#### Vector Calculus – 2

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#### Derivatives in Vector Calculus

In calculus, derivatives map a scalar field onto another scalar field.

In vector calculus, derivatives are use to create the following mappings:

- ▶ The gradient maps a scalar field onto a vector field.
- ► The divergence maps a vector field onto a scalar field.
- ► The curl maps a vector field onto a vector field.

#### The Gradient of a Scalar Field

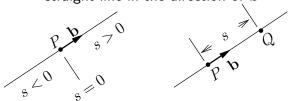
Given: f(x,y,z) :a scalar function defined over a domain in 3D space  $(\mathbb{R}^3)$ . f is differentiable. Then

$$\operatorname{grad} f = \nabla f \equiv \underbrace{\frac{\partial f}{\partial x} \mathbf{i} + \frac{\partial f}{\partial y} \mathbf{j} + \frac{\partial f}{\partial z} \mathbf{k}}_{\text{This is a vector.}}$$

What is the utility of the gradient?

#### Directional Derivative

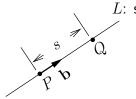
straight line in the direction of  $\ensuremath{\mathbf{b}}$ 



s is measured from P in the direction of  $\mathbf{b}$ .

$$D_{\mathbf{b}}f = \left. \frac{df}{ds} \right|_{\mathbf{b}} \equiv \lim_{s \to 0} \frac{f(Q) - f(P)}{s}$$

#### Directional Derivative and Gradient



L: st line through P

$$L$$
: st line through  $P$ 

Equation of  $L$ 
 $\mathbf{r}(s) = \mathbf{p}_0 + \mathbf{b}s$ 
 $= x(s)\mathbf{i} + y(s)\mathbf{i} + z(s)\mathbf{k}$ 

Assume f has partial derivatives which are continuous.

Then 
$$\frac{df}{ds} = \frac{\partial f}{\partial x} \frac{dx}{ds} + \frac{\partial f}{\partial y} \frac{dy}{ds} + \frac{\partial f}{\partial z} \frac{dz}{ds},$$
but 
$$\frac{d\mathbf{r}}{ds} = \frac{dx}{ds} \mathbf{i} + \frac{dy}{ds} \mathbf{j} + \frac{dz}{ds} \mathbf{k} = \mathbf{b} \quad \text{(why?)}$$

$$\therefore D_{\mathbf{b}} f = \frac{df}{ds} \Big|_{\mathbf{b}} = \nabla f \cdot \mathbf{b}$$

#### Direction of Maximum Increase

#### **Thorem**

Let f(P)=f(x,y,z) be a scalar function, with first partial derivatives which are continuous, defined over some domain B. Then  $\nabla f$  exists,  $\nabla f$  is a vector, its length independent of the choice of coordinates, and if  $\nabla f \neq 0$ , then the direction of  $\nabla f$  is the direction of the maximum rate of increase of f.

#### **Proof**

$$D_{\mathbf{b}}f = \nabla f \cdot \mathbf{b} = |\nabla f||\mathbf{b}|\cos\gamma$$

where  $\gamma$  is the included  $\angle$  between  $\nabla f$  and  $\mathbf{b}$ .

But f and s are scalar functions of the position P.

: their values do not depend on the choice of the coordinate system.

# Direction of Maximum Increase (cont)

#### Proof (cont)

$$\therefore D_{\mathbf{b}}f = \frac{df}{ds}\Big|_{\mathbf{b}}$$
: independent of choice of coordinate system.

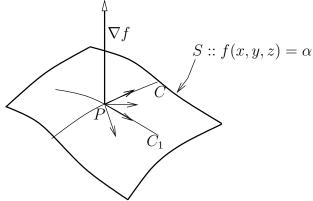
Hence, the length and direction of  $\nabla f$  does not depend on the choice of the coordinate system; this makes  $\nabla f$  a proper vector.

 $D_{\mathbf{b}}f = \nabla f$  will be maximum when  $\cos \gamma = 1$ , *i.e.* when  $\mathbf{b}$  is chosen in the direction of  $\nabla f$ .

 $\therefore$ , direction of  $\nabla f$ : direction of maximum rate of increase, and  $|\nabla f|$ : magnitude of the maximum rate of increase in f.

#### Gradients and Surfaces

Let f(x,y,z): a scalar function, differentiable. Then  $f(x,y,z)=\alpha$  (a constant) represents a surface, called a "level surface of f".



# Gradients and Surfaces (cont)

P: a point on S.

C: a curve on S passing through P, represented by  $\mathbf{r}(t)=x(t)\mathbf{i}+y(t)\mathbf{j}+z(t)\mathbf{k}.$ 

C lies on S, components of  $\mathbf{r}$  satisfy the relation for S: [1]  $f(x(t), y(t), z(T)) = \alpha$ .

A tangent vector of C will be  $\mathbf{r}'(t) = x'(t)\mathbf{i} + y'(t)\mathbf{j} + z'(t)\mathbf{k}$ .

For different curves on S passing through P, the set of tangent vectors at P forms a surface tangent to S at P. A normal to this plane at P will then be the surface-normal of S at P.

# Gradients and Surfaces (cont)

Differentiating [1] w.r.t. t:

$$\frac{\partial f}{\partial x}x'(t) + \frac{\partial f}{\partial y}y'(t) + \frac{\partial f}{\partial z}z'(t) = 0$$
$$\therefore \nabla f \cdot \mathbf{r}' = 0$$

 $\therefore \nabla f$  is normal to all vectors  $\mathbf{r}'$  in the tangent plane of S.  $\implies \nabla f$  is the normal vector of S at P.

#### ⇒ Theorem:

Let f(x,y,z): a differentiable function in space. Let  $f(x,y,z)=\alpha=(const)$ : a surface in that space. Then  $\nabla f$  at a point P on S is either 0 or a normal vector of S at P.

#### Potential

Nomenclature in science and engineering:

If f(P) = f(x, y, z) is a scalar field, and  $\mathbf{V}(P) = \mathbf{V}(x, y, z) = \nabla f$  is a vector field, then, often f is known as the **potential** for the vector field  $\mathbf{V}(=\nabla f)$ .

Sometimes, a negative sign and a scale factor is included.

### The Divergence of a Vector Field

Let  $\mathbf{V}(x, y, z) = v_1 \mathbf{i} + v_2 \mathbf{j} + v_3 \mathbf{k}$  be a differentiable vector field, with  $v_1, v_2, v_3$  functions of (x, y, z). Then

$$\operatorname{div} \mathbf{V} = \nabla \cdot \mathbf{V} \equiv \frac{\partial v_1}{\partial x} + \frac{\partial v_2}{\partial y} + \frac{\partial v_3}{\partial z},$$

where  $\nabla$  is the same scalar-to-vector operator used in the definition of the gradient.

## Divergence (cont)

#### Theorem

Given a vector function  $\mathbf{V}(x,y,z)$  on a domain in space,  $\mathrm{div}\mathbf{V}(=\nabla\cdot\mathbf{V})$  is a scalar function, whose values depend on  $\mathbf{V}$  and the location in space, but <u>not</u> on the choice of the coordinate system.

What is the physical significance of divergence?

#### Divergence – Significance

In many physical situations, divergence has a meaning:

- ▶ In fluid mechanics,  $\rho V$  represents the mass flux.  $\nabla \cdot (\rho V)$  equals the net outflow of mass per unit volume.
- ▶ In heat conduction, q represents the heat flux.
  - $\nabla \cdot (\mathbf{q})$  equals the net heat outflow per unit volume.

#### The Laplacian

Let f(x, y, z) be a scalar function in space. Its gradient

$$\nabla f = \frac{\partial f}{\partial x}\mathbf{i} + \frac{\partial f}{\partial y}\mathbf{j} + \frac{\partial f}{\partial z}\mathbf{k}$$

is a vector function. If we take the divergence of this, we will get another scalar function:

$$\nabla \cdot (\nabla f) = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2} = \nabla^2 f.$$

This function is known as the Laplacian of f, and is denoted by  $\nabla^2 f$ .

 $abla^2$  is known as the Laplacian operator.

#### Use of the Laplacian

Fourier's Law of heat conduction states that  $\mathbf{q}=-k\nabla T$ , where T is the temperature and k the thermal conductivity. The First Law states that "rate of storage of energy equals the net inflow of energy". Hence, we get the heat conduction equation:

$$\rho c_p \frac{\partial T}{\partial t} = -\nabla \cdot \mathbf{q} = -\nabla \cdot (-k\nabla T) = \nabla \cdot (k\nabla T)$$

If k is constant, this reduces to the 'heat conduction equation':

$$\rho c_p \frac{\partial T}{\partial t} = k \nabla^2 T$$

#### Curl of a Vector Field

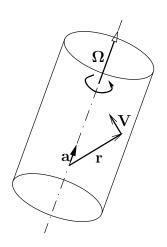
Let  $\mathbf{V}(x,y,z) = v_1 \mathbf{i} + v_2 \mathbf{j} + v_3 \mathbf{k}$  be a differentiable vector field, with  $v_1, v_2, v_3$  functions of (x, y, z). Then

$$\operatorname{curl} \mathbf{V} \equiv \nabla \times \mathbf{V}$$

$$= \left(\frac{\partial v_3}{\partial y} - \frac{\partial v_2}{\partial z}\right) \mathbf{i} + \left(\frac{\partial v_1}{\partial z} - \frac{\partial v_3}{\partial x}\right) \mathbf{j} + \left(\frac{\partial v_2}{\partial x} - \frac{\partial v_1}{\partial y}\right) \mathbf{k}$$

$$= \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ v_1 & v_2 & v_3 \end{vmatrix}$$

## Rotation of a Solid Body



$$\begin{array}{rcl} \boldsymbol{\Omega} &=& \omega \mathbf{a} \\ \mathbf{V} &=& \boldsymbol{\Omega} \times \mathbf{r} \\ \text{If } \mathbf{a} = \mathbf{k} \text{ (choice of direction),} \\ \text{then } \boldsymbol{\Omega} &=& [0,\ 0,\ \omega]. \\ \text{If } \mathbf{r} &=& [x,\ y,\ z], \\ \text{then } \mathbf{V} &=& \boldsymbol{\Omega} \times \mathbf{r} \\ &=& [-\omega y,\ \omega x,\ 0] \\ &=& -\omega y \mathbf{i} + \omega x \mathbf{j}. \end{array}$$

# Rotation of a Solid Body (cont)

Thus, 
$$\operatorname{curl} \mathbf{V} \equiv \nabla \times \mathbf{V} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ -\omega y & \omega x & 0 \end{vmatrix}$$
$$= 2\omega \mathbf{k} = 2\Omega$$

Thus, solid body rotation leads to a V-field such that  $\nabla \times \mathbf{V}$  is in the direction of  $\Omega$  and its magnitude =  $|2\Omega| = 2\omega$ .

### Properties of Curl

#### 1. Curl of a gradient

Let f be a scalar field, differentiable.

Let  $\nabla f$ , a vector field, be also differentiable.

Then  $\nabla \times (\nabla f) = \mathbf{0}$  (vector).

That is, the gradient field of a scalar field has zero curl.

We say that that the gradient field is irrotational.

Note:  $\nabla \cdot (\nabla f) = \nabla^2 f$ , is a scalar field, and need not be zero.

# Properties of Curl (cont)

#### 2. Divergence of a curl

Let V be a vector field, twice differentiable.

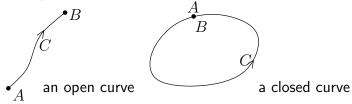
Then  $\nabla \times \mathbf{V}$  is another vector field.

Then  $\nabla \cdot (\nabla \times \mathbf{V}) = 0$ .

We say that the divergence of the curl of a vector field is zero.

### Integral Calculus of Vectors

Integration of vectors is a generalisation of ordinary integrals like  $\int_a^b f(x)dx$ . Line integrals are defined over a curve in space:



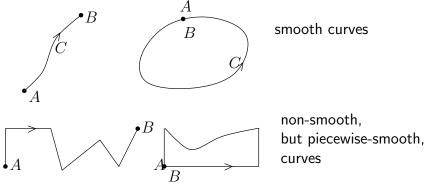
The curve C can be parametrically represented:  $\mathbf{r}(t) = x(t)\mathbf{i} + y(t)\mathbf{j} + z(t)\mathbf{k}; \quad a \leq t \leq b,$  where A at  $\mathbf{r}(a)$  is the initial point, and B at  $\mathbf{r}(b)$  is the terminal point. t increases from a to b: the positive direction of C. If points A and B coincide, then we have a closed curve.

#### Smooth Curves

C is called a smooth curve if:

1: at each point, a unique tangent exists, and

2: the direction of this tangent varies continuously.



Henceforth, all curves are assumed piecewise smooth.

#### Line Integral

Line integral of a vector function:

$$\int_{C} \mathbf{F}(\mathbf{r}) \cdot d\mathbf{r} = \int_{a}^{b} \mathbf{F}(\mathbf{r}(t)) \cdot \underbrace{\frac{d\mathbf{r}}{dt}}_{\mathbf{r}'(t)} dt$$

$$= \int_{C} (F_{1}dx + F_{2}dy + F_{3}dz)$$

$$= \int_{C} (F_{1}x' + F_{2}y' + F_{3}z')dt$$
If C is closed, 
$$\int_{C} \text{may be written as } \oint_{C} .$$

# Line Integral (cont)

The integral is scalar (: dot product).

$$\int_{C} \mathbf{F}(\mathbf{r}) \cdot d\mathbf{r} = \int_{a}^{b} \underbrace{\mathbf{F}(\mathbf{r}(t)) \cdot \frac{\mathbf{r}'}{|\mathbf{r}'|}}_{F_{\mathsf{tan}}} \underbrace{|\mathbf{r}'| dt}_{ds}$$

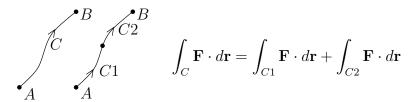
Rules for line integrals:

$$\int_{C} \alpha \mathbf{F}(\mathbf{r}) \cdot d\mathbf{r} = \alpha \int_{C} \mathbf{F}(\mathbf{r}) \cdot d\mathbf{r}, \quad (\alpha : \mathsf{const})$$

$$\int_{C} (\mathbf{F}(\mathbf{r}) + \mathbf{G}(\mathbf{r})) \cdot d\mathbf{r} = \int_{C} \mathbf{F}(\mathbf{r}) \cdot d\mathbf{r} + \int_{C} \mathbf{G}(\mathbf{r}) \cdot d\mathbf{r}$$

## Line Integral (cont)

If C is partitioned into C1 and C2, and all three have the same direction, then:

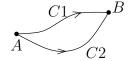


If the direction of traverse (orientation) is reversed, then the value of the integral changes sign.

## Vector Line Integral

$$\int_{C} \mathbf{F}(\mathbf{r})dt = \left[ \int_{C} F_{1}(\mathbf{r})dt \right] \mathbf{i} + \left[ \int_{C} F_{2}(\mathbf{r})dt \right] \mathbf{j}$$
$$+ \left[ \int_{C} F_{3}(\mathbf{r})dt \right] \mathbf{k}$$

### Path Dependence

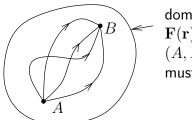


Generally,  $\int_A^B \mathbf{F} \cdot d\mathbf{r}$  depends on (1) the end points A and B, and (2) the path from A to B.

Hence, different paths (curves) from A to B may lead to different values of the integral.

Some integrals may be path-independent.

These are important items in physics.



domain D over which  $\mathbf{F}(\mathbf{r})$  is defined (A,B, any curve, must be inside D)

#### Theorem

The following statements are equivalent.

- 1.  $\int_A^B \mathbf{F}(\mathbf{r}) \cdot d\mathbf{r}$  is path-independent.
- 2.  $\mathbf{F} = \nabla f$ ; i.e. a potential (f) exists.
- 3.  $\oint \mathbf{F}(\mathbf{r}) \cdot d\mathbf{r} = 0$  for any closed curve in D.
- 4. If D is simply-connected,  $\nabla \times \mathbf{F} = 0$ .

From [2.], we can deduce:

$$\int_{B}^{A} \mathbf{F} \cdot d\mathbf{r} = f(B) - f(A), \text{ where } \mathbf{F} = \nabla f.$$

Note the similarity to definite integrals.

[3.] is used interchangeably with [1.], as in Thermodynamics (where it is used twice — once for the First Law, and once for the Second Law).

[4.] is based on the Pfaffian<sup>1</sup> form:

$$\mathbf{F} \cdot d\mathbf{r} = F_1 dx + F_2 dy + F_3 dz.$$

This form is called **exact**; the RHS is called an **exact differential** (or Pfaffian differential), if we do have a differentiable function f(x, y, z) s.t.

$$\mathbf{F} \cdot d\mathbf{r} = F_1 dx + F_2 dy + F_3 dz$$

$$= df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz$$

$$= \nabla f \cdot d\mathbf{r}$$

That means,  $\mathbf{F} = \nabla f$ ,

$$\therefore F_1 = \frac{\partial f}{\partial x}, F_2 = \frac{\partial f}{\partial y}, F_3 = \frac{\partial f}{\partial z}.$$

<sup>&</sup>lt;sup>1</sup>Johann Friedrich Pfaff (1765–1825), German mathematician.

For simply connected domains<sup>2</sup>

$$\label{eq:F} \text{if } \mathbf{F} = \nabla f$$
 then  $\nabla \times \mathbf{F} = \nabla \times (\nabla f) = 0$ 

By this, or by cross-differentiation,

$$\frac{\partial^2 f}{\partial x \partial y} = \frac{\partial^2 f}{\partial y \partial x} \implies \frac{\partial F_2}{\partial x} = \frac{\partial F_1}{\partial x} \implies \frac{\partial F_2}{\partial x} - \frac{\partial F_1}{\partial x} = 0,$$

etc.. : each component of  $\nabla \times F = 0$ . :  $\nabla \times F = 0$ .

$$\therefore V \times F = 0.$$

 $<sup>^2\</sup>text{A}$  SCD is one in which any closed curve can be continuously shruck to a point without leaving that domain.

A spacial case: Flatland, which is a domain in (x, y)-plane.

$$\int_{C} \mathbf{F}(\mathbf{r}) \cdot d\mathbf{r} = \int_{C} (F_1 dx + F_2 dy)$$

will be path independent

iff  $F_1dx + F_2dy$  is an exact differential,

i.e. iff 
$$\frac{\partial F_1}{\partial y} = \frac{\partial F_2}{\partial x}$$
,

i.e. iff 
$$\frac{\partial F_1}{\partial y} - \frac{\partial F_2}{\partial x} = 0$$
,

i.e. iff 
$$\nabla \times \mathbf{F} = \left(\frac{\partial F_1}{\partial y} - \frac{\partial F_2}{\partial x}\right) \mathbf{k} = 0.$$

#### Vector Integral Theorems

- ▶ Help us convert integrals from one form to another.
- ▶ volume integral ←→ surface integral.
- ▶ surface integral ←→ volume integral.

#### These are useful

- for their convenience often lead to ease of evaluation, and
- by making all integrals in an equation of the same type, often help convert integral equations into differential equations.

## Green's Theorem in the Plane (x, y)

R: closed, bounded region in (x,y) plane (need not be simply-connected).

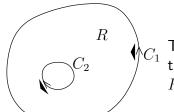
C: boundary of R, consisting of finitely many smooth curves.

 $F_1(x,y)$  and  $F_2(x,y)$ : continuous functions, s.t. the partial derivatives  $\partial F_1/\partial y$  and  $\partial F_2/\partial x$  exist and are continuous everywhere in R.

Then

$$\iint\limits_{R} \left(\partial F_2/\partial x - \partial F_1/\partial y\right) dx \, dy = \oint\limits_{C} (F_1 dx + F_2 dy).$$

## Green's Theorem (cont)

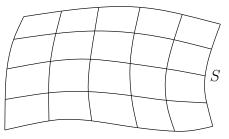


The line-integral on the RHS is on the entire  $C,\,$  traversed in such a way that R is on the left as C is traversed.

If  $\mathbf{F} \equiv F_1 \mathbf{i} + F_2 \mathbf{j}$ , then we get the vector form of Green's Theorem:

$$\iint\limits_{\mathcal{D}} (\nabla \times \mathbf{F}) \cdot \mathbf{k} \, dx \, dy = \oint_{C} \mathbf{F} \cdot d\mathbf{r}.$$

### Surface Integrals in 3-D



Surface in 3-D (assumed piecewise smooth)  $\mathbf{r}(u,v)$ :definition of the surface

 $\mathbf{r}(u,v)=x(u,v)\mathbf{i}+y(u,v)\mathbf{j}+z(u,v)\mathbf{k}$ , where (u,v) : defined over R in the u,v plane.

At any point on S,

$$\mathbf{r}_u = \left(\frac{\partial \mathbf{r}}{\partial u}\right)_v \text{ :tangent vector along a const-}u \text{ line,}$$
 
$$\mathbf{r}_v = \left(\frac{\partial \mathbf{r}}{\partial v}\right)_y \text{ :tangent vector along a const-}v \text{ line}$$
 
$$\therefore \ \mathbf{N} = \mathbf{r}_u \times \mathbf{r}_v \text{ is a normal vector, and}$$
 and 
$$\mathbf{n} = \frac{\mathbf{N}}{|\mathbf{N}|} \text{ is a unit normal vector.}$$

All of these will exist, except perhaps at edges and corners.

For F: a vector function,

$$\iint\limits_{S} \mathbf{F} \cdot \mathbf{n} \, dA \equiv \iint\limits_{R} \mathbf{F}(\mathbf{r}(u, v)) \cdot \mathbf{N}(u, v) \, du \, dv$$

is the definitition of the surface integral of  $\mathbf{F}$  over S.

 $|\mathbf{r}_u imes \mathbf{r}_v| = |\mathbf{N}| =$  area of parallelogram of sides  $\mathbf{r}_u$  and  $\mathbf{r}_v$ .

 $\therefore \mathbf{n} \, dA = \mathbf{n} |\mathbf{N}| \, du \, dv = \mathbf{N} \, du \, dv.$ 

 $\mathbf{F} \cdot \mathbf{n} = \text{normal component of } \mathbf{F} \text{ (normal to surface } S).$ 

$$\mathbf{F} = F_1 \mathbf{i} + F_2 \mathbf{j} + F_3 \mathbf{k},$$

$$\mathbf{N} = N_1 \mathbf{i} + N_2 \mathbf{j} + N_3 \mathbf{k},$$

$$\mathbf{n} = \cos \alpha \mathbf{i} + \cos \beta \mathbf{j} + \cos \gamma \mathbf{k}.$$

 $\cos \alpha = \mathbf{n} \cdot \mathbf{i}, \cos \beta = \mathbf{n} \cdot \mathbf{j}, \cos \gamma = \mathbf{n} \cdot \mathbf{k}$ , are the direction cosines for  $\mathbf{n}$ .

Thus,

$$\iint_{S} \mathbf{F} \cdot \mathbf{n} \, dA = \iint_{S} (F_{1}N_{1} + F_{2}N_{2} + F_{3}N_{3}) \, du \, dv$$

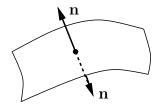
$$= \iint_{S} (F_{1}\cos\alpha + F_{2}\cos\beta + F_{3}\cos\gamma) \, dA$$
But  $\cos\alpha \, dA = dy \, dz$ 

$$\cos\beta \, dA = dz \, dx$$

$$\cos\gamma \, dA = dx \, dy$$

$$\iint_{S} \mathbf{F} \cdot \mathbf{n} \, dA = \iint_{S} (F_{1} \, dy dz + F_{2} \, dz dx + F_{3} \, dx dy)$$

#### Effect of Orientation



We have a choice of normals  $\mathbf{n}$  or  $-\mathbf{n}$  ( $\mathbf{N}$  or  $-\mathbf{N}$ ). This changes the value of the integral by a factor of (-1)

A smooth surface S is **orientable** if the positive normal direction at some point P on S can be moved in a continuous and unique way over the entire surface.

We assume our surfaces to be smooth and orientable.

### Triple Integrals

- ▶ These are integrals over a volume (dV = dx dy dz)
- $\blacktriangleright$  We consider a 3-D region T

#### Divergence Theorem of Gauss

Perhaps the most important and celebrated theorem in Vector Calculus.

Transforms a **triple integral over a volume** into a **surface integral over the surface bounding the volume**, and *vice-versa*.

### Divergence Theorem of Gauss (cont)

T: a closed, bounded region in space.

 ${\cal S}$  : the bounding surface of  ${\cal T}$  , piecewise smooth and orientable.

 $\mathbf{F}(x,y,z)$ : a vector function, continuous, with continuous first partial derivatives all over T (and S, and maybe beyond).

Then

$$\iiint\limits_{T}\nabla\cdot\mathbf{F}\,dV=\iint\limits_{S}\mathbf{F}\cdot\underbrace{\mathbf{n}\,dA}_{d\mathbf{S}}=\iint\limits_{S}\mathbf{F}\cdot d\mathbf{S}$$

## Divergence Theorem of Gauss (cont)

If  $\mathbf{F} = [F_1, F_2, F_3]$  and  $\mathbf{n} = [\cos \alpha, \cos \beta, \cos \gamma]$ , then the theorem becomes:

$$\iiint_{T} \left( \frac{\partial F_{1}}{\partial x} + \frac{\partial F_{2}}{\partial y} + \frac{\partial F_{3}}{\partial z} \right) dV$$

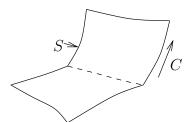
$$= \iint_{S} (F_{1} \cos \alpha + F_{2} \cos \beta + F_{3} \cos \gamma) dA$$

$$= \iint_{S} (F_{1} dy dz + F_{2} dz dx + F_{3} dx dy)$$

#### Stokes' Theorem

Stokes' Theorem is the 3-dimensional version of Green's Theorem.

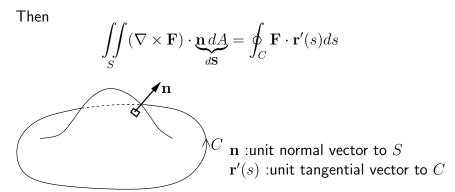
It relates a surface-integral over a 3-D surface to a line-integral over the bounding line in 3-D.



S :surface in 3-D, piecewise smooth, and oriented C :boundary of S, piecewise smooth, simple, closed

 $\mathbf{F}(x,y,z)$  :vector function, defined over S++, continuous, with continuous (first) partial derivatives.

# Stokes' Theorem (cont)



Direction of  $\oint$  : s.t. topologically, a vertically upwards  ${\bf n}$  is to the left of the traverse.

### Stokes' Theorem (cont)

Stokes' Theorem relates directly to path independence:

If the RHS equals zero everywhere (: path independent), then the LHS equals zero everywhere.

 $\implies \nabla \times F = 0$  everywhere.