

PHYS 425 - w5l1

Road Map to 10mK Understanding
Dilution Refrigerators.

Mixtures of ^3He - ^4He

1. Understand the difference between the behaviour of systems of identical particles at low temps. \Rightarrow Fermions & Bosons.

- ☒ Some QM review
- ☒ Develop w.f. for a pair of identical particles
- ☒ Deduce Pauli Exclusion principle
- ☒ Write down dist'n fns for Fermions & Bosons (hope to derive these later in the course)

2. Systems of Fermions at Low temp.

- Fermi Energy
- Density of States
- Low-temp heat capacity
- Low-temp Entropy

3. Mixtures of ^3He - ^4He

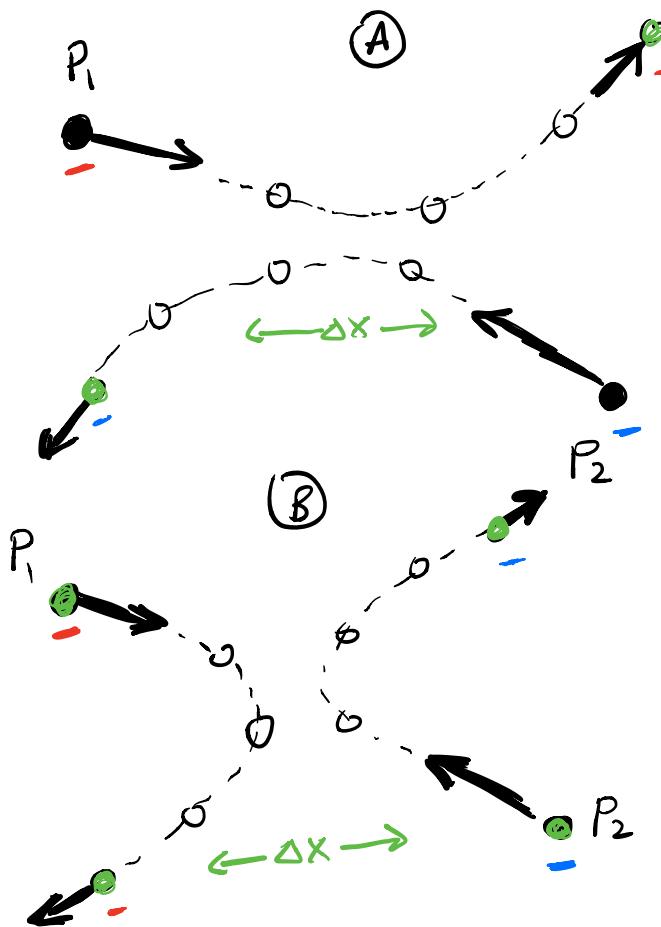
- Interpretation of Chemical Potential
- Energetics of ^3He - ^4He mixtures
- Gibbs Free Energy & Chemical Pot.

4. Dilution Refrigerator

- Analogy w/ Evaporative Cooling
- Cooling Power
- Base temp.
- Operation of a dilution Refrigerator

Identical Particles

Classically, can always distinguish between particles. However, this is not always possible in case of Quantum Mechanics.



Classically, can track position & momentum of P₁ & P₂ w/ arbitrary precision & can always distinguish between scenario A & B even if P₁ & P₂ are identical.

However, b/c of Heisenberg uncertainty principle

$$(\text{HUP}) \quad \Delta x \Delta p \geq \frac{\hbar}{2} \quad \text{in Q.M.}$$

cannot always tell which particle is which after scattering event. If particles come within Δx of one another lose ability to distinguish between the two \therefore between scenarios

(A) $\&$ (B).

This implies that, for identical particles, exchanging positions of $p_1 \& p_2$ leaves system unchanged.

Physically measurable quantities in Q.M.
determined from abs. square of wavefn:

E.g. $|\psi(x)|^2 dx$ gives prob. of finding particle
 \sim between $x \& x+dx$

$$\begin{aligned}
 \langle \psi | \hat{H} | \psi \rangle &= E \langle \psi | \psi \rangle \\
 &= E \int \psi^*(x) \psi(x) dx \\
 &= E \underbrace{\int | \psi(x) |^2 dx}_{\text{is expectation value of system energy.}}
 \end{aligned}$$

For a system of two identical particles require

$$\Rightarrow |\psi(x_1, x_2)|^2 dx_1 dx_2 = |\psi(x_2, x_1)|^2 dx_1 dx_2$$

↑ ↑ ↑ ↑
 P₁ at x₁ P₂ at x₂ P₁ at x₂ P₂ at x₁

i.e. since particles are identical, prob. of finding
 P₁ at x₁, P₂ at x₂ same as prob. of finding
 P₂ at x₁, P₁ at x₂.

This condition is satisfied if

$$\textcircled{A} \quad \boxed{\psi(x_2, x_1) = e^{i\delta} \psi(x_1, x_2)}$$

where $e^{i\delta}$ is
a phase factor.

$$|\psi(x_2, x_1)|^2 = e^{-i\delta} \psi^*(x_1, x_2) e^{+i\delta} \psi(x_1, x_2)$$

$$= |\psi(x_1, x_2)|^2$$

Define exchange operator

$$P_{12} \psi(x_1, x_2) = \psi(x_2, x_1) = e^{i\delta} \psi(x_1, x_2)$$

$$\begin{aligned} & \vdash P_{12} \left(\underbrace{P_{12} \psi(x_1, x_2)}_{\psi(x_2, x_1)} \right) = P_{12} \underbrace{\psi(x_2, x_1)}_{e^{i\delta} \psi(x_1, x_2)} \\ & = e^{i\delta} \left(\underbrace{P_{12} \psi(x_1, x_2)}_{\psi(x_2, x_1)} \right) = e^{i\delta} \underbrace{\psi(x_2, x_1)}_{e^{i\delta} \psi(x_1, x_2)} \\ & \qquad \qquad \qquad = \boxed{e^{2i\delta} \psi(x_1, x_2)} \end{aligned}$$

Must also be the case that

$$P_{12} \left(P_{12} \psi(x_1, x_2) \right) = \boxed{\psi(x_1, x_2)}$$

$$\therefore \Psi(x_1, x_2) = e^{2i\delta} \Psi(x_1, x_2)$$

Holds only if $e^{2i\delta} = 1$

or
$$e^{i\delta} = \sqrt{1} = \pm 1$$

$$e^{i\delta} = \cos \delta + i \sin \delta = \pm 1$$

true if $\delta = 0, \pi$

Returning to ④, we find

$$\Psi(x_2, x_1) = \pm \Psi(x_1, x_2)$$

If $\Psi(x_1, x_2) = +\Psi(x_2, x_1)$ w.f. is said to
be symmetric

If $\Psi(x_1, x_2) = -\Psi(x_2, x_1)$ w.f. is said to
antisymmetric

Suppose that our system is two identical non-interacting particles.

Eg. Particles in a box.

Hamiltonian for P_1 is \mathcal{H}_1 s.t.

$$\underbrace{\mathcal{H}_1 \psi_A(x_1)}_{\sim} = E_A \psi_A(x_1)$$

P_1 is state A w/ energy E_A

Hamiltonian for P_2 is \mathcal{H}_2 s.t.

$$\underbrace{\mathcal{H}_2 \psi_B(x_2)}_{\sim} = E_B \psi_B(x_2)$$

P_2 in state B w/ energy E_B .

For example for free particles $\mathcal{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2}$

ψ_A & ψ_B are single-particle wavefns of P_1 & P_2

Eg. particles in 1-D box

$$\phi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

\therefore Requirements for two-particle w.f. for net system are:

$$\textcircled{1} \quad \psi(x_1, x_2) = \pm \psi(x_2, x_1) \quad \text{exchange property.}$$

$$\textcircled{2} \quad (\underbrace{f_{l_1} + f_{l_2}}_{\substack{\text{not hamiltonian} \\ \text{of non-interacting} \\ \text{system.}}} \psi(x_1, x_2) = (E_1 + E_2) \psi(x_1, x_2)$$

w.f. must be
an eigenstate
of net Hamiltonian.

Try a guess

$$\psi(x_1, x_2) = \phi_A(x_1) + \phi_B(x_2)$$

P_1 in state A P_2 in state B

test the exchange property

$$\psi(x_2, x_1) = \phi_A(x_2) + \phi_B(x_1) \neq \pm \psi(x_1, x_2)$$

P_2 in state A P_1 in state B

Fails exchange test.

Try another guess

$$\Psi(x_1, x_2) = \mathcal{D}_A(x_1) \mathcal{D}_B(x_2)$$

exchange $\Psi(x_2, x_1) = \phi_A(x_2) \phi_B(x_1)$ fails
 exchange requirement.

Finally try

$$\Psi(x_1, x_2) = \underbrace{\mathcal{D}_A(x_1) \mathcal{D}_B(x_2) + \mathcal{D}_A(x_2) \mathcal{D}_B(x_1)}_{\text{Superposition of } p_1 \text{ in A, } p_2 \text{ in B}} + \underbrace{\frac{1}{\sqrt{2}} [P_1(x_1) P_2(x_2) - P_2(x_1) P_1(x_2)]}_{\text{Antisymmetric part}}$$

Makes some sense b/c w/ two identical particles one in state & other in state B, can't tell which is in A & which is in B.

Try exchange:

$$\Psi(x_2, x_1) = \phi_A(x_2)\phi_B(x_1) + \phi_A(x_1)\phi_B(x_2)$$

$$= \pm (\cancel{\phi_A(x_1)\phi_B(x_2)} \pm \cancel{\phi_A(x_2)\phi_B(x_1)})$$

$$\Psi(x_1, x_2)$$

$$\therefore \Psi(x_2, x_1) = \pm \Psi(x_1, x_2) \quad \text{Exchange requirement satisfied!}$$

Check that is two-particle w.f. is an eigenstate of $\hat{H}_1 + \hat{H}_2$.

$$(\hat{H}_1 + \hat{H}_2) \Psi(x_1, x_2) = (\hat{H}_1 + \hat{H}_2) [\psi_A(x_1) \psi_B(x_2) \pm \psi_A(x_2) \psi_B(x_1)]$$

$$= (\underbrace{\hat{H}_1 \psi_A(x_1)}_{E_A \psi_A(x_1)}) \psi_B(x_2) \pm \psi_A(x_2) (\underbrace{\hat{H}_1 \psi_B(x_1)}_{E_B \psi_B(x_1)}) \\ + \psi_A(x_1) (\underbrace{\hat{H}_2 \psi_B(x_2)}_{E_B \psi_B(x_2)}) \pm (\underbrace{\hat{H}_2 \psi_A(x_2)}_{E_A \psi_A(x_2)}) \psi_B(x_1)$$

$$= E_A \psi_A(x_1) \psi_B(x_2) \pm E_B \psi_A(x_2) \psi_B(x_1)$$

$$+ E_B \psi_A(x_1) \psi_B(x_2) \pm E_A \psi_A(x_2) \psi_B(x_1)$$

$$= E_A (\underbrace{\psi_A(x_1) \psi_B(x_2) \pm \psi_A(x_2) \psi_B(x_1)}_{\Psi(x_1, x_2)})$$

$$+ E_B (\underbrace{\psi_A(x_1) \psi_B(x_2) \pm \psi_A(x_2) \psi_B(x_1)}_{\Psi(x_1, x_2)})$$

$$\therefore (\mathcal{H}_1 + \mathcal{H}_2) \Psi(x_1, x_2) = (E_A + E_B) \Psi(x_1, x_2)$$

\therefore our two-particle w.f. is an eigenstate
of the total Hamiltonian.

Symmetric 2-particle w.f.

$$\Psi_s(x_1, x_2) = \mathcal{D}_A(x_1) \mathcal{D}_B(x_2) + \mathcal{D}_A(x_2) \mathcal{D}_B(x_1)$$

Antisymmetric 2-particle w.f.

$$\Psi_a(x_1, x_2) = \mathcal{D}_A(x_1) \mathcal{D}_B(x_2) - \mathcal{D}_A(x_2) \mathcal{D}_B(x_1)$$