PHYS 5321 Mechanics Fall 2021 8/24/21

What is mechanics? . What is classical mechanics?

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. Is the universe quahtum or classical?

. Is classical mechanics "complete"? Has everything been discount.

. What are the main concepts of classical mechanics?

. What are the formulations of classical mechanics?

There are "levels of theory." For example, Newtonian gravitation and Coulomb's law are empirical laws. They describe how masses behave in the presence of other masses and charges in the presence of other charges, but their applicability is limited to masses and charges, respectively.

A How would a "higher level" theory look like? Higher level of abstraction? More "fundamental"?

A higher level theory would tell you, e.g., the properties that empirical laws must have, rules that they must follow, the structure of the theory. The more fundamental a theory is, the more general the principles it is based on must be what is the most fundamental principle you can think of;

Conservation laws (and some definitions)

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HG Eq. 1.3

Consider Newton's Second law F= ma = dp = p with $\vec{p} = m\vec{v}$, $\vec{a} = \frac{d}{dt}\vec{v}$, $\vec{v} = \frac{d}{dt}\vec{r}$. For a particle.

 $f \vec{F} = 0 \Rightarrow \frac{d\vec{p}}{dt} = 0$

p' 15 fre 1, near momentum

 $d\vec{p} = 0$

 $\int d\vec{p} = 0$

 $\vec{p} + C = 0$

 $\vec{p} = constant$

Also known as Newton's first law

Conservation theorem for the Inear momentum of a particle: if the total force is zero, linear momentum is conserved

Let's know define the angular momentum as $\vec{L} = \vec{r} \times \vec{p}$ about point \vec{R} about point o

And the torque as $\vec{N} = \vec{r} \times \vec{p} \vec{F} = \vec{r} \times \frac{d}{dt} (m\vec{v})$ $d(u \cdot v) = udv + vdu \times product rule$

d (rxmv) = rxd (mv) + mvxd rxd dt latell

 $\frac{d}{dt} = \vec{r} \times \frac{d}{dt} (m\vec{v}) + m\vec{v} \times \vec{v}$ vanishes since angle is zero

 $\frac{d\vec{L}}{dt} = \vec{L} = \vec{N}$

Eg. 1.11

$$f \vec{N} = 0 \Rightarrow \frac{d}{dt} \vec{L} = 0$$

$$d\vec{L} = 0 \cdot dt = 0$$

$$\int d\vec{L} = 0$$

$$\vec{L} + C = 0$$

$$\vec{L} = constant$$

Conservation theorem for the angular momentum of a particle: if the total torque is zero, angular momentum is conserved /

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Consider, again, Newton's Second Law F= ma and let's define the work done by an external force F on a particle going from point 1 to point 2 as W= SF. ds A BALLAGO BALLAGO

For constant mass, $W = m \int_{1}^{2} \frac{d\vec{v}}{dt} \cdot d\vec{s}$

since $\vec{v} = \frac{d\vec{s}'}{dt} \Rightarrow d\vec{s}' = \vec{v}dt$,

product rule

W= m $\int_{1}^{2} \frac{d\vec{v}}{dt} \cdot \vec{v} dt$ $\int_{0}^{1} \frac{d\vec{v}}{dt} \cdot \vec{v} dt = \int_{0}^{1} \frac{d\vec{v}}{dt} = \int_{0}^{1} \frac{d\vec{v}}{dt} + \int_{0}^{1} \frac{d\vec{v}}{dt} = 2\vec{v} d\vec{v}$

 $\Rightarrow \frac{1}{2} \frac{d}{dt} (v^2) = \frac{d\vec{v}}{dt} \cdot \vec{v}$ So $W = \frac{m}{2} \int_{-\infty}^{\infty} \frac{d}{dt} (v^2) dt$

 $W = \frac{m}{2} \int_{1}^{2} dv^{2} = \frac{m}{2} v^{2} \Big|_{1}^{2} = \frac{m}{2} \left(v_{z}^{2} - v_{1}^{2} \right) = \frac{1.13}{2}$

The quantity my is called the Kinetic energy, we will use T. W= Tz-T, The work done is the change in kinetic energy

The value of a line integral SF(r). dr = S (Fx dx + Fy dy + Fz dz) #

Over a path (from point to point 2) in general

'energy int depends not only on (1) and (2), but also on the path (along which we integrate. What we are the conditions for path independence? Define a line integral as independent of path in a domain D in space if for every pair of points D and D in D, integral * has the same value for all paths in D that start at 10 and end at 20. A "well-known" theorem of vector analysis says that a line integral is independent of the path in D if and only if $\vec{F} = \nabla f$, so $\vec{F}_{X} = \frac{\partial f}{\partial x}$; $\vec{F}_{Y} = \frac{\partial f}{\partial y}$; $\vec{F}_{Z} = \frac{\partial f}{\partial z}$ I will prove the if, but not fre only if.

Let C be any path in D from point O to point O given by $\vec{r}(t) = \vec{x}(t)\hat{i} + \vec{y}(t)\hat{j} + \vec{z}(t); a \le t \le b$ $d\vec{r}(t) = d\vec{x}(t)\hat{i} + d\vec{y}(t)\hat{j} + d\vec{z}(t)\hat{k}$

$$\int_{C} \vec{F}(\vec{r}) \cdot d\vec{r} = \int_{C} F_{x} dx + F_{y} dy + F_{z} dz \qquad \text{we make time-}$$

$$= \int_{a}^{b} \left(F_{x} \frac{dx}{dt} + F_{y} \frac{dy}{dt} + F_{z} \frac{dz}{dt} \right) dt \qquad \text{explicit}$$

Let
$$\vec{F}(\vec{r}) = \nabla f$$
, then

$$\int_{C} \vec{F}(\vec{r}) \cdot d\vec{r} = \int_{\partial x}^{2} \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz$$

$$= \int_{a}^{b} \left(\frac{\partial f}{\partial x} \frac{dx}{dt} + \frac{\partial f}{\partial y} \frac{dy}{dt} + \frac{\partial f}{\partial z} \frac{dz}{dt} \right) dt$$

$$= \int_{a}^{b} \left(\frac{\partial f}{\partial x} + \frac{\partial f}{\partial t} + \frac{\partial f}{\partial z} \right) dt = \int_{a}^{b} \frac{df}{dt} dt \quad v_{1} - v_{2}$$

$$= \int_{a}^{b} \left(\frac{\partial f}{\partial t} + \frac{\partial f}{\partial t} + \frac{\partial f}{\partial z} \right) dt = \int_{a}^{b} \frac{df}{dt} dt \quad v_{1} - v_{2}$$

$$\Rightarrow \int df = f[x(t), y(t), z(t)] = f[x(t), z(t)]$$

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$$\Rightarrow \int_{a}^{b} df = f[x(t), y(t), z(t)] = \begin{cases} t=b \\ t=a \end{cases}$$

$$= \begin{cases} f[x(t), y(t), z(t)] \\ t=a \end{cases}$$

For convenience, let $f = -V(\vec{r})$, then $\vec{F} = -\nabla V(\vec{r})$.

Which the work done by an external force on a particle is independent of the path taken is called a conservative force field

Notice that this would not happen in general and friction is a particularly important case. But when it holds, then Tz-T, = V,-V2

$$\Rightarrow T_1 + V_1 = T_2 + V_2$$

=> T,+V, = Tz+Vz Conservation theorem for the energy of a particle: if the forces acting on a particle are conservative, then the total energy of the particle, T+V, is conserved.