x}(L, t)$$

## Bessel Functions

Arise in problems with cylindrical symmetry.

**Bessel's Equation:**

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} + (x^2 - n^2)y = 0$$

**General Solution:**

$$y(x) = c_1 J_n(x) + c_2 Y_n(x)$$

where  $J_n(x)$  is Bessel function of the first kind and  $Y_n(x)$  is Bessel function of the second kind.

**Properties:**

$$\bullet \quad J_n(0) = \begin{cases} 1 & n = 0 \\ 0 & n > 0 \end{cases}$$

- $Y_n(x) \rightarrow -\infty$  as  $x \rightarrow 0^+$
- $J_n(x)$  and  $Y_n(x)$  oscillate with decreasing amplitude as  $x \rightarrow \infty$

**Orthogonality:**

$$\int_0^a x J_n(\alpha_m x) J_n(\alpha_k x) dx = 0 \quad (m \neq k)$$

where  $\alpha_m$  and  $\alpha_k$  are zeros of  $J_n$ .

## Legendre Polynomials

Arise in problems with spherical symmetry.

**Legendre's Equation:**

$$(1 - x^2) \frac{d^2 y}{dx^2} - 2x \frac{dy}{dx} + n(n + 1)y = 0$$

**Solutions:** Legendre polynomials  $P_n(x)$  for integer  $n$

**First Few Legendre Polynomials:**

$$P_0(x) = 1$$

$$P_1(x) = x$$

$$P_2(x) = \frac{1}{2}(3x^2 - 1)$$

$$P_3(x) = \frac{1}{2}(5x^3 - 3x)$$

**Orthogonality:**

$$\int_{-1}^1 P_m(x) P_n(x) dx = \begin{cases} 0 & m \neq n \\ \frac{2}{2n+1} & m = n \end{cases}$$

## Green's Functions

Solution to PDE can be expressed using Green's function  $G(\mathbf{x}, \mathbf{x}_0)$ :

For Poisson equation  $\nabla^2 u = f$ :

$$u(\mathbf{x}) = \int_{\Omega} G(\mathbf{x}, \mathbf{x}_0) f(\mathbf{x}_0) d\mathbf{x}_0 + \text{boundary terms}$$

**Green's Function satisfies:**

$$\nabla^2 G(\mathbf{x}, \mathbf{x}_0) = \delta(\mathbf{x} - \mathbf{x}_0)$$

**3D Free Space:**

$$G(\mathbf{x}, \mathbf{x}_0) = -\frac{1}{4\pi|\mathbf{x} - \mathbf{x}_0|}$$



## Useful Formulas

### Energy Method (for uniqueness):

Define energy:  $E(t) = \frac{1}{2} \int_{\Omega} u^2 dx$

Show  $\frac{dE}{dt} \leq 0$  to prove uniqueness.

### Separation Constant:

When separating variables, if spatial part gives negative eigenvalues, use  $-\lambda$  to get:

$$X'' + \lambda X = 0$$

This gives sinusoidal solutions for  $\lambda > 0$ .