Eigenvalue equation et oscillations 11/18/21 (148) Eq. 100 leaves us in a familiar place: Tin = -V.M.

It is thooke's law But in n dimensions Luckily, we derived in pgs. 142-143 F=ma=-KAX what the solution looks like : $\chi = e^{i\sqrt{E/m}t+C}$ n. = Ca; e timaginary, not index with w= JK/m Hence, $n_j = \frac{d}{dt} n_j = \frac{d}{dt} \left[Ca_j e^{i\omega t} \right] = Ca_j \frac{d}{dt} e^{i\omega t}$ n; = Caj eint. in = in des n; and $\vec{n}_{j} = \frac{d}{dt}\vec{n}_{j} = \frac{d}{dt}\left(Ca_{j} iwe^{iwt}\right) = iwCa_{j}\frac{d}{dt}e^{iwt}$ nj = Cajeiwt, izw = -wzCajeiwt=-wznj sartifle Ewilling so Tij (-w² (aj eint) + Vij (Caj eint) = 0 Ceiwt [to Vijaj - w Tijaj] = 0 Eq. 6.12 Vij aj - w2 Tij aj = 0 System of n homogeneous equation s Good old matrix multiplication Actually Vis and Tis are rank-2 tensors

Eq. 6.12 leaves us in a familiar place

 $\tilde{V}\vec{a} - \omega^2 \tilde{T}\vec{a} = 0$

 $\left(\tilde{V} - \omega^2 \tilde{T}\right) \vec{a} = 0$

either a = 0, or more interestingly 1 v-w27 = 0

Luckily , we derived in pgs. 89-90 what the solution looks like.

 $\begin{vmatrix} V_{11} - \omega^{2} T_{11} & V_{12} - \omega^{2} T_{12} & \cdots \\ V_{21} - \omega^{2} T_{21} & V_{22} - \omega^{2} T_{22} & \cdots \\ V_{31} - \omega^{2} T_{3i} & \cdots & \cdots \end{vmatrix} = 0$

The roots of the characteristic equation, which is an algebraic equation of degree n for wz, are solutions to the equations of motion. They are the eigenfrequencies We can rewrite Eq. 6.12 as $\sqrt{\vec{a}} = \omega^2 \vec{T} \vec{a}$ Va = λTa Eq. 6.14

Each eigenfrequency WK is related to the eigenvalue ZK and to the eigenvector ax. The matrix of eigenvectors is A. More generally, $\vec{v}\vec{a_k} = \lambda_k \vec{T}\vec{a_k}$

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In the most general case, ω_k , λ_k , $\vec{\lambda}_k$ are complex Take the transpose on both sides of Eq. 6.14

 $(\vec{\nabla} \vec{a_k})^T = (\lambda_k \vec{T} \vec{a_k})^T$ Properties of transpose $\vec{a_k}^T \vec{\nabla}^T = \lambda_k (\vec{T} \vec{a_k})^T = \lambda_k \vec{a_k}^T \vec{\nabla}^T (cA)^T = cA^T$

Now take the conjugate, this is the adjoint or conjugate transpose

 $\vec{a}_{k}^{T} \vec{v}^{T} = \lambda_{k} \vec{a}_{k}^{T} \vec{T}^{T}$

we know that the elements of \tilde{V} are force constants which must be real, and \tilde{T} are masses which must also be real. Further more, we showed before that \tilde{V} and \tilde{T} are symmetric. As in the case of the moment of inertia tensor (pg. 115), \tilde{V} and \tilde{V} are Hermitian. Hence,

 $\vec{a}_{k}^{T} \vec{v} = \overline{\lambda_{k}} \vec{a}_{k}^{T} \vec{T} \Rightarrow \vec{a}_{k}^{*} \vec{v} = \overline{\lambda_{k}} \vec{a}_{k}^{*} \vec{T} = Eq. 6.16$

* denotes the adjoint, also called the conjugate transpose
The bar denotes the complex conjugate.

K and L are columns of \widetilde{A} , in principle they could be the same

Eq. 6.16 Multiply from the right times ax, so

$$\vec{a}_{\ell}^* \vec{v} \vec{a}_{k} = \vec{\lambda}_{\ell} \vec{a}_{\ell}^* \vec{T} \vec{a}_{k}$$

Multiply Eq. 6.15 from the left times at, so

$$\vec{a}_{\ell}^* \hat{V} \vec{a}_{k} = \vec{a}_{\ell}^* \hat{\lambda}_{k} \hat{T} \vec{a}_{k}$$

Algebra says that if two quantities are equal to a third one, the first two are equal. Hence,

$$\lambda_{k} \vec{a}_{k}^{*} \vec{T} \vec{a}_{k} = \bar{\lambda}_{e} \vec{a}_{e}^{*} \vec{T} \vec{a}_{k}$$

$$\Rightarrow (\lambda_k - \overline{\lambda}_\ell) \vec{a}_\ell^* \vec{T} \vec{a}_k = 0$$

When
$$l=K$$
, $\left[\left(\lambda_{k}-\overline{\lambda_{k}}\right)\overrightarrow{a}_{k}^{*}\overrightarrow{\tau}\overrightarrow{a}_{k}^{*}\right]=0$

take the transpose $(\lambda_k - \overline{\lambda}_k)[\vec{a}_k^* \vec{T} \vec{a}_k]^T = (\lambda_k - \overline{\lambda}_k)(\vec{T} \vec{a}_k^*)^T \vec{a}_k^T$

$$= \frac{(\lambda_k - \overline{\lambda}_k) \vec{a}_k^T \vec{T} \vec{a}_k^{*T}}{(\lambda_k - \overline{\lambda}_k) \vec{a}_k^T \vec{T} \vec{a}_k^{*T}}$$

then the conjugate $(\overline{\lambda}_{k} - \overline{\lambda}_{k}) \vec{a}_{k}^{T} \vec{\gamma}^{T} \vec{a}_{k}^{T} = (\overline{\lambda}_{k} - \lambda_{k}) \vec{a}_{k}^{*} \vec{\gamma}^{*} \vec{a}_{k}^{T}$

Both quantities are equal to zero, so they are (52) equal; \tilde{T} is Hermitian, so $\tilde{T}^* = T$; Hence $(\lambda_k - \lambda_k) \vec{a}_k^* \vec{T} \vec{a}_k = (\lambda_k - \lambda_k) \vec{a}_k^* \vec{T} \vec{a}_k$

we can see that the matrix product $\vec{a_k} + \vec{r} \vec{a_k}$ is Hermitian, so it is real. Also

ak = aktibk; akt = xkt + ipkt; akt

 $\frac{1}{a_{k}} = \alpha_{k}^{T} + i\beta_{k}^{T} = \alpha_{k}^{T} - i\beta_{k}^{T} = \alpha_{k}^{T} - i\beta_{k}^{T} = \bar{\alpha}_{k}^{*}$

so & the matrix product is

 $\left(\alpha_{k}^{T}-i\beta_{k}^{T}\right)\left(\tilde{T}\right)\left(\alpha_{k}+i\beta_{k}\right)=\left(\alpha_{k}^{T}\tilde{T}-i\beta_{k}^{T}\tilde{T}\right)\left(\alpha_{k}+i\beta_{k}\right)$

= XxTXx - iBxTxxx + i xxTTBx - i 2 BxTTBx

= XKTXK + BKT TBK + I (XKTBK BKTXK) Eq. 6.19

So yeah, the matrix product

Since T=T

Symmetric,

[(Tak)TBK]T

[Tak)TBK

[Tak)TBK

BrTTXK

Since $T = \frac{1}{2} \vec{n} T \vec{n}$, if $\vec{n}_{k} = \alpha_{k}$ $\vec{n}_{k} = \beta_{k}$ $T = \frac{1}{2} \alpha_{k} T \alpha_{k} T = \frac{1}{2} \beta_{k} T \beta_{k}$ So the matrix product $\vec{a}_{k} T \vec{a}_{k} = 2T + 2T = 4T$ The kinetic energy can't be negative, so the matrix product is real and positive greater than a equal to zero.

If the amplitudes are not zero, the kinetic energy is not zero, and it can't be negative. Hence the matrix product is real and positive.