

# **Pseudospectral Time-Domain Method in EM Simulations**

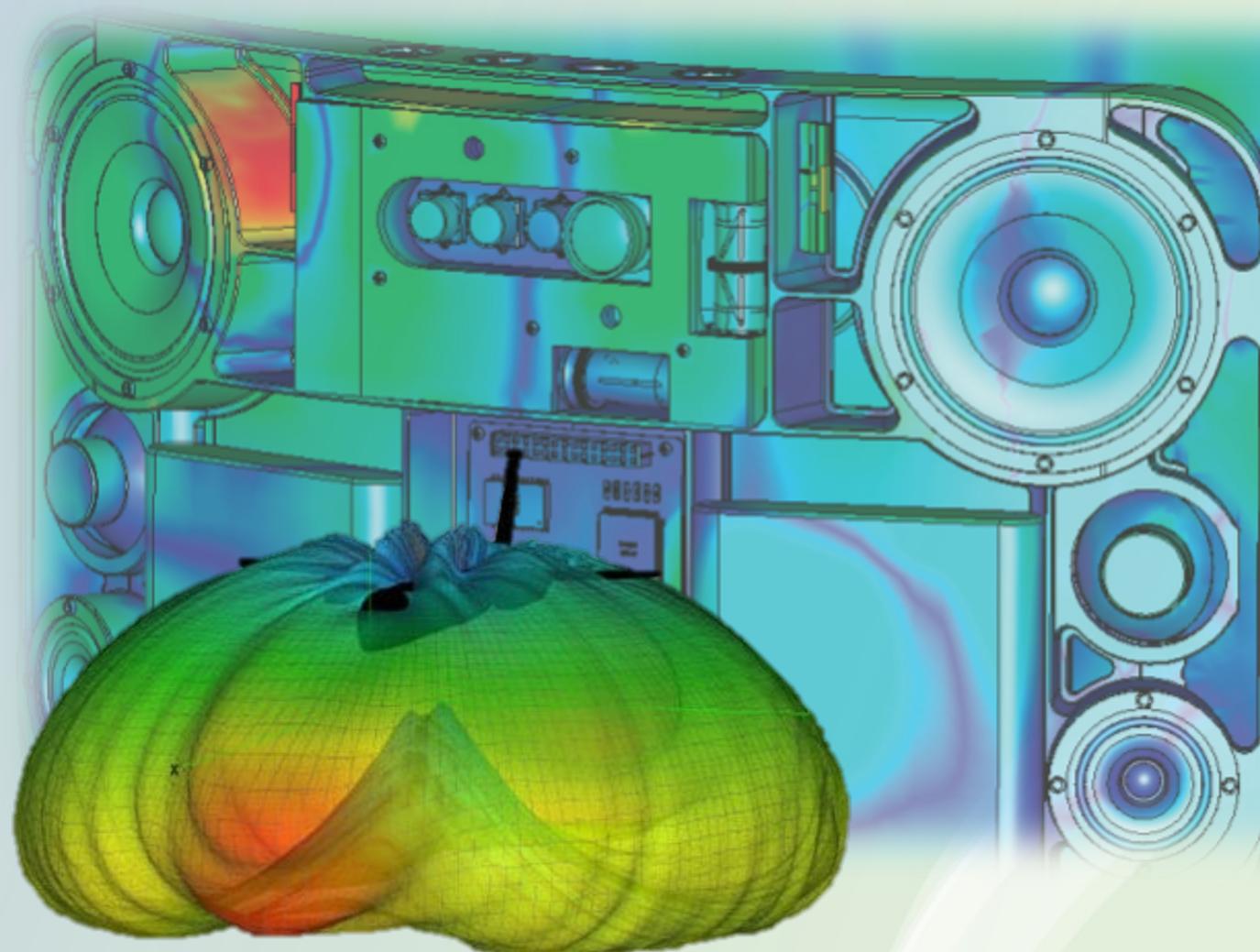
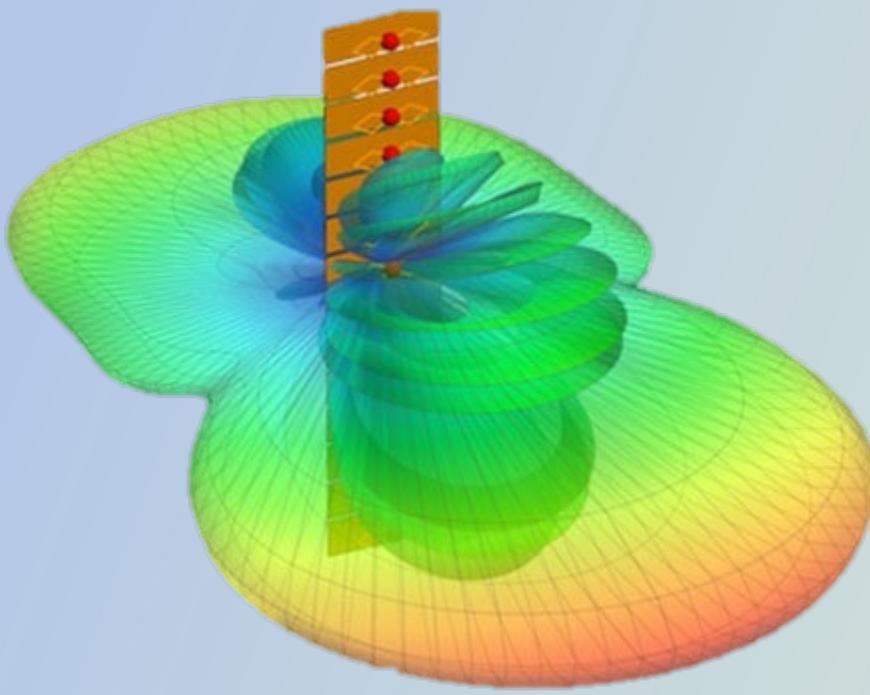
Speaker: Jake W. Liu

# Outline

- Basics of Computational Electromagnetics
- Recap of FDTD
- Introduction to PSTD
- Advanced topics in PSTD

# Numerical Simulations in EM

What are the applications?



- <https://altair.com/feko-applications>
- [https://www.3ds.com/fileadmin/PRODUCTS-SERVICES/SIMULIA/PRODUCTS/CST/CST\\_Brochure\\_A4.pdf](https://www.3ds.com/fileadmin/PRODUCTS-SERVICES/SIMULIA/PRODUCTS/CST/CST_Brochure_A4.pdf)

# Some Applications

- Antenna radiation modeling
- Device modeling (SIP / EMC)
- Metamaterials and nanostructures
- Wave propagation and scattering
- ...

Based on the application requirements, we select an appropriate computational electromagnetics (CEM) method for modeling.

	<b>IE</b>	<b>DE</b>
<b>TD</b>	$\oint \mathbf{E} \cdot d\mathbf{l} = -\frac{d}{dt} \int \mathbf{B} \cdot d\mathbf{s}$ $\oint \mathbf{H} \cdot d\mathbf{l} = \int (\mathbf{J} + \frac{d}{dt} \mathbf{D}) \cdot d\mathbf{s}$ $\oint \mathbf{D} \cdot d\mathbf{s} = \int q \, dv$ $\oint \mathbf{D} \cdot d\mathbf{s} = 0$	$\nabla \times \mathbf{E} = -\partial \mathbf{B} / \partial t$ $\nabla \times \mathbf{H} = \mathbf{J} + \partial \mathbf{D} / \partial t$ $\nabla \cdot \mathbf{D} = q$ $\nabla \cdot \mathbf{B} = 0$
<b>FD</b>	$\oint \mathbf{E} \cdot d\mathbf{l} = -i\omega \int \mathbf{B} \cdot d\mathbf{s}$ $\oint \mathbf{H} \cdot d\mathbf{l} = \int (\mathbf{J} + i\omega \mathbf{D}) \cdot d\mathbf{s}$ $\oint \mathbf{D} \cdot d\mathbf{s} = \int q \, dv$ $\oint \mathbf{D} \cdot d\mathbf{s} = 0$	$\nabla \times \mathbf{E} = -i\omega \mathbf{B}$ $\nabla \times \mathbf{H} = \mathbf{J} + i\omega \mathbf{D}$ $\nabla \cdot \mathbf{D} = q$ $\nabla \cdot \mathbf{B} = 0$

# Computational Electromagnetics

- Well known numerical methods are listed

	<b>IE</b>	<b>DE</b>
<b>TD</b>	<b>TDIE</b>	<b>FDTD</b>
<b>FD</b>	<b>MOM</b>	<b>FEM</b>

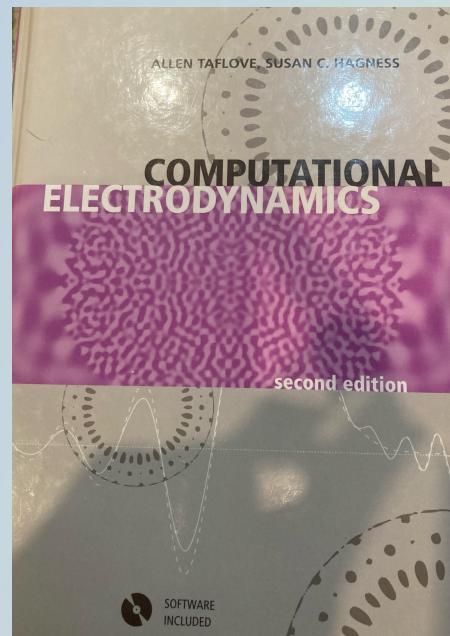
These are called full-wave methods

# Asymptotic Methods

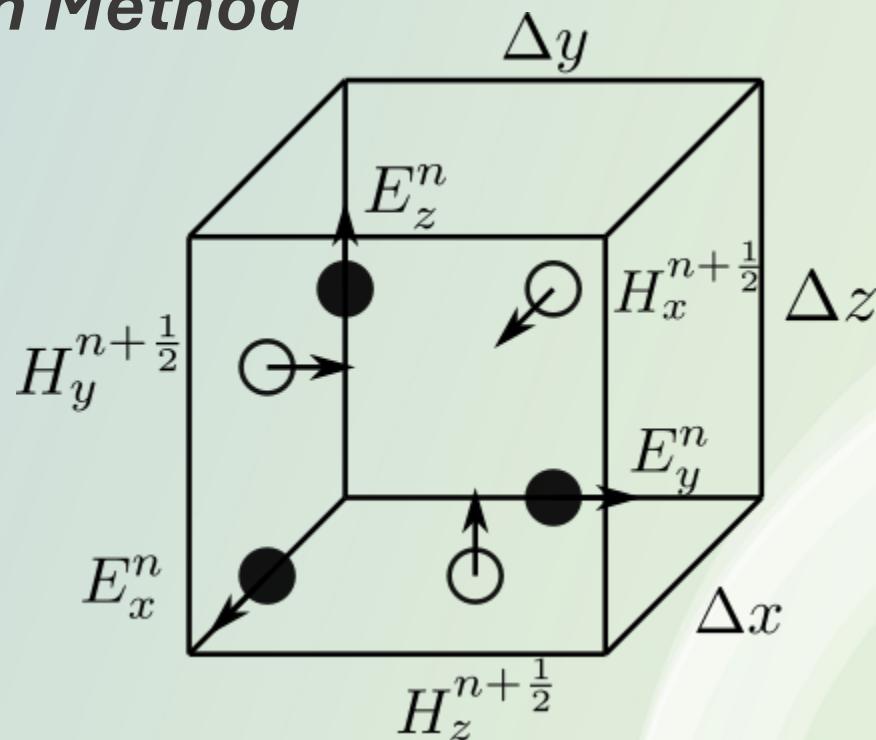
- Asymptotic methods can be applied for the **electrically-large cases**:
  - ⇒ Geometrical optics (GO)
  - ⇒ Physical Optics (PO)
- And their advanced version:
  - ⇒ Geometrical Theory of Diffraction (GTD)
  - ⇒ Physical Theory of Diffraction (PTD)

# Finite-Difference Time-Domain Method

- Kane Yee, 1966, "*Numerical solution of initial boundary value problems involving Maxwell's equations in isotropic media*"
- Allen Taflove, 1995, **Computational Electrodynamics: The Finite-Difference Time-Domain Method**



This is my 2<sup>nd</sup> edition.  
I got my 3<sup>rd</sup> edition  
recently 😊



# FDTD Formulation

- Discretization in time => FD
- Discretization in space => FD

2<sup>nd</sup> order accuracy

$$\frac{E_{x,k}^{n+1/2} - E_{x,k}^{n-1/2}}{\Delta t} = - \frac{H_{y,k+1/2}^n - H_{y,k-1/2}^n}{\epsilon \Delta z}$$

2<sup>nd</sup> order accuracy

# Topics in FDTD

- Discretization
- Material
- Perfectly Matching Layer (PML)
- Hard Source / Soft Source
- Pure Scattered-Field (SF) / Total-Field Scattered-Field (TFSF)
- Near-to-Far Field Transformation (NTFF)

# Pseudospectral Time-Domain Method

- Q. H. Liu, 1997, “*The PSTD algorithm: A time-domain method requiring only two cells per wavelength*”
- Discretization in time => **FD**
- Discretization in space => **PS**
  - using global basis function to approximate differential operations
    - **Fourier** => uniform collocated points
    - **Multi-domain** => non-uniform collocated points

# Fourier PSTD Formulation

- Denoting the Fourier transform operator as

$$\begin{aligned}\Psi(\mathbf{k}) &= \mathfrak{F}_\eta(\psi) = \int_{-\infty}^{\infty} \psi(\mathbf{r}) e^{-ik_\eta \eta} d\eta \\ \psi(\mathbf{r}) &= \mathfrak{F}_\eta^{-1}(\Psi) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \Psi(\mathbf{k}) e^{ik_\eta \eta} dk_\eta\end{aligned}$$

- Then

$$\begin{aligned}\partial_\eta \psi &= \partial_\eta \left[ \frac{1}{2\pi} \int_{-\infty}^{\infty} \Psi(\mathbf{k}) e^{ik_\eta \eta} dk_\eta \right] = \frac{i k_\eta}{2\pi} \int_{-\infty}^{\infty} \Psi(\mathbf{k}) e^{ik_\eta \eta} dk_\eta = \\ &\quad \mathfrak{F}_\eta^{-1}(i k_\eta \Psi) = \mathfrak{F}_\eta^{-1}[i k_\eta \mathfrak{F}_\eta(\psi)]\end{aligned}$$

$$\frac{E_{x,\hbar}^{n+1/2} - E_{x,\hbar}^{n-1/2}}{\Delta t} = -\frac{1}{\varepsilon} \mathfrak{F}_z^{-1}[i k_z \mathfrak{F}_z(H_y^n)]|_\hbar$$

# Implementation of PSTD

- FFT can be applied to compute the derivatives.
- Be aware of the use of **fftshift/ifftshift** if k vector is defined as zero-centered.
- To avoid numerical errors, taking **the real part** of the computed result is necessary!

```
dF = real.(ifft(im .* K .* fft(F, dim), dim))
```

# Advantages and Disadvantages of PSTD

- Advantages
  - Requires only two cells per wavelength according to Nyquist theorem (theoretically speaking...)
  - Computations for E and H fields are collocated\*
- Disadvantages
  - Difficult to deal with field discontinuities
    - Source
    - TFSF formulation
    - Metallic surfaces
    - High-contrast dielectrics

# Topics in PSTD (vs. FDTD)

- Discretization
- Material
- Perfectly Matching Layer (PML)
- Hard Source / Soft Source
- Pure Scattered-Field / Total-Field Scattered-Field (TFSF)
- Near-to-Far Field Transformation (NTFF)

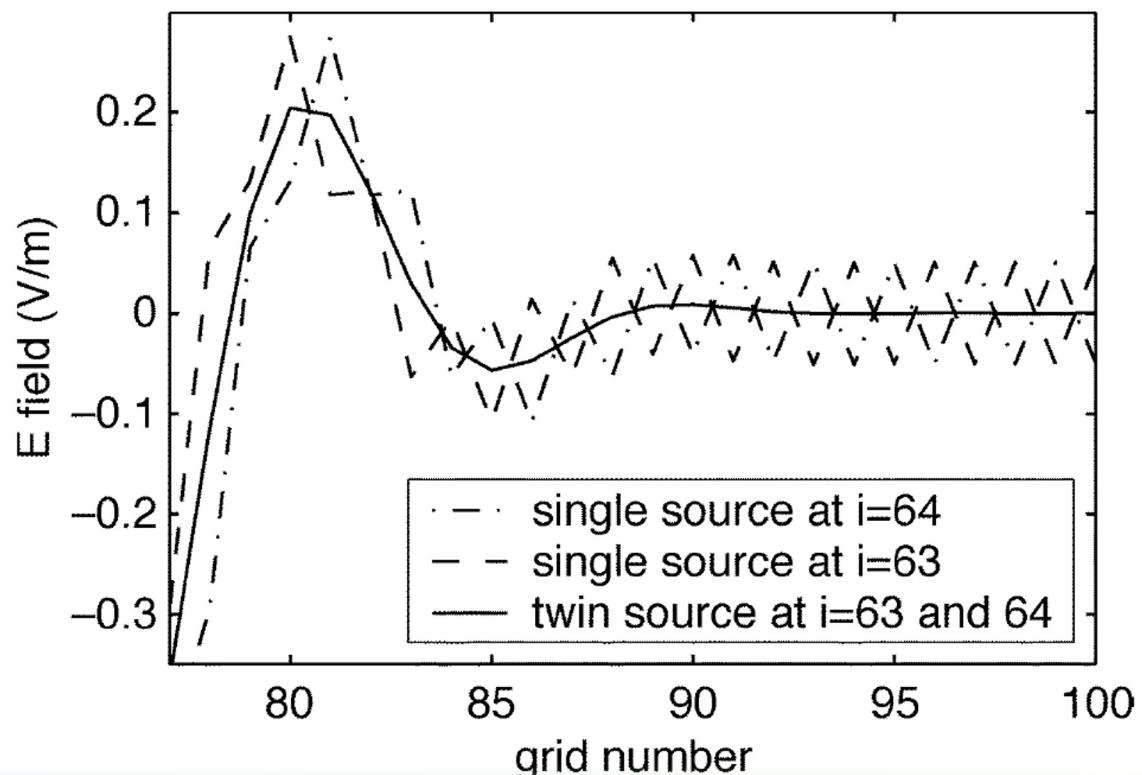
How does the topic differ from that of FDTD?

- Few
- Medium
- High

# Compact Wave Source Condition

- Tae-Woo Lee and Susan C. Hagness, 2004, “**A Compact Wave Source Condition for the Pseudospectral Time-Domain Method**”

Discontinuity  
introduced by  
point sources



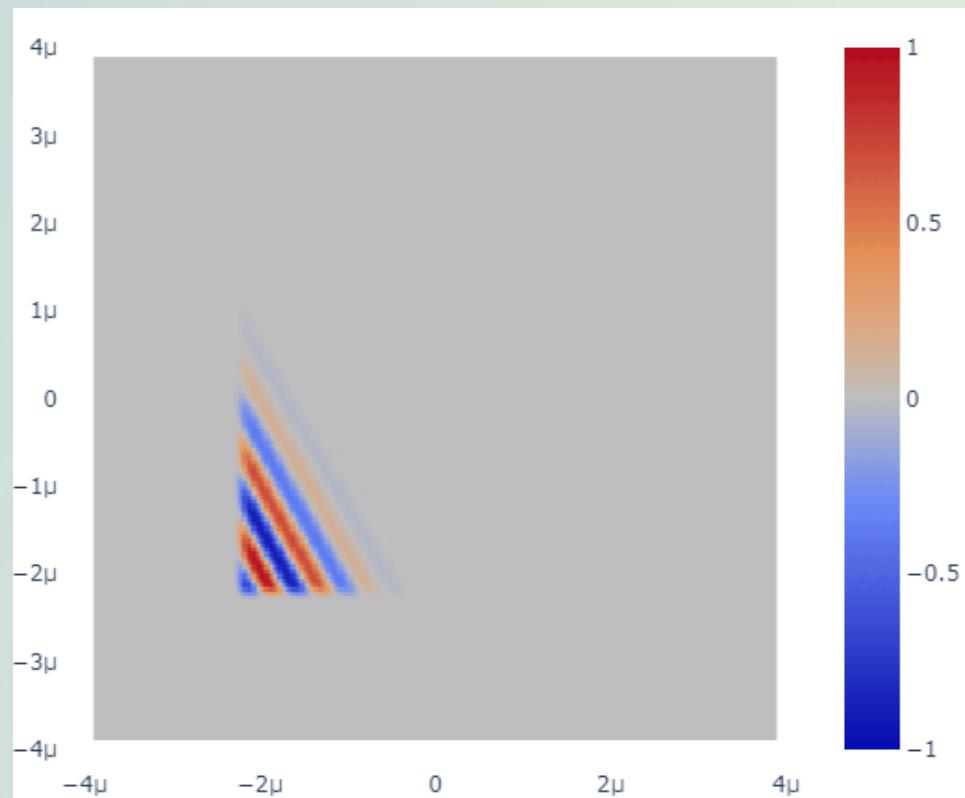
# SF / TFSF Formulation

- SF / TFSF are techniques to introduce plane wave into the space
- Decompose the total field into incident field and scattered field as

$$E_{tot} = E_{scat} + E_{inc}$$

→ assumed known

Illustration of the  
TFSF technique

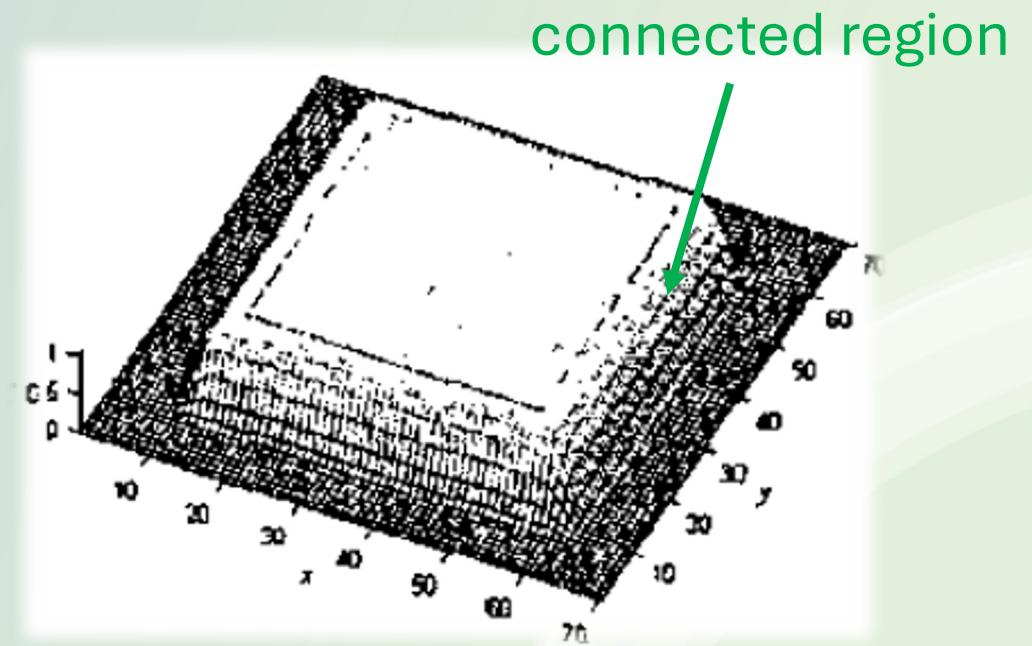


# TFSF Formulation in PSTD

- Xiang Gao, Mark S. Mirochnik and Dennis W. Prather, 2004, “**Soft Source Generation in the Fourier PSTD Algorithm**”

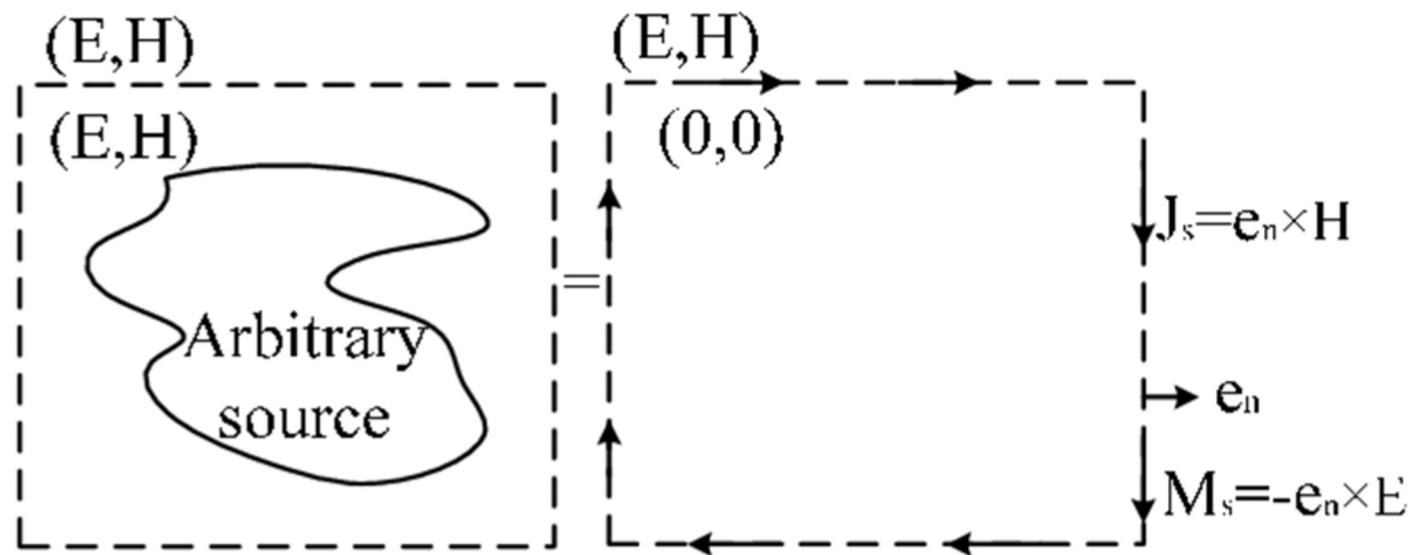
Discontinuity introduced  
by TFSF boundaries

$$\begin{cases} \hat{\vec{E}}_{tot} = \hat{\vec{E}}_{inc} + \vec{E}_{scat} = \zeta \vec{E}_{inc} + \vec{E}_{scat} \\ \hat{\vec{H}}_{tot} = \hat{\vec{H}}_{inc} + \vec{H}_{scat} = \zeta \vec{H}_{inc} + \vec{H}_{scat} \end{cases}$$



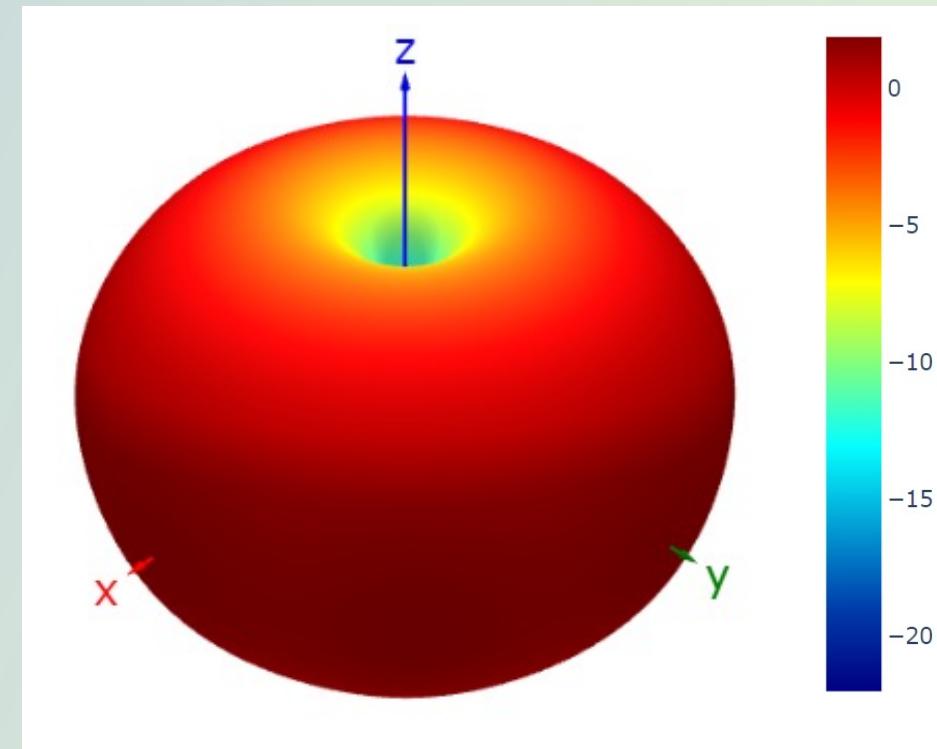
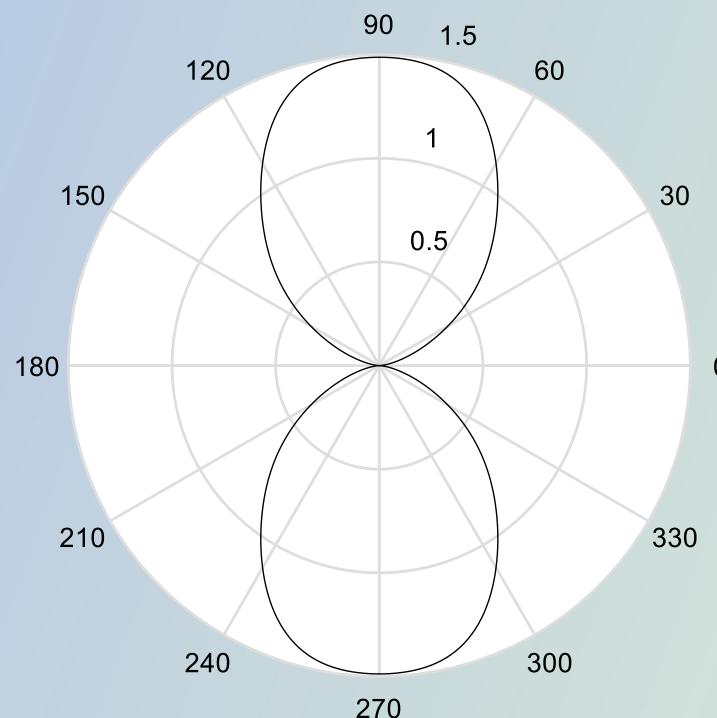
# NTFF Transformation

- From equivalence principle, when the surface currents on a selected imaginary surface are known, the fields inside the surface or outside the surface can be deduced from the imaginary currents.



# NTFF Transformation

- From the imaginary currents, one can compute the far field of the source with NTFF transformation.



# Staggered PSTD

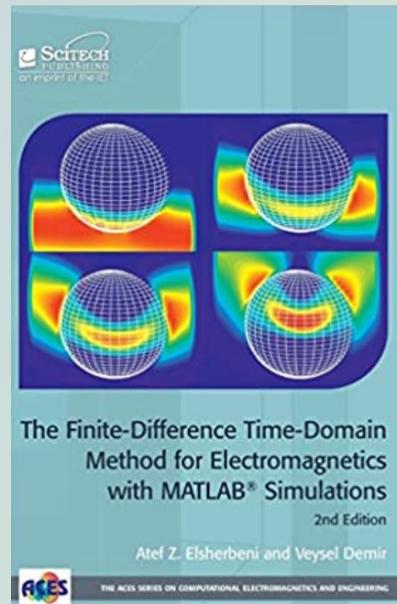
- M. Ding and K. Chen, 2010, “***Staggered-grid PSTD on local Fourier basis and its applications to surface tissue modeling***”
- Set PSTD grid to Yee grid.
- Advantages include mitigation of discontinuities and least modification of original FDTD code.

# Some Advices on Implementing Your Own Code...

- FDTD (or PSTD) is relatively easy to implement as compared to other CEM methods (there are lots of “toy” codes out there...)
- Start with the basics, progressing from 1D to 2D to 3D. This approach helps avoid getting bogged down in details.
- Find some simple problems to model (dipole radiation, reflection from a plane, etc.) → open your college EM book

# Some Useful References For You to Start

- John B. Schneider, *Understanding the FDTD Method*
- Atef Z. Elsherbeni and Veysel Demir, *The Finite-Difference Time-Domain Method for Electromagnetics with MATLAB® Simulations*



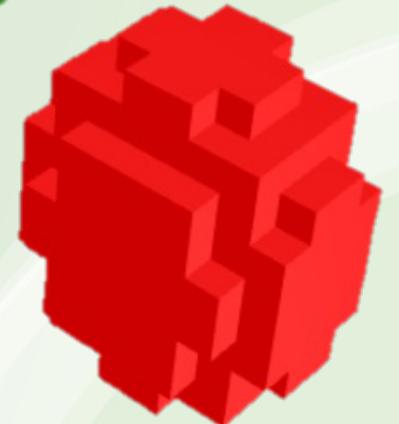
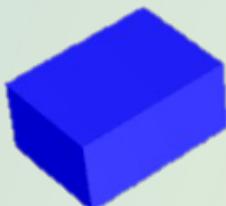
# Some Interesting Tools I Made (in Julia)

- **VoxelModel.jl**

- Create voxel models with simple geometries

- **RadiationPatterns.jl**

- Plot 2D/3D radiation patterns using PlotlyJS.jl



# Exercises

- Try to modify your 1D FDTD code into PSTD formulation
- Find out the differences in the implementation
- Compare results of both methods with different simulation scenarios

# Thank You!