

Basic Background / motivation for Bifurcation analysis

Consider ODEs $\dot{x}(t) = f(x(t), p)$ $x(t) \in \mathbb{R}^n$, $f: \mathbb{R}^n \times \mathbb{R}^m \rightarrow \mathbb{R}^n$

Initial-value problems can be safely solved numerically

$x(0) = x_0$, $\varphi^T(x_0)$:= solution of ODE at time T when starting from x_0 at time 0,

$x_0 \rightarrow \varphi^T(x_0)$ is diffeomorphism of \mathbb{R}^n , is approximated by

x_N where $x_{k+h} = x_k + h f(x_k, p)$, $\|x_N - \varphi^T(x_0)\| = G(N^{-1})$
 $(N \gg 1, h = \frac{T}{N})$

Trouble arises when studying long-time behaviour:

$$\|x_N - \varphi^T(x_0)\| \leq \underbrace{\alpha_p(L, T)}_{\text{if } T \text{ is large}} \frac{T}{N}, L \sim \|\partial f\|$$

if T is large, this is uncontrollably large

Observation often $\mathbf{C}^T(\mathbf{x}_*)$ converges to invariant sets.

Simple invariant sets can be computed directly:

- **Equilibria**: solve $\mathbf{0} = \mathbf{f}(\mathbf{x}, \mathbf{p})$ nonlinear system
→ $\mathbf{x}_{\text{eq}}(\mathbf{p})$
- If $\partial_{\mathbf{x}} \mathbf{f}(\mathbf{x}_{\text{eq}}, \mathbf{p}) \in \mathbb{R}^{n \times n}$ has no eigenvalues on imaginary axis
behaviour of $y = \mathbf{x} - \mathbf{x}_{\text{eq}}(\mathbf{p})$ for small y is governed
by $\dot{y} = \partial_{\mathbf{x}} \mathbf{f}(\mathbf{x}_{\text{eq}}, \mathbf{p}) y$ (linear ODE)
 - all eigenvalues λ have $\operatorname{Re} \lambda < 0 \Rightarrow \mathbf{x}_{\text{eq}}$ stable
 - some $\operatorname{Re} \lambda > 0 \Rightarrow \mathbf{x}_{\text{eq}}$ unstable
- Eigenvalues on imaginary axis \Rightarrow "local bifurcation"
 - $\lambda = 0 \Rightarrow$ Saddle-node } \hookrightarrow soc Kurzschluss
 - $\lambda = \pm i\omega \Rightarrow$ Hopf bifurcation } Applied Bifurcation Theory

Periodic orbits are trajectories $\varphi^{(\cdot)}(x)$ such that $\varphi^T(x) = x$ for some $T > 0$, but not all $T > 0$

They are found by solving periodic boundary-value problems

$$x'(t) = T f(x(t), p) \quad \leftarrow \text{rescaled time}$$

$$\theta = x(1) - x(0)$$

$$(\text{e.g.}) \quad \theta = m^T x(0) - c \quad (\text{Poincaré / return section})$$

where $f(x(0), p)$, m are linearly independent.

- periodic orbits branch off near equilibria at Hopf bifurcation
- their stability is determined by eigenvalues of $\partial_x \varphi^T(x)$