

Some derivation of the axon problem

March 6, 2024

This is about Nestor's diffusion advection problem. We have a finite axon of length L , with some influx j_0 at the soma ($x = 0$). We also know that the dendrite is closed, so no net flux will be possible at $x = L$. Our chemical species $u : (x, t) \mapsto u(x, t)$ undergoes diffusion, advection, and decay. If I am not mistaken, the problem can be formulated as the initial boundary problem:

$$\begin{cases} u_t = D u_{xx} - v u_x - k u \\ [D u_x - v u]_{(0,t)} = j_0 \\ [D u_x - v u]_{(L,t)} = 0 \\ u(x, 0) = \varphi(x) \end{cases} \quad (1)$$

on the domain $x \in D = [0, L]$ and $t \in [0, +\infty]$. This is very classic stuff and I sure hope there exists a solution of some sort. Let us start first with the steady-state solution

Steady-state In order to find the steady-state solution $s(x) = u(x, t \rightarrow \infty)$, we impose $u_t = 0$. That leaves us with the Sturm-Liouville problem

$$\begin{cases} D s'' - v s' - k s = 0 \\ D s'(0) - v s(0) = j_0 \\ D s'(L) - v s(L) = 0 \end{cases} \quad (2)$$

which admits the general solution

$$s(x) = a e^{\lambda_1 x} + b e^{\lambda_2 x}, \quad a, b \in \mathbb{R} \quad (3)$$

Since $D, v, k > 0$ (things diffuse, move rightwards and degrade in time), also the eigenvalues λ_1 and λ_2 will be real. We can now proceed and find the exact values of a and b by direct substitution in the boundary conditions. Before going ahead, we rewrite the expressions for $s(x)$ and $s'(x)$ in the more convenient form

$$s(x) = e^{\frac{v}{2D}x} \left[a e^{\frac{\Delta}{2D}x} + b e^{-\frac{\Delta}{2D}x} \right] \quad (4)$$

$$s'(x) = \frac{v}{2D} s(x) + \frac{\Delta}{2D} e^{\frac{v}{2D}x} \left[a e^{\frac{\Delta}{2D}x} - b e^{-\frac{\Delta}{2D}x} \right] \quad (5)$$

where $\Delta = \sqrt{v^2 + 4Dk}$. By inserting these expressions in the boundary conditions, we obtain

$$\begin{cases} b = \frac{\Delta - v}{\Delta + v} a - \frac{2j_0}{\Delta - v} \\ b = \exp\left\{\frac{\Delta}{D}L\right\} \frac{\Delta - v}{\Delta + v} a \end{cases}$$

This yields the constant values

$$a = \frac{2j_0}{(\Delta - v)(1 - e^{\frac{\Delta}{D}L})}, \quad b = \frac{2j_0 e^{\frac{\Delta}{D}L}}{(\Delta + v)(1 - e^{\frac{\Delta}{D}L})} \quad (6)$$

and the final expression for $s(x)$ as

$$s(x) = \frac{2j_0}{1 - e^{\frac{\Delta}{D}L}} e^{\frac{v}{2D}x} \left[\frac{e^{\frac{\Delta}{2D}x}}{\Delta - v} + \frac{e^{-\frac{\Delta}{2D}(x-2L)}}{\Delta + v} \right] \quad (7)$$

Interestingly, the flux j_0 acts as a scaling parameter.

Full solution I don't really know how to do this for now but soon. Below are some failed attempts. We start with the substitution

$$u(x, t) = w(x, t) e^{\eta t - \mu x} \quad (8)$$

where

$$\eta = k + \frac{v^2}{4D}, \quad \mu = \frac{v}{2D} \quad (9)$$

by direct substitution we get

$$\begin{cases} w_t = D w_{xx} \\ D w_x(0, t) - (D\mu + v)w(0, t) = j_0 e^{-\eta t} \\ D w_x(L, t) - (D\mu + v)w(L, t) = 0 \\ w(x, 0) = \varphi(x) e^{\mu x - \eta t} \end{cases} \quad (10)$$

We have, therefore, a initial boundary problem for the diffusion of $w(x, t)$, with Robin boundary conditions that depend on time. To try and tackle this problem, we will follow the recipe in [1] and integrating it with [2].

We start by separating our solution $w(x, t)$ in two parts, one denominated $S(x, t)$, representing the steady-state of (10), and the other $U(x, t)$ representing the transient perturbation of $w(x, t)$ from its asymptotic steady state

$$w(x, t) = S(x, t) + U(x, t) \quad (11)$$

By construction, $S(x, t)$ has to respect the boundary value problem (10) for $t \rightarrow \infty$. Following, without any pretensions of originality, the method illustrated in [1], we posit for $S(x, t)$ the following “quasi-steady-state” form

$$S(x, t) = A(t) \left(1 - \frac{x}{l}\right) + B(t) \frac{x}{l} \quad (12)$$

The two unknown functions can be promptly found by substituting $S(x, t)$ in the boundary conditions

$$\begin{cases} \frac{D}{l} (B - A) - \frac{3}{2} v A = j_0 e^{-\eta t} \\ \frac{D}{l} (B - A) - \frac{3}{2} v B = 0 \end{cases} \quad (13)$$

obtaining

$$A(t) = \left[\frac{2D}{3lv} - 1 \right] \frac{2j_0}{3v} e^{-\eta t}, \quad B(t) = \frac{4Dj_0}{9lv^2} e^{-\eta t} \quad (14)$$

and therefore

$$S(x, t) = \left[\frac{x}{l} + \frac{2D}{3lv} - 1 \right] \frac{2j_0}{3v} e^{-\eta t} \quad (15)$$

After finding $S(x, t)$, we can correctly define the problem for the transient solution $U(x, t)$, subject to Dirichlet boundary conditions:

$$\begin{cases} U_t = DU_{xx} \\ D U_x(0, t) - \frac{3}{2} U(0, t) = 0 \\ D U_x(L, t) - \frac{3}{2} U(L, t) = 0 \\ U(x, 0) = \varphi(x) e^{\mu x - \eta t} - S(x, 0) =: \tilde{\varphi}(x) \end{cases} \quad (16)$$

We can proceed with separation of variables, and posit that $U(x, t) = X(x)T(t)$. We first solve the spatial boundary value problem

$$\begin{cases} X'' + \lambda X = 0 \\ D X'(0) - \frac{3}{2} X(0) = 0 \\ D X'(L) - \frac{3}{2} X(L) = 0 \end{cases} \quad (17)$$

The only non-trivial solution of this BVP can be written, assuming $\lambda > 0$, as

$$X(x) = a \cos \sqrt{\lambda} x + b \sin \sqrt{\lambda} x \quad (18)$$

and obtain the values for λ and the relationship between a and b by substituting in the boundary conditions

$$\begin{cases} Db\sqrt{\lambda} - \frac{3}{2}va = 0 \\ D\sqrt{\lambda} \left[b \cos \sqrt{\lambda}l - a \sin \sqrt{\lambda}l \right] - \frac{3}{2}v \left[a \cos \sqrt{\lambda}l + b \sin \sqrt{\lambda}l \right] = 0 \end{cases} \quad (19)$$

$$\begin{cases} Db\sqrt{\lambda} - \frac{3}{2}va = 0 \\ \cos \sqrt{\lambda}l \left[Db\sqrt{\lambda} - \frac{3}{2}va \right] + \sin \sqrt{\lambda}l \left[Da\sqrt{\lambda} + \frac{3}{2}vb \right] = 0 \end{cases} \quad (20)$$

$$\begin{cases} a = \frac{2Db}{3v}\sqrt{\lambda} \\ \frac{b}{6v} \sin \sqrt{\lambda}l [4D^2\lambda + 9v^2] = 0 \end{cases} \quad (21)$$

Since $\lambda > 0$, its only possible values are $\lambda = \frac{n^2\pi^2}{l^2}$, $n \in \mathbb{N}$, yielding the final form for $X(x)$ is

$$X(x) = C_n \left[\frac{2Dn\pi}{3vl} \cos \frac{n\pi}{l}x + \sin \frac{n\pi}{l}x \right] \quad (22)$$

References

1. Drábek, P. & Holubová, G. *Elements of Partial Differential Equations* in (2007). <https://api.semanticscholar.org/CorpusID:118449101>.
2. Salsa, S. *Partial Differential Equations in Action. UNITEXT*. <https://api.semanticscholar.org/CorpusID:124525275> (2022).