Recall the Grand Partition fon:

$$= = \sum_{\substack{\text{all} \\ \text{states}}} -(\epsilon_i - \mu_i)/k_B T$$

For a single energy level that can be occupied by any number of Bosons:

$$\frac{1}{1-e^{-(\xi(k)-\mu)/k_{\rm B}T}} = \left[1-e^{-(\xi(k)-\mu)/k_{\rm B}T}\right]^{-1}$$
Single energy level

If consider Grand Canonical Ensemble of entire system (many levels of different & values):

$$= \prod_{k} = \prod_{k} (1-e^{-(\epsilon(k)-\mu)/k_BT})^{-1}$$

Eq. Consider 2 levels each 
$$\omega/0,1$$
, or 2 particles
$$\frac{\epsilon_2}{\epsilon_1} = \frac{1}{-(\epsilon_1-\mu)(k_3T)} = -(2\epsilon_1-2\mu)(k_3T)$$

$$= \frac{1 + e^{-(\xi_1 - \mu)/k_B T} - (2\xi_1 - 2\mu)/k_B T}{+ e^{-(\xi_2 - 2\mu)/k_B T}}$$

$$+ \frac{e^{-(\xi_2 - \mu)/k_B T} - (\xi_1 + \xi_2 - 2\mu)/k_B T}{+ e^{-(2\xi_1 + \xi_2 - 3\mu)/k_B T}}$$

$$+ \frac{e^{-(2\xi_1 + \xi_2 - 3\mu)/k_B T}}{+ e^{-(2\xi_1 + 2\xi_2 - 4\mu)/k_B T}}$$

$$+ \frac{e^{-(2\xi_1 + 2\xi_2 - 4\mu)/k_B T}}{+ e^{-(2\xi_1 + 2\xi_2 - 4\mu)/k_B T}}$$

c.t. 
$$(1+e^{-(\xi_1-\mu)/k_BT} - (2\xi_1-2\mu)/k_BT)$$
.  
 $(1+e^{-(\xi_2-\mu)/k_BT} - (2\xi_2-2\mu)/k_BT)$   
 $(1+e^{-(\xi_2-\mu)/k_BT} + e^{-(\xi_2-2\mu)/k_BT})$ 

$$:= (=,)(=_z)$$

In general,

$$= \frac{1}{k} \left( 1 - e^{-\left( \mathcal{E}(k) - \mu \right) / k_B T} \right)^{-1}$$

.: grand Potential for many-level system is

$$\overline{Q}_{q} = -k_{B}T \ln \overline{\Xi}$$

$$= -k_{B}T \ln \left[ TT \left( 1 - e^{-(\xi(k) - \mu)/k_{B}T} \right)^{-1} \right]$$

$$= -k_{B}T \overline{\Xi} \ln \left[ 1 - e^{-(\xi(k) - \mu)/k_{B}T} \right]^{-1}$$

Recall density of states depends only on mag. of k

 $\int_{k}^{\infty} f(k) = \frac{V}{2\pi^{2}} \int_{0}^{\infty} k^{2} f(k) dk$ 4 any for of mag. of  $\vec{k}$ 

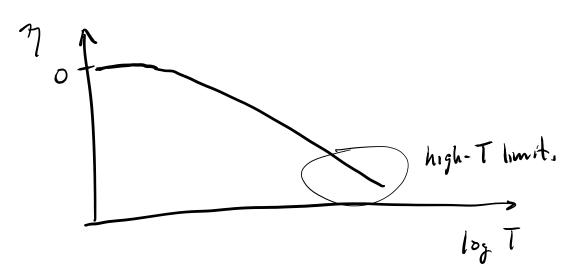
$$\frac{1}{2\pi^2} \sum_{k_B} \left\{ \int_{0}^{\infty} k^2 \ln \left[ 1 - e^{-\left( \mathcal{E}(k) - \mu_0 \right) / k_B T} \right] dk \right\}$$

General expression for grand potential of system of Bosons

- valid for any temp T - valid for any form of Elk)

Low-temp Limit:

know 
$$7 = \frac{u}{k_p T}$$
 goes to zero as  $T \to 0$ 



$$\overline{\Psi}_{a} \approx \frac{V}{\lambda \pi^{2}} k_{B} T \int_{0}^{\infty} k^{2} \ln \left[ 1 - e^{-\varepsilon(k)/k_{B}T} \right] dk$$

only valid for Bosons for which

u > 0 as T > 0

$$\frac{\mu}{k_n T} \rightarrow 0 \text{ as } T \rightarrow 0$$

high-temp limit:

For a gas of particles (Fermions or Bosons)

in dilute high-temp limit

$$\eta = \frac{m}{k_B T} = \ln \left( \frac{n}{n_Q} \right)$$
 is large of neg.

$$N_{Q} = \left(\frac{m k_{B}T}{2\pi h^{2}}\right)^{3/2} \qquad N = \frac{N}{V}$$

$$\overline{N}_{BF} = \frac{1}{e^{(\mathcal{E}(k) - \mu)/k_B T}} \approx e^{-(\mathcal{E}(k) - \mu)/k_B T}$$

Recall In (1±x) ≈ ±x

for |x| << |

$$\approx -e^{-(\xi(k)-\mu J/k_n T)}$$

: In high-T (dilute) limit
$$\Phi_{a} \approx -\frac{V}{2\pi^{2}} k_{B} T e^{\eta} \int_{0}^{\infty} \frac{1}{k^{2}} e^{-\epsilon(k)/k_{B}T} dk \qquad (2)$$

## valid for Bosons & Fermions.

Recall 
$$\Phi_{\alpha} = U - TS - \mu N$$

$$d\Phi_{\alpha} = dU - TdS - SdT - \mu dN - Nd\mu$$

$$= TdS - PdV + \mu dN - TdS - SdT - \mu dN - Nd\mu$$

$$d\Phi_{\alpha} = -PdV - SdT - Nd\mu$$

$$S = -\left(\frac{\partial \Phi_{\alpha}}{\partial V}\right)_{T,\mu}$$

$$S = -\left(\frac{\partial \Phi_{\alpha}}{\partial T}\right)_{V,\mu} \longrightarrow Cv = T\frac{\partial S}{\partial T}$$

$$N = -\left(\frac{\partial \Phi_{\alpha}}{\partial \mu}\right)_{V,T} \longrightarrow con^{14} use in low-T limit blocs set  $\frac{\mu}{k_{B}T} = 0$$$

For gas of free Bosons 
$$(E(k) = \frac{\hbar^2 k^2}{am})$$

of use 3 to find in high-T limit:

$$S = Nk_B \left[ \ln \left( \frac{n_Q}{n} \right) + \frac{5}{2} \right]$$
 Sackur-Tetrode eq.m.

leave for homework

Recovered classical results starting from Quantum theory of identical particles.

For gas of free Bosons use 1) to find in low-Tlimit;

$$P \propto T^{5/2}$$
  
 $S \propto T^{3/2}$   
 $C_{V} \propto T^{3/2}$   
homework.