

UNIVERSITY OF MANNHEIM

MASTER THESIS

Quasi-Monte Carlo Methods and Applications

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Research Group Name
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Abstract

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x

RITA Rational Inverse Transform with Aliasing

CDF Cumulative Distribution Function

QMC Quasi-Monte Carlo

MC Monte Carlo

HU Hounsfield Unit

CT Computed Tomography

CBCT Cone-Beam Computed Tomography

FFD Forced Fixed Detection

Part I

Fundamentals of Quasi-Monte Carlo Methods

Chapter 1

Monte Carlo and Quasi-Monte Carlo Integration

1.1 Motivation and Problem Setting

The numerical evaluation of high-dimensional integrals is a fundamental task in modern applied mathematics and scientific computing. Applications range from Bayesian inference and financial mathematics to machine learning and computational physics. In particular, contemporary domains such as deep neural network training and medical imaging often require the estimation of integrals of the form

$$I = \int_{[0,1]^s} f(x) dx, \quad (1.1)$$

where s denotes the dimensionality of the problem and f is a function that may be expensive or impractical to evaluate analytically.

One of the most widely used approaches to compute such integrals is the classical Monte Carlo (**MC**) method, which estimates the expectation based on averages over randomly sampled points. The primary appeal of **MC** integration lies in its dimension-independent convergence rate and minimal assumptions on the integrand. However, its asymptotic error rate of $\mathcal{O}(N^{-1/2})$ limits its efficiency – especially when high precision is required or function evaluations are computationally expensive.

Quasi-Monte Carlo (**QMC**) methods offer an alternative paradigm: instead of random samples, they employ deterministic sequences – so-called low-discrepancy sequences—that aim to fill the integration domain more uniformly. This structured sampling allows for faster convergence under certain smoothness conditions and forms the basis for many state-of-the-art techniques in high-dimensional numerical integration.

The goal of this part is to develop a rigorous mathematical foundation for **QMC** methods and to understand how they improve upon classical Monte Carlo integration. The key concepts introduced here – particularly discrepancy theory and low-discrepancy sequences – serve as theoretical tools that will be revisited in later parts of this thesis. Chapter 2 will delve into the measurement of uniformity via star discrepancy and provide concrete constructions of low-discrepancy sequences such as Sobol and Halton, which are central to the numerical methods applied in Part II

and Part III, dedicated to neural network training and CT-based photon transport simulation.

1.1.1 High-Dimensional Integration in Applications

High-dimensional integration problems arise naturally in a wide range of scientific and engineering disciplines. Whenever expectations with respect to multivariate distributions must be computed numerically, they are typically formulated as integrals over the s -dimensional unit cube – such as the integral in Equation (1.1).

Prominent examples include:

- **Bayesian statistics:** Computing posterior expectations, marginal likelihoods, or predictive distributions.
- **Financial mathematics:** Pricing complex financial derivatives and evaluating risk measures under stochastic models.
- **Machine learning:** Estimating expectations in variational inference, training neural networks using randomized optimization techniques, or evaluating generalization bounds.
- **Medical imaging:** Simulating photon transport and estimating physical quantities based on noisy measurements, especially in Computed Tomography (CT) and magnetic resonance imaging (MRI).

In all these cases, the dimensionality s can be moderate to very high – sometimes even exceeding hundreds or thousands of variables. This introduces significant challenges for numerical integration methods, which must balance accuracy, computational cost, and robustness with respect to the structure of the integrand.

The remainder of this chapter explores how MC and QMC methods address these challenges. Before doing so, we briefly review the Monte Carlo method and its fundamental convergence properties in the next section.

1.1.2 Monte Carlo Integration: Principle and Convergence

The classical MC method is a probabilistic approach to numerical integration that relies on random sampling. It is particularly suited for high-dimensional settings, as its convergence rate does not deteriorate with increasing dimension.

Definition 1.1.1 (Monte Carlo Estimator).

Let $f \in L^2([0, 1]^s)$ and let $X_0, \dots, X_{N-1} \sim \mathcal{U}([0, 1]^s)$ be independent and identically distributed random samples. The Monte Carlo estimator of the integral I is given by

$$I_N^{\text{MC}} := \frac{1}{N} \sum_{n=0}^{N-1} f(X_n). \quad (1.2)$$

This estimator is unbiased and converges almost surely to the true integral as $N \rightarrow \infty$, as established by the Strong Law of Large Numbers:

Theorem 1.1.2 (Strong Law of Large Numbers).

Let $f \in L^2([0, 1]^s)$. Then

$$\mathbb{P} \left[\lim_{N \rightarrow \infty} Q_N(f) = \int_{[0,1]^s} f(x) dx \right] = 1. \quad (1.3)$$

Beyond this almost sure convergence, the expected integration error of the Monte Carlo estimator can be quantified using the root-mean-square error:

Theorem 1.1.3 (Monte Carlo Convergence Rate, [12, Satz 1.2]).

Let $f \in L^2([0, 1]^s)$. Then

$$\mathbb{E} [|Q_N(f) - I|] \leq \frac{\sigma[f]}{\sqrt{N}}, \quad (1.4)$$

where $\sigma[f] := \sqrt{\text{Var}[f]}$.

Sketch of Proof. Using independence and linearity of expectation, the variance of $Q_N(f)$ is

$$\text{Var}[Q_N(f)] = \frac{\text{Var}[f]}{N}. \quad (1.5)$$

Applying Jensen's inequality yields the stated bound on the expected absolute error. \square

Remark 1.1.1. The convergence rate $\mathcal{O}(N^{-1/2})$ is independent of the integration dimension s , which is a key advantage of the Monte Carlo method. In contrast, classical grid-based methods often suffer from the curse of dimensionality, where the number of required samples grows exponentially with s .

Example 1.1.1. Let $f \in C^1([0, 1]^s)$ be a Lipschitz-continuous function. A uniform grid with $N = m^s$ points has a worst-case error of $\mathcal{O}(N^{-1/s})$. In high dimensions, this becomes prohibitively inefficient, while MC integration retains the dimension-agnostic rate $\mathcal{O}(N^{-1/2})$. [5, Section 1.1]

Despite its robustness and simplicity, MC integration has several well-known limitations:

- The convergence rate is relatively slow, especially for smooth integrands.
- Error bounds are probabilistic rather than deterministic.
- The method does not exploit structural properties of the integrand, such as smoothness or sparsity.

These limitations motivate the development of quasi-Monte Carlo methods, which will be introduced in the next section.

1.2 Quasi-Monte Carlo Methods

1.2.1 Deterministic Sampling and Intuition

Classical MC methods approximate integrals over the unit cube $[0, 1]^s$ using randomly sampled points. In contrast, QMC methods replace stochasticity with a deterministic strategy, aiming to cover the integration domain in a more uniform and structured way.

Definition 1.2.1 (Quasi-Monte Carlo Estimator).

Let $f: [0, 1]^s \rightarrow \mathbb{R}$ be a measurable function and let $\{x_n\}_{n=1}^N \subset [0, 1]^s$ be a deterministic point set. Then the *quasi-Monte Carlo estimator* for the integral

$$I(f) = \int_{[0,1]^s} f(\mathbf{x}) d\mathbf{x}$$

is defined as

$$I_N^{\text{QMC}}(f) := \frac{1}{N} \sum_{n=1}^N f(\mathbf{x}_n). \quad (1.6)$$

Unlike in the Monte Carlo setting, these points are not drawn from a probability distribution, but are generated deterministically — typically by rules that aim to avoid clustering and oversampling.

Figure 1.1 illustrates this contrast by comparing 300 randomly sampled points to 300 QMC points from the Sobol' sequence in two dimensions. The quasi-Monte Carlo points distribute more evenly, avoiding both gaps and clusters.

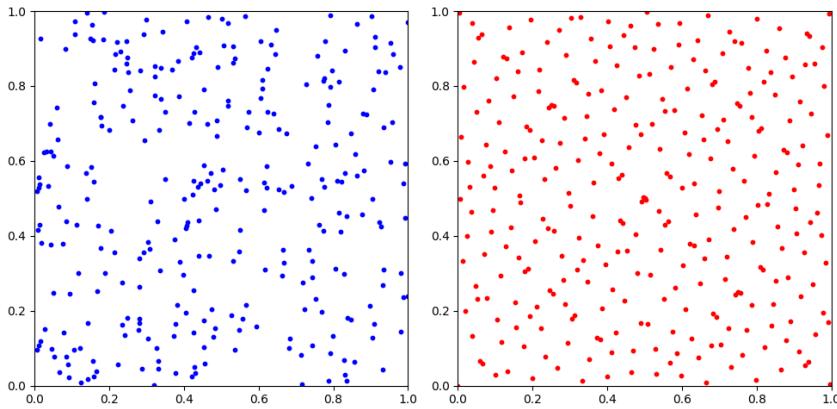


FIGURE 1.1: Comparison of 300 sample points in $[0, 1]^2$ using (left) standard Monte Carlo and (right) a Sobol' sequence. QMC points avoid clustering and fill the domain more evenly.

A natural alternative to MC sampling is the use of regular tensor-product grids. However, these grids suffer from the *curse of dimensionality*: placing m points per dimension leads to a total of $N = m^s$ points, which becomes infeasible even for moderate s . Furthermore, grids are not easily extensible – adding new points requires global recompilation and destroys nesting.

By contrast, many **QMC** sequences, such as Sobol' and Halton (introduced in Chapter 2), are designed to be *incremental*: each new point can be added without changing the existing ones. This makes **QMC** particularly suitable for adaptive and anytime algorithms. However, not all **QMC** constructions share this feature such as lattice rules and some optimized designs are fixed-size by nature.

QMC methods thus offer the best of both worlds: they combine the space-filling structure of grids with the flexibility and scalability of sampling methods. Low-discrepancy sequences fill the space more evenly than random points while avoiding the redundancy of regular grids. This often leads to significantly lower integration errors, especially for smooth functions or problems with low effective dimension.

This shift from probabilistic to deterministic sampling also changes the way we analyze error: instead of using statistical bounds, we rely on *discrepancy measures*, which quantify the uniformity of the point set. These concepts, along with the notion of function variation, will be introduced in the following sections.

Remark 1.2.1. **QMC** estimators retain the same algebraic structure as their **MC** counterparts, but their convergence behavior is governed by entirely different theoretical tools – namely discrepancy theory and function variation.

1.2.2 Monte Carlo vs. Quasi-Monte Carlo: A First Comparison

MC and **QMC** methods share the same high-level goal: estimating an integral by averaging function evaluations at selected sample points. The difference lies in the sampling strategy – stochastic versus deterministic – and the consequences this has for accuracy, convergence and theoretical guarantees.

Table 1.1 summarizes key conceptual distinctions between the two paradigms.

Aspect	Monte Carlo (MC)	Quasi-Monte Carlo (QMC)
Sampling	Independent random points	Deterministic low-discrepancy points
Error bounds	Probabilistic (in expectation)	Deterministic (worst-case)
Regularity assumptions on f	Square integrability (L^2)	Bounded variation (HK)
Theoretical convergence	$\mathcal{O}(N^{-1/2})$	Up to $\mathcal{O}(N^{-1})$ (heuristic)
Point extensibility	Trivial	Often supported (e.g., Sobol)
Robustness to noise	High	Low

TABLE 1.1: Conceptual comparison of Monte Carlo and quasi-Monte Carlo integration methods.

While **MC** estimators are unbiased and robust even under minimal assumptions, their convergence is slow and does not improve when the integrand is smooth. **QMC** methods, on the other hand, exploit structure in the integrand – such as smoothness or low effective dimension – and can yield significantly lower integration errors.

However, **QMC** methods lack probabilistic guarantees and depend more strongly on the careful construction of the point set. Their performance can degrade if the integrand exhibits high variability along many input dimensions or if the chosen sequence is not well matched to the function.

The next subsection discusses the convergence rates of both methods in more detail and highlights the interplay between dimensionality and integration error.

1.2.3 Convergence Rates and the Curse of Dimensionality

A major appeal of the Monte Carlo method is its dimension-agnostic convergence behavior from Theorem 1.1.3. Under the assumption that $f \in L^2([0, 1]^s)$, the root-mean-square error of the MC estimator satisfies

$$\mathbb{E}[|I_N^{\text{MC}} - I|] = \mathcal{O}(N^{-1/2}), \quad (1.7)$$

regardless of the dimension s . This makes MC a go-to method for very high-dimensional problems, even when no structure in the integrand is known or exploitable.

QMC methods, while lacking this universality, offer a fundamentally different convergence profile. For sufficiently regular integrands (specifically, those of bounded variation in the sense of Hardy and Krause), the integration error of QMC estimators is bounded by

$$|I_N^{\text{QMC}} - I| \leq D_N^* \cdot V_{\text{HK}}(f), \quad (1.8)$$

where D_N^* is the star discrepancy of the point set and $V_{\text{HK}}(f)$ is the Hardy–Krause variation of the integrand (both formalized later in TODO). For well-constructed low-discrepancy sequences, it is known that

$$D_N^* = \mathcal{O}\left(\frac{(\log N)^s}{N}\right). \quad (1.9)$$

This bound implies a near-linear convergence rate in N – much faster than MC – but introduces a logarithmic dependence on the dimension s . As s increases, the factor $(\log N)^s$ grows rapidly, leading to the curse of dimensionality in worst-case QMC analysis.

Fortunately, in many practical problems, the *effective dimension* of the integrand is much lower than its nominal input dimension. That is, most of the function’s variation is concentrated along a small number of directions. When this is the case, QMC methods frequently outperform MC in practice, even in problems with moderately high s .

Example 1.2.1. In deep learning applications, such as training neural networks, stochastic gradients often lie near low-dimensional manifolds. Similarly, in CT-based simulation models, the system’s physical response depends more strongly on a few dominant input parameters. These scenarios are highly favorable for QMC.

We will revisit the precise meaning of discrepancy and variation – and how they control the QMC error – in Chapters 2 and 3.

1.3 Uniformity Concepts

The efficiency of QMC methods hinges on how well the employed point sets "fill" the integration domain $[0, 1]^s$. This motivates the need for a precise mathematical understanding of *uniformity*. The present section introduces two key concepts in this regard – *uniform distribution modulo one* and *equidistribution* – and clarifies their relevance to numerical integration.

1.3.1 Uniform Distribution vs. Equidistribution

While the terms *uniform distribution* and *equidistribution* are sometimes used interchangeably in informal contexts, they carry distinct meanings in the context of numerical integration. In quasi-Monte Carlo methods, the notion of *uniform distribution modulo one* provides a rigorous criterion for how evenly a sequence covers the unit cube. This concept forms the foundation for analyzing the convergence behavior of QMC estimators.

Definition 1.3.1 (Uniform Distribution Modulo One).

Let $(\mathbf{x}_n)_{n \in \mathbb{N}_0} \subset [0, 1]^s$ be a sequence of sample points. The sequence is said to be *uniformly distributed modulo one* (u.d. mod 1) if, for every axis-aligned box $[\mathbf{a}, \mathbf{b}] \subset [0, 1]^s$, the proportion of points falling into this box converges to its Lebesgue measure:

$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \chi_{[\mathbf{a}, \mathbf{b}]}(\mathbf{x}_n) = \lambda_s([\mathbf{a}, \mathbf{b}]),$$

where λ_s denotes the s -dimensional Lebesgue measure and χ_A is the indicator function of the set A .

This definition emphasizes asymptotic spatial coverage: in the limit, every subregion of the domain is sampled proportionally to its volume. Importantly, this property is purely deterministic and does not rely on any probabilistic assumptions — in contrast to Monte Carlo methods, which only guarantee uniformity in expectation.

In contrast, when speaking of a random sample $\{\mathbf{X}_1, \dots, \mathbf{X}_N\} \subset [0, 1]^s$ drawn i.i.d. from the uniform distribution, we refer to a probabilistic concept of uniformity: each point is independently drawn according to the uniform measure, but the empirical distribution may not be uniformly spread in finite samples. Thus, equidistribution refers to the ideal uniform coverage that we seek deterministically, while uniform random sampling only ensures this behavior in expectation or with high probability.

Remark 1.3.1. Uniform distribution modulo one is a necessary condition for the convergence of quasi-Monte Carlo estimators. If a point sequence is not u.d. mod 1, then there exists at least one Riemann-integrable function for which the sample average fails to converge to the integral.

The following result, known as Weyl's criterion, provides a necessary and sufficient condition for uniform distribution in terms of exponential sums.

Theorem 1.3.2 (Weyl's Criterion).

A sequence $(\mathbf{x}_n)_{n \geq 0} \subset [0, 1]^s$ is uniformly distributed modulo one if and only if, for all nonzero $\mathbf{h} \in \mathbb{Z}^s \setminus \{\mathbf{0}\}$, we have

$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \exp(2\pi i \mathbf{h} \cdot \mathbf{x}_n) = 0.$$

Example 1.3.1. Let $\alpha \in \mathbb{R}$ be irrational. Then the Kronecker sequence $(n\alpha \bmod 1)_{n \geq 0}$ is uniformly distributed in $[0, 1]$. This is a consequence of Weyl's criterion and illustrates the existence of simple, deterministic sequences with excellent uniformity

properties.

1.3.2 Why Uniformity Matters in Numerical Integration

In QMC integration, the objective is to approximate the integral defined in Equation (1.1) by a finite average over a deterministic point set $\{x_0, \dots, x_{N-1}\} \subset [0, 1]^s$. The accuracy of this approximation depends crucially on how uniformly the point set samples the domain.

Unlike MC methods, which rely on probabilistic guarantees and variance-based error estimates, QMC methods exploit deterministic structure: the integration error is directly influenced by how well the point set fills the domain. This insight gives rise to one of the most fundamental theoretical tools in QMC analysis: the *Koksma–Hlawka inequality*.

Remark 1.3.2. The Koksma–Hlawka inequality bounds the integration error of a QMC estimator by the product of two quantities: the *star discrepancy* of the point set and the *variation* of the integrand in the sense of Hardy and Krause. In short,

$$\left| \frac{1}{N} \sum_{n=0}^{N-1} f(x_n) - I \right| \leq D_N^*(\{x_n\}) \cdot V_{HK}(f).$$

A rigorous statement and detailed discussion of this inequality will follow in Chapter 3.

This inequality highlights why uniformity is essential: if the integrand has bounded variation, then smaller discrepancy directly implies smaller integration error. In this sense, discrepancy theory becomes a cornerstone of effective QMC integration.

Remark 1.3.3. The star discrepancy quantifies the maximal deviation between the empirical distribution of the point set and the uniform distribution over $[0, 1]^s$. It vanishes asymptotically for uniformly distributed sequences and governs the convergence behavior of QMC estimators.

The next chapter provides a formal introduction to discrepancy theory. We will define extreme and star discrepancy, study their geometric interpretation, and analyze the behavior of structured low-discrepancy sequences such as Halton and Sobol'. These sequences, due to their uniform space-filling properties, play a central role in modern high-dimensional numerical integration.

Chapter 2

Discrepancy Theory and Low-Discrepancy Sequences

2.1 Measuring Uniformity: Discrepancy

One of the cornerstones of QMC theory is the concept of *discrepancy*, which quantifies how uniformly a finite point set samples the s -dimensional unit cube $[0, 1]^s$. While uniform distribution modulo one describes the asymptotic behavior of infinite sequences, discrepancy provides a finite-sample measure of deviation from perfect uniformity. Low-discrepancy sequences are the foundation of QMC methods because they minimize this deviation and thus yield more accurate numerical integration.

This section introduces formal definitions, geometric intuition and empirical insights into discrepancy measures. It lays the theoretical groundwork for understanding how the structure of point sets affects QMC integration performance.

2.1.1 Definition of Discrepancy and Star Discrepancy

First, we define the more general *extreme discrepancy*:

Definition 2.1.1 (Extreme Discrepancy).

Let $\mathcal{P} \subset [0, 1]^s$, with $|\mathcal{P}| = N$, being a finite point set. Then the extreme discrepancy $D_N(\mathcal{P})$ is defined as

$$D_N(\mathcal{P}) = \sup_{\substack{\mathbf{a}, \mathbf{b} \in [0, 1]^s \\ \mathbf{a} \leq \mathbf{b}}} \left| \frac{A([\mathbf{a}, \mathbf{b}], \mathcal{P}, N)}{N} - \lambda_s([\mathbf{a}, \mathbf{b}]) \right|.$$

Hereby $A([\mathbf{a}, \mathbf{b}], \mathcal{P}, N)$ denotes the number of points in \mathcal{P} that fall into the box $[\mathbf{a}, \mathbf{b}]$, and $\lambda_s([\mathbf{a}, \mathbf{b}])$ is the Lebesgue measure of that box, given by $\prod_{i=1}^s (b_i - a_i)$.

In practice, a common and slightly more tractable variant is the *star discrepancy*, which restricts the test boxes to be anchored at the origin.

Definition 2.1.2 (Star Discrepancy).

Let $\mathcal{P} \subset [0, 1]^s$, with $|\mathcal{P}| = N$, being a finite point set. Then the star discrepancy D_N^* is defined as

$$D_N^*(\mathcal{P}) = \sup_{t \in [0, 1]^s} \left| \frac{A([\mathbf{0}, t], \mathcal{P}, N)}{N} - \lambda_s([\mathbf{0}, t]) \right|.$$

This form is used in the Koksma–Hlawka inequality and serves as a central measure in QMC analysis. Note that both discrepancy definitions are deterministic and depend solely on the point configuration.

Remark 2.1.1. A low star discrepancy implies that the empirical distribution of the points approximates the uniform distribution well over all axis-aligned subrectangles anchored at the origin.

2.1.2 Geometric Interpretation and Examples

To build geometric intuition, consider the 1-dimensional case. Given N sample points $x_0, \dots, x_{N-1} \in [0, 1]$, we can visualize discrepancy as the maximum vertical deviation between the empirical distribution function

$$F_N(t) := \frac{1}{N} \sum_{n=0}^{N-1} \chi_{[0, t)}(x_n)$$

and the uniform cumulative distribution function $F(t) = t$. The discrepancy corresponds to the largest gap between $F_N(t)$ and $F(t)$.

In higher dimensions, the idea generalizes: the star discrepancy measures the maximal difference in the number of points falling into an axis-aligned box $[\mathbf{0}, t)$ versus the volume of that box. Intuitively, it quantifies whether the point set "overpopulates" or "underpopulates" certain regions.

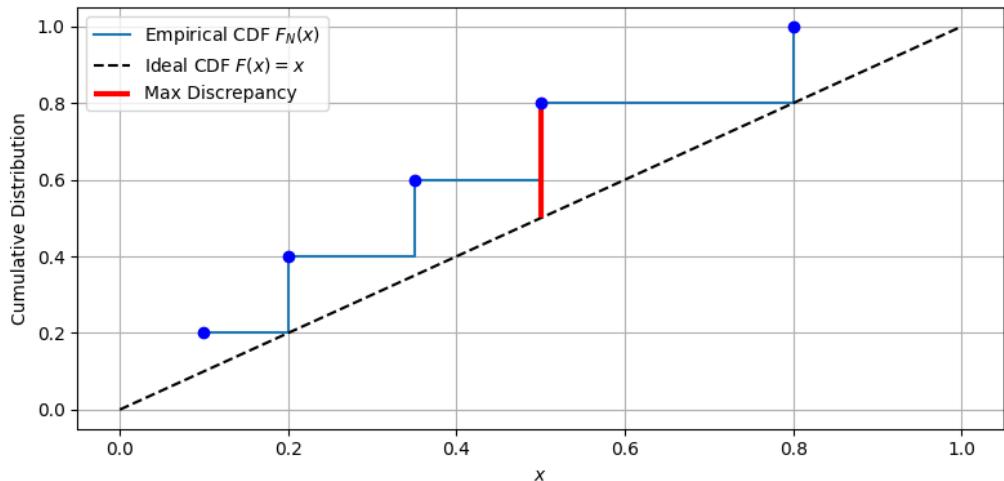


FIGURE 2.1: Visualization of 1D discrepancy: the maximum vertical distance between the empirical CDF $F_N(t)$ and the uniform CDF $F(t) = t$.

Example 2.1.1. Let $x_n = \frac{n}{N}$ for $n = 0, \dots, N - 1$. Then $F_N(t)$ is a piecewise constant staircase function, and the discrepancy can be shown to be $\mathcal{O}(1/N)$. This is an optimal rate in 1D.

Remark 2.1.2. Discrepancy provides a worst-case error metric over all subintervals. Thus, even a point set that appears visually well-distributed may have large discrepancy due to subtle gaps or clustering in certain regions.

2.1.3 Empirical Observation of Discrepancy in QMC

Discrepancy is not just a theoretical construct – its practical relevance becomes evident when comparing the behavior of **QMC** sequences with **MC** samples. In particular, structured point sets like Halton and Sobol sequences achieve much lower discrepancy than random samples of the same size.

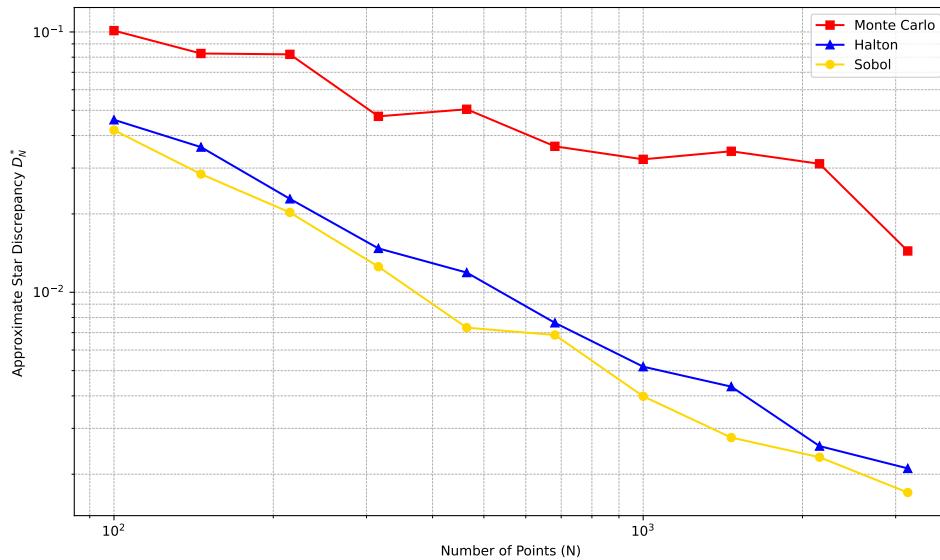


FIGURE 2.2: Comparison of discrepancy growth for Sobol, Halton and random (**MC**) point sets in 2D. Low-discrepancy sequences exhibit significantly slower discrepancy growth.

Remark 2.1.3. The expected star discrepancy of i.i.d. Monte Carlo samples decreases at the rate $\mathcal{O}(N^{-1/2})$, whereas for low-discrepancy sequences, provable upper bounds of order $\mathcal{O}((\log N)^s / N)$ exist. [5, Section 2.2]

The next section will introduce specific constructions of low-discrepancy sequences and analyze their dimensional performance and implementation details.

2.2 Construction of Low-Discrepancy Sequences

In this section, we will explore several well-known constructions of low-discrepancy sequences, focusing on their mathematical foundations and properties.

2.2.1 The Halton Sequence

The Halton sequence is one of the earliest and most widely used constructions of low-discrepancy sequences in arbitrary dimensions. It generalizes the one-dimensional van der Corput sequence to multiple dimensions by using mutually prime bases.

Definition 2.2.1 (Halton Sequence). Let b_1, \dots, b_s be pairwise coprime integers greater than 1, typically chosen as the first s prime numbers. The n -th point $\mathbf{x}_n \in [0, 1]^s$ of the s -dimensional Halton sequence is defined componentwise by

$$\mathbf{x}_n = (\phi_{b_1}(n), \dots, \phi_{b_s}(n)),$$

where $\phi_b(n)$ denotes the van der Corput radical-inverse function in base b . The Halton sequence is then given by $\mathcal{S}_{b_1, \dots, b_s} = (\mathbf{x}_n)_{n \in \mathbb{N}}$.

Example 2.2.1 (First Elements of the Halton Sequence in 2D).

Consider the two-dimensional Halton sequence based on the first two prime bases $b_1 = 2$ and $b_2 = 3$. The first three elements are obtained by computing the radical-inverse values $\phi_{b_1}(n)$ and $\phi_{b_2}(n)$ for $n = 1, 2, 3$:

- $n = 1: \mathbf{x}_1 = (\phi_2(1), \phi_3(1)) = (\frac{1}{2}, \frac{1}{3})$
- $n = 2: \mathbf{x}_2 = (\phi_2(2), \phi_3(2)) = (\frac{1}{4}, \frac{2}{3})$
- $n = 3: \mathbf{x}_3 = (\phi_2(3), \phi_3(3)) = (\frac{3}{4}, \frac{1}{9})$

This illustrates how the Halton sequence fills the unit square in a low-discrepancy manner even for small n .

The radical-inverse function $\phi_b(n)$ maps an integer n to a real number in $[0, 1)$ by reflecting its base- b representation about the decimal point:

$$\phi_b(n) = \sum_{k=0}^{\infty} d_k b^{-k-1}, \quad \text{where } n = \sum_{k=0}^{\infty} d_k b^k.$$

Intuitively, this construction ensures that each component of the sequence explores the unit interval in a structured, non-redundant way. By combining several one-dimensional van der Corput sequences in different coprime bases, the resulting multi-dimensional point set avoids regular grid-like patterns and achieves asymptotic uniformity.

The Halton sequence has star discrepancy of order

$$D_N^*(\{\mathbf{x}_0, \dots, \mathbf{x}_{N-1}\}) = \mathcal{O}\left(\frac{(\log N)^s}{N}\right),$$

making it a prototypical example of a low-discrepancy sequence. However, for larger dimensions, correlation effects between the different base components can lead to degraded uniformity and higher discrepancy. Variants such as scrambling or leaping are commonly employed to mitigate this issue.

Remark 2.2.1. The Halton sequence is extensible in N and s , making it suitable for applications that require growing or adaptive point sets. However, its performance

in high dimensions is often inferior to more modern constructions like the Sobol sequence.

2.2.2 The Sobol Sequence

Direction Numbers and Gray Code

Scrambling Techniques for Sobol

2.2.3 Dimensional Performance and Sequence Comparison

Chapter 3

Function Variation and Error Bounds in QMC

3.1 Function Variation in Multiple Dimensions

- 3.1.1 Variation in the Sense of Vitali
- 3.1.2 Definition of the Hardy–Krause Variation
- 3.1.3 Examples and Interpretation of Hardy–Krause Variation

3.2 The Koksma–Hlawka Inequality

- 3.2.1 Formal Theorem and Preconditions
- 3.2.2 Role of Discrepancy and Function Variation
- 3.2.3 Interpretation: Bounding the Integration Error

3.3 Convergence Implications

- 3.3.1 From Error Bound to Convergence Rate
- 3.3.2 Comparison to Monte Carlo: $O(N^{-1/2})$ vs. $O\left(\frac{(\log N)^s}{N}\right)$
- 3.3.3 Limitations in High Dimensions and Practical Workarounds

Part II

Quasi-Monte Carlo Sampling for Neural Network Training

Chapter 4

Motivation and Problem Formulation

- Many-query problems and the need for surrogate models
- Limitations of randomly sampled training data in DNNs
- Formal definition of the supervised learning setup
- Motivation for using low-discrepancy sampling in training

Chapter 5

Theoretical Foundations of QMC-Based Learning

- Adapting QMC theory to supervised learning
- Hardy-Krause variation of the loss function
- Generalization error bounds using low-discrepancy sampling
-
- Comparison to MC-based bounds and practical implication
-
- Validity conditions and smoothness assumptions

Chapter 6

QMC-Based Deep Learning Algorithm

- Algorithm description (based on Mishra & Rusch Algirthm 2.2)
- Network Architecture, activation functions and regularization
- QMC vs. MC sampling: sample complexity and efficiency
- Discussion of smooth vs. non-smooth activations (e.g. ReLU vs. tanh)
- Practical considerations: dimensionality, optimal choices

Chapter 7

Empirical Evaluation and Discussion

- Overview of benchmark tasks:
 - Synthetic function approximation
 - ODE simulation (projectile motion)
 - Computational finance (option pricing)
 - CFD (airfoil flow)
- Comparison of DLSob (QMC) and DLrand (MC) Performance
- Error metrics and convergence plots
- Limitations, edge cases and future directions

Part III

X-ray Simulation using QMC Methods

Chapter 8

Motivation and Problem Statement

8.1 Computed Tomography Imaging

CT imaging is a powerful medical imaging technique that provides detailed cross-sectional images of the body by combining multiple X-ray projections taken from different angles. This technique is widely used for diagnostic purposes, allowing clinicians to visualize internal structures with high spatial resolution.

The process involves rotating an X-ray source around the patient while simultaneously capturing X-ray projections on a detector array. Each projection represents the cumulative attenuation of X-rays as they pass through various tissues, influenced by the density and composition of the materials they encounter.

The collected data is then reconstructed into a three-dimensional volume using advanced algorithms, such as filtered backprojection or iterative reconstruction methods. These algorithms convert the raw projection data into cross-sectional images, which can be further processed to enhance contrast, reduce noise, and improve overall image quality.

To achieve accurate reconstructions, it is crucial to acquire high-quality X-ray projection images. However, one of the most significant challenges in CT imaging is the presence of scattered radiation - such as Compton and Rayleigh scattering - which can lead to artifacts and distortions in the reconstructed images.

Throughout this thesis, the term *scatter reduction in CT* refers to the mitigation of scatter-induced artifacts in the two-dimensional X-ray projection images acquired by the detector array during a CT scan. These projections serve as the input for reconstructing the final three-dimensional volume. Given the complexity of the full reconstruction process, the focus of this work is limited to analyzing and correcting scatter effects at the level of the two-dimensional projection data.

8.2 X-ray Imaging and the Challenge of Scatter

X-ray **CT** relies on measuring the attenuation of X-rays along straight-line paths through a phantom. Here, the X-ray beam is further abstracted as individual photons with defined initial energies and directions. The photon trajectories typically extend from an X-ray source \mathcal{S} to a detector element \mathcal{D} , forming the path:

$$l = \overrightarrow{\mathcal{S}\mathcal{D}}$$

Under idealized conditions, the attenuation process is governed by the *Lambert-Beer law*, which describes an exponential decay in intensity I as the beam interacts with the material. As stated in [8, Chap. 7], this relationship is given by:

$$I = I_0(E) \cdot \int_0^{E_{\max}} \exp \left(- \int_S^D \mu(x, E) dx \right) dE \quad (8.1)$$

Here, $I_0(E)$ denotes the incident intensity of X-rays with initial energy E at the source S , and $\mu(x, E)$ represents the linear attenuation coefficient at spatial location x along the path l , for energy E . The resulting intensity I is measured at the detector element D .

Although exponential attenuation along straight-line paths such as \overrightarrow{SD} is the idealized model for X-ray imaging, this assumption is systematically violated in practice. As X-rays traverse the scanned phantom, many photons undergo scattering interactions - such as Compton or Rayleigh scattering - that alter their trajectories. Despite deviating from the primary path, these scattered photons may still reach the detector, adding unintended signal components. As a result, the measured intensities no longer represent pure line integrals of the attenuation map. This discrepancy introduces nonlinear errors and visible artifacts in the reconstructed image, ultimately degrading both visual quality and quantitative accuracy.

Scattered radiation is a major source of image artifacts - such as cupping and streaks - and reduces both spatial and contrast resolution. These effects not only degrade visual image quality but also compromise the accuracy of quantitative measurements, such as Hounsfield units μ_* , which are used for clinical interpretation, such as tissue characterization [8, Chap. 8]. Hounsfield units represent a normalized attenuation relative to the attenuation of water:

$$\mu_* = \left(\frac{\mu}{\mu_{\text{water}}} - 1 \right) \cdot 1000 \quad (8.2)$$

The impact of scatter becomes especially pronounced in modern CT systems using high-energy X-rays or large-area flat-panel detectors, where scatter may dominate the measured signal. As such, accurate modeling and correction of scatter are essential for achieving high-fidelity CT images, particularly in clinical applications where precision and reliability are paramount [8].

8.3 Scatter Correction Methods for X-ray Imaging

8.3.1 Overview of Scatter Correction Techniques

In X-ray imaging, scattered photons are a major source of image artifacts and quantitative inaccuracies. To mitigate these effects, a range of computational scatter correction methods has been developed. These approaches can be broadly classified into the following categories:

- **Empirical and Analytical Methods:**

These include techniques such as primary modulation, convolution-based correction, and energy windowing. Scatter is typically estimated using simplified

models or empirical kernels, often assuming a smooth background distribution, and then subtracted from the measured signal. While computationally efficient, these methods rely on assumptions regarding the spatial and energy distribution of scattered photons. As a result, they may fail to accurately model complex scatter phenomena in heterogeneous anatomical structures.

- **Physics-Based Models:**

These methods aim to provide a more accurate representation of the underlying photon transport physics, including scattering phenomena. Among them, **MC** simulation is regarded as the most rigorous and comprehensive technique due to its ability to statistically model complex photon interactions without relying on simplifying assumptions. In such simulations, primary and scattered photon contributions are detected separately, allowing the estimated scatter signal to be subtracted from measured data in order to restore image fidelity.

In 2011, Sisniega et al. [16] already demonstrated promising results using computationally expensive **MC** simulations of CT scans of a cylindrical phantom with two bone inserts. The results showed significant improvements in image quality, as illustrated by Sisniega et al. in Figure 8.1.

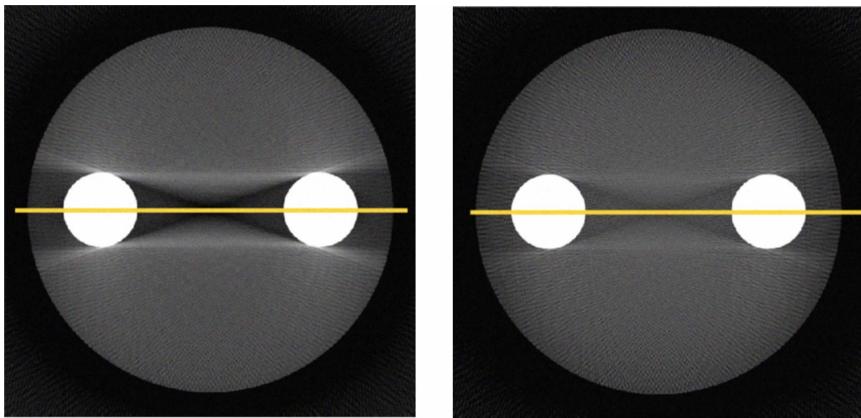


FIGURE 8.1: Cone-Beam Computed Tomography (**CBCT**) of a 70 mm soft-tissue cylinder with two bone inserts. Right: Simulation of a X-ray image with scattering, exhibiting typical cupping and streak artifacts. Right: Image after scatter correction applying **MC** techniques for the simulation of 5×10^6 photons, showing improved uniformity and quantitative accuracy. Right: Image without correction, exhibiting typical cupping and streak artifacts. (Figure adapted from [16], ©2011 IEEE)

8.3.2 Monte Carlo Simulation

To achieve scatter correction using **MC** simulation, a three-dimensional phantom is first estimated by analyzing the X-ray projections. Then, an **MC** simulation is performed to estimate the scatter signal $I_{\text{scattered}}$ for each detector pixel. This estimated scatter signal is then subtracted from the measured intensity I_{measured} :

$$I_{\text{corrected}} = I_{\text{measured}} - I_{\text{scattered}} \quad (8.3)$$

The MC simulation is a powerful computational technique that follows physics-based principles to model the transport of photons through matter. It is widely recognized as the gold standard for scatter correction in CT and related imaging modalities. This status is attributed to several key factors:

- **Physical Accuracy:**

MC methods simulate the stochastic nature of photon interactions - including Compton and Rayleigh scattering, photoelectric absorption, and multiple scattering events - based on fundamental physical cross-sections and material properties.

- **Comprehensive Modeling:**

Unlike analytical or empirical methods, MC simulations can account for complex geometries, heterogeneous materials and realistic X-ray spectra, providing highly accurate estimates of the scatter signal.

- **Validation Benchmark:**

Due to their accuracy, MC-based scatter estimates are routinely used as reference standards for validating and benchmarking faster, approximate correction methods.

8.4 High-Level Overview of the Monte Carlo Simulation

The MC simulation of photon transport for scatter correction typically involves the following key steps [16]:

1. **Photon Emission:**

Photons are emitted from a virtual X-ray source, with their energies sampled from a predefined source spectrum.

2. **Photon Tracking:**

Each photon is tracked as it propagates through the object. At every step, the probability of interaction (either scattering or absorption) is determined by the local material properties and corresponding interaction cross-sections.

3. **Interaction Sampling:**

Upon interaction, the event type (e.g., Compton scattering, Rayleigh scattering, or photoelectric absorption) and the resulting change in photon direction and energy are sampled from the relevant probability distributions.

4. **Detection:**

Photons that reach the detector—either unscattered or after one or more scattering events—are recorded. The simulation distinguishes between primary and scattered photons, thereby enabling an accurate estimation of the scatter contribution at each detector element.

5. **Statistical Averaging:**

By simulating a large number of photons, the method builds a statistically robust estimate of the scatter distribution. The accuracy of the result increases with the number of simulated photons and ultimately converges toward a stable intensity distribution.

The schematic flow of the Monte Carlo simulation process is summarized in Figure 8.2.

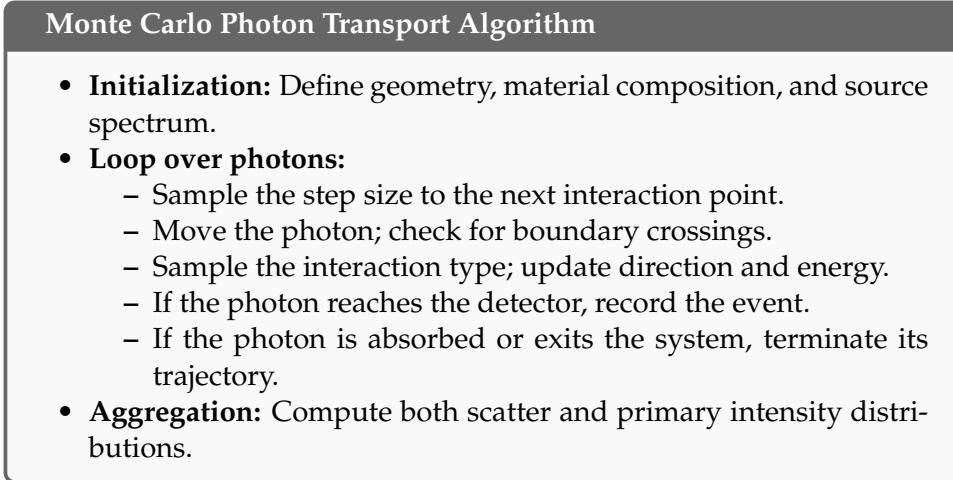


FIGURE 8.2: Pseudocode representation of the Monte Carlo algorithm for photon transport in scatter modeling.

8.5 Monte Carlo Methods: Benefits & Drawbacks

MC simulations are widely regarded as the gold standard for scatter correction in computed tomography due to their ability to model photon-matter interactions from first principles. These simulations accurately reproduce all relevant physical scattering phenomena - including Compton and Rayleigh scattering - and can accommodate arbitrarily complex geometries and heterogeneous material compositions. Owing to this high degree of physical fidelity, MC-based methods yield highly accurate scatter estimates and are commonly employed as reference standards for evaluating and validating alternative correction approaches.

However, the high accuracy of Monte Carlo simulations comes at the cost of substantial computational effort. Producing low-noise scatter estimates requires the simulation of a large number of photon histories to ensure statistical convergence. Consequently, the runtime scales with the desired level of accuracy and may range from several minutes to multiple hours or even days, depending on the complexity of the scene and the available computational resources. This computational burden poses a major limitation, particularly in applications involving large datasets or iterative reconstruction workflows.

To address the computational inefficiency of slow-converging MC methods, recent research has explored strategies to accelerate physically accurate photon transport simulations. Among these, the application of QMC methods has gained increasing attention. By replacing random sampling with deterministic low-discrepancy sequences, QMC techniques can achieve significantly faster convergence while maintaining comparable accuracy. Recent QMC-based scatter correction algorithms have demonstrated runtime reductions of multiple orders of magnitude, thereby enabling high-fidelity simulations even in time-sensitive or resource-constrained environments [6, 7, 1].

Chapter 9

Physical laws of Photon Simulation

This chapter introduces the physical and mathematical principles underlying the simulation of X-ray imaging. Central to this is the modeling of individual photon interactions with matter, including attenuation and scattering processes. These interactions are inherently stochastic due to the quantum nature of radiation-matter interactions. The stochastic nature of photons and their interactions with matter is modelled using monte Carlo methods of uniformly distributed values between 0 and 1, which are then transformed to follow the physical probability distributions relevant to the specific interactions being simulated. This approach allows for a realistic representation of photon transport and interaction within the imaging system.

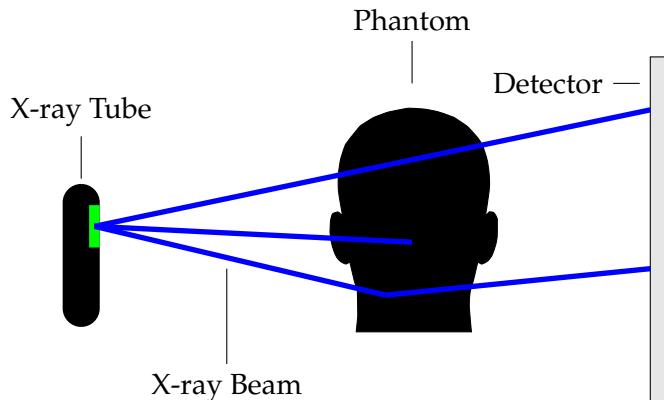


FIGURE 9.1: Schematic illustration of photon transport in X-ray imaging.

In a typical simulation workflow, individual photons are emitted from the X-ray tube and propagate through air before reaching the phantom. Upon entering the phantom, they may interact with the material through scattering or absorption, depending on the local composition of the tissue and photon energy. Only a fraction of the photons will exit the phantom, continue their path through air, and ultimately reach the detector, contributing to the final image formation.

The simulation framework described here forms the basis for all subsequent analyses presented in this thesis.

9.1 X-Ray Tube

X-ray beams are generated within evacuated glass tubes containing several critical components that convert electrical energy into X-ray photons and heat. The X-ray tube acts as an energy converter, where the vast majority of the input energy is transformed into heat and only a small fraction becomes useful radiation. The following Figure 9.2 illustrates the basic construction of an X-ray tube as in [14]

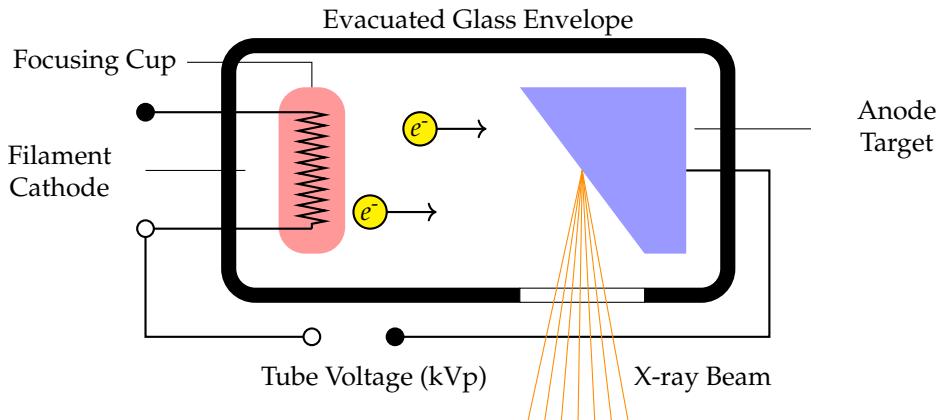


FIGURE 9.2: Schematic of an X-ray tube.

9.1.1 Photon Generation

To capture this randomness, Monte Carlo methods are employed. In such simulations, random numbers—typically sampled uniformly from the interval $[0, 1]$ are transformed into physically meaningful quantities according to the relevant probability distributions. This probabilistic modeling enables realistic simulation of photon transport and forms the basis for the analyses presented in this thesis.

The key components are:

- **Filament (Cathode):** A tungsten wire filament serves as the electron source through thermionic emission. When a current of approximately $3\text{-}6 \text{ A}$ passes through it, the filament reaches incandescence, releasing electrons from its surface. These free electrons form a cloud near the cathode until they are accelerated toward the anode by the applied high voltage.
- **Focusing Cup:** The filament is embedded in a negatively charged, nickel-made focusing cup. Its function is to electrostatically shape and direct the electron stream toward the anode's focal spot, thereby influencing the resolution and size of the resulting X-ray beam.
- **Anode Target:** The anode consists of a tungsten target, often embedded in a copper support. Tungsten is used for its high atomic number and melting point, enhancing X-ray production and durability. The copper base improves heat dissipation. Typically, less than 1% of the electron energy is converted into X-rays, with the remainder generating heat that must be effectively managed.
- **Evacuated Glass Envelope:** All components are sealed within a borosilicate glass or metal-ceramic housing evacuated to a low pressure (typically 10^{-5} to 10^{-7} hPa). The vacuum allows unimpeded electron flow and prevents arcing.

The housing is usually immersed in insulating oil to provide thermal and electrical isolation.

When high-energy electrons strike the tungsten target, X-ray photons are produced through two primary mechanisms:

- **Bremsstrahlung (Braking Radiation):** This accounts for approximately 80% of X-ray production. When electrons pass close to tungsten nuclei, they are decelerated by the electrostatic attraction, causing them to lose kinetic energy that is emitted as X-ray photons. This process produces a continuous spectrum of X-ray energies from near zero up to the maximum electron energy.
- **Characteristic Radiation:** This occurs when high-energy electrons knock inner shell electrons from tungsten atoms. When outer shell electrons drop down to fill these vacancies, they emit X-rays with discrete, characteristic energies specific to tungsten. Therefore peaks are occurring at the difference of the binding energies of the electron shells. For reference, the atomic model of tungsten is given in Figure 9.7 and the approximate binding energies of the electron shells in Table 9.1.

Shell / Subshell	Binding Energy (keV)
K	69.5
L ₁	12.1
L ₂	11.5
L ₃	10.2
M ₁	2.82
M ₂	2.30
M ₃	2.15
N ₁	0.43
N ₂	0.32
N ₃	0.22

TABLE 9.1: Approximate Electron Binding Energies of Tungsten (W, Z = 74)

In Table 9.1 only binding energies of the K, L and M shells are given, since these are the most relevant for X-ray production. The binding energies of the N shell and higher shells are negligible in this context due to their low binding energies.

Although it is not necessary to understand every technical detail of the X-ray apparatus for the purposes of simulation, I found it important to briefly present the fundamental working principles of the X-ray tube. This background allows one to appreciate how key simulation parameters for photon generation - specifically the tube voltage and the cathode material - influence the resulting X-ray spectrum and photon behavior.

Figure 9.4 below shows an resulting energy spectrum for an X-ray tube with a tungsten cathode operated at 100 kVp. It illustrates the resulting distribution of photon energies, which is shaped by both the material and the applied voltage. The intensity of the different photon energies is measured in *spectral fluence* in $\text{cm}^{-2} \text{ keV}^{-1}$,

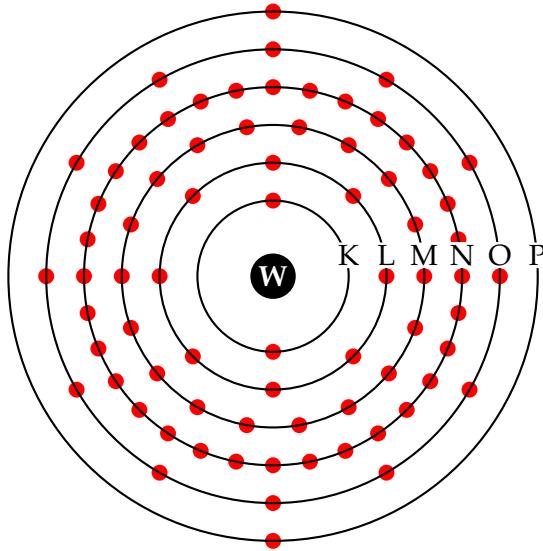


FIGURE 9.3: Bohr model of the tungsten atom (W , $Z = 74$) with electron shells and subshells.

which describes the number of X-ray photons per unit area per unit energy interval. It certainly shows the two components of the X-ray spectrum: the continuous bremsstrahlung spectrum and the characteristic radiation peaks:

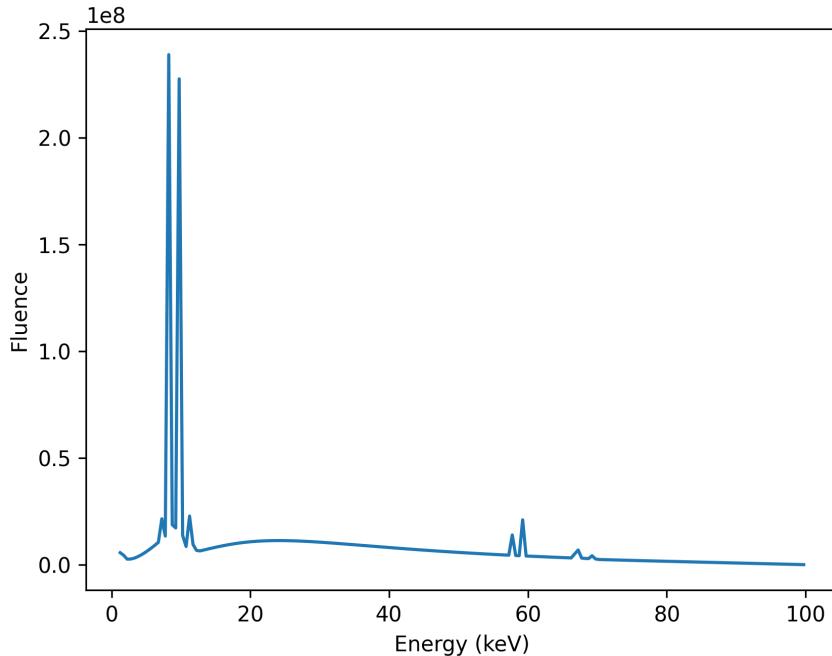


FIGURE 9.4: X-ray spectrum for a tungsten cathode at 100 kVp showing the continuous bremsstrahlung spectrum and characteristic peaks. Build with *SpekPy* [13]

9.1.2 Filter

Low energy X-ray photons contribute little to image formation but significantly increase patient dose. Therefore, X-ray tubes are often equipped with filters to selectively attenuate these low-energy photons while allowing higher-energy photons to pass through. This process, known as beam hardening, improves image quality by reducing scatter and enhancing contrast. Common filter materials include aluminum, copper and molybdenum, which are chosen based on their atomic number and thickness to effectively absorb low-energy photons while minimizing the impact on higher-energy photons [14].

As can be seen in Figure 9.5, the filter is placed at the X-ray tube margin such that the photons hit the filter after the anode target and before leaving the X-ray tube. Followingly X-rays pass through the filter before reaching the patient.

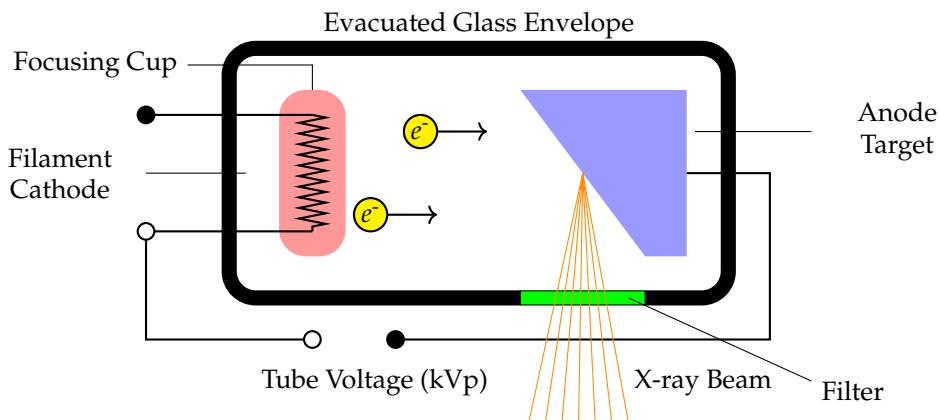


FIGURE 9.5: Schematic of an X-ray tube with Filter.

Throughout this thesis, a filter consisting of 0.4 mm Tin (Sn) and 0.1 mm Copper (Cu) is used for the simulations. This filter is chosen to represent a realistic X-ray tube filter that is commonly used in clinical practice [17]. The resulting spectrum is shown in Figure 9.6.

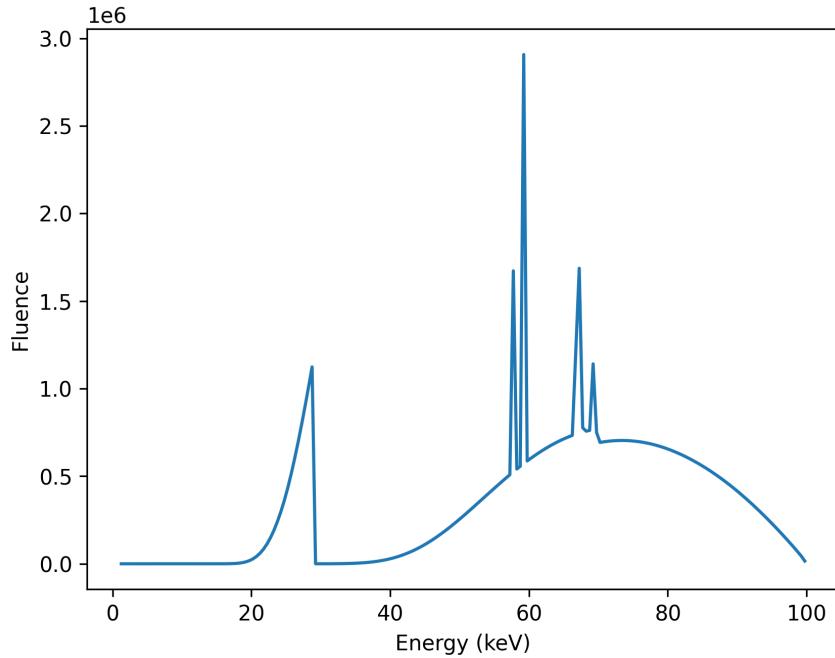


FIGURE 9.6: X-ray spectrum for a tungsten cathode at 100 kVp with a 0.4 mm showing the continuous bremsstrahlung spectrum and characteristic peaks. Build with *SpekPy* [13]

9.2 Photon Generation

In the context of Monte Carlo simulations, photons are generated from the X-ray tube source, which is typically modeled as a point source emitting photons within a conical beam. The emission cone is defined by a half-angle Θ , which determines the angular distribution of emitted photons. – The emitted photons are characterized by two key properties: their direction and energy.

9.2.1 Photon Energy

The photon energy is sampled from the X-ray tube spectrum as described in Section 9.1.2. In the simulations referenced in this thesis the spectra are generated utilizing *SpekPy* [13, 15]. The energies of the resulting photons are sampled via inverse transform sampling from the normalized spectrum.

Inverse transform sampling [9] is a method to generate random samples from a target distribution using uniformly distributed random variables, such as those generated by a Quasi-Monte Carlo sequence.

Definition 9.2.1 (Inverse Transform Sampling).

Let $F_X : \mathbb{R} \rightarrow [0, 1]$ be the cumulative distribution function (CDF) of a continuous, strictly increasing random variable X . The method of *inverse transform sampling* generates a realization of X by the following procedure:

1. Generate a sample U from the uniform distribution on the unit interval, i.e., $U \sim \mathcal{U}(0, 1)$.
2. Compute the value $X := F_X^{-1}(U)$, where F_X^{-1} denotes the inverse of the CDF F_X .

Theorem 9.2.2.

Let F_X be a continuous and strictly increasing cumulative distribution function and let $U \sim \mathcal{U}(0, 1)$. Then the random variable

$$X := F_X^{-1}(U)$$

has cumulative distribution function F_X , i.e.,

$$\mathbb{P}(X \leq x) = F_X(x), \quad \text{for all } x \in \mathbb{R}.$$

Proof.

Since F_X is continuous and strictly increasing, its inverse F_X^{-1} exists. For any $x \in \mathbb{R}$, we compute:

$$\mathbb{P}(X \leq x) = \mathbb{P}(F_X^{-1}(U) \leq x) = \mathbb{P}(U \leq F_X(x)) \quad \begin{matrix} \text{since } F_X \text{ is strictly} \\ \text{decreasing} \end{matrix} \quad F_X(x),$$

because $U \sim \mathcal{U}(0, 1)$ and thus $\mathbb{P}(U \leq u) = u$ for all $u \in [0, 1]$. This shows that X has CDF F_X . \square

9.2.2 Photon Direction

The direction of each emitted photon is sampled uniformly within a conical emission cone defined by the half-angle Θ . For the simulation a spherical alignment of the X-ray tube is assumed, such that the beam is oriented along the vector $\vec{v} = (v_1, v_2, v_3)$ in the Cartesian coordinate system.

Followingly, the direction of the emitted photon is sampled based on two random values $u_1, u_2 \in [0, 1)$ as follows:

1. Sample a random angle θ uniformly from the interval $[0, \Theta]$ with u_1 :

$$\theta = u_1 \cdot \Theta$$

2. Sample a random azimuthal angle ϕ uniformly from the interval $[0, 2\pi)$ with u_2 :

$$\phi = u_2 \cdot 2\pi$$

3. Compute the direction vector \vec{d} of the photon as:

$$\vec{d} = (\sin(\theta) \cos(\phi), \sin(\theta) \sin(\phi), \cos(\theta))$$

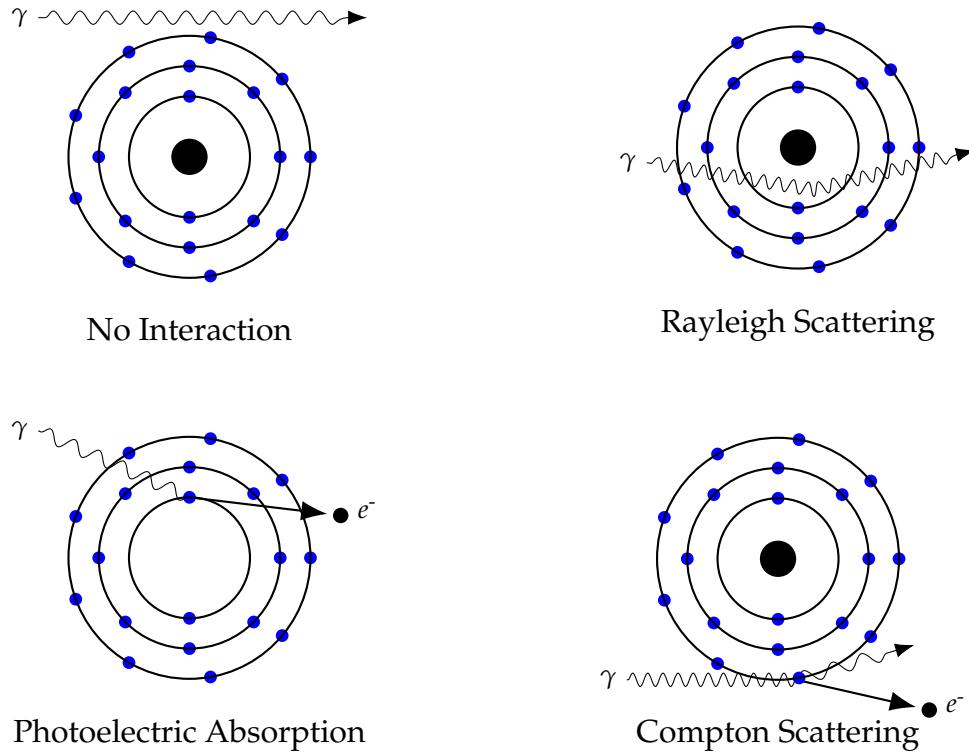


FIGURE 9.7: Principles from photon interaction with matter similar to [8, Chap. 7]

9.3 Scattering and Attenuation

This section provides the basic concepts of photon interaction with matter. The interaction relevant for medical imaging results in a reduction of radiation intensity, which corresponds to a decreased number of photons reaching the detector. Hereby X-ray photons may be fully absorbed by *photoelectric absorption* or undergo either *elastic scattering* (Rayleigh) or *inelastic scattering* (Compton) as they interact with matter.

The attenuation of X-ray photons arises from physical processes that alter their number, direction, or energy as they interact with matter. These interactions occur at the level of individual photons and are highly dependent on the photon energy. This section presents an overview of the primary interaction mechanisms relevant to attenuation like in [8].

For correctness, in the simulations it is further assumed that the X-ray photons are propagating through vacuum before entering and after exiting the phantom. Usually the tissues are surrounded by air, which has a negligible effect on the photon transport.

9.3.1 Free Path Length

All mentioned interaction processes - the photoelectric effect, Compton scattering and Rayleigh scattering - are probabilistic in nature. The according attenuation coefficients μ characterizes the extend of the beam being reduced by its according effect, when passing through the material, in $\text{cm}^2 \text{g}^{-1}$. For the simulation of photons, the

attenuation coefficients are dependend on the according energy E of the photon and the material at spherical location x of the photon. The coefficients are summarized in Table 9.2.

Interaction Type	Attenuation Coefficient
Photoelectric Effect	$\mu_{\text{PE}}(x, E)$
Rayleigh Scattering	$\mu_{\text{RS}}(x, E)$
Compton Scattering	$\mu_{\text{CS}}(x, E)$

TABLE 9.2: Attenuation coefficients for different photon interaction mechanisms.

The linear attenuation coefficient $\mu(x, E)$ is the sum of the individual attenuation coefficients of the interaction coefficients:

$$\mu(x, E) = \mu_{\text{PE}}(x, E) + \mu_{\text{RS}}(x, E) + \mu_{\text{CS}}(x, E)$$

When an X-ray photon enters the phantom two main effects are considered:

- **Attenuation:** The photon may be absorbed or Scattered.
- **No Interaction:** The photon may pass through the phantom without any interaction.

The *free path length* t of a photon describes the distance a photon takes to traverse through the phantom before it interacts with matter. Supposing the phantom's location is x and the photon is traveling in the unit direction \vec{v} , as in [6] the free path length t follows the distribution given by:

$$t \sim \mu(x, E) \exp \left[- \int_0^t \mu(x + s \cdot \vec{v}) dt \right]$$

Check: nach dem paper müsste μ abhängig vom Endpunkt sein.

The free path length t is constrained by the maximum distance c ($0 \leq t \leq c$) to the next exit point of the phantom along the ray with direction \vec{v} . After leaving the phantom, we assume the photon is reaching the detector or leaving the simulation domain.

The probability of the photon to leave the phantom unaltered is described by the *escape probability* \mathcal{P} of a photon at a given position A with unit direction \vec{v} and energy E as in [6]:

$$\mathcal{P}(x, \vec{v}, E) = \exp \left[- \int_0^c \mu(x + t\vec{v}, E) dt \right],$$

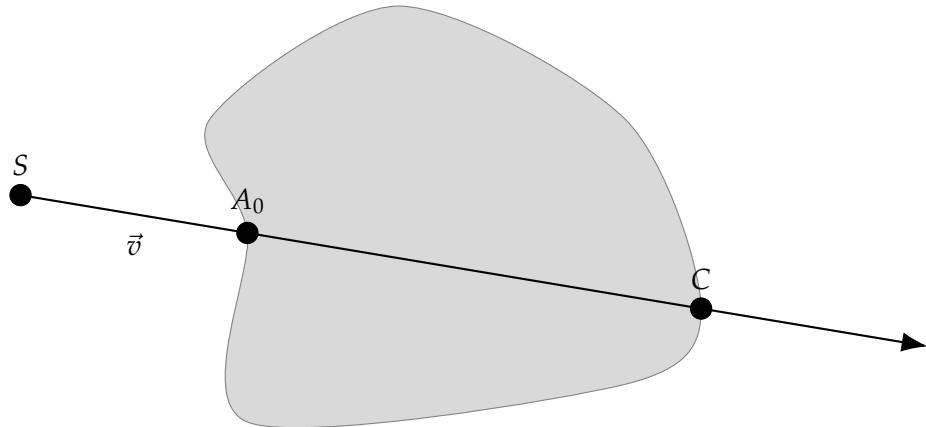


FIGURE 9.8: Illustration of the entry and exit point of the ray of a photon without interaction.

9.3.2 Compton Scattering

Compton scattering is the most dominant interaction mechanism for X-ray photons in tissue [8, Chap. 7]. It occurs when a X-ray photon with considerably higher energy than the binding energy of an outer shell electron collides with this electron. This interaction results in a transfer of energy and momentum. The electron is ejected from the atom, while the incoming photon is scattered at an angle and its energy is reduced.

A portion of the incident photon energy is transferred to the electron, which is referred to as the "recoil electron" or Compton electron. The interaction produces a positive ion, the "recoil electron" and a scattered photon. If the deflection angle of the scattered photon is small, most of the energy is retained by the scattered photon. The deflection angle can vary from 0 to 180 degrees, depending on the energy transfer during the interaction.

For the distribution of the scattering angle θ of the scattered photons in the differential cross section, we apply the formula derived by Klein and Nishina in 1928. This equation describes the distribution of the scattering angle of X-ray photons in Compton scattering assuming the electron is initially free and stationary, the relation between energy loss and photon deflection is determined from conservation of momentum and energy between the photon and the recoiling electron [2, 4].

$$\frac{d\sigma}{d\Omega} = \frac{1}{2} r_e^2 \frac{1}{(1 + k(1 - \cos \theta))^2} \left(1 + \cos^2 \theta + \frac{k^2 (1 - \cos \theta)^2}{1 + k(1 - \cos \theta)} \right) \quad (9.1)$$

where $k = \frac{e}{mc^2} \approx \frac{E}{0.511}$ with mc^2 being the rest mass energy of the electron, E being the energy of the incident photon and r_e being the classical electron radius.

For sampling from the distribution of the scattering angle we cannot apply the same procedure as for the photon energy in Section 9.2.1, since the Klein-Nishina formula is not a cumulative distribution function (CDF). Instead the method of *rejection sampling* is used.

Rejection sampling is a technique to generate samples from a target distribution by using an envelope distribution that is easy to sample from. The envelope distribution is an upper bound to the original distribution. The basic idea is to sample from this proposal distribution and accept or reject the samples based on a criterion that relates the envelope distribution to the target distribution from Klein-Nishina [9, Chap. 4].

In the context of the distribution of Klein-Nishina, the convergence behavior for $E \rightarrow \infty$ and $k \rightarrow \infty$ consequently is:

$$f(k, x) \in \mathcal{O}\left(\frac{1}{k(1-x)}\right) \quad (9.2)$$

A suitable envelope distribution is of the form

$$h(k, x) = \frac{1}{k(1-x)}$$

with the given conditions

$$\begin{aligned} f(k, 1) &= h(k, 1) = 2 \\ f(k, -1) &= h(k, -1) = \zeta(k) \\ \zeta(k) &= \frac{2(2k^2 + 2k + 1)}{(2k + 1)^3} \end{aligned}$$

By a straightforward calculation, this leads to the following values[11]:

$$\begin{aligned} a(k) &= 2(b(k) - 1) \\ b(k) &= \frac{1 + \frac{\zeta(k)}{2}}{1 - \frac{\zeta(k)}{2}} \end{aligned}$$

Now by using the rejection sampling method, $x = \cos \theta$ is sampled from the distribution h by generating a uniform random variable $r \sim \mathcal{U}(0, 1)$ and calculating the value x as follows:

$$x = b(k) - (b(k) - 1) \left(\frac{\zeta}{2}\right)^r \quad (9.3)$$

We will accept this value if $\frac{f(k,x)}{h(k,x)} \geq u_1$, where u_1 denotes the random input variable. To derive an exact direction of the scattered photon, we need to further sample an azimuthal angle ϕ , which is done by the second random input variable:

$$\phi = 2\pi u_2 \quad (9.4)$$

Once the scattering angle θ is sampled, the energy of the scattered photon E' can be calculated using the Compton formula:

$$E' = \frac{E}{1 + \frac{E}{m_e c^2} (1 - \cos \theta)} \quad (9.5)$$

9.3.3 Rayleigh Scattering

Rayleigh scattering is a type of elastic scattering that occurs at low X-ray energies [8, Chap. 7]. The incident photon interacts with several electrons that are usually bound in the outer shells of the atom. In this process, the low energy photon is not ejected, but rather the electrons and in turn the whole atom is set to vibration with respect to the incident photon's wavelength. The vibrating photon transfers its excess energy to an electromagnetic photon with the same wavelength but probably a different direction than the incident photon. The majority of these scattered photons are emitted in a forward direction. This interaction does not result in the ejection of electrons from the atom and no ionization occurs as no energy is converted into kinetic energy.

Although Rayleigh scattering is taking place in the X-ray tube, it can be excluded from the linear attenuation coefficient, as it does not contribute to the attenuation of the X-ray beam in the same way as Compton scattering or photoelectric absorption. As described in [14], the exclusion of Rayleigh scattering from the linear attenuation coefficient is based on the assumption that the characteristic radiation emitted by the target is isotropic, meaning it is emitted uniformly in all directions. As the X-ray is not attenuated by Rayleigh scattering, it is a common assumption to exclude Rayleigh scattering from the attenuation coefficient.

Chapter 10

The Simulation Algorithm

This chapter presents the algorithm used for X-ray image simulation. It incorporates elements of *Forced Detection* as introduced in [3], in order to accelerate convergence.

Unlike standard Monte Carlo simulations, where a random variable determines whether a photon escapes the phantom or undergoes another scattering event, the forced detection approach employed, calculates the probability of escaping the phantom at each interaction point and reflects this distribution in accordingly updated photon intensities for the case of escaping the phantom and the case of scattering. Instead of probabilistically terminating the photon trajectory, the algorithm tracks it through a predefined number of scattering events N . At each interaction, the remaining intensity is updated to reflect both the probability of Compton scattering and the conditional escape probability. This ensures that both potential outcomes—escape or continued scattering—are implicitly accounted for in the evolving photon intensity. As a result, the simulation maintains physical accuracy while achieving a significant reduction in variance.

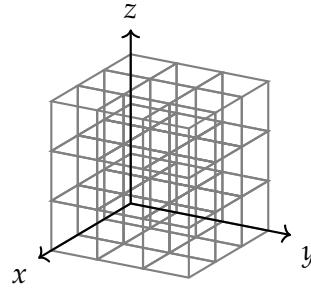
When a photon eventually escapes the phantom, it contributes to the corresponding detector pixel, weighted by its current intensity, photon energy, and escape probability. Conversely, the distance to the next interaction is sampled based on the probability of not escaping the phantom. Further implementation details are discussed in Section 10.1.

This variance reduction strategy enables the simulation to converge more rapidly toward high-resolution images with suppressed noise and well-preserved structural detail.

10.1 Algorithm overview

The algorithm begins by initializing a set of photons, each assigned an initial energy E_0 , sampled from the X-ray tube spectrum and an initial direction $\vec{\omega}_0$ sampled uniformly within a cone centered around the principal beam axis. A total of three random variables are drawn for each photon to determine its energy and direction.

Given a fixed source position A representing the location of the X-ray tube, each photon's initial origin is placed within a voxel grid composed of cubic voxels that define the geometry of the scene. Each voxel encodes an integer value identifying a specific material or tissue type. Furthermore, each photon is initialized with an intensity of $W_0 = 1$, which is iteratively updated throughout the simulation to account for attenuation and scattering processes.



Then an algorithm is applied to traverse the photon through the voxel grid until it reaches the first material which is not air. This is done with a ray traversal algorithm. Once the phantom at a point A_0 is reached, the photon transport is simulated by the physical laws described in Chapter 8.

Hereby, as in all other scatter points within the phantom, the free *free path length* t_i is sampled from the exponential distribution based on *Beer-Lambert's law* (Eq. 8.1) and the *total attenuation coefficient* $\mu(A_i, E_i)$ accordingly. A_i hereby denotes the current position and E_i the current energy. First, the *escape probability* depending on the current position A_i and the unit direction $\vec{\omega}_i$ and the energy E_i is computed. The escape probability in Equation 10.1 is the probability of the photon escaping the phantom at the current position A_i without being further scattered.

$$p(A_i, \omega_i, E_i) = \exp \left(- \int_{\overrightarrow{A_i C_i}} \mu(x, E_i) dx \right) \quad (10.1)$$

Hereby C_i denotes the exit point of the phantom in the direction of the photon $\vec{\omega}_i$. The escape probability is used to determine the intensity of the photon without the $(i+1)$ -th scatter event.

Using one random variable u_{3i+1} , the free path length is sampled. Hereby both cases are being considered:

1. Photon escapes the phantom without further scattering:

This happens with the portion matching the escape probability $p(A_i, \omega_i, E_i)$. Accordingly

$$W_{i+1}^{\text{escape}} = W_i \cdot p(A_i, \omega_i, E_i) \quad (10.2)$$

The photon contributes to the detector signal in direction $\vec{\omega}_i$ with a final intensity of $E_i \cdot W_{i+1}^{\text{escape}}$. In case $i = 0$, this intensity is accounted as primary intensity.

2. Photon does not escape the phantom and scatters:

This happens with the portion matching the inverse of the escape probability $1 - p(A_i, \omega_i, E_i)$. Accordingly, the free path length is sampled from the exponential distribution by solving Equation 10.3 for t_i :

$$\exp \left(- \int_0^{t_i} \mu(A_i + s \cdot \vec{\omega}_i, E_i) ds \right) = u_{3i+4} \quad (10.3)$$

Accordingly, the next scatter point is being computed as:

$$A_{i+1} = A_i + t_i \cdot \vec{\omega}_i \quad (10.4)$$

The photon intensity is then updated with:

$$W_{i+1} = W_i \cdot (1 - p(A_i, \omega_i, E_i)) \cdot \frac{\mu_{CS}(A_{i+1}, E_i)}{\mu(A_{i+1}, E_i)} \quad (10.5)$$

In the next step, the photon energy and opening angle after the Compton scatter event is sampled with a second random variable u_{3i+2} . A third variable is used to apply a azimuthal rotation around the direction of the photon $\vec{\omega}_i$ to sample the new direction $\vec{\omega}_{i+1}$ of the photon after the scatter event.

This process is repeated according to the maximum scatter order N .

TODO: what happens with the leftover intensity? It is forwarded without any further scattering!

10.2 Sub-Algorithms

To model the relevant physical processes in a structured manner, the main simulation is decomposed into several sub-algorithms. This section describes the purpose and implementation of each sub-algorithm in detail.

10.2.1 Photon Generation

For the photon generation step, two fundamental properties are sampled for each photon:

- *Photon energy*, sampled using a single random variable from the X-ray spectrum.
- *Photon direction*, sampled using two random variables uniformly within a cone defined by the beam axis \vec{d} and opening angle α .

Photon Energy Sampling

The photon energy is sampled from the X-ray tube spectrum, which is represented as a discrete set of energy values with corresponding fluence values. The spectrum was generated using *Spekpy* [13, 15], based on an X-ray tube configured with the parameters listed in Table 10.1.

Parameter	Value
Tube voltage	120 kV
Anode material	Tungsten
Filtration	0.4 mm Tin (Sn), 0.1 mm Copper (Cu)
Target angle	12.5°

TABLE 10.1: Parameters used to generate the X-ray spectrum with *Spekpy*.

Photon energies are sampled using inverse transform sampling. The algorithm normalizes the fluence values to obtain a probability density function (PDF), computes the corresponding cumulative distribution function (CDF) and uses a uniformly distributed random variable to select an energy according to the CDF.

Algorithm 1 Photon Energy Sampling from Spectrum

Require: Array of energies $E = [E_1, E_2, \dots, E_n]$
Require: Corresponding fluence values $\Phi = [\phi_1, \phi_2, \dots, \phi_n]$
Require: Number of samples N
Require: Random variable $u \sim \mathcal{U}(0, 1)$
Ensure: Sampled photon energies $S = [s_1, s_2, \dots, s_N]$

// Normalize fluence values:

- 1:
$$T \leftarrow \sum_{i=1}^n \phi_i$$
- 2:
$$\text{PDF}[i] \leftarrow \frac{\phi_i}{T}$$
-
- // Compute cumulative distribution function:
- 3:
$$\text{CDF}[1] \leftarrow \text{PDF}[1]$$
- 4: **for** $i = 2$ to n **do**
- 5: $\text{CDF}[i] \leftarrow \text{CDF}[i - 1] + \text{PDF}[i]$
- 6: **end for**
- // Note: $\text{PDF}[i] > 0$ in the spectrum, therefore CDF is strictly increasing.
- 7: Create interpolating function $\text{InverseCDF}(u)$ from $(\text{CDF}[i], E[i])$
- 8: **return** $\text{InverseCDF}(u)$

Photon Direction Sampling

The direction of each photon is sampled uniformly within a cone defined by the beam axis \vec{d} and opening angle α . This requires two independent random variables $u_1, u_2 \sim \mathcal{U}(0, 1)$.

Algorithm 2, adapted from [18], generates a random unit vector \vec{v} uniformly distributed within a cone of half-angle α around the direction \vec{d} :

1. **Sampling the polar angle θ :**
Draw $u_1 \sim \mathcal{U}(0, 1)$ and compute

$$\theta = \arccos(1 - u_1(1 - \cos \alpha)).$$

This corresponds to inverse transform sampling such that $\cos \theta$ is uniformly distributed on $[\cos \alpha, 1]$, resulting in uniform sampling over the spherical cap.

2. **Sampling the azimuthal angle ϕ :**
Draw $u_2 \sim \mathcal{U}(0, 1)$ and compute

$$\phi = 2\pi u_2,$$

ensuring a uniform distribution around the cone axis.

3. Constructing the local direction vector:

In a local spherical coordinate system with the cone axis aligned along the z -axis, the sampled unit vector is

$$\vec{v}_{\text{local}} = \begin{pmatrix} \sin \theta \cos \phi \\ \sin \theta \sin \phi \\ \cos \theta \end{pmatrix}.$$

4. Rotation into global coordinates:

To align the cone axis from the local z -axis to an arbitrary unit vector \vec{d} , an orthonormal basis $(\vec{u}, \vec{v}, \vec{d})$ is constructed via the Gram–Schmidt process:

- Choose a helper vector $\vec{a} = (0, 0, 1)$ if $|d_3| < 0.999$, otherwise $\vec{a} = (1, 0, 0)$.
- Compute $\vec{u} = \frac{\vec{a} \times \vec{d}}{\|\vec{a} \times \vec{d}\|}$.
- Set $\vec{v} = \vec{d} \times \vec{u}$ to complete a right-handed orthonormal basis.

Finally, rotate \vec{v}_{local} into global coordinates via the transformation

$$\vec{v}_{\text{global}} = (\vec{v}_{\text{local}})_x \vec{u} + (\vec{v}_{\text{local}})_y \vec{v} + (\vec{v}_{\text{local}})_z \vec{d}.$$

Algorithm 2 Uniform Direction Sampling Within a Cone

Require: Cone angle α

Require: Unit beam direction vector $\vec{d} = (d_1, d_2, d_3)$

Require: Random variables $u_1, u_2 \sim \mathcal{U}(0, 1)$

Ensure: Sampled direction vector \vec{v} uniformly within cone around \vec{d}

// Calculate angles according to samples

1: $\theta \leftarrow \arccos(1 - u_1(1 - \cos \alpha))$ // Polar angle

2: $\phi \leftarrow 2\pi u_2$ // Azimuthal angle

// Calculate local direction vector in spherical coordinates

3:

$$\vec{v}_{\text{local}} \leftarrow \begin{pmatrix} \sin \theta \cos \phi \\ \sin \theta \sin \phi \\ \cos \theta \end{pmatrix}$$

// Orthonormal basis construction (Gram-Schmidt)

4: **if** $|d_3| < 0.999$ **then**

5: $\vec{a} \leftarrow (0, 0, 1)$

6: **else**

7: $\vec{a} \leftarrow (1, 0, 0)$

8: **end if**

9: $\vec{u} \leftarrow \frac{\vec{a} \times \vec{d}}{\|\vec{a} \times \vec{d}\|}$ // Orthogonal vector

10: $\vec{v} \leftarrow \vec{d} \times \vec{u}$ // Complete right-handed basis

// Rotate local vector into global coordinates

11: $\vec{v}_{\text{global}} \leftarrow \vec{u}(\vec{v}_{\text{local}})_x + \vec{v}(\vec{v}_{\text{local}})_y + \vec{d}(\vec{v}_{\text{local}})_z$

12: **return** \vec{v}_{global}

10.2.2 Ray Traversal

The ray traversal algorithm is responsible for simulating the propagation of a photon through the voxel grid. It is used to forward the photon until it reaches the first chemical compound, which is not air (or is air). It is also used to simulate the photon transport within the phantom until it reaches the exit point of the phantom.

The Input values of the algorithm are:

- The voxel grid *materialGrid* representing the geometry of the scene.
- The initial position of the photon *A*.
- The (unit) direction of the photon $\vec{\omega}$.

The algorithm traverses the photon through the voxel grid, voxel by voxel while checking the material of the next voxel and iterrupts in case the next voxel is not air (or is air).

The return values of the algorithm are:

- The coordinates of the final position: *C*
- An array of all crossed voxel indices: *crossedVoxels*
- An array of all crossed voxel materials: *crossedMaterials*
- An array of al entry points of each crossed voxel: *entryPoints*
- An array of the distances traverses in each voxel: *distances*
- A boolean mask indicating whether the photon exited the grid: *exitGrid*

The sequence of the algorithm can be summarized as follows:

- I. Initialize empty arrays for *crossedVoxels*, *crossedMaterials*, *entryPoints* and *distances*.
- II. Convert the photon position *A* into voxel coordinates *pos*.
- III. Determine the indices of the next voxel *voxelIdx* based on the *pos* and the direction $\vec{\omega}$.
- IV. Set *exitGrid* to false.
- V. If the *voxelIdx* of the next voxel is out of bounds, set *exitGrid* to true and exit the algorithm.
- VI. Determine *material* of the next *voxelIdx*
- VII. If the *material* is not air, exit the algorithm.
- VIII. While *material* is air:
 1. Append *voxelIdx* to *crossedVoxels*.
 2. Append *material* to *crossedMaterials*.
 3. Append *pos* to *entryPoints*.
 4. Determine *maxDist* to the next voxel boundary.
 5. Append *maxDist* to *distances*.

6. Update pos by adding $\vec{\omega} \cdot maxDist$.
7. Determine $voxelIdx$ of the next voxel based on the updated pos .
8. If the $voxelIdx$ is out of bounds, set $exitGrid$ to true and exit the algorithm.
9. Determine $material$ of the next $voxelIdx$.

When the algorithm finishes, the final Position C is set to the last value of $position$.

To use this algorithm for traversing the photons through air until they reach the phantom and to traverse the photons through the phantom until they reach the exit point of the phantom, the algorithm is called with the boolean value $throughAir$. In case $throughAir = true$, the algorithm will traverse the photons through air until they reach the first material which is not air. In case $throughAir = false$, the algorithm will traverse the photons through the phantom until they reach the exit point of the phantom.

Therefore we define a short helper function in Algorithm 3 beforehand, which generates the specific expression depending on the boolean value of $throughAir$.

Algorithm 3 Ray Traversal Helper Function

Require: Boolean $throughAir$
Require: $material$
Ensure: Boolean value

```

1: if  $throughAir$  then
2:   return  $material = \text{air}$ 
3: else
4:   return  $material \neq \text{air}$ 
5: end if

```

Algorithm 4 implements the process of ray traversing through the voxel grid in detail.

Algorithm 4 Ray Traversal Algorithm

Require: Boolean *throughAir*
Require: Material grid $materialGrid \in \mathbb{Z}^{X \times Y \times Z}$, voxel size $s \in \mathbb{R}^+$
Require: Initial pos $A \in \mathbb{R}^3$, direction $\vec{\omega} \in \mathbb{R}^3$
Ensure: Final pos $C \in \mathbb{R}^3$
Ensure: Arrays *crossedVoxels*, *crossedMaterials*, *entryPoints*, *distances*
Ensure: Boolean *exitGrid* indicating whether the photon exited the grid

```

1:  $pos \leftarrow O/s$ 
2:  $exitGrid \leftarrow \text{false}$ 
3:  $currentVoxelIdx \leftarrow \lfloor pos \rfloor$ 
4:  $negDir \leftarrow \vec{\omega} < 0$ 
5:  $onB \leftarrow (pos = \lfloor currentVoxelIdx \rfloor)$ 
6:  $nextVoxelIdx \leftarrow currentVoxelIdx$ 
7:  $nextVoxelIdx[negDir \wedge onB] \leftarrow nextVoxelIdx[negDir \wedge onB] - 1$ 
8: if any( $nextVoxelIdx < 0$ ) or any( $nextVoxelIdx \geq \text{shape}(materialGrid)$ ) then
9:    $exitGrid \leftarrow \text{true}$ 
10:   $C \leftarrow pos$ 
11:  return  $C, crossedVoxels, crossedMaterials, entryPoints, distances, exitGrid$ 
12: end if
13:  $material \leftarrow materialGrid[nextVoxelIdx]$ 
14: while Algorithm 3(throughAir, material) do
15:    $crossedVoxels.append(nextVoxelIdx)$ 
16:    $crossedMaterials.append(material)$ 
17:    $entryPoints.append(pos)$ 
18:    $fracPos \leftarrow pos - \lfloor pos \rfloor$ 
19:    $maxDist \leftarrow [\infty, \infty, \infty]$ 
20:    $negDir \leftarrow \vec{\omega} < 0$ 
21:    $posDir \leftarrow \vec{\omega} > 0$ 
22:    $onB \leftarrow (pos = \lfloor currentVoxelIdx \rfloor)$ 
23:    $maxDist[negDir \wedge onB] \leftarrow -1/\vec{\omega}[negDir \wedge onB]$ 
24:    $maxDist[negDir \wedge \neg onB] \leftarrow -fracPos[negDir \wedge \neg onB]/\vec{\omega}[negDir \wedge \neg onB]$ 
25:    $maxDist[posDir] \leftarrow (1 - fracPos[posDir])/\vec{\omega}[posDir]$ 
26:    $distance = \min(maxDist)$ 
27:    $distances.append(maxDist)$ 
28:    $pos \leftarrow pos + \vec{\omega} \cdot distance$ 
29:    $currentVoxelIdx \leftarrow \lfloor pos \rfloor$ 
30:    $nextVoxelIdx \leftarrow currentVoxelIdx$ 
31:    $nextVoxelIdx[negDir \wedge onB] \leftarrow nextVoxelIdx[negDir \wedge onB] - 1$ 
32:   if any( $nextVoxelIdx < 0$ ) or any( $nextVoxelIdx \geq \text{shape}(materialGrid)$ ) then
33:      $exitGrid \leftarrow \text{true}$ 
34:      $C \leftarrow pos$ 
35:     return  $C, crossedVoxels, crossedMaterials, entryPoints, distances, exitGrid$ 
36:   end if
37:    $material \leftarrow materialGrid[nextVoxelIdx]$ 
38: end while
39:  $C \leftarrow pos$ 
40: return  $C, crossedVoxels, crossedMaterials, entryPoints, distances, exitGrid$ 
```

10.2.3 Forced Detection

The forced detection algorithm is responsible to apply one iteration of the forced detection process to a photon within the phantom. Hereby the algorithm accounts:

- The voxel grid *materialGrid* representing the geometry of the scene.
- The voxel grid *totalAttenuationGrid* representing the attenuation coefficients of the materials in the voxel grid *materialGrid*.
- The voxel grid *comptonAttenuationGrid* representing the Compton scattering coefficients of the materials in the voxel grid *materialGrid*.
- The voxel grid *absorptionGrid* representing the absorption coefficients of the materials in the voxel grid *materialGrid*.
- The initial position of the photon A_i .
- The (unit) direction of the photon $\vec{\omega}_i$.
- The initial energy of the photon E_i .
- The initial intensity of the photon W_i .
- A random variable $u \sim \mathcal{U}(0, 1)$ to sample the free path length.
- The voxel size s .

The algorithm then applies the ray traversal algorithm (Algorithm 4) to traverse the photon through the phantom until it reaches the exit point of the phantom C_i . Obtaining the exit point C_i , the arrays of *crossedVoxels* and *distances*, now the attenuation coefficients of the materials along the ray can be extracted from the voxel grids for the crossed voxels:

$$\begin{aligned} \text{totalAttenuationCoefficients} &= \text{totalAttenuationGrid}[\text{crossedVoxels}] \\ \text{comptonScatteringCoefficients} &= \text{comptonAttenuationGrid}[\text{crossedVoxels}] \end{aligned}$$

With the following helper algorithm (Algorithm 5), the *escapeProbability* is computed. Further the last distance index k and last interpolation factor f are returned, to solve Equation 10.3 for the free path length t_i , which can then be simply computed by:

$$t_i = \sum_{j=0}^{k-1} \text{distances}[j] + f \quad (10.6)$$

Now the next scatter point $A_{i+1} = A_i + t_i \cdot s \cdot \vec{\omega}_i$ can be determined and the new intensities together with the exit point C_i can be computed:

Algorithm 5 Partial Product Sum for Free Path Sampling in Forced Detection

Require: Arrays $distances, totalAttenuationCoefficients \in \mathbb{R}^n$, weight $u \in [0, 1]$
Ensure: Last distance index k , last interpolation factor f , $escapeProbability$

```

1:  $S_{\text{total}} \leftarrow 0$ 
2: Initialize array  $P[0 \dots n - 1]$ 
3: for  $i \leftarrow 0$  to  $n - 1$  do
4:    $P[i] \leftarrow distances[i] \cdot totalAttenuationCoefficients[i]$ 
5:    $S_{\text{total}} \leftarrow S_{\text{total}} + P[i]$ 
6: end for
7:  $T \leftarrow u \cdot S_{\text{total}}$ 
8:  $S \leftarrow 0$ 
9: for  $i \leftarrow 0$  to  $n - 1$  do
10:   if  $S + P[i] > T$  then
11:      $f \leftarrow \frac{T-S}{totalAttenuationCoefficients[i]}$ 
12:     return  $(i, f, S_{\text{total}})$ 
13:   end if
14:    $S \leftarrow S + P[i]$ 
15: end for
16: return  $(n, 1.0, S_{\text{total}})$  //  $u = 1.0$  or exact fit

```

$$voxelIdx = \lfloor A_{i+1}/s \rfloor$$

$$W_{i+1} = W_i \cdot (1 - escapeProbability) \cdot \frac{comptonAttenuationGrid[voxelIdx]}{totalAttenuationCoefficients[voxelIdx]}$$

$$W_{i+1}^{\text{escape}} = W_i \cdot escapeProbability$$

And accordingly the forced detection algorithm returns the following values:

- The new position of the photon A_{i+1} .
- The new intensity of the photon after the scatter event W_{i+1} .
- The intensity of the photon for the case of escaping the phantom W_{i+1}^{escape} .
- The exit point of the phantom C_i .
- The Compton scattering coefficient μ_{CS} at the scatter point A_{i+1} .

Algorithm 6 Forced Detection Algorithm

Require: Material grid $materialGrid \in \mathbb{Z}^{X \times Y \times Z}$
Require: Total attenuation grid $totalAttenuationGrid \in \mathbb{R}^{X \times Y \times Z}$
Require: Compton scattering grid $comptonAttenuationGrid \in \mathbb{R}^{X \times Y \times Z}$
Require: Absorption grid $absorptionGrid \in \mathbb{R}^{X \times Y \times Z}$
Require: Initial position $A_i \in \mathbb{R}^3$
Require: Unit direction $\vec{\omega}_i \in \mathbb{R}^3$
Require: Initial energy $E_i \in \mathbb{R}^+$
Require: Initial intensity $W_i \in \mathbb{R}^+$
Require: Random variable $u \sim \mathcal{U}(0, 1)$
Require: Voxel size $s \in \mathbb{R}^+$
Ensure: New position $A_{i+1} \in \mathbb{R}^3$
Ensure: New intensity $W_{i+1} \in \mathbb{R}^+$
Ensure: Escape intensity $W_{i+1}^{\text{escape}} \in \mathbb{R}^+$
Ensure: Exit point $C_i \in \mathbb{R}^3$
Ensure: Compton Attenuation coefficient μ_{CS}

```

1:  $C_i, crossedVoxels, crossedMaterials, entryPoints, distances, exitGrid \leftarrow$ 
   Algorithm 4( $throughAir = true, materialGrid, s, A_i$ )
2: if  $exitGrid$  then
3:   return  $(C_i, W_i, 0, C_i)$  // Photon escaped the phantom
4: end if
5:  $totalAttenuationCoefficients \leftarrow totalAttenuationGrid[crossedVoxels]$ 
6:  $comptonScatteringCoefficients \leftarrow comptonAttenuationGrid[crossedVoxels]$ 
7:  $absorptionCoefficients \leftarrow absorptionGrid[crossedVoxels]$ 
8:  $k, f, escapeProbability \leftarrow$  Algorithm 5( $distances, totalAttenuationCoefficients, u$ )
9:  $escapeProbability \leftarrow \frac{escapeProbability}{\sum_{j=0}^{k-1} distances[j] \cdot totalAttenuationCoefficients[j]}$ 
10:  $t_i \leftarrow \sum_{j=0}^{k-1} distances[j] + f$ 
11:  $A_{i+1} \leftarrow A_i + t_i \cdot s \cdot \vec{\omega}_i$ 
12:  $voxelIdx \leftarrow \lfloor A_{i+1} / s \rfloor$ 
13:  $\mu_{\text{CS}} \leftarrow comptonScatteringCoefficients[voxelIdx]$ 
14:  $W_{i+1} \leftarrow W_i \cdot (1 - escapeProbability) \cdot \frac{\mu_{\text{CS}}}{totalAttenuationCoefficients[voxelIdx]}$ 
15:  $W_{i+1}^{\text{escape}} \leftarrow W_i \cdot escapeProbability$ 
16: return  $(A_{i+1}, W_{i+1}, W_{i+1}^{\text{escape}}, C_i, \mu_{\text{CS}})$ 

```

10.2.4 Photon Exit Point Determination

To determine the exit point of the photon in the world after exiting the phantom, a more efficient algorithm than the ray traversal algorithm from Section 10.2.2 is used can be applied. This algorithm yields the exit point of the phantom which is used to determin the detector pixel the photon contributes to.

The algorithm is initialized with the following parameters:

- The voxel grid shape (N_x, N_y, N_z) of the $materialGrid$ representing the geometry of the scene.
- The initial position of the photon C^{phantom} .
- The (unit) direction of the photon $\vec{\omega}$.

- The voxel size s .

The algorithm then computes the point where the photon exits the voxel grid efficiently and returns the exit point C^{grid} and the according Coordinates $C^{\text{gridCoords}}$. The algorithm is implemented as follows:

Algorithm 7 Compute Exit Point of Ray from Voxel Grid

Require: Ray origin $C^{\text{phantom}} \in \mathbb{R}^3$, direction $\vec{\omega} \in \mathbb{R}^3$
Require: Grid shape (N_x, N_y, N_z) , voxel size $s \in \mathbb{R}^+$
Ensure: Exit point C^{grid} , $C^{\text{gridCoords}}$

```

1:  $C^{\text{phantomCoords}} \leftarrow C^{\text{phantom}} / s$                                 // Convert to voxel coordinates
2:  $x_{\min} \leftarrow 0, x_{\max} \leftarrow N_x$ 
3:  $y_{\min} \leftarrow 0, y_{\max} \leftarrow N_y$ 
4:  $z_{\min} \leftarrow 0, z_{\max} \leftarrow N_z$ 
5: for axis in {x, y, z} do
6:    $o \leftarrow C^{\text{phantomCoords}}_{\text{axis}}$ 
7:    $d \leftarrow \vec{\omega}_{\text{axis}}$ 
8:   if  $d > 0$  then
9:      $t^{(\text{axis})} \leftarrow \frac{x_{\max}-o}{d}$ 
10:    else if  $d < 0$  then
11:       $t^{(\text{axis})} \leftarrow \frac{x_{\min}-o}{d}$ 
12:    else if  $d = 0$  then
13:       $t^{(\text{axis})} \leftarrow -\infty$ 
14:    end if
15:  end for
16:   $t_{\text{exit}} \leftarrow \min(t^{(x)}, t^{(y)}, t^{(z)})$                                 // Exit time
17:   $C^{\text{gridCoords}} \leftarrow C^{\text{phantomCoords}} + t_{\text{exit}} \cdot \vec{\omega}$ 
18:   $C^{\text{grid}} \leftarrow C^{\text{gridCoords}} \cdot s \cdot \vec{\omega}$ 
19: return  $C^{\text{grid}}, C^{\text{gridCoords}}$ 

```

10.2.5 Compton Scattering

The Compton scattering algorithm is responsible for simulating the physical process for the event of Compton scattering. Based on the Klein-Nishina formula (Equation 9.1) and the rejection sampling method from Section 9.3.2, the algorithm samples the scatter angle and computes the new photon energy and direction after the scattering event.

The algorithm initializes with the following parameters:

- The initial energy of the photon E_i .
- The (unit) direction of the photon $\vec{\omega}_i$.
- Two random variables $u_1, u_2 \sim \mathcal{U}(0, 1)$ to sample the new photon energy and direction.

The algorithm applies the Klein-Nishina procedure to compute the scatter angle and therefore utilizes the electron rest energy E_{rest} and follow this procedure:

1. Initialize

$$k = E_i/mc^2, \quad \zeta(k) = \frac{2(2k^2 + 2k + 1)}{(2k + 1)^3}, \quad b(k) = \frac{1 + \frac{\zeta}{2}}{1 - \frac{\zeta}{2}}, \quad a(k) = 2(b - 1).$$

2. Rejection Sampling loop:

2.1. Generate random number $r \sim \mathcal{U}(0, 1)$.

2.2. Calculate

$$\begin{aligned} x &= b - (b + 1) \left(\frac{\zeta}{2} \right)^r \\ f(k, x) &= \frac{1 + x^2 + \frac{k^2(1-x)^2}{1+k(1-x)}}{(1 + k(1 - x))^2} \\ h(k, x) &= \frac{a}{b - x} \end{aligned}$$

2.3. If $u_1 \leq \frac{f(k, x)}{h(k, x)}$ then: $\theta = \arccos(x)$

3. Calculate photon energy with Equation 9.5 and direction using the azimuthal angle ϕ from Equation 9.4.

Given the scatter angle θ , the new photon energy E_{i+1} is computed as in [10]:

$$E_{i+1} = \frac{E_i}{1 + k(1 - \cos \theta)}. \quad (10.7)$$

With the random variable u_2 an azimuthal angle ϕ is sampled to determine the new direction.

$$\phi = 2\pi u_2 \quad (10.8)$$

Assuming the vector \vec{u} is an orthogonal unit vector $\vec{u} \perp \vec{\omega}_i$ and $\vec{v} = \vec{u} \times \vec{\omega}_i$, then the new direction $\vec{\omega}_{i+1}$ is as follows:

$$\vec{\omega}_{i+1} = \sin(\theta)\cos(\phi) \cdot u + \sin(\theta)\sin(\phi) \cdot v + \cos(\theta) \cdot \vec{\omega}_i \quad (10.9)$$

This leads to the following return values of the algorithm:

- The new photon energy E_{i+1} after the Compton scattering event.
- The new (unit) direction $\vec{\omega}_{i+1}$ of the photon after the Compton scattering event.

In the pseudoalgorithm is more explicitly describing this process in detail in Algorithm 8 by implementing omitted tweaks and details.

Algorithm 8 Compton Scattering Algorithm

Require: Initial photon energy E_i , direction $\vec{\omega}_i$
Require: Electron rest energy E_{rest}
Require: Random variables $u_1, u_2 \sim \mathcal{U}(0, 1)$
Ensure: New photon energy E_{i+1} , direction $\vec{\omega}_{i+1}$

// Scatter angle sampling

- 1: $k \leftarrow E_i / E_{\text{rest}}$
- 2: $\varepsilon_0 \leftarrow 1/(2k + 1)$
- 3: **repeat**
- 4: Generate $r \sim \mathcal{U}(0, 1)$
- 5: **if** $r < 0.5$ **then**
- 6: $\varepsilon \leftarrow \varepsilon_0 + (1 - \varepsilon_0) \cdot 2r$
- 7: **else**
- 8: $\varepsilon \leftarrow \varepsilon_0 + (1 - \varepsilon_0) \cdot 2(1 - r)$
- 9: **end if**
- 10: $\cos \theta \leftarrow 1 + \frac{1}{k} (1 - \frac{1}{\varepsilon})$
- 11: **if** $|\cos \theta| \leq 1$ **then**
- 12: $\sin^2 \theta \leftarrow 1 - \cos^2 \theta$
- 13: **if** $u_1 \leq \frac{1}{2} (\varepsilon + \frac{1}{\varepsilon} - \sin^2 \theta)$ **then**
- 14: $\theta \leftarrow \arccos(\cos \theta)$
- 15: **break**
- 16: **end if**
- 17: **end if**
- 18: **until** accepted

// New photon energy calculation

- 19: $E_{i+1} \leftarrow \frac{E_i}{1+k(1-\cos \theta)}$

// Orthonormal basis construction

- 20: $\phi \leftarrow 2\pi u_2$
- 21: **if** $|(\vec{\omega}_i)_z| < 0.999$ **then**
- 22: $\vec{a} \leftarrow (0, 0, 1)$
- 23: **else**
- 24: $\vec{a} \leftarrow (1, 0, 0)$
- 25: **end if**
- 26: $\vec{u} \leftarrow \frac{\vec{a} \times \vec{\omega}_i}{\|\vec{a} \times \vec{\omega}_i\|}$
- 27: $\vec{v} \leftarrow \vec{\omega}_i \times \vec{u}$

// New direction calculation

- 28: $\vec{\omega}_{i+1} \leftarrow \sin \theta \cos \phi \cdot \vec{u} + \sin \theta \sin \phi \cdot \vec{v} + \cos \theta \cdot \vec{\omega}_i$
- 29: **return** $E_{i+1}, \vec{\omega}_{i+1}$

10.3 Algorithm Composition

The main algorithm is composing the sub-algorithms from Section 10.2.

The algorithm is initialized with the following parameters:

- The voxel grid $materialGrid$ representing the geometry of the scene.
- The voxel grid $totalAttenuationGrid$ representing the total attenuation coefficients of the materials in the voxel grid $materialGrid$.

- The voxel grid $\text{comptonAttenuationGrid}$ representing the Compton scattering coefficients of the materials in the voxel grid materialGrid .
- The voxel grid absorptionGrid representing the absorption coefficients of the materials in the voxel grid materialGrid .
- The source position S of the X-ray tube.
- The number of scatter events N to simulate.
- The opening angle α of the X-ray beam.
- The voxel size s of the voxel grid.
- A sequence of random variables $u \sim \mathcal{U}(0, 1)^{3(N+1)}$ to sample the photon energies, directions and free path lengths.

Hereby, the photon generation algorithm is called to generate the photon energies and directions.

Then the photon is initialized with the an initial energy E_0 , an initial direction $\vec{\omega}_0$ and an initial intensity $W_0 = 1$ by applying Algorithm 1 and Algorithm 2. Based on the initial position $A_0 = S$, direction $\vec{\omega}$, the voxel grid and voxel size s , the ray traversal algorithm (Algorithm 4) is applied to traverse the photon through the voxel grid until it reaches the first material which is not air. The exit point of the phantom is denoted as C_0 .

The loop over the scatter events:

Later N iterations of the forced detection algorithm (Algorithm 6) are applied together with one random variable to simulate the photon transport through the phantom. In each iteration, the photon position and intensity are updated. With Algorithm 7, the exit point C_i^{grid} of the photon is computed and captured together with the relevant intensity $E_i \cdot W_{i+1}^{\text{escape}}$.

Then the Compton scatter event is simulated, taking into account the photon energy E_i , the direction $\vec{\omega}_i$ and the Compton scattering coefficient μ_{CS} at A_{i+1} . Together with two random variables, the new photon energy E_{i+1} with an according direction $\vec{\omega}_{i+1}$ is sampled.

Algorithm 9 Main Simulation Algorithm Composition

Require: Material grid $materialGrid$
Require: Total attenuation grid $totalAttenuationGrid$
Require: Compton attenuation grid $comptonAttenuationGrid$
Require: Absorption grid $absorptionGrid$
Require: Source position S
Require: Number of scatter events N
Require: Beam opening angle α
Require: Voxel size s
Require: Random variables $u \sim \mathcal{U}(0, 1)^{3(N+1)}$

Ensure: Detector contributions (primary and scatter signals)

// Photon generation

- 1: $E_0 \leftarrow$ Algorithm 1(spectrum, u_1)
- 2: $\vec{\omega}_0 \leftarrow$ Algorithm 2(α , beam axis, u_2, u_3)
- 3: $W_0 \leftarrow 1$
 // Traverse photon through air to phantom
- 4: $A_0, \dots \leftarrow$ Algorithm 4(throughAir=true, $materialGrid, s, S, \vec{\omega}_0$)
- 5: **for** $i = 0$ to $N - 1$ **do**
- // Forced detection step
- 6: $A_{i+1}, W_{i+1}, W_i^{\text{escape}}, C_i^{\text{phantom}}, \mu_{\text{CS}} \leftarrow$ Algorithm 6($materialGrid, totalAttenuationGrid, comptonAttenuationGrid, absorptionGrid, A_i, \vec{\omega}_i, E_i, W_i, u_{3i+1}, s$)
 // Record escaped photon contribution
- 7: **if** $W_i^{\text{escape}} > 0$ **then**
- 8: $C_i^{\text{world}}, \dots \leftarrow$ Algorithm 7($C_i^{\text{phantom}}, \vec{\omega}_i, \text{grid shape}, s$)
- 9: Add $E \cdot W_i^{\text{escape}}$ to detector pixel at C_i^{world}
- 10: **end if**
 // Compton scatter step
- 11: $E_{i+1}, \vec{\omega}_{i+1} \leftarrow$ Algorithm 8($E, \vec{\omega}_i, \mu_{\text{CS}}, u_{3i+2}, u_{3i+3}$)
- 12: **end for**
 // Handle leftover intensity
- 13: $C_N^{\text{world}}, \dots \leftarrow$ Algorithm 7($A_N, \vec{\omega}_N, \text{grid shape}, s$)
- 14: Add $E \cdot W$ to detector pixel intensity at C_N^{world}

Appendix A

Frequently Asked Questions

A.1 How do I change the colors of links?

The color of links can be changed to your liking using:

```
\hypersetup{urlcolor=red}, or  
\hypersetup{citecolor=green}, or  
\hypersetup{allcolor=blue}.
```

If you want to completely hide the links, you can use:

```
\hypersetup{allcolors=.}, or even better:  
\hypersetup{hidelinks}.
```

If you want to have obvious links in the PDF but not the printed text, use:

```
\hypersetup{colorlinks=false}.
```


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