

Identification of Power BJT Operating Stages Based on Experimental Excess Charge Estimation

1st Jan Miklas

*Department of Power Electrical and Electronic Engineering
Faculty of Electrical Engineering and Communication
Brno University of Technology
Brno, Czech Republic
jan.miklas@vutbr.cz*

2nd Petr Prochazka

*Department of Power Electrical and Electronic Engineering
Faculty of Electrical Engineering and Communication
Brno University of Technology
Brno, Czech Republic
prochazkap@feec.vutbr.cz*

Abstract—The paper proposes and demonstrates an experimental way of estimating the amount of stored charge of excess minority carriers within power BJT base and collector. This consequently allows a detailed identification of transistor operating stage.

A brief device operation analysis is provided as a clear support of the measured characteristics.

The method is based on determining the steady-state stored charge at various operating conditions by integration of negative transient base current during turn-off event which deflates the stored charge.

An ultimate objective of these and future experiments is an accurate interpretation and modelling of various device stages during IGBT switching process. Most of the observed phenomena are common among a power BJT and IGBT's intrinsic BJT. As IGBT doesn't provide access to the internal base current, it is advantageous to measure and interpret the relations between stored charge and switching waveforms of power BJT first and further generalize the observations to IGBT measurements.

Index Terms—IGBT, power BJT, bipolar semiconductor devices saturation, minority carriers concentration, excess charge storage, transistor switching measurement,

I. INTRODUCTION

Devices utilizing bipolar current transport used as a fast power switches such as IGBTs are not going to be fully replaced by unipolar counterparts in the near future thanks to their high current density capability at optimum manufacturing price. Notorious downside of bipolar devices is charge storage within forward-biased junctions, limiting the switching speed.

High current density is allowed by two phenomena unique for bipolar devices physics (i.e. including the minority carriers action) - the saturation and conductivity modulation in lightly doped region. Both are related to minority carriers charge storage and dependent on the base and collector current levels. As a result, the good on-state performance can be achieved at the cost of degrading the switching performance and vice-versa. In addition, the charge storage undesired side-effects are clearly reflected in the switching waveforms [1].

Saturation mode is essential for steady on-state operation as described by a static current - voltage characteristic. It

is defined by excess minority carriers concentration within the base and collector during forward bias of both BJT junctions. While there is only a little practical significance in determining the amount of excess charge stored in the collector for static saturation characteristic analysis of already-manufactured device, estimating the boundaries of individual operational modes is highly beneficial for accurate identification and analysis of particular stages during switching action. However, determining the amount and geometric profile of excess charge stored in the collector N-drift region is not trivially possible by means of static saturation characteristic measurement. On the other hand, the excess charge - continuously recombined with majority carriers - is supplied by the base current. Likewise, negative base current can be utilized to deflate the excess charge. So, one only need the access to the base current transient switching waveform and to eliminate the recombination impact on measured waveform to gain a quantitative information about the steady-state minority carriers stored within transistor right before the switch-off event. Unfortunately IGBT doesn't provide access to the internal base current, but most of the phenomena in question are common among a power BJT and intrinsic BJT within IGBT. It is thus beneficial to use an "archaic" power BJT as a qualitative testing reference with the advantage of direct control of base current parameters, waveform shape and its accurate measurement.

The gained observations can be further almost-directly applied to the IGBT on-state and switching waveforms interpretation [2].

Measurements presented in next sections were performed to experimentally validate the possibility to interpret the power BJT operational stages by estimating the amount of base and collector stored charge of excess minority carriers based on base current transient waveform measurements. In other words, to validate the following presumptions:

- Recombination of minority carriers is at least one order of magnitude slower than times required for external deflation of the excess charge (i.e. the turn-off time),
- Basic function of power BJT is given primarily by the desired "theoretical" phenomena prior to the undesired side-effects.

II. A BRIEF THEORY

A. Structure of Power BJT

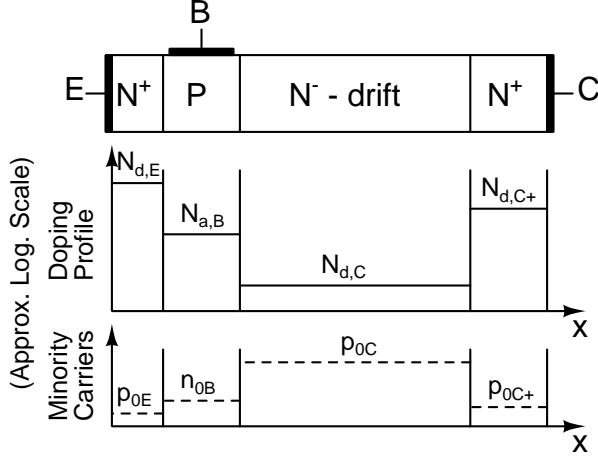


Fig. 1. Structure of power BJT. Dimensions and concentrations scale are illustrative i.e. not proportional to real scales.

The basic one-dimensional structure of power BJT is depicted on Fig. 1. It approximates the cross-section of a real device right beneath the emitter contact. Unlike a low voltage BJT, the structure involves a very lightly doped so-called drift region in the collector to support high collector voltage by the depletion region while reverse biased. The current flow through drift region creates additional voltage drop unless being flooded with injected minority carriers exceeding the equilibrium concentration multiple times (conductivity modulation [2]). The minority carriers concentration in equilibrium within doped semiconductor with intrinsic carrier concentration n_i and doping concentration N_D or N_A is given by:

$$n_{p0} = \frac{n_i^2}{N_D}, \quad p_{n0} = \frac{n_i^2}{N_A} \quad (1)$$

The higher doping, the lower minority carriers concentration; as depicted on Fig. 1.

B. Operation of power BJT

Bipolar transistor operates in 3 basic modes, corresponding to 3 combinations of biasing 2 junctions within transistor, as summarized in Tab. I. This can be considered as opera-

Mode	B-E bias	B-C bias
Cutoff	R	R
Active	F	R
Saturation	F	F

TABLE I
THREE BASIC OPERATION MODES OF BJT.

tion mode definitions. Corresponding illustrative (not properly scaled) minority carrier concentrations and potentials within device structure are depicted on Fig. 2. The minority carrier concentrations on the junction (depletion region) boundaries - let us denote it $p_n(0)$ and $n_p(0)$ - adhere the following principles:

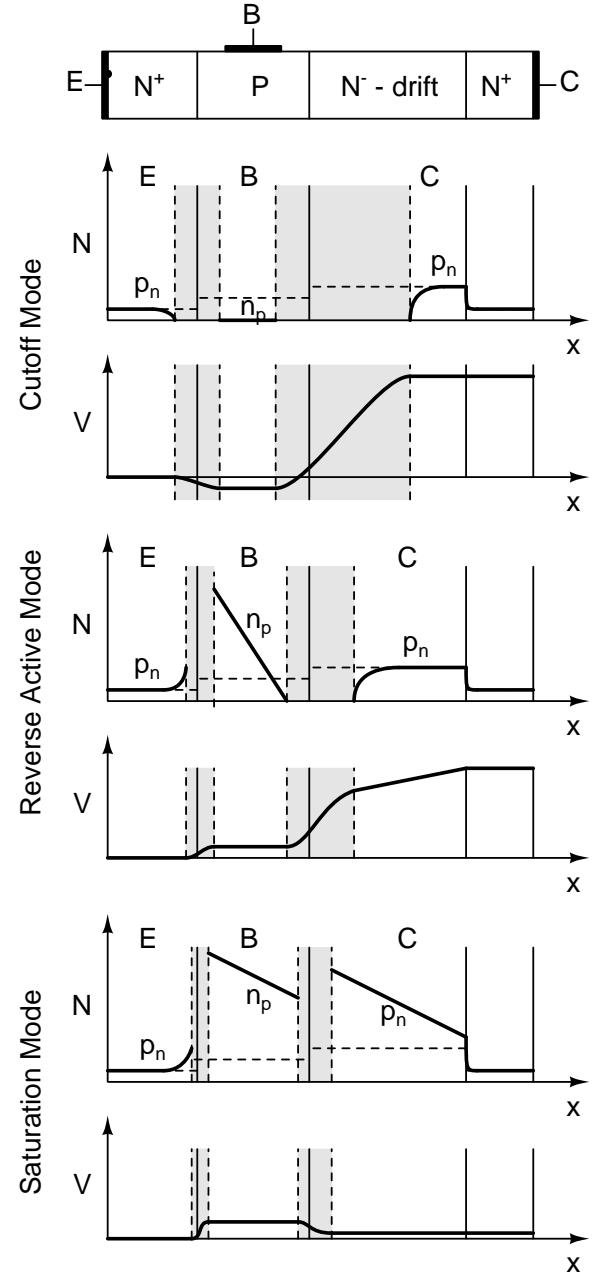


Fig. 2. Illustrative power BJT minority carrier and potential profiles for 3 basic operation modes.

- reverse biased junction:

$$p_n(0) = 0, \quad n_p(0) = 0 \quad (2)$$

- forward biased junction: “The Law of Junction” [3]:

$$p_p(0)n_p(0) = n_i^2 e^{qV_{PN}/kT} \quad (3)$$

$$n_p(0) = n_{p0} \cdot e^{qV_{PN}/kT} \quad (4)$$

$$p_n(0) = p_{n0} \cdot e^{qV_{PN}/kT} \quad (5)$$

Shaded area represents the depletion region.

Carriers concentration gradient signifies the diffusion current while the potential (voltage) gradient signifies the drift current. Drift current (and thus the voltage gradient - i.e. also the capability to support voltage) is typical for depletion regions and lightly doped region without conductivity modulation. In fact, the drift current principally signifies the current flow through the resistance in general.

In active mode, the B-E junction is forward biased which results in minority carriers injection from emitter into the base region (the “Law of Junction”) and B-C junction is reverse biased resulting in zero electron concentration at the junction. The current flows through the collector despite the reverse bias, due to minority carriers diffusion at the depletion region boundary (carriers gradient within base) and drift through the depletion region (potential gradient within depletion region). The minority carriers in collector does not exceed the equilibrium concentration.

Collector voltage (with respect to the emitter) can be varied by varying the depletion width. In other words, varying the resistivity of the current path.

In saturation mode, both junctions are forward biased. N-drift region is flooded with minority carriers injected from the base, exceeding an equilibrium concentration, typically exceeding even the majority carriers equilibrium concentration. Consequently, the majority carriers concentration is increased as well to maintain charge neutrality [2]. As a result, low doping causing additional voltage drop while proportional to the collector current is effectively increased and voltage drop disappears. This is called the *conductivity modulation*.

Thus, there is no significant drift current, the collector current is maintained by diffusion. As stated before, drift signifies the resistive material, while transistor in saturation exhibits the low-resistivity (on-state) performance. Carrier concentration gradient is equal in base and collector region as guaranteed by the same (diffusion) current I_C and depicted on Fig. 2.

C. Active and Saturation Mode at Various Drive and Load Conditions

Figure 3 illustrates the minority carrier concentrations within base and N-drift region at variable drive (I_B) and load (I_C) conditions. It is important to understand the effect of each parameter as both constitute the amount of stored charge and alternations of measured charge must correspond to these parameters alternations if the used theory is correct. Operator “++” applied on any parameter on the figure designates the increasing trend of that parameter in the direction of an arrow.

I_B level determines the forward voltage across B-E junction. This in turn determines the injection level of minority carriers into the base as stands in the “Law of Junction” (4). So increasing I_B increases the minority electron concentration in the base at B-E junction. This rule applies on both active and saturation mode, as indicated on the picture.

As for I_C , there is not a common rule for both modes. In the active mode with B-C junction reverse biased, the *transistor effect* is valid so I_C is proportional to I_B through the current gain β . In saturation, however, both junctions are forward

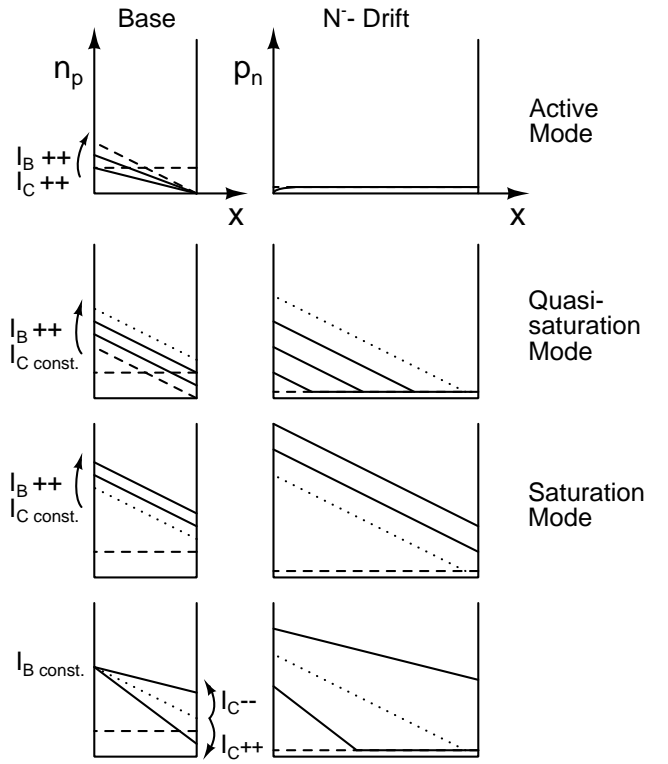


Fig. 3. Minority carrier concentration dependence on I_B and I_C variations. The geometry and concentration scales are illustrative, not proportional to real device values. Also the base and collector diagrams use the different scale to depict the same concentration gradient.

biased, transistor effect is not valid any more, significant amount of charge storage within collector is introduced and current transport is determined by diffusion. That means the value of I_C is given by the gradient of carrier concentration (both electrons and holes to maintain the charge neutrality) is constant in base and collector region. The rule for explanation of I_C variations impact on the stored charge thus is to alternate the concentration gradient (see the bottom picture on Fig. 3). The higher current, the higher concentration gradient; while the “left boundary” concentration in the base is firmly determined by I_B . As a result, increasing I_C causes the decrease of stored charge and vice versa, as opposed to possible intuitive expectation.

Increasing I_C at constant I_B can lead to quasi-saturation region (see the bottom picture on the figure) and additional voltage drop across the unmodulated drift region.

Increasing I_B at constant I_C is possible (see the second and third picture on the figure), this only magnifies the depth of saturation and reduces the value of current gain β .

In addition, retaining in saturation mode, $Q = f(I_C)$ should exhibit the linear dependence because all carrier profiles remain linear.

$Q = f(I_B)$ is also expected to be linear. However, the amount of stored charge rapidly increases after transition from active to the saturation mode due to charge storage within collector. $Q = f(I_B)$ thus exhibits a double-slope

linear dependency with apparent boundary between active and (quasi)saturation region. The boundary between quasi-saturation and saturation region is not apparent by means of stored charge volume; one would need to observe the collector voltage additionally. It is good to note that charge stored in collector is much higher than that stored in the base (depicted charge profiles are of different scale). As a result, any alternations i.e. storing and deflating the charge in active mode is a way much faster than alternations during saturation.

D. Charge Control Relation

Qualitative considerations in section II-C regarding the stored charge within base and collector can be supported by simple yet accurate enough *charge control* relation [4], [5] as follows:

$$\frac{dQ}{dt} = i_B(t) - \frac{Q(t)}{\tau} \quad (6)$$

where Q denotes the total stored charge, i_B is the base current supplying (or deflating) the charge and τ stands for the recombination time constant - i.e. Q/τ represents the “recombination current”, proportional to amount of stored charge [3].

All excess minority carriers charge is supplied by the base current and it is not necessary to partition i_B into individual parts (i_{pC} etc.) as it is not needed to identify individual portions of total stored charge.

In steady state, I_B is just sufficient to supply the recombined charge (proportional to total stored charge volume):

$$I_B = \frac{Q}{\tau} \quad (7)$$

which matches all 3 modes in Tab. I. $I_B = 0$ in cutoff mode so stored charge equals zero; $I_B > 0$ which signifies the forward bias of B-E junction in active and saturation mode further defined by collector bias voltage.

Negative base current is not possible in steady state.

III. METHODOLOGY

A. Negative Base Drive Transient Current Relation with Stored Charge

As stated before, negative base current is not possible in steady state. However, it is common to apply a negative transient base current during turn-off event to speed up the charge deflation from base and collector. Time integral of this current provides an information about the steady state charge Q_{ss} stored within transistor right before the event.

Integration of the “discharging” analogy to (6) yields:

$$\int i_{B-}(t) dt = Q_{ss} - \frac{1}{\tau} \int Q(t) dt \quad (8)$$

That is, the obtained steady-state charge projection $Q_{sst,m}$ (“m” stands for “measured”) is not necessarily equal to an actual stored charge Q_{ss} . The difference between those two is the portion of charge that already got recombined during deflation time.

$$Q_{ss,m} = Q_{ss} - Q_{recomb} \quad (9)$$

One of the objectives of proposed tests is to find such conditions that Q_{recomb} is certainly negligible compared to total charge Q_{ss} , prior to other measurements.

B. Test Summary

There are two quantities, the actual steady-state Q_{ss} stored in the collector and its measured projection $Q_{ss,m}$. They are both dependent on varying test conditions like I_{B+} , I_{B-} , I_C , either directly (Q_{ss} , according to (7) and Fig. 3) or indirectly ($Q_{ss,m}$) through the deflation time allowing slow recombination (8).

The following experimental tests were proposed based on the analysis in section II-C. To be noted, all of the tests are based on steady-state stored charge quantification at varying conditions.

- Q vs. I_{B+} at constant I_{B-} , I_C .

Linearly growing dependency with double slope is expected as described in the analysis of Fig. 3. Linearity is predicted also by (7)

The purpose of the measurement is verification of provided analysis and verification of assumption that the amount of charge stored within base is negligible compared to charge stored within collector during saturation.

- Q vs. I_{B-} at constant I_{B+} , I_C .

Constant dependency is expected, because constant I_B yields a constant Q_{ss} according to (7). However, the charge-deflation time for constant Q_{ss} increases with decreasing I_{B-} , so the recombined charge portion becomes non-negligible and the resulting characteristic $Q_{ss,m} = f(I_{B-})$ should tend to zero at low I_{B-} levels.

The purpose of this measurement is to find such conditions where recombination effect is negligible and $Q_{ss,m} = Q_{ss}$ is assured. In other words, to assure the measurement (reflected charge) credibility.

- Q vs. I_C at constant I_{B+} , I_{B-} .

Linearly decreasing dependency is expected in saturation as described in the analysis of Fig. 3. The charge should remain approximately constant after transition from saturation to the active region, with slightly decreasing trend due to increasing collector voltage and consequent base width modulation (this mode is forced by a constant I_{B+} and increasing I_C , so current gain β needs to decrease)

C. Test Bench

All testing was performed on test bench built for fast switching measurement of power semiconductor devices with minimized parasitic influences as described in [6] and [7] in detail.

The device under test - power BJT BUV48A - was driven by discrete base driver through a variable base resistor to provide a constant I_{B+} during whole “charging” on-state period and constant I_{B-} during most of the “discharging” turn-off event. Separated current path and base resistors were ensured for I_{B+} and I_{B-} via diodes.

IV. MEASUREMENT AND INTERPRETATION OF RESULTS

The integration time and thus the deflated charge was divided into two parts:

$$Q = Q_1 + Q_2 \quad (10)$$

Q_1 represents the excess charge stored in collector until transition from saturation to an active region. The boundary is determined by the decrease of collector current below 90 % of its steady value I_C .

Q_2 represents mainly the charge stored within base. It is expected to be constant at given I_B according to (7).

Q_1 and Q_2 can be considered with somewhat daring approximation as collector and base stored charge respectively.

A. Q vs. I_{B+}

The measured characteristic on Fig. 5 exhibits the **linear dependency with double slope**, exactly as expected by the test proposal in section III-B.

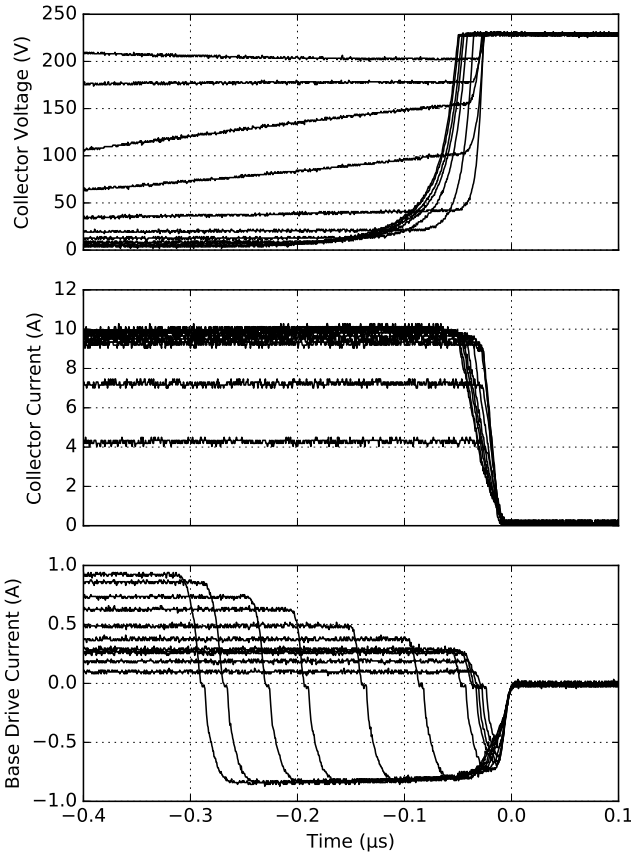


Fig. 4. Transient waveforms at varying I_{B+} .

The great slope of characteristic in saturation region corresponds to charge storage within collector as illustrated in Fig. 3.

Linear dependency is valid in the active region as well, expressing the injection level dependency on I_B . However, the charge stored within base is small compared to collector, hence

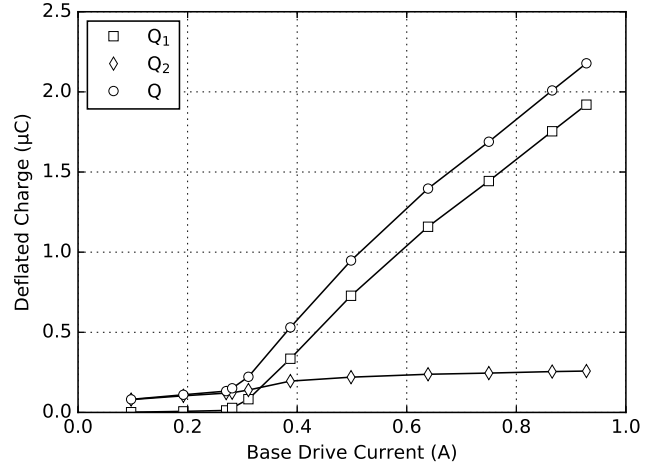


Fig. 5. $Q = f(I_{B+})$ characteristic, test results. The boundary between active and saturation region is clearly apparent. Conditions: $I_{B-} = 0.8$ A, $I_C = 10$ A.

the small slope of the characteristic. Charge stored within the base ranges from zero to approx. $0.2 \mu C$.

The boundary between active and saturation region at given test conditions is clearly apparent. It can be verified looking at the base and collector current transient waveforms on Fig. 4. Collector steady current is proportional to the base current in the active mode according to $I_C = \beta \cdot I_B$. $\beta = 4A/0.1A = 40$ according to the picture.

B. Q vs I_{B-}

The measured characteristic on Fig. 7 fulfills the expectation of **constant value** at reasonably high I_{B-} levels (> 0.3 A) and decreasing trend as I_{B-} decreases.

An important observation arises when comparing Q_1 with Q_2 decreasing trend with increasing deflation time (decreasing I_{B-}). Q_1 remains constant while Q_1 drops below $I_{B-} < 0.3$ A. The explanation is the nature of recombination - it is proportional to total stored charge - equation (6). As Q_2 is deflated after Q_1 deflation, the remaining amount of charge (equal to Q_1) is really small, but deflating current level I_{B-} remains constant and thus the recombined charge proportion to deflated charge becomes minimal and recombination effect does not distort the measured charge any more.

The practical outcome of this result is the knowledge of usable value of $I_{B-} > 0.3$ A for further measurements at given conditions.

C. Q vs I_C

The measured characteristic on Fig. 9 exhibits the **linear dependency** with negative slope in saturation mode, i.e. until approaching the active region as expected by the test proposal in section III-B.

V. CONCLUSION AND FUTURE WORK

The proposed method summarized in section III-B was verified by multiple measurements listed in section IV. All of

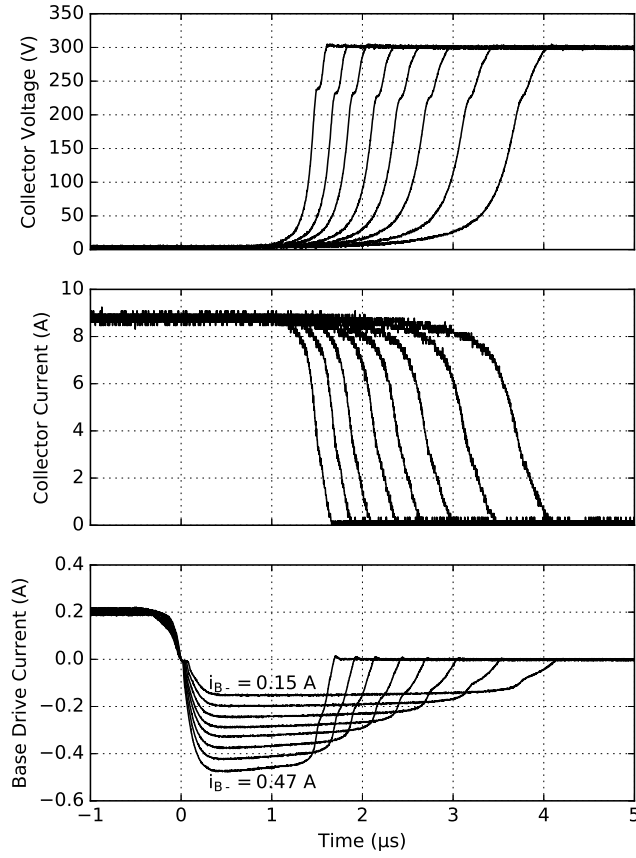


Fig. 6. Transient waveforms at varying I_{B-} .

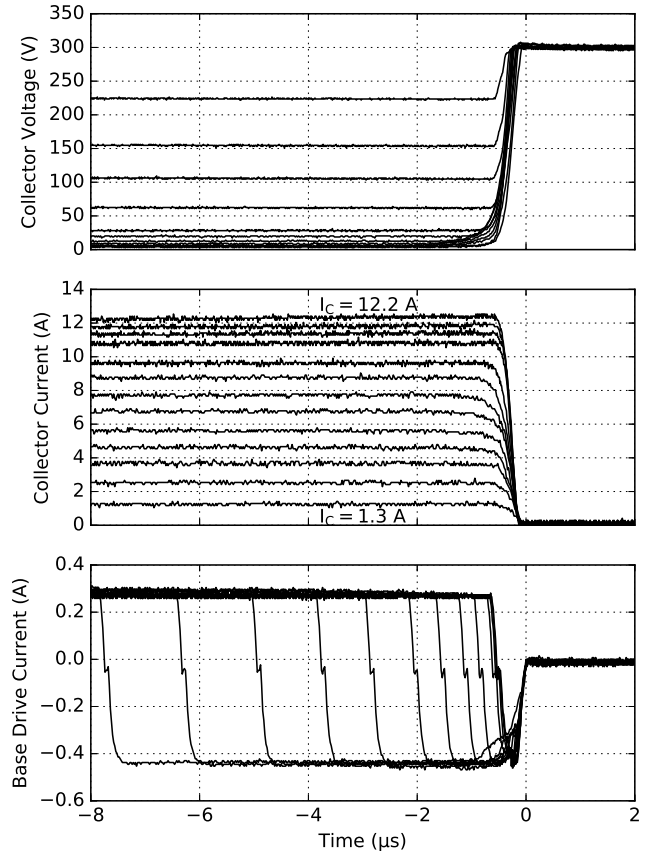


Fig. 8. Transient waveforms at varying I_{C-} .

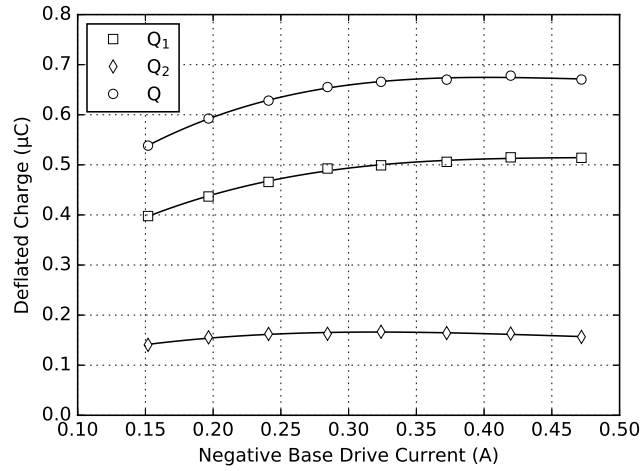


Fig. 7. $Q = f(I_{B-})$ characteristic, test results. Conditions: $I_{B+} = 0.2$ A, $I_C = 9$ A.

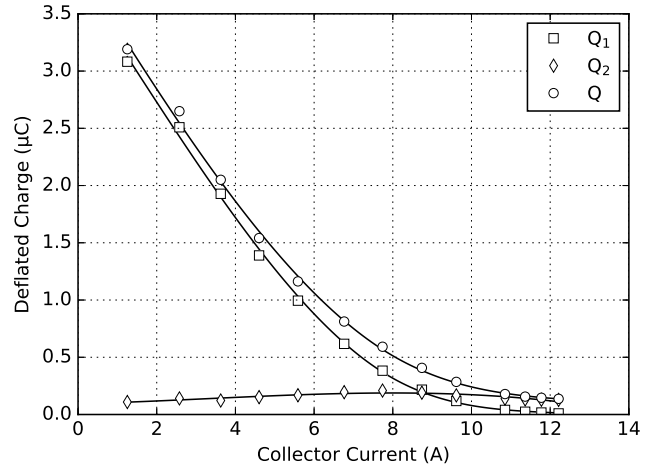


Fig. 9. $Q = f(I_C)$ characteristic, test results. Transition from saturation to the active region is apparent. Conditions: $I_{B+} = 0.3$ A, $I_{B-} = 0.4$ A.

the expected dependences outlined in the analysis in section II-C were clearly verified. That validates the possibility of using the excess charge deflated via base as full-featured indicator of any operation stage of power BJT.

The method will be further used for identification of var-

ious stages and impacts on switching waveforms and further applied to IGBT switching analysis.

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REFERENCES

- [1] C. Hu and M. J. Model, "A model of power transistor turn-off dynamics," 1980 IEEE Power Electronics Specialists Conference, Atlanta, Georgia, USA, 1980, pp. 91-96.
- [2] B. J. Baliga, Fundamentals of Power Semiconductor Devices, New York: Springer. 2008.
- [3] R. F. Pierret, Semiconductor Device Fundamentals, Addison-Wesley Publishing Company, 1996, ISBN 0-201-54393-1
- [4] H. K. Gummel, H. K., "A Charge Control Relation for Bipolar Transistors", Bell Syst. Tech. J., vol. 49, 1970
- [5] S. M. Sze and M. K. Lee, Semiconductor Devices: Physics and Technology, 3rd ed., New York: John Wiley & Sons, Inc. 2012.
- [6] J. Miklas, Power Switching Transistors, Brno: Brno University of Technology, Faculty of Electrical Engineering and Communication. 2016. Head of Diploma Thesis doc. Dr. Ing. Miroslav Patocka
- [7] P. Prochazka, J. Miklas, I. Pazdera, M. Patocka, J. Knobloch, R. Cipin, "Measurement of Power Transistors Dynamic Parameters", Mechatronics 2017, pp.571-577, ISBN 978-3-319-65959-6